

Scalable Linear Solvers for Next Generation Weather and Climate

Christopher Maynard

- Met Office, Exeter, UK and

University of Reading

Thomas Melvin

- Met Office, Exeter, UK

Eike Müller

- University of Bath



Gungho, LFRic and PSyclone Time-stepping and the solver Performance Analysis

1. *PSyclone and its uses in LFRic* I. Kavcic **MS02** Wed 1300-1330
2. *On using a DSL Approach Performance Portability of the LFRic Weather and Climate Model* CMM **MS08** Wed 1630-1700
3. *Building a Performance Portable Software System for the Met Office's Weather and Climate Model, LFRic* **CSM07** (Poster session)

Acknowledgments

Nerc: Gungho partner Universities

STFC: PSyclone development

Met Office: LFRic development team

Dynamics Research

Subgrid Physics developers

工合

Gungho: Mixed finite element dynamical core



LFRic: Model infrastructure for next generation modelling



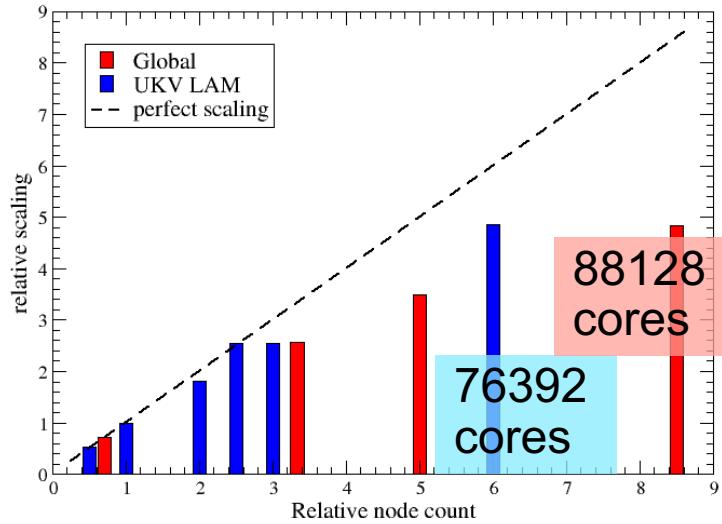
PSyClone: Parallel Systems code generation used in LFRic and Gungho



UM: Current modelling environment (UM parametrisations are being reused in LFRic

Uses Lon-Lat grid
Scientifically very good
Good computational performance

Very High Resolution scaling
6.5 Km resolution



The finger of blame ...
Lon-lat grid is preventing scaling
10km resolution (mid-latitudes) → 10m at poles

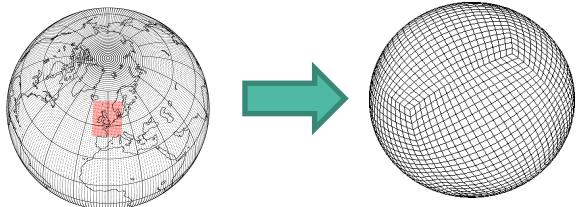
Move to quasi-uniform mesh to remove polar singularity

Maintain 'good' aspects of current model

- No computational modes
- Accurate dispersion
- Semi-Implicit timestepping
- Reuse subgrid parametrizations

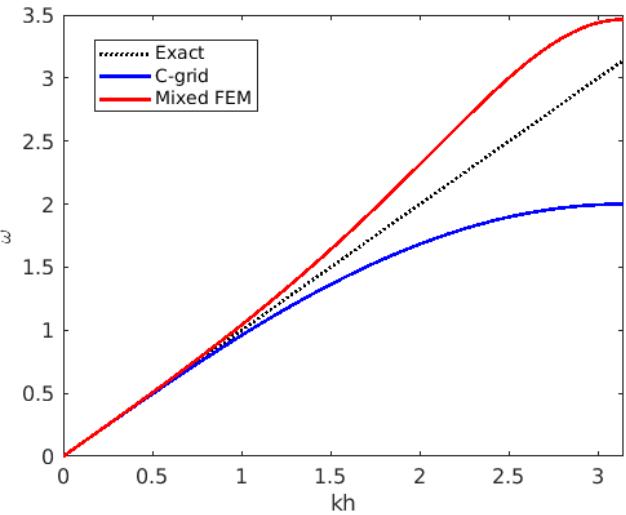
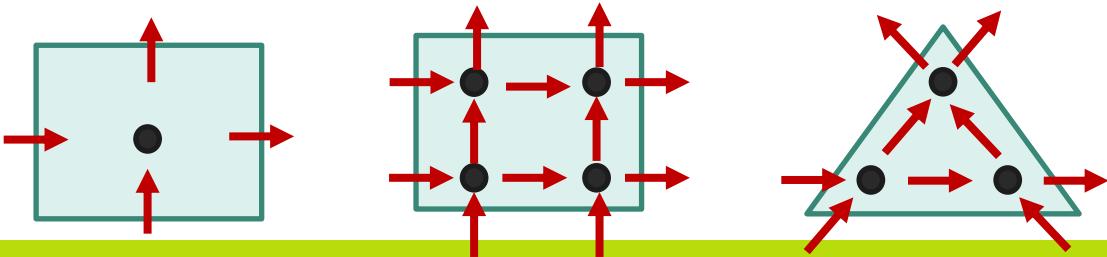
Improve inherent conservation

Improve scalability



Mixed Finite Element method gives

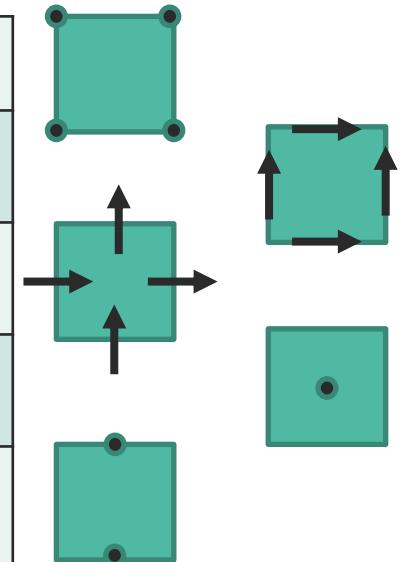
- Compatibility: $\nabla \times \nabla \varphi = 0, \nabla \cdot \nabla \times \boldsymbol{v} = 0$
- Accurate balance and adjustment properties
- No orthogonality constraints on the mesh
- Flexibility of choice mesh (quads, triangles) and accuracy (polynomial order)



