EEM 323

ELECTROMAGNETIC WAVE THEORY II

NORMAL INCIDENCE

AT PLANAR BOUNDARY

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Önemli not: Ders notlarındaki şekillerin hazırlanmasında internet ortamından faydalanılmıştır. Özellikle belirtilmeyen tüm şekil, tablo, eşitlik ve denklemler vb. "D. K, Fundamentals of Engineering Electromagnetics, Addison-Wesley Inc." ile "D. K, Field and Wave Electromagnetics, Mc-Graw Hill Inc." kitabından taranarak elde edilmiştir. Alıntıların kaynağına kolay ulaşılabilmesi maksadıyla numarası ve altyazıları da gösterilmektedir.

DERS KİTABI

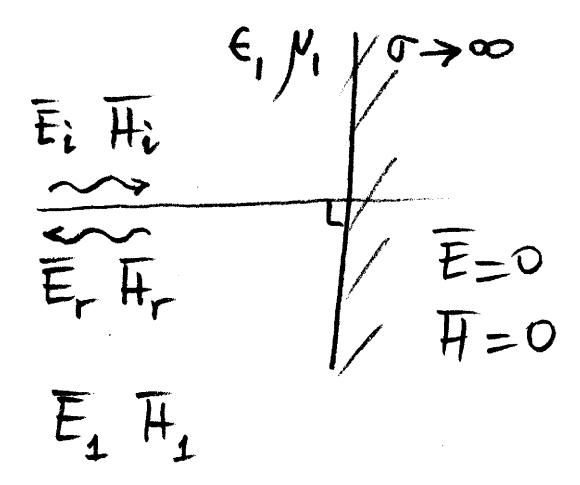
[1] David Keun Cheng, Fundamentals of Engineering Electromagnetics, Addison-Wesley Publishing, Inc., 1993. veya David Keun Cheng, Çeviri: Adnan Köksal, Birsen Saka, Mühendislik Elektromanyetiğinin Temelleri – Fundamentals of Engineering Electromagnetics, Palme Yayınları.

KAYNAK / YARDIMCI KİTAPLAR:

- [2] David Keun Cheng, Field and Wave Electromagnetics, Addison-Wesley Publishing, Inc. veya David Keun Cheng, Çeviri: Mithat İdemen, Elektromanyetik Alan Teorisinin Temelleri Field and Wave Electromagnetics, Literatür Yayıncılık.
- [3] Stanley V. Marshall, Richard E. DuBroff, Gabriel G. Skitek, *Electromagnetic Concepts and Applications*, Dördüncü Basım, Prentice Hall International, Inc., 1996.
- [4] Joseph A. Edminister, Elektromanyetik, 2. Baskıdan çeviri, Çevirenler: M. Timur Aydemir, E. Afacan, K. C. Nakipoğlu, Schaum's Outlines, McGraw Hill Inc., Nobel Yayın Dağıtım, Ankara, 2000.

NORMAL INCIDENCE AT PLANAR BOUNDARY OF

Dielectric – Perfect conductor boundary:



Problem:

Given the incident E field, calculate all the other fields

* NORMAL

INCIDENCE

AT A COND.

BOUNDARY

$$E_i(2) = \hat{q}_x \quad E_{io} \quad e^{-i\beta_1 2}$$

$$H_i(a) = \hat{a}_y \frac{E_{io}}{\eta_i} e^{-j\beta_i a}$$

$$\overline{H}$$
 \widehat{a}_{ni} \widehat{q}_{ni} $\widehat{\sigma}_{2} = \infty$

reflected waves? $E_{f}(2) = \hat{q}_{x} E_{f} e^{i\beta_{x}^{2}}$

$$E_{i}(2) = E_{i}(2) + E_{r}(2) = \hat{q}_{x}(E_{io} e^{-ip_{i}2} + E_{ro} e^{-ip_{i}2})$$

Conde Surface: Boundary Condition

 $E_{r}(z=0) = a_{x}(E_{i0} + E_{r0}) = E_{2}(0) = 0$

=) \(\bar{\xi_{ro}} = -\bar{\xi_{io}}\)

Toloma

Ero (7) = -ax Eio = +ip, 2

1,2:

MW(BONUS)

Plot | E | VS. 7

in time.

(multiplot)

 $E_{1}(2) = \hat{q}_{x} E_{io}(e^{-j\beta_{i}2} - e^{-2j\sin\beta_{i}2}$

TOTAL FIELD .

