

Note: inverting a Laplace transform is often very hard. We usually use a table of transforms and hope to find the relevant function / transform pair.

Inverting Laplace transforms

Example Using the facts that $\mathcal{L}[1] = 1/s$ and $\mathcal{L}[e^{at}] = \frac{1}{s-a}$,

where a is a constant, find $f(t)$ if

$$\tilde{f}(s) = \frac{s+3}{s(s+1)}.$$

- To bring $\tilde{f}(s)$ into a form where we can use the known transforms, use partial fractions.

$$\tilde{f}(s) = \frac{s+3}{s(s+1)} = \frac{A}{s} + \frac{B}{s+1} \Rightarrow s+3 = A(s+1) + Bs = (A+B)s + A$$

Equating coefficients:

$$A = 3 \quad \text{and} \quad A + B = 1 \Rightarrow B = 1 - A = -2$$

$$\Rightarrow \tilde{f}(s) = \frac{3}{s} - \frac{2}{s+1}$$

Then, the inverse Laplace transform of $3/s$ is given by

$$\mathcal{L}^{-1}[3/s] = 3 \mathcal{L}^{-1}[1/s] = 3$$

and that of $2/(s+1)$ by

$$\mathcal{L}^{-1}[2/(s+1)] = 2 \mathcal{L}^{-1}[1/(s+1)] = 2e^{-t}$$

so that $f(t) = 3 - 2e^{-t}$.

Properties of Laplace transforms

① The shift theorem

$$\mathcal{L}[e^{-at} f(t)] = \tilde{f}(s+a)$$

Proof $\mathcal{L}[e^{-at} f(t)] = \int_0^{\infty} e^{-at} e^{-st} f(t) dt$

$$= \int_0^{\infty} e^{-(s+a)t} f(t) dt = \tilde{f}(s+a)$$

definition of the Laplace transform
with s replaced by $s+a$

Example Since $\mathcal{L}[\sin \omega t] = \frac{\omega}{s^2 + \omega^2}$,

the Laplace transform of $\exp(-\gamma t) \sin \omega t$
(a damped oscillation) is

$$\mathcal{L}[e^{-\gamma t} \sin \omega t] = \frac{\omega}{(s+\gamma)^2 + \omega^2}$$

② Scaling $\mathcal{L}[f(at)] = \frac{1}{a} \tilde{f}(s/a)$

proved by a change of variable as for Fourier transforms.

③ $\mathcal{L}[t^n f(t)] = (-1)^n \frac{d^n}{ds^n} \tilde{f}(s)$

Proof

$$\frac{d^n}{ds^n} \tilde{f}(s) = \frac{d^n}{ds^n} \int_0^\infty e^{-st} f(t) dt = \int_0^\infty \frac{d^n}{ds^n} (e^{-st}) f(t) dt$$

$$= \int_0^\infty (-1)^n t^n e^{-st} f(t) dt = (-1)^n \int_0^\infty e^{-st} [t^n f(t)] dt$$

$$= (-1)^n \mathcal{L}[t^n f(t)] \Rightarrow \mathcal{L}[t^n f(t)] = (-1)^n \frac{d^n}{ds^n} \tilde{f}(s)$$

Example Given that $\mathcal{L}[1] = 1/s$, calculate the Laplace transform of $f(t) = t^n$, $n \in \mathbb{Z}_+$.

$$\mathcal{L}[t^n \cdot 1] = (-1)^n \frac{d^n}{ds^n} \mathcal{L}[1] = (-1)^n \frac{d^n}{ds^n} \frac{1}{s}$$

$$= (-1)^n (-1)^n \frac{n!}{s^{n+1}} = \frac{n!}{s^{n+1}}$$

④ The convolution theorem

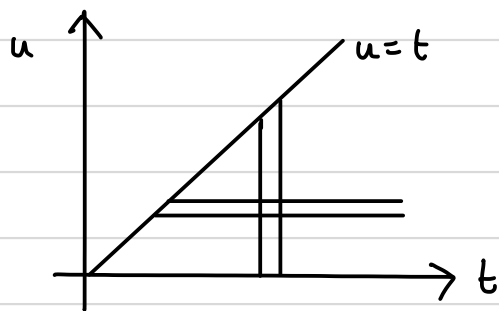
$$\mathcal{L} \left[\int_0^t f(t-u) g(u) du \right] = \tilde{f}(s) \tilde{g}(s)$$

- The integral $\int_0^t f(t-u) g(u) du$ is the convolution of f and g , denoted by $f * g$.
- Note the difference in integration limits from the Fourier transform case.

Proof

- By definition,

$$\mathcal{L} \left[\int_0^t f(t-u) g(u) du \right] = \int_0^\infty e^{-st} \int_0^t f(t-u) g(u) du dt$$



- The region of integration is an infinite wedge

- The inner integral over u from $u=0$ to $u=t$ is performed first. It runs over the vertical strip shown above.
- The outer integral over t sums the integrals over the vertical strips to give the integral over the whole infinite wedge.

- We can also integrate over horizontal strips running from $t = u$ to ∞ and then integrate over these from $u = 0$ to ∞ . This gives

$$\int_0^{\infty} g(u) \int_{t=u}^{\infty} e^{-ts} f(t-u) dt du$$

- The inner integral is performed at a fixed value of u , so we can write $t' = t - u$, so that

$$\begin{aligned} & \int_0^{\infty} g(u) \int_{t'=0}^{\infty} e^{-s(u+t')} f(t') dt' du \\ &= \int_0^{\infty} g(u) e^{-su} du \int_0^{\infty} f(t') e^{-st'} dt' \end{aligned}$$

- Then,

$$\begin{aligned} \mathcal{L} \left[\int_0^t f(t-u) g(u) du \right] &= \int_0^{\infty} f(t') e^{-st'} dt' \int_0^{\infty} g(u) e^{-su} du \\ &= \tilde{f}(s) \tilde{g}(s) \end{aligned}$$

Example Given that

$$\mathcal{L} [\cos \omega t] = \frac{s}{s^2 + \omega^2} \quad \text{and} \quad \mathcal{L} [e^{at}] = \frac{1}{s-a},$$

calculate $\mathcal{L} \left[\int_0^t e^{au} \cos(b(t-u)) du \right]$.