Mixed Finite Element Method

$$W_0 \xrightarrow{\nabla} W_1 \xrightarrow{\nabla \times} W_2 \xrightarrow{\nabla \cdot} W_3.$$

W_0	Pointwise scalars
W_1	Circulation Vectors
W_2	Flux Vectors
W_3	Volume integrated Scalars
W_θ	Pointwise scalars



Inspired by iterative-semi-implicit semi-Lagrangian scheme used in UM

Scalar transport uses high-order, upwind, explicit Eulerian FV scheme

Wave dynamics (and momentum transport) use iterative-semi-implicit, lowest order mixed finite element method (equivalent to C-grid/Charney-Phillips staggering)

$$\bar{F}^\alpha \equiv \alpha F^{n+1} + (1 - \alpha) F^n$$

$$\delta_t \mathbf{u} = -\overline{(2\Omega + \nabla \times \mathbf{u}) \times \mathbf{u} + \nabla (K + \Phi)} + c_p \theta \nabla \Pi^\alpha$$

$$\delta_t \rho = -\nabla \cdot [\mathcal{F}(\rho^n, \bar{\mathbf{u}}^{1/2})]$$

$$\delta_t \theta = -\mathcal{A}(\theta^n, \bar{\mathbf{u}}^{1/2})$$

Quasi-Newton Method: $\mathcal{L}(\mathbf{x}^*) \mathbf{x}' = -\mathcal{R}(\mathbf{x}^{(k)})$.

Linearized around reference state (previous time-step state) $\mathbf{x}^* \equiv \mathbf{x}^n$

Solve for increments on latest state: $\mathbf{x}' \equiv \mathbf{x}^{(k+1)} - \mathbf{x}^{(k)}$

Semi-Implicit system contains terms needed for acoustic and buoyancy terms

$$\mathcal{L}(\mathbf{x}_{\text{phys}}^*) \mathbf{x}'_{\text{phys}} = \begin{cases} \mathbf{u}' - \mu \left(\frac{\mathbf{n}_b \cdot \mathbf{u}'}{\mathbf{n}_b \cdot \mathbf{z}_b} \right) \mathbf{z}_b \\ + \tau_u \Delta t c_p (\theta' \nabla \Pi^* + \theta^* \nabla \Pi'), \\ \rho' + \tau_\rho \Delta t \nabla \cdot (\rho^* \mathbf{u}'), \\ \theta' + \tau_\theta \Delta t \mathbf{u}' \cdot \nabla \theta^*, \\ \frac{1 - \kappa}{\kappa} \frac{\Pi'}{\Pi^*} - \frac{\rho'}{\rho^*} - \frac{\theta'}{\theta^*}, \end{cases}$$

Solver Outer system with Iterative (GCR) solver

$$\begin{pmatrix} M_2^{\mu,C} & -P_{2\theta}^{\Pi^*} & -G^{\theta^*} \\ D^{\rho^*} & M_3 & \\ P_{\theta 2}^{\theta^*} & M_\theta & \\ -M_3^{\rho^*} & -P_{3\theta}^* & M_3^{\Pi^*} \end{pmatrix} \begin{pmatrix} \tilde{u}' \\ \tilde{\rho}' \\ \tilde{\theta}' \\ \tilde{\Pi}' \end{pmatrix} = \begin{pmatrix} -\mathcal{R}_u \\ -\mathcal{R}_\rho \\ -\mathcal{R}_\theta \\ -\mathcal{R}_\Pi \end{pmatrix}$$

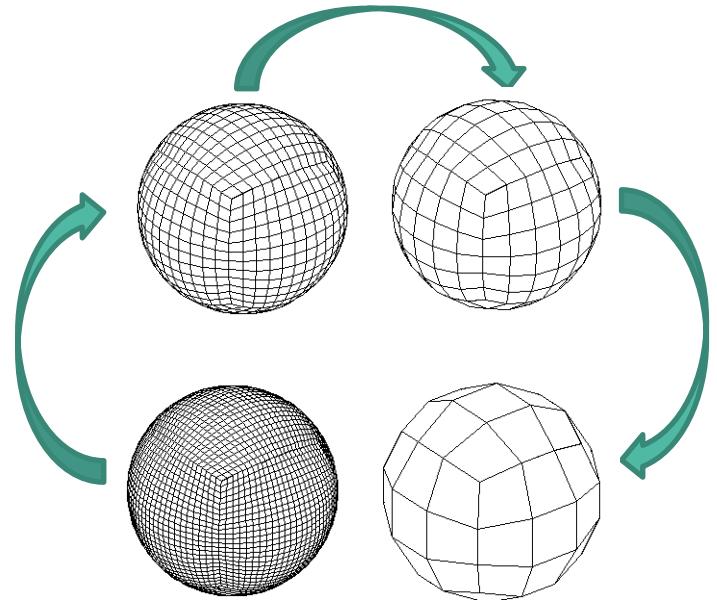
- Contains all couplings
- Preconditioned by approximate Schur complement for the pressure increment
- Velocity and potential temperature mass matrices are lumped

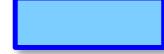
- Helmholtz system $H\Pi' = R$ solved using a single Geometric-Multi-Grid V-cycle with block-Jacobi smoother