 $E_{i}(t) = -\frac{\partial}{\partial x}$ il E_{io} singt

Standing Wave

 \overline{EX} Y-polarized uniform plane wave $(\overline{E_i} \ \overline{H_i})$ frequency 100 MHz., in air, normal incidence to x=0.

X=0; perfectly conducting plane boundary. $|\overline{E_i}| = 6 \, (mV/m)$

(a) Phasor and Instantaneous expressions Ei Hi

(b) Fr Hr reflected wave.

(c) \(\varE_1 \) \(\varH_1 \) total work

(d) determine The location recovert to the conducting plane where $\bar{E}_1 = 0$.

$$\omega = 2\pi f = 2\pi. 10^{8} (rad/s.)$$

$$\beta_{1} = k = \frac{\omega}{c} = \frac{2\pi. 10^{8}}{3.10^{8}} = \frac{2\pi}{3} (rad/s.)$$

$$\gamma_{1} = \gamma_{0} = \sqrt{\frac{\omega}{\epsilon_{0}}} = 120\pi. (2)$$

$$\gamma_{2} = \gamma_{0} = \sqrt{\frac{\omega}{\epsilon_{0}}} = 120\pi. (2)$$

(a)
$$\{\overline{t}_{i}(x) = \widehat{a}_{i} = 0.70^{3} \cdot e^{-\frac{2\pi}{3}x} \}$$
 $\{v/m\}$
 $\overline{H}_{i}(x) = \frac{1}{2} \hat{a}_{i} \times \overline{E}_{i}(x) = \hat{a}_{i} = \frac{10^{-4}}{2\pi} \cdot e^{-\frac{2\pi}{3}x} \}$ $\{A/m\}$

Phaier

Instructaneous Expressions (Arth Begerless)

$$\bar{E}_{i}(x,t) = \Re\{\bar{E}_{i}(x) e^{j\omega t}\}
= \hat{a}_{i}(x,t) = \Re\{\bar$$

$$\vec{\mathcal{H}}_{i}(x;t) = \hat{a}_{j} \frac{10^{-4}}{2\pi} \cdot \cos\left(2\pi \cdot 10^{8} \cdot t - \frac{2\pi}{3} \times\right) (A/m)$$

(b)
$$\bar{t}_{i}(x) = -\hat{a}_{i} 6.10^{-3} \cdot e^{i\frac{2\pi}{3}x}$$
 (V/m)

$$\overline{H}_r(x) = \frac{1}{\eta_r} (-\hat{a}_x) \times \overline{E}_r(x) = \hat{a}_y \frac{10^{-4}}{2\pi} e^{-\frac{1}{3}x}$$
 (A/m)

Instantaneous Expressions

26 blumper

$$\overline{E}(x,t) = Re \left\{ \overline{E}_r(x) e^{i\omega t} \right\} = -\hat{a}_y 6.10^3 \cdot \cos \left(2\pi.10^8 t + \frac{2\pi}{3} x \right)$$

$$\frac{1}{2\pi} (x,t) = \frac{1}{2\pi} \left\{ \frac{1}{2\pi} (x) e^{i\omega t} \right\} = \frac{1}{2\pi} \left(\frac{10^{-4}}{2\pi} \cos(2\pi t, 10^8 t + \frac{2\pi}{3} x) \right)$$

$$\overline{E}_{1}(x) = \overline{E}_{i}(x) + \overline{E}_{i}(x) = +\hat{a}_{y} 6.10^{3} \left(e^{-j\frac{2\pi}{3}x} - e^{-j\frac{2\pi}{3}x}\right)$$

$$\overline{E}_{1}(x) = -\hat{a}_{y} j12.10^{-3}. sin\left(\frac{2\pi}{3}x\right) \left(v_{h}\right)$$

$$\overline{H}_{4}(x) = \overline{H}_{i}(x) + \overline{H}_{r}(x) = \hat{q}_{1} \frac{10^{4}}{\pi} \cos(\frac{2\pi}{3}x)$$
 (A/m)

Instantaneous Expressions

$$\bar{E}_{i}(x,t) = 2e \left\{ \bar{E}_{i}(x) e^{i\omega t} = \hat{a}_{i} 12.10^{3} \sin(\frac{2\pi}{3}x) . \sin(2\pi.10^{8}t) \right\}$$

$$\overline{\mathcal{H}}_{i}(x,t) = 2e \left\{ \overline{H}_{i}(x)e^{iwt} \right\} = \hat{a}_{i} \frac{10^{-4}}{11} \cos\left(\frac{2\pi}{3}x\right) \cdot \cos(2\pi \cdot 10^{-6}t)$$