- We identify $f(t-u) = \cos(b(t-u))$
and $g(u) = e^{au}$

so that $f(t) = \cos(bt)$
and $g(t) = e^{at}$

- By the convolution theorem,

$$\mathcal{L} \left[\int_0^t e^{au} \cos(b(t-u)) du \right] = \tilde{f}(s) \tilde{g}(s) = \frac{s}{s^2 + b^2} \frac{1}{s-a}$$

$$= \frac{s}{(s^2 + b^2)(s-a)}$$

Inversion using the convolution theorem

- Since $\mathcal{L} \left[\int_0^t f(t-u)g(u) du \right] = \tilde{f}(s) \tilde{g}(s)$

we have that $\mathcal{L}^{-1} [\tilde{f}(s) \tilde{g}(s)] = \int_0^t f(t-u)g(u) du$

- If a function $\tilde{F}(s)$ is given as the product of two functions $\tilde{f}(s)$ and $\tilde{g}(s)$, the inverse transforms of which are known, we can use the above to calculate $F(t)$, the inverse transform of $\tilde{F}(s)$.

Example Evaluate $\mathcal{L}^{-1} \left[\frac{4}{s^2 (s+2)^2} \right]$

- We write $\tilde{f}(s) = \frac{4}{s^2}$ and $\tilde{g}(s) = \frac{1}{(s+2)^2}$

- We saw before that $\mathcal{L}[t^n] = \frac{n!}{s^{n+1}}$
- Then, $\mathcal{L}^{-1}\left[\frac{n!}{s^{n+1}}\right] = t^n$, so $\mathcal{L}^{-1}\left[\frac{1}{s^2}\right] = t$

(setting $n=1$) and $f(t) = \mathcal{L}^{-1}\left[\frac{4}{s^2}\right] = 4t$

- The other function, $\tilde{g}(s) = 1/(s+2)^2$, has the same form, but with s shifted to $s+2$.
- We therefore use the shift theorem

$$\mathcal{L}[e^{-at} f(t)] = \tilde{f}(s+a) \text{ or } \mathcal{L}^{-1}[\tilde{f}(s+a)] = e^{-at} f(t)$$

to see that $g(t) = \mathcal{L}^{-1}\left[\frac{1}{(s+2)^2}\right] = e^{-2t} \cdot t = te^{-2t}$

- From the convolution theorem,

$$\mathcal{L}^{-1}[\tilde{f}(s)\tilde{g}(s)] = \int_0^t f(t-u)g(u)du$$

and $\mathcal{L}^{-1}\left[\frac{4}{s^2} \frac{1}{(s+2)^2}\right] = 4 \int_0^t (t-u) u e^{-2u} du$

$$= 4 \left\{ \left[\frac{(t-u) u e^{-2u}}{-2} \right]_0^t + \frac{1}{2} \int_0^t (t-2u) e^{-2u} du \right\}$$

$$= 2 \left\{ \left[\frac{(t-2u) e^{-2u}}{-2} \right]_0^t - \int_0^t \frac{(-2) e^{-2u}}{-2} du \right\}$$

$$= \cancel{2} \left\{ \frac{t}{\cancel{2}} e^{-2t} + \frac{t}{\cancel{2}} - \left[\frac{e^{-2u}}{\cancel{-2}} \right]_0^t \right\} = t e^{-2t} + t + e^{-2t} - 1$$

$$= (t+1) e^{-2t} + t - 1.$$

The gamma function

- We wish to calculate the Laplace transform of $f(t) = t^p$ for $p > -1$, or

$$\mathcal{L}[t^p] = \int_0^\infty e^{-st} t^p dt$$

- Putting $u = st$,

$$\mathcal{L}[t^p] = \int_0^\infty e^{-u} \left(\frac{u}{s} \right)^p \frac{du}{s} = \frac{1}{s^{p+1}} \int_0^\infty e^{-u} u^p du$$

- The integral on the right is independent of s .
- It is a function of p called the gamma function. Usually, we define

$$\Gamma(p) = \int_0^\infty e^{-u} u^{p-1} du, \quad p > 0.$$

- We can then write

$$\mathcal{L}[t^p] = \frac{\Gamma(p+1)}{s^{p+1}} \quad p > -1$$

- We saw earlier for positive integers n that

$$\mathcal{L}[t^n] = \frac{n!}{s^{n+1}} \quad n \in \mathbb{Z}_+$$

- This implies that $\Gamma(n+1) = n!$ for $n \in \mathbb{Z}_+$, and the gamma function is a generalisation of the factorial to non-integer arguments.

Properties of the gamma function

① For every $p > 0$, $\Gamma(p+1) = p \Gamma(p)$

Proof:
$$\begin{aligned}\Gamma(p+1) &= \int_0^{\infty} e^{-u} u^p du \\ &= \left[-e^{-u} u^p \right]_0^{\infty} - \int_0^{\infty} (-e^{-u})(p u^{p-1}) du \\ &= p \int_0^{\infty} e^{-u} u^{p-1} du \\ &= p \Gamma(p)\end{aligned}$$

② $\Gamma(1) = 1$

Proof:
$$\int_0^{\infty} e^{-u} du = \left[-e^{-u} \right]_0^{\infty} = 1$$