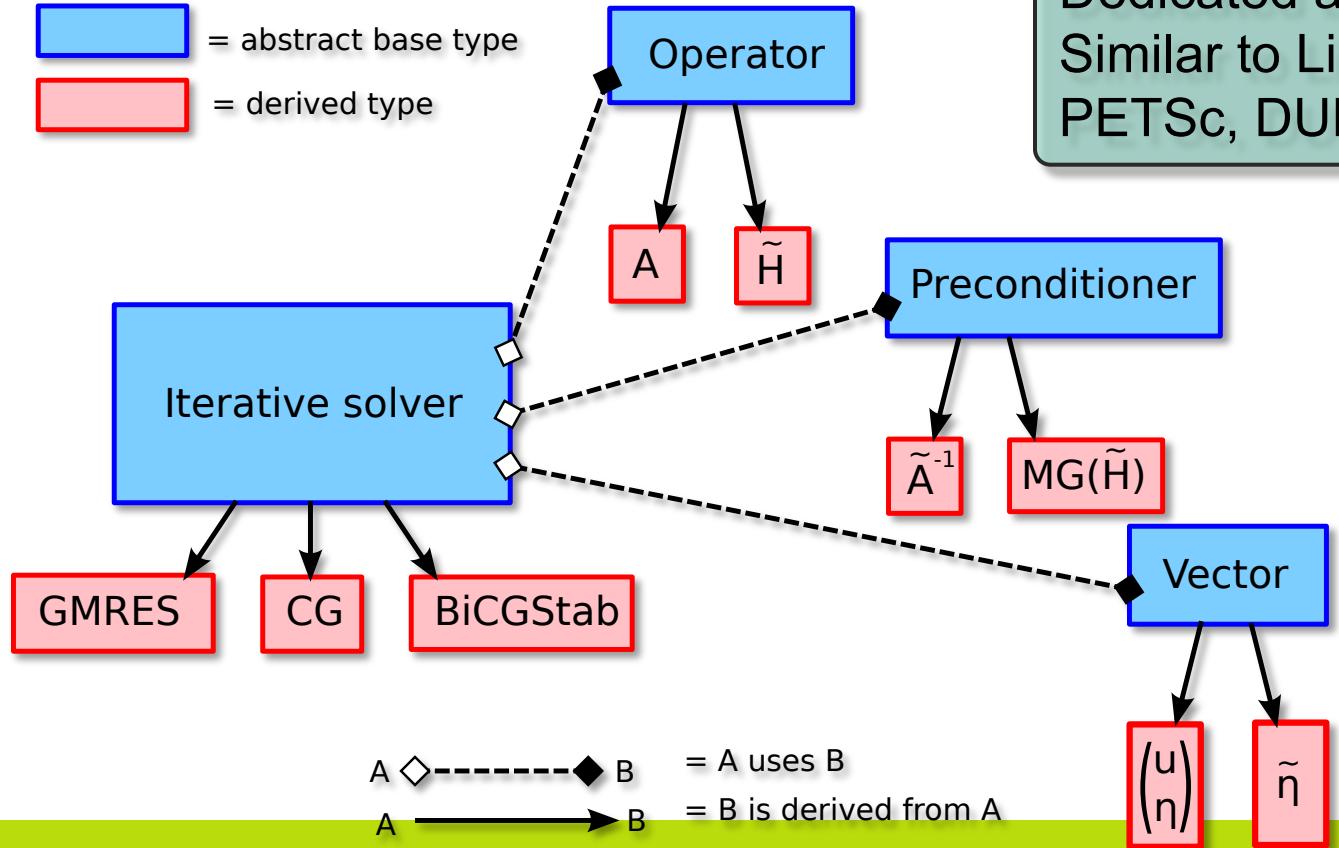
$$H = M_3^{\Pi^*} + \left(P_{3\theta}^* \mathring{M}_{\theta}^{-1} P_{\theta 2}^{\theta^*, z} + M_3^{\rho^*} M_3^{-1} D^{\rho^*} \right) \left(\mathring{M}_2^{\mu, C} \right)^{-1} G^{\theta^*}.$$

- Block-Jacobi smoother with small number (2) of iterations on each level
- Exact (tridiagonal) vertical solve: \widehat{H}_z^{-1}

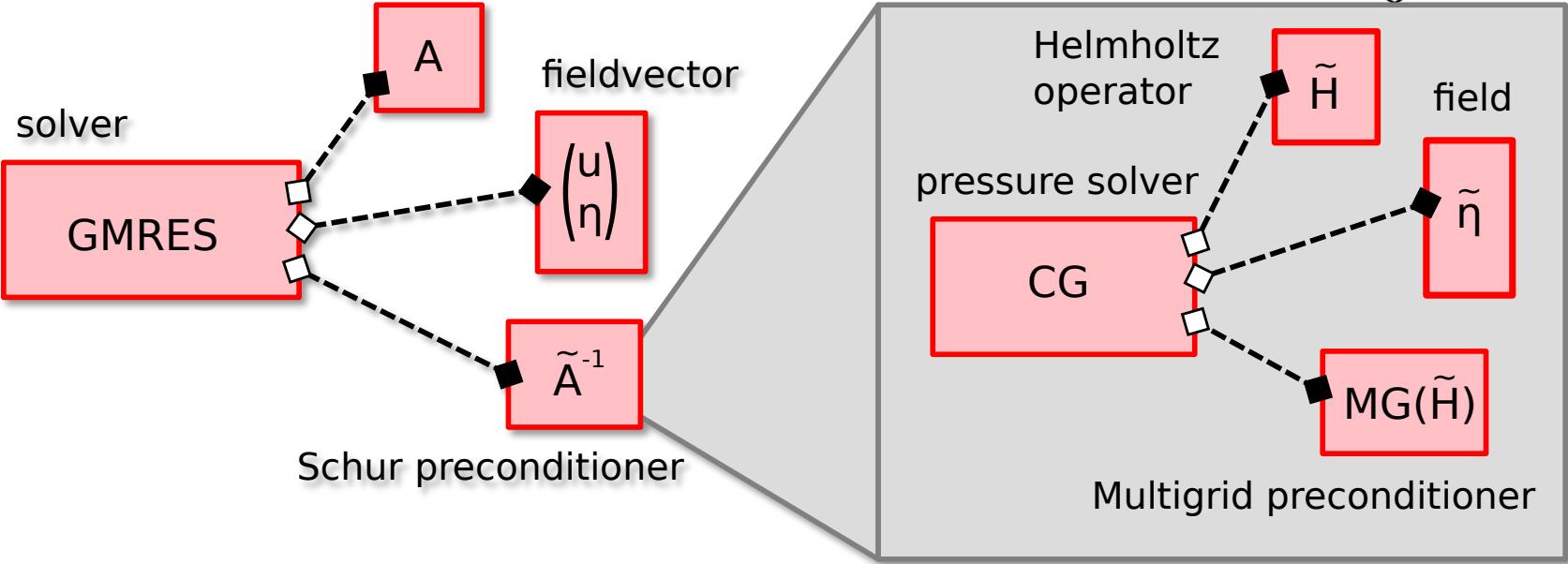
$$\widetilde{\Pi}' \leftarrow \widetilde{\Pi}' + \omega \widehat{H}_z^{-1} \left(\mathcal{B} - H \widetilde{\Pi}' \right)$$