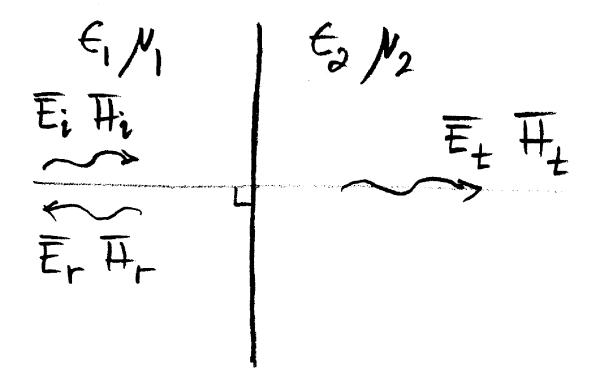
$$\chi = -\frac{\chi_1}{2} = -\frac{\pi}{\beta_1} = -\frac{3}{5}$$
 (m)

ÖDEV:

Durağan dalga çözümünü Manyetik alan için hesaplayınız.

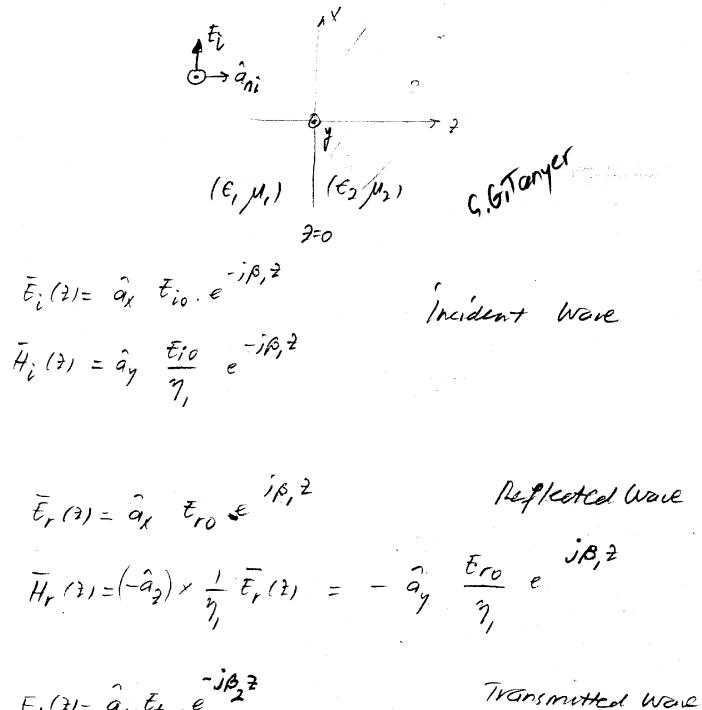
NORMAL INCIDENCE AT PLANAR BOUNDARY OF

Dielectric – Conductor boundary:



Problem:

Given the incident E field, calculate all the other fields



$$E_{t}(t) = \hat{a}_{x} E_{to} e^{-j\beta_{2}^{2}}$$

$$H_{t}(t) = \hat{a}_{x} \times \frac{1}{\eta_{2}} E_{t}(t) = \hat{a}_{y} \frac{E_{to}}{\eta_{2}} e^{-j\beta_{2}^{2}}$$

$$H_{t}(t) = \hat{a}_{x} \times \frac{1}{\eta_{2}} E_{t}(t) = \hat{a}_{y} \frac{E_{to}}{\eta_{2}} e^{-j\beta_{2}^{2}}$$

QUESTION:

- Given only the incident Electric field, what are the fields in both regions?
- How many equations do we need?

Boundary Condition

$$\exists \quad \exists_{io} + \exists_{ro} = \exists_{to}$$

$$=) \frac{1}{\eta_{i}} (E_{io} - E_{ro}) = \frac{E_{to}}{\eta_{2}}$$

$$E_{ro} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} E_{io}$$

$$E_{to} = \frac{2\eta_2}{\eta_2 + \eta_1} E_{io}$$

$$\Gamma = \frac{\mathcal{E}_{ro}}{\mathcal{E}_{io}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}$$

$$T = \frac{\mathcal{E}_{to}}{\mathcal{E}_{io}} = \frac{2\eta_2}{\eta_2 + \eta_1}$$

If Medium 2 is perfect conductor, $\Rightarrow \Gamma = -1$, T = 0 $\Rightarrow E_{ro} = -E_{io}$, $E_{to} = 0$

