RIGID BODY | INERTIA TENSORS

1. Carbon monoxide

Calculating the inertia tensor of a molecule requires knowing the distance between the atoms and their masses.

The distance, or bodn length, can be found at the Computational Chemistry Comparison and Benchmark DataBase of the National Institute of Standards and Technology, NIST, U.S.A. Querying about the chemical formula, CO in this case, we can find the distances, or the positions in a cartesian frame of reference, expressed in angstrom, equivalent to 1×10^{-10} m.

The mass for each chemical element is expressed in unified atomic mass units, u, in the periodic table published by the International Union of Pure and Applied Chemistry, IUPAC. This is the mass expressed in grams of one mole of atoms with the proportion of isotopes found in nature. To obtain the mass in grams of just one atom, simply divide by the number of atoms in this mole, Avogadro's constant, $N_A = 6{,}022 \times 10^{23} \,\mathrm{mol}^{-1}$.

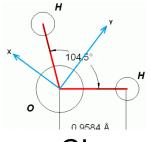
Express the inertia tensor in SI units (kg² m). Result:

$$\overline{\overline{I}} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1.45 \cdot 10^{-46} & 0 \\ 0 & 0 & 1.45 \cdot 10^{-46} \end{bmatrix}$$

2. Water

Express the inertia tensor in SI units. Result:

$$\overline{\overline{I}} = \begin{bmatrix} 1.02 \cdot 10^{-47} & 0 & 0\\ 0 & 1.92 \cdot 10^{-47} & 0\\ 0 & 0 & 2.95 \cdot 10^{-47} \end{bmatrix}$$



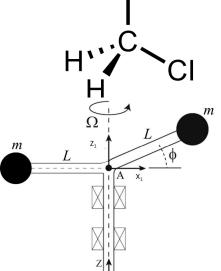
3. Dichloromethane

The chemical formula for this molecule is CH₂Cl₂. Express the inertia tensor in SI

$$\overline{\overline{I}} = \begin{bmatrix} 2.69 \cdot 10^{-46} & 0 & 0\\ 0 & 2.56 \cdot 10^{-45} & 0\\ 0 & 0 & 1.04 \cdot 10^{-36} \end{bmatrix}$$



The figure shows a system at t = 0 with weights at the ends of two arms. The vertical rod rotates without friction at a constant angular velocity Ω with respect to the inertial frame of reference O_{xyz} . The rod and the arms have negligible masses, compared to the mass m of each weight. Calculate:



- a) the inertia tensor with respect to A as a function of time $\overline{I}_A(t)$
- b) the angular momentum $\vec{L}_A(t) = \overline{\bar{I}}_A(t)\vec{\Omega}$ and torque $\vec{\tau}(t) = \dot{\vec{L}}(t)$.

Results:

$$\begin{split} \overline{\overline{I}}_A = \begin{bmatrix} \ell^2 m \left(-\cos^2\left(\phi\right)\cos^2\left(\Omega t\right) - \cos^2\left(\Omega t\right) + 2\right) & -\ell^2 m \left(\cos^2\left(\phi\right) + 1\right)\sin\left(\Omega t\right)\cos\left(\Omega t\right) & \frac{\ell^2 m (\sin\left(\Omega t - 2\phi\right) - \sin\left(\Omega t + 2\phi\right))}{4} \\ -\ell^2 m \left(\cos^2\left(\phi\right) + 1\right)\sin\left(\Omega t\right)\cos\left(\Omega t\right) & \ell^2 m \left(\sin^2\left(\phi\right)\sin^2\left(\Omega t\right) - 2\sin^2\left(\Omega t\right) + 2\right) & -\frac{\ell^2 m (\cos\left(\Omega t - 2\phi\right) - \cos\left(\Omega t + 2\phi\right))}{4} \\ \frac{\ell^2 m (\sin\left(\Omega t - 2\phi\right) - \sin\left(\Omega t + 2\phi\right))}{4} & -\frac{\ell^2 m (\cos\left(\Omega t - 2\phi\right) - \cos\left(\Omega t + 2\phi\right))}{4} & \ell^2 m \left(\cos^2\left(\phi\right) + 1\right) \end{bmatrix} \\ \vec{L}_A = \begin{bmatrix} \frac{\Omega \ell^2 m (\sin\left(\Omega t - 2\phi\right) - \sin\left(\Omega t + 2\phi\right))}{4} \\ -\frac{\Omega \ell^2 m (\cos\left(\Omega t - 2\phi\right) - \cos\left(\Omega t + 2\phi\right))}{4} \end{bmatrix} & \vec{\tau}_A = \begin{bmatrix} \frac{\Omega^2 \ell^2 m (\cos\left(\Omega t - 2\phi\right) - \cos\left(\Omega t + 2\phi\right))}{4} \\ \frac{\Omega^2 \ell^2 m (\sin\left(\Omega t - 2\phi\right) - \sin\left(\Omega t + 2\phi\right))}{4} \\ 0 \end{bmatrix} \end{split}$$