 = abstract base type
 = derived type



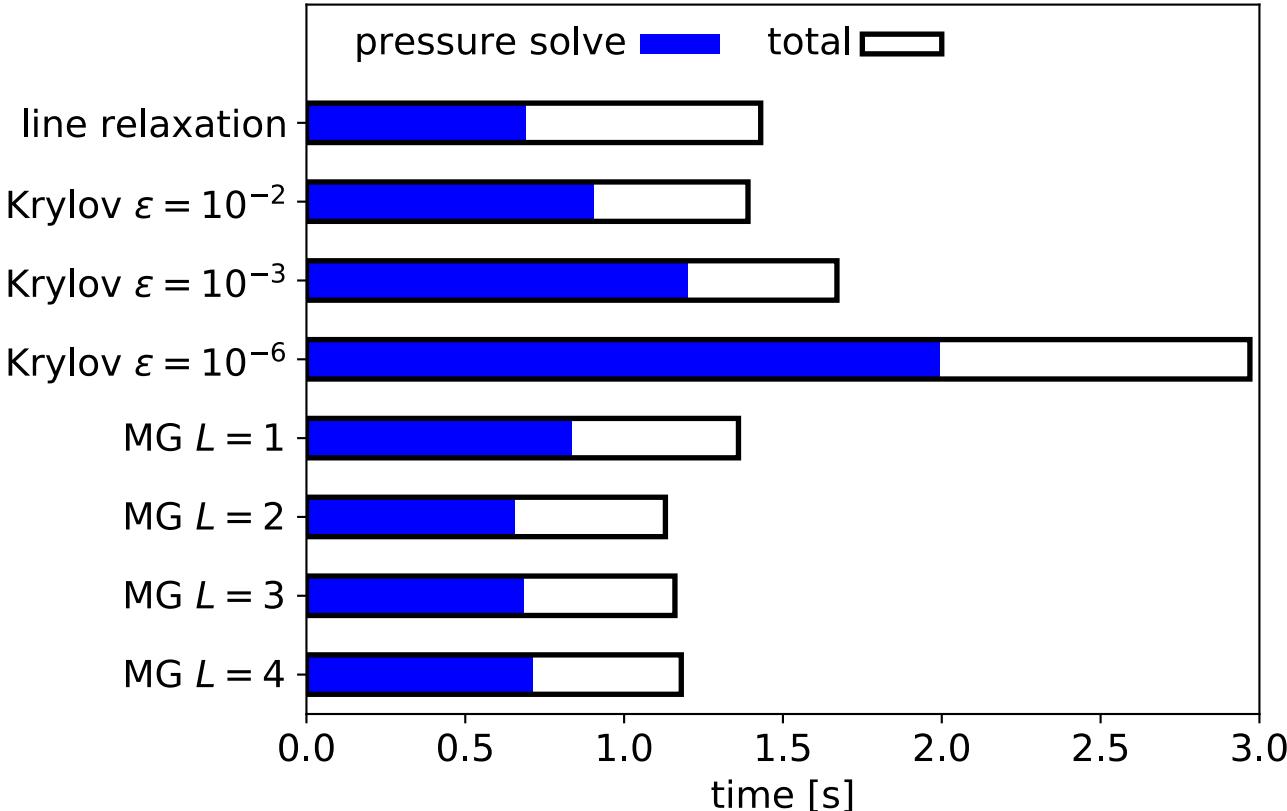
Dedicated abstraction in F2K3 OO
 Similar to Lin. Alg Libs e.g.
 PETSc, DUNE-ISTL, Trillinos



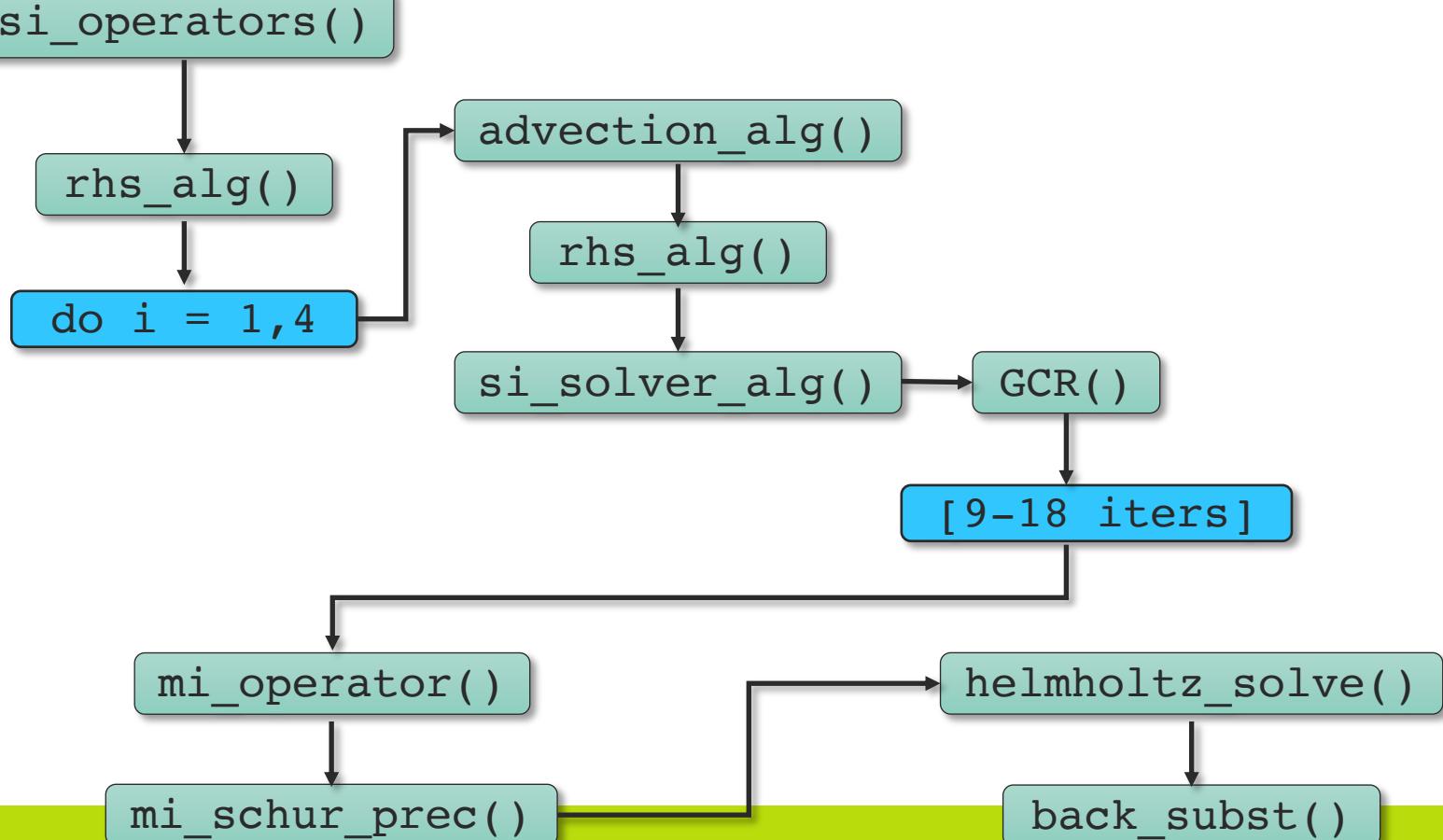
Allows for easy implementation of sophisticated nested solver
Multigrid preconditioner - reduce work for iterative solver
- faster and less global sums (better scaling)