=> Total reflection and Standing wave

If medium a is not a perfect conductor,

=> Portial reflection, partial transmission

Note: Given \overline{E}_{i} \Rightarrow Find \overline{E}_{i} \overline{E}_{t} \Rightarrow find \overline{F}_{i} \overline{F}_{t} \overline{F}_{t} \Rightarrow find \overline{E}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i} \overline{F}_{i}

$$\overline{H}_{i}(z) = \frac{\widehat{a}_{z} \times \overline{E}_{io}}{\eta_{i}} e^{-j\beta_{i}z}$$

$$\overline{H}_{r}(z) = \frac{(-\hat{a}_{z} \times \overline{E}_{io})}{\eta_{1}} (\Gamma) e^{j\beta_{1}z}$$

$$\overline{H}_{t}(z) = \frac{(\hat{a}_{2} \times E_{io})}{\eta_{2}} (z) e^{-i\beta_{0}z}$$

REVIEW:

Assume
$$|\Gamma| \neq 0$$
 yam' $0 < |\Gamma| < 1$

$$|T| \neq 0$$
 and $0 < |T| < 1$

$$|T| \neq 0$$
 and $0 < |T| < 1$

$$\frac{E_{1}(z)}{E_{1}(z)} = \frac{E_{1}(z)}{E_{1}(z)} + \frac{E_{1}(z)}{E_{1}(z)} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} + \frac{e^{-i\beta_{1}z}}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} + \frac{e^{-i\beta_{1}z}}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})} = \frac{a_{1}}{a_{1}} \frac{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z})}{E_{10}(e^{-i\beta_{1}z} + \Gamma e^{-i\beta_{1}z}$$

$$\overline{E}_{1}(2) = \hat{q}_{\chi} \overline{E}_{io} \left[\overline{Z} \cdot e^{-i\beta,2} + \Gamma(j2\sin\beta,2) \right]$$

Standing Wave Ratio

$$S = \frac{|E|_{max}}{|E|_{min}} = \frac{|+|\Gamma|}{|-|\Gamma|}$$

$$|\Gamma| = \frac{S-1}{S+1}$$

$$\Gamma > 0 \left(\eta_1 > \eta_1\right)$$

$$\begin{array}{c|c}
\hline
0 & \overleftarrow{k} \\
\hline
2=0
\end{array}$$

QUESTION: If the reflection coefficient is positive and +1, what would be the standing wave function?

$$\frac{2}{max} = -\frac{n\pi}{\beta_1} = -\frac{n\pi}{2}$$

minimum occurs when
$$|\bar{t}_{i}(t)| = \bar{t}_{io}(Lr)$$

$$2_{min} = -\frac{(2n+1)\pi}{2\beta} = -\frac{(2n+1)\pi}{4}$$

QUESTION: If the reflection coefficient is negative and -1, what would be the standing wave function?

$$\max \left\{ |\bar{t}_{i}(t)| \right\} = \bar{t}_{io} (1-\Gamma)$$

 $\min \left\{ |\bar{t}_{i}(t)| \right\} = \bar{t}_{io} (1+\Gamma)$

locations of min and max are Interchanged.

$$\overline{P}_{av} = \frac{1}{2} \operatorname{Re} \left\{ \overline{E} \times \overline{H}^{*} \right\}^{2} \qquad \overline{E}_{i} \alpha_{i} = \hat{a}_{i} \operatorname{Eio} e^{-j\beta_{i} t} \left(H \operatorname{Fe}^{j\beta_{i} t} \right) \\
\overline{H}_{i}(t) = \hat{a}_{i} \operatorname{Eio} e^{-j\beta_{i} t} \left(H \operatorname{Fe}^{j\beta_{i} t} \right) \\
\overline{P}_{av}_{i} = \hat{a}_{i} \frac{\operatorname{Eio}}{2\eta_{i}} \operatorname{Re} \left\{ \left(1 - \operatorname{Fe}^{-j2\beta_{i} t} \right) \left(1 - \operatorname{Fe}^{-j2\beta_{i} t} \right) \right\} \\
= \hat{a}_{i} \frac{\operatorname{Eio}}{2\eta_{i}} \operatorname{Re} \left\{ \left(1 - \operatorname{Fe}^{-j2\beta_{i} t} \right) \left(1 - \operatorname{Fe}^{-j2\beta_{i} t} \right) \right\}$$

$$=\hat{q}_{2}\frac{\xi_{io}^{2}}{2\eta_{i}}Re^{\left\{ (1-\Gamma^{2})+\frac{1}{2}2\Gamma.\sin 2\rho_{i}+\frac{3}{2}\right\}}$$

$$\overline{P_{av,1}} = \hat{q}_{3} \frac{\overline{E_{io}}}{2\eta_{i}} (1-\Gamma^{2})$$

