

## 44(1) : Analysis of New Accelerations.

### The a Type Accelerations

This is given by the orbital type of acceleration,  $\underline{\underline{a}}^a$ . These are  $\underline{\underline{a}}^a$ , and the spin type of acceleration,  $\underline{\underline{a}}^s$ . There are also relativistic accelerations defined by:

$$\underline{\underline{a}}^a = \frac{d\underline{\underline{v}}^a}{dt} + c \underline{\underline{\omega}}^a \underline{\underline{v}}^a + c \underline{\underline{\omega}}^a b \underline{\underline{v}}^b - c v^b \underline{\underline{\omega}}^a b \quad (1)$$

$$\text{and } \underline{\underline{a}}^s = \underline{\underline{\Omega}} \times \underline{\underline{v}}^a - \underline{\underline{\omega}}^a b \times \underline{\underline{v}}^b. \quad (2)$$

The spin correction terms in eq. (1) represent the difference between a Newtonian orbit and a relativistic orbit. The terms in eq. (2) are:

- i) The relativistic velocity,  $\underline{\underline{\Omega}} \times \underline{\underline{v}}$ ;
- ii) The relativistic Coriolis acceleration,  $c \underline{\underline{\omega}}^a b \times \underline{\underline{v}}$ .

Now define the force:

$$\underline{\underline{F}}^a = m \underline{\underline{a}}^a \quad (3)$$

$$\frac{1}{c} \underline{\underline{F}}^s = m \underline{\underline{a}}^s \quad (4)$$

The force  $\underline{\underline{F}}^s$  is  $c$  times smaller than the force  $\underline{\underline{F}}^a$ . They are both relativistic forces.

2) and are analogous respectively to the electric field strength  $\underline{E}^a$  and magnetic flux density  $\underline{B}^a$  in generally covariant electrodynamics (ECE theory):

$$\underline{E}^a = -\frac{\partial A^a}{\partial t} - c \nabla A_0^a - c \omega^a_b A^b + c A^b \omega^a_b \quad (5)$$

$$\underline{B}^a = \underline{\Omega} \times \underline{A}^a - \underline{\omega}^a_b \times \underline{A}^b \quad (6)$$

The sign change between eq. (1) and eq. (5) is a matter of convention. In electrodynamics (5) is a matter of convention. It is S.I. units.

The acceleration in eq. (5) is the gravitomagnetic type. It was used to explain the process of the expression. It may be denoted:

$$\underline{s}^a = \underline{a}_s^a / c^2 \quad (7)$$

and

$$\underline{s}^a = \underline{\Omega} \times \underline{v}^a - \underline{\omega}^a_b \times \underline{v}^b \quad (8)$$

3) In non-relativistic dynamics the velocity is defined by:

$$\underline{\Omega} = \underline{\nabla} \times \underline{v} \quad - (9)$$

but this is not generally covariant by definition.

Limit of zero Sp. i Connection

In this limit, eqs. (1) and (2) reduce to Newtonian dynamics and classical vorticity respectively:

$$\underline{a}_0^a \rightarrow \frac{d\underline{v}^a}{dt} + c \underline{\nabla} \times \underline{v}_0^a, \quad - (10)$$

$$\frac{\underline{a}^a}{c} = \underline{\Omega}^a \rightarrow \underline{\nabla} \times \underline{v}^a \quad - (11)$$

Now sum up over:

$$a = (1), (2), (3) \quad - (12)$$

so:

$$\underline{a}_0 = \underline{a}_0^{(1)} + \underline{a}_0^{(2)} + \underline{a}_0^{(3)} \quad - (13)$$

$$\underline{\Omega} = \underline{\Omega}^{(1)} + \underline{\Omega}^{(2)} + \underline{\Omega}^{(3)} \quad - (14)$$

$$\underline{a}_0 = \frac{d\underline{v}}{dt} + c \underline{\nabla} \times \underline{v}_0 \quad - (15)$$

and

$$\underline{\Omega} = \underline{\nabla} \times \underline{v} \quad - (16)$$

4) Define the gravitational potential by:

$$\Phi = CV_0 \quad \text{--- (17)}$$

$$\text{then: } \underline{a}_0 = \frac{\partial V}{\partial t} + \nabla \Phi \quad \text{--- (18)}$$