C192 cubed sphere
with 30 L (~50Km)
Baroclinic wave test
Met Office Cray
XC40 64 nodes
(2304 cores) Mixed
mode 6 MPI/6 OMP
threads

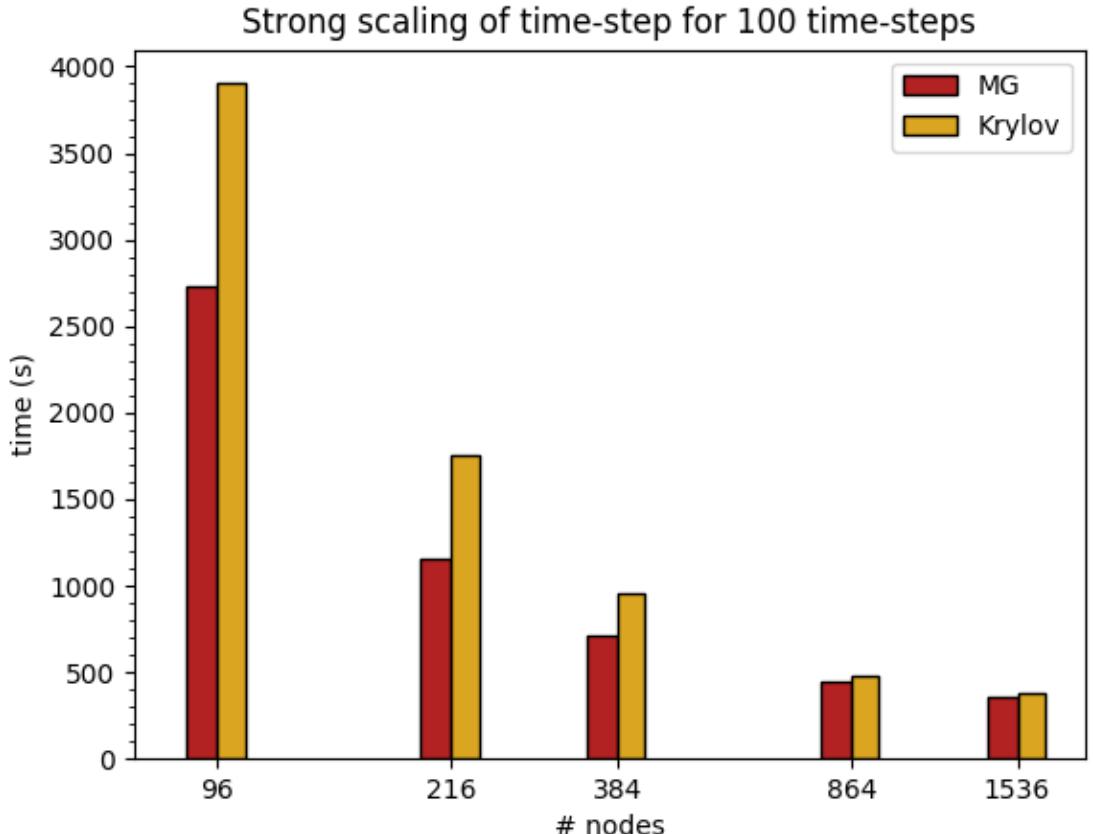
c.f. $\|r\| = \|\mathbf{A}x - b\|$ Of
Krylov 10^{-2}
Before and after MG
3-level V-cycle