$$\Gamma : real Since OD lossless.$$

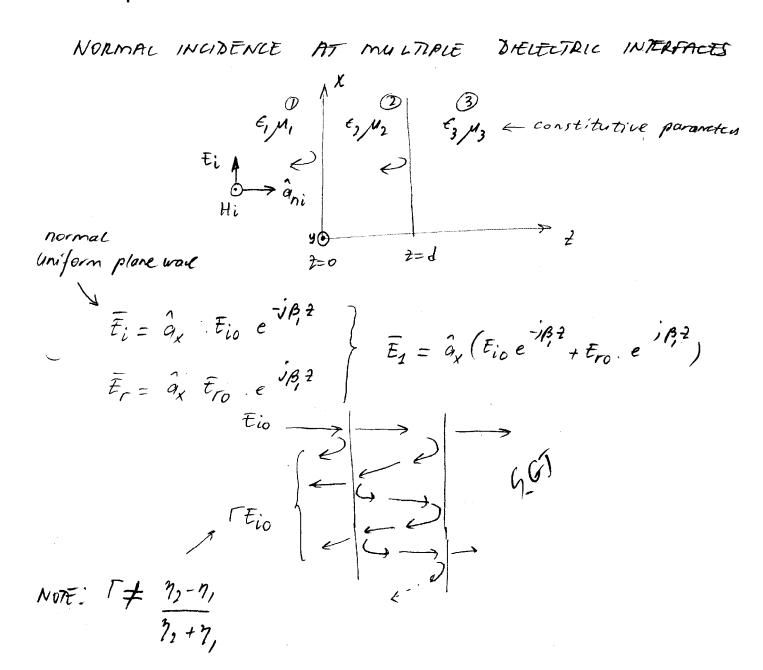
Similarly by
$$\frac{2}{9av_{12}} = \frac{3}{9} + \frac{tio}{2\eta_{2}} = \frac{7}{2}$$

$$\overline{Q}_{av_{j1}} = \overline{Q}_{av_{j2}}$$
 $\Rightarrow \sqrt{-\Gamma^2 - \frac{\eta_j}{\eta_2}}$ $\Rightarrow \sqrt{-\Gamma^2 - \frac{\eta_j}{\eta_2}}$ Show itis

QUESTION:

NORMAL INCIDENCE AT PLANAR BOUNDARY OF MULTIPLE DIELECTRIC INTERFACES

- How would you approach the problem?
- How many unknowns are there and how many do you need?
- How many equations are there?
- If there are enough equations, how would you solve the problem?



QUESTION: How many methods/approaches can you think of to solve this problem?

1. Calculate all recursive reflections and transmissions, one by one.

Que: How many steps do you think you need to solve?

Que: If it is infinite, can you define a method for

convergence? Can you define the constraints for

the right time to stop?

2. Write down the total E and H fields and apply the boundary conditions, and see what you can do..

Let us skip the first method (you think of it later).

Now let us solve the fields in all three media.

Let us number the fields according to the media number; 1, 2 and 3.

For media (1):

$$\overline{E}_{1} = \widehat{q}_{x}(E_{io} e^{-i\beta_{i}^{2}} + \overline{t}_{ro} e^{-i\beta_{i}^{2}})$$

$$\overline{H}_{i} = \widehat{q}_{y} \frac{1}{\eta}(E_{io} e^{-i\beta_{i}^{2}} - \overline{t}_{ro} e^{-i\beta_{i}^{2}})$$

For media (2):

Medium 2: Forward and Back word Ware

$$\overline{E}_{3} = \hat{q}_{x} \left(\overline{E}_{3}^{\dagger} e^{-i\beta_{3}^{2}} + \overline{E}_{3}^{-} e^{-i\beta_{3}^{2}} \right)$$

For media (3):

Medium 3: Only a forward wave in (+2) direction $\overline{E}_3 = \hat{a}_x \, \overline{E}_3^{\dagger} \, e^{-j\beta_3 Z}$ $\overline{H}_3 = \hat{a}_y \, \frac{\overline{E}_3^{\dagger}}{\gamma_2} \, e^{-j\beta_3 Z}$

Question:

Define the known and the unknown variables (fields).

Unknowns Fro, Et, E, , Est

Boundary Conditions

$$Q_{2=0}$$
 $E_{1}(0) = E_{2}(0)$
 $H_{1}(0) = H_{2}(0)$

(a)
$$z=d$$
 $E_{3}(d) = E_{3}(d)$
 $H_{3}(d) = H_{3}(d)$