The antisymmetry law of ECE means that:

$$\frac{\partial V}{\partial t} = -\nabla \Phi \quad \text{--- (19)}$$

and this is the weak equivalence principle.

$$\underline{F} = m \frac{\partial V}{\partial t} = m \nabla \Phi \quad \text{--- (20)}$$

This gives Newtonian dynamics and adds.

$$\underline{R}_0^a = C \omega^a b \nabla^b - C V^b \omega^a \quad \text{--- (21)}$$

$$\underline{\Omega}_S^a = -\omega^a b \times \nabla^b \quad \text{--- (22)}$$

so  $\underline{R}_0^a$  gives the relativistic correction of Newtonian, and  $\underline{\Omega}_S^a$  gives the relativistic correction of vorticity.

144(2): The 5 Types Velocities and 6 Accelerations

The spin velocity is defined by:

$$\underline{\omega}^a = c(\underline{\nabla} \times \underline{r}^a - \underline{\omega}^a b \times \underline{r}^b) \quad (1)$$

and the acceleration by:

$$\underline{d}_{\text{orb}}^a = \frac{d\underline{\omega}^a}{dt} + c \nabla \omega^a + c \omega^a b \underline{\omega}^b - c \omega^b \underline{\omega}^a b \quad (2)$$

and  $\underline{d}_{\text{spin}}^a = c(\underline{\nabla} \times \underline{\omega}^a - \underline{\omega}^a b \times \underline{r}^b) \quad (3)$

In the limit of vanishing spin correction:

$$\underline{\omega}^a \rightarrow c \nabla \times \underline{r}^a \quad (4)$$

$$\underline{d}_{\text{spin}}^a \rightarrow c \nabla \times \underline{\omega}^a \quad (5)$$

Summing up over the  $a$  indices:

$$\underline{\omega} \rightarrow c \nabla \times \underline{r} \quad (6)$$

$$\underline{d}_{\text{spin}} \rightarrow c \nabla \times \underline{\omega} \quad (7)$$

So:

$$\underline{d}_{\text{spin}} = c^2 \nabla \times (\nabla \times \underline{r}) \quad (8)$$

Now we use vector identity ("Vedda Analysis Problem Solver", problem 10. 22, page 442):

$$\nabla \times (\nabla \times \underline{r}) = \nabla (\nabla \cdot \underline{r}) - \nabla^2 \underline{r} \quad (9)$$

2) To find that:

$$\underline{\underline{a}}_{\text{Spin}} = c^2 (\nabla (\underline{\underline{\gamma}} \cdot \underline{\underline{\gamma}}) - \nabla^2 \underline{\underline{\gamma}}) \quad (10)$$

It is also useful to use the result (VAPS problem 11-8).

But if

$$\underline{\underline{\gamma}} = \underline{\omega}_{av} \times \underline{\underline{\gamma}} \quad (11)$$

$$\underline{\omega}_{av} = \frac{1}{2} \nabla \times \underline{\underline{\gamma}} \quad (12)$$

then

The spin  $\underline{\underline{\gamma}}$  is

$$\underline{\underline{a}}_{\text{Spin}} = c (\nabla \times \underline{\underline{\gamma}}^a - \underline{\omega}^a b \times \underline{\underline{\gamma}}^b) \quad (13)$$

and in the limit of vanishing acceleration spin conversion:

$$\underline{\underline{a}}_{\text{Spin}}^a \rightarrow c (\nabla \times \underline{\underline{\gamma}}^a) \quad (14)$$

Summing over  $a$ :

$$\underline{\underline{a}}_{\text{Spin}} = c \nabla \times \underline{\underline{\gamma}}$$

$$\underline{\underline{a}}_{\text{Spin}} = 2c \underline{\omega}_{av} \quad (15)$$

$$-\quad(16)$$

In summary:

$$\underline{\underline{a}}_{\text{Spin}} = 2c \underline{\omega}_{av}$$

$$\underline{\underline{a}}_{\text{Spin}} = c^2 (\nabla (\nabla \cdot \underline{\underline{\gamma}}) - \nabla^2 \underline{\underline{\gamma}})$$

"In the limit of vanishing spin conversion Eqs. (16)  
are the result of spacetime torsion.