Anatomy of a time-step



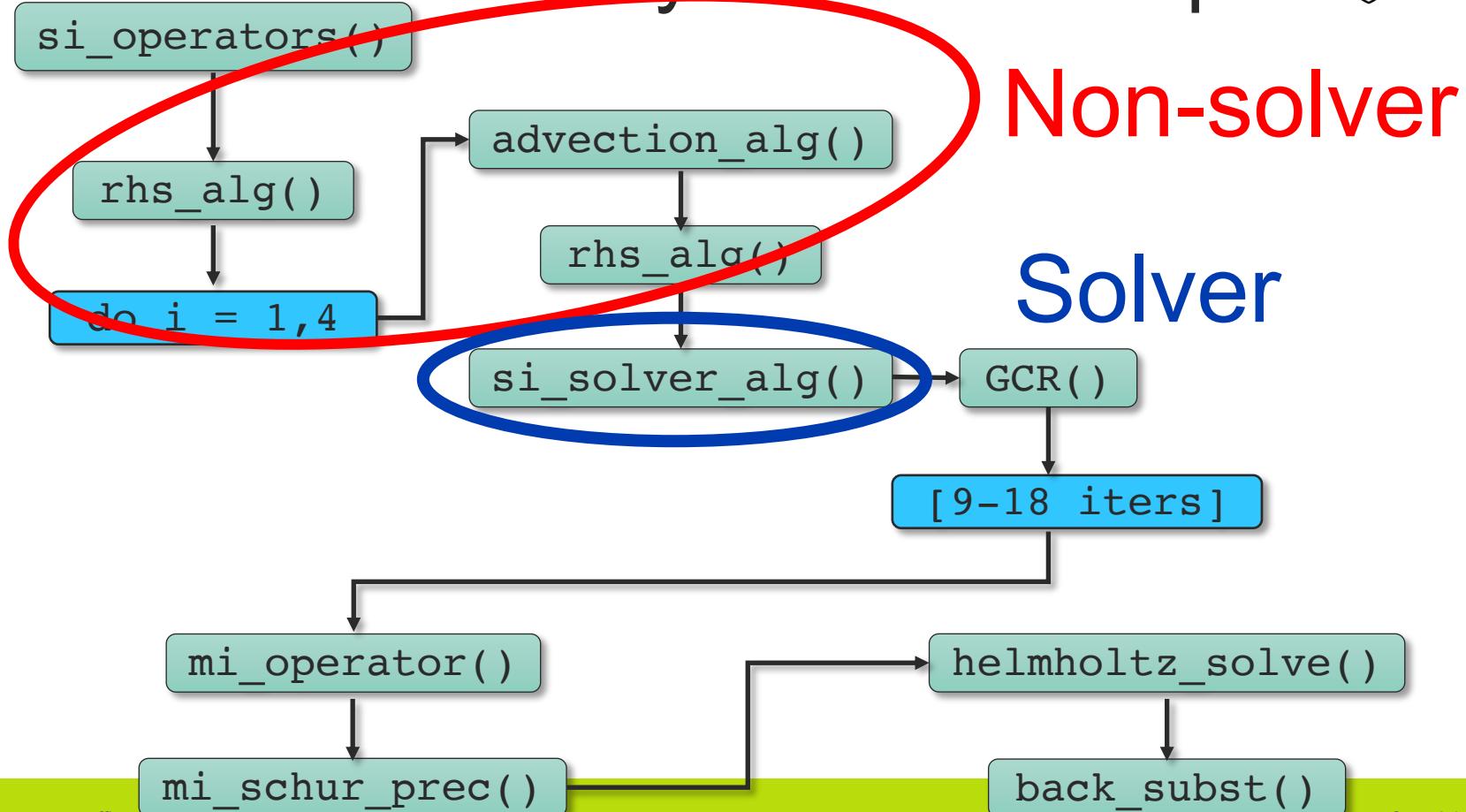
Strong Scaling

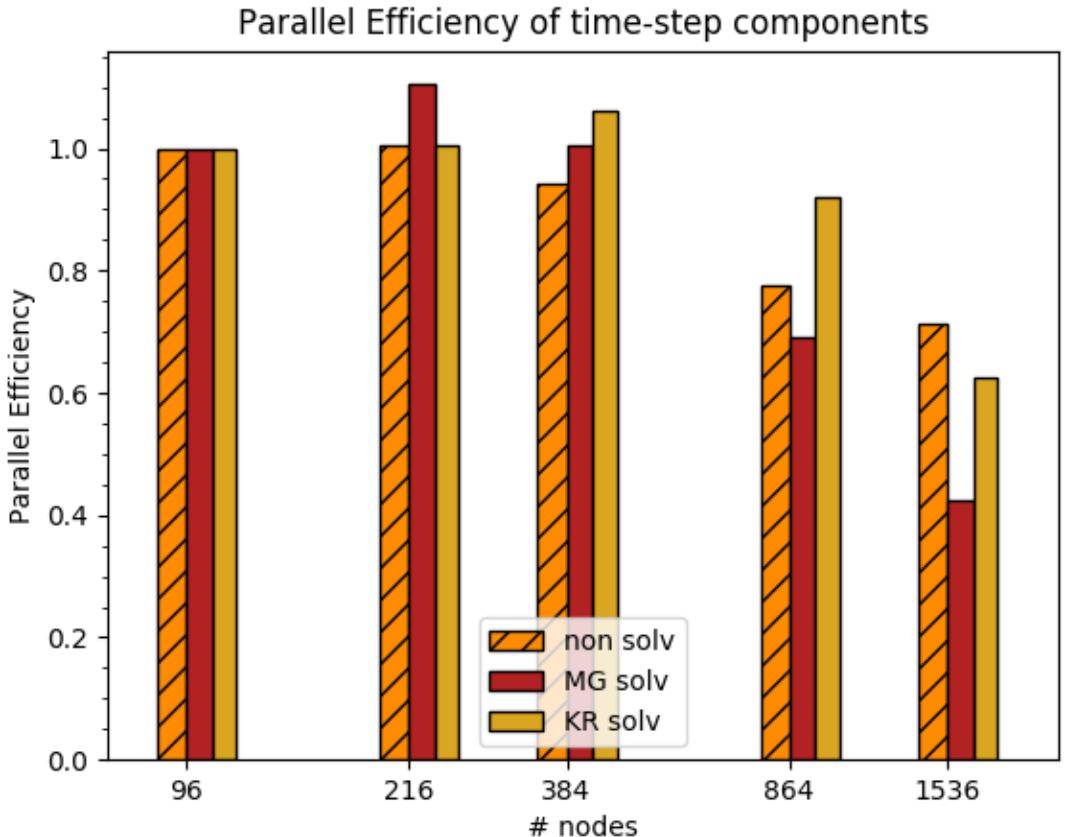


C1152 cubed sphere ~ 9 Km
Mixed mode 6 MPI/6 OMP

96 nodes LV = 48x48 x30L ~
75 K dofs
1536 nodes LV = 12x12 x30L
~4300 dofs

Multigrid is faster, but ...
Krylov subspace solver
scales better but always
slower ...





C.F. 96 node run
Non-solve is same for both,
scales OK, not great

MG and KR solver start to
drop over

MG scaling looks worse
because starts better, always
faster

si_operators()

rhs_alg()

do i = 1,4

advection_alg()

rhs_alg()

si_solver_alg()

Mixed op

Mixed Schur preconditioner

GCR()

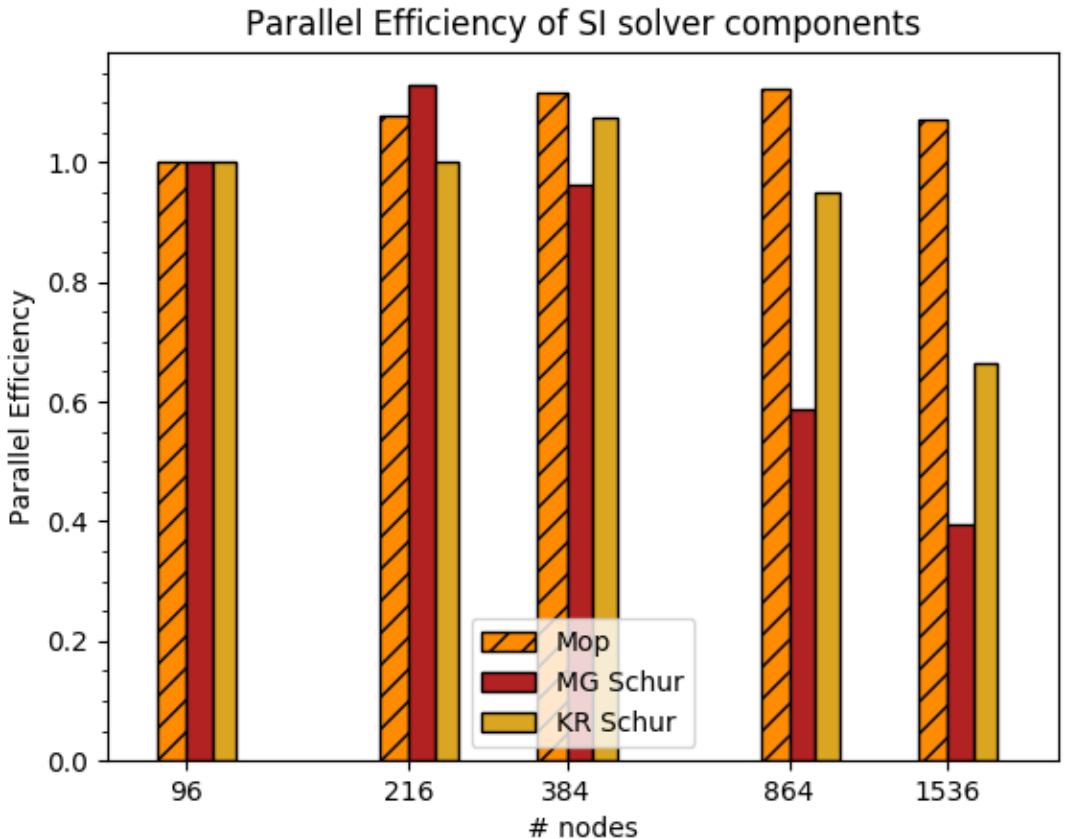
[9-18 iters]

mi_operator()

mi_schur_prec()

helmholtz_solve()

back_subst()



Mixed operator is same for both shows perfect scaling

MG and KR Mixed Schur precon, MG looks to be scaling much worse.

MG scaling looks worse because starts better, always faster

si_operators()

rhs_alg()

do i = 1,4

advection_alg()

rhs_alg()

si_solver_alg()

[9-18 iters]

mi_operator()

mi_schur_prec()

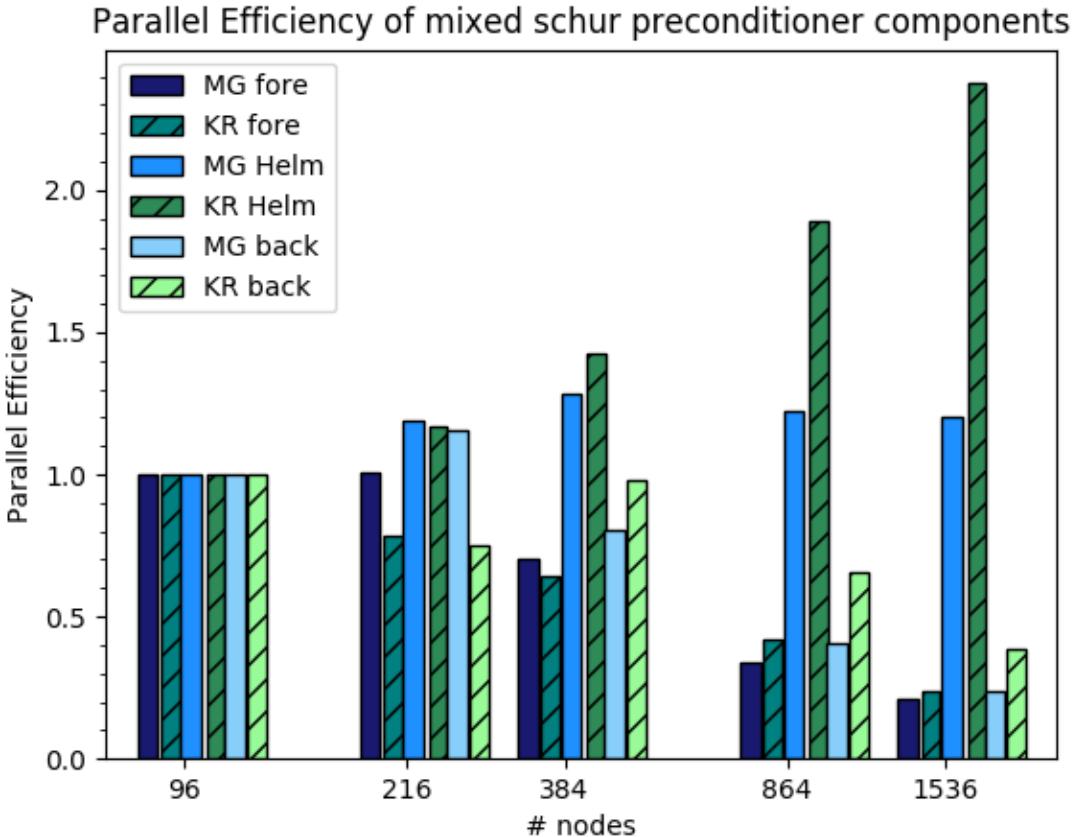
Pressure solve

Back substitution

GCR()

helmholtz_solve()

back_subst()



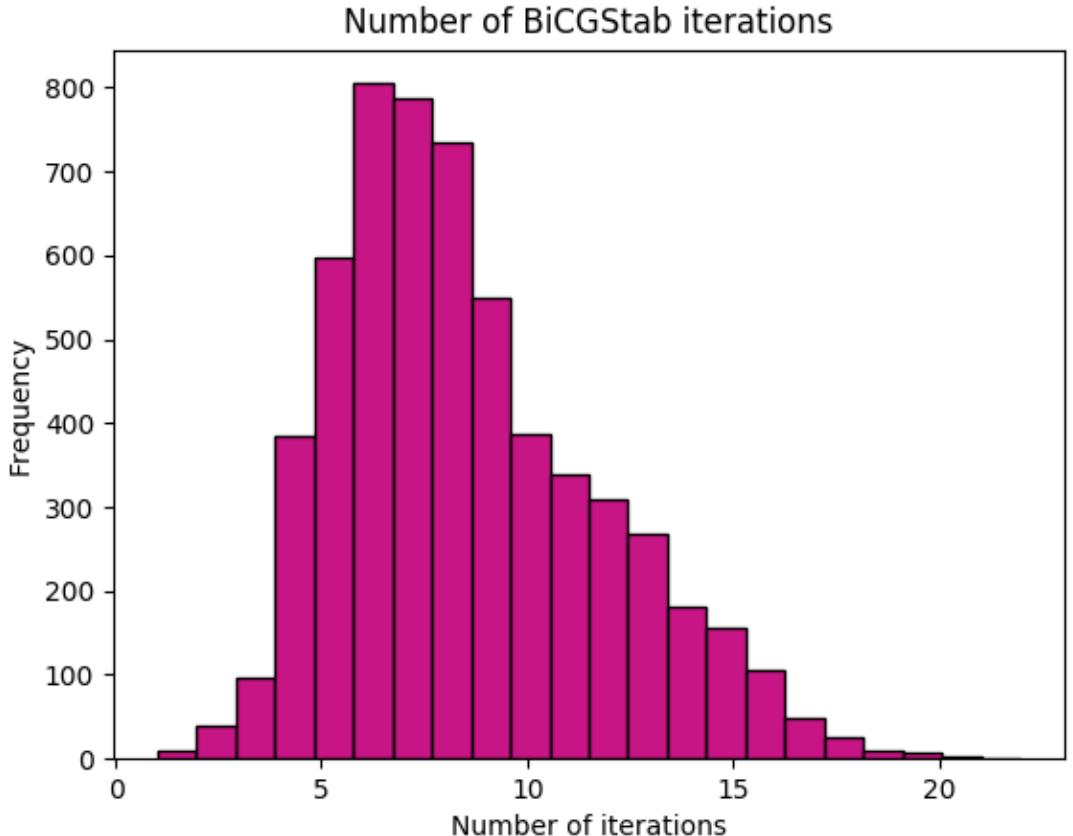
Forward and backward substitution for both KR and MG show poor scaling

KR helm (solver) shows super linear scaling (very slow on 96 nodes)

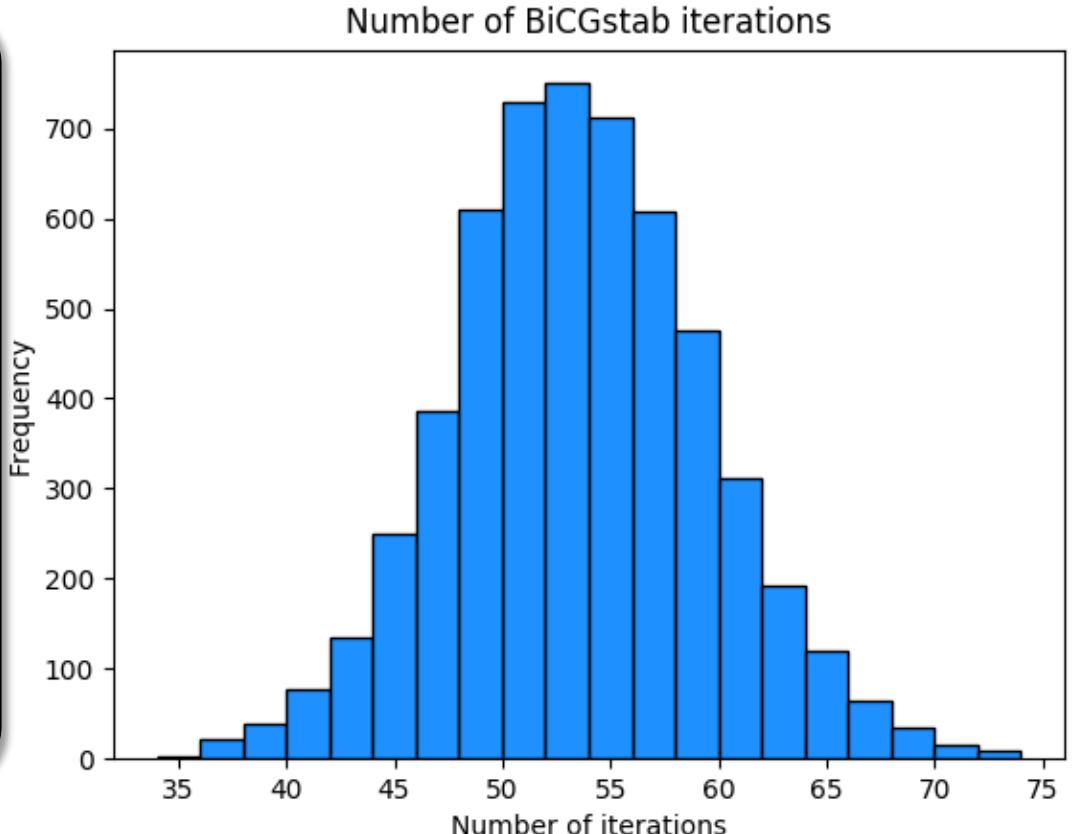
MG shows perfect scaling
→ Efficient everywhere.
Always fastest

Run-time on 1536 nodes is tiny for solver – scaled away
What happened to global sums?

100 ts x 4 sub ts = 400 calls to GCR outer solver
9-18 iters → 5800 calls to BiCGstab
Solve to 10^{-2} ave 7 iters
Total number of BiCGstab iters ~ 50K
5 Global sums per iter
~250K GS
At 55K cores, GS latency < $10\mu\text{s}$ → 2.5 s
Best case scenario,
adaptive network slow that down



Solve 10^{-6} ave is 50-55 iters
Total BiCGStab iters ~
300,000
→ 1.5 million Global sums
5x more than 10^{-2}
Here, multigrid has massive
scaling advantage. 5-10x
reduction in cost of GS
GCR still has GS



Implemented sophisticated solver framework
Allows for ease-of-use algorithmic changes
Including multigrid solver
Compared BiCGstab 10^{-2} to 3-level MG

MG always faster
BiCGStab super-scales – because dominated by computation on few nodes for large residual
Both solvers scale away, and become small in the profile → local comms then dominate

Is $\|r\| = 10^{-2}$ realistic? With Physics & Orography?
Less ideal problem – multigrid has much bigger advantage as Kr solver will see many more Global Sums.