### 3) Applications

#### 1) Whirlpool Galaxy

The whirlpool galaxy is characterised by a constant  $\omega$  in eq. (12), so  $\omega$  and  $\vec{\alpha}_{\text{spin}}$  are constant. The angular velocity of the whirlpool galaxy is:

$$\omega_{\text{av}} = \frac{1}{2c} \vec{\alpha}_{\text{spin}} - (17)$$

(TBC)

$\omega_{\text{av}} = \text{constant}$

#### 2) Viscous Fluid

As in VAPS p. 478, the most general form of second derivative that can occur is a vector equation is a linear combination of terms  $\nabla^2 \underline{\Sigma}$  and  $\underline{\Sigma} (\underline{\nabla} \cdot \underline{\Sigma})$ . So  $\vec{\alpha}_{\text{spin}}$  is the most general  $\underline{\Sigma} (\underline{\nabla} \cdot \underline{\Sigma})$ . Analogously, the second derivative of  $\underline{\Sigma}$  can be used to build up a spin acceleration can be used to build up a viscous force. Usually, the latter is expressed as

$$\underline{\vec{f}_v} = \mu \nabla^2 \underline{\Sigma} + (\mu' \underline{\nabla}) \underline{\Sigma} (\underline{\nabla} \cdot \underline{\Sigma}) - (18)$$

where  $\mu$  and  $\mu'$  are coefficients.

From eq. (15)

$$\frac{1}{c} \underline{\nabla} \times \vec{\alpha}_{\text{spin}} = \underline{\Sigma} \times (\underline{\nabla} \times \underline{\Sigma})$$

$$4) \quad = \nabla (\underline{\nabla} \cdot \underline{v}) - \nabla^2 \underline{v} \quad -(19)$$

This is the structure of viscous force. The vorticity is:

$$\underline{\Omega} = \frac{a_{\text{spin}}}{c} = \underline{\nabla} \times \underline{v} \quad -(20)$$

### The Effect of Spin (Coriolis)

This is to <sup>(19)</sup> introduce non-inertial acceleration such as centripetal and Coriolis. For example:

$$\underline{w}^a = c(\underline{\nabla} \times \underline{\zeta}^a - \underline{\omega}_b^a \times \underline{\zeta}^b) \quad -(21)$$

Summing over <sup>(20)</sup>  $a$ :

$$\underline{w} = c(\underline{\nabla} \times \underline{\zeta} - \underline{\omega}_b \times \underline{\zeta}^b) \quad -(22)$$

Defining the spin conversion vector

$$\underline{\omega} = \omega_{23} \underline{i} + \omega_{31} \underline{j} + \omega_{12} \underline{k} \quad -(23)$$

After

$$r\underline{\omega} = -\underline{\omega}_b \times \underline{\zeta} \quad -(24)$$

and

$$\boxed{\underline{w} = c(\underline{\nabla} \times \underline{\zeta} + r\underline{\omega})} \quad -(25)$$

The term  $r\underline{\omega}$  comes from a spinning frame of reference. This is used in the analysis of a gyroscope as in the levator.

5) Note that the units of  $\underline{\omega}$  are invers metres, and it should not be confused with the angular velocity  $\omega_{av}$ .

Similarly :

$$\underline{a}_{\text{spin}} = c (\underline{\nabla} \times \underline{v} + \underline{\nu \omega}) - (26)$$

$$\underline{d}_{\text{spin}} = c (\underline{\nabla} \times \underline{w} + \underline{w \omega}). - (27)$$

From eq. (12), the angular velocity in general is :

$$\boxed{\underline{\omega}_{av} = \frac{1}{2} (\underline{\nabla} \times \underline{v} + \underline{\nu \omega})} - (28)$$

and consists of a variety  $\underline{\nabla} \times \underline{v}$  and a contribution  $\underline{\nu \omega}$  due to its spin conversion. The spin conversion vector is proportional to angular velocity :

$$\underline{\omega} = \frac{2}{\sqrt{3}} \underline{\omega}_{av} - (29)$$

If  $v$  is constant as in <sup>a</sup> Andromeda galaxy, then the angular velocity and spin conversion are directly proportional:

$$\underline{\omega} = d \underline{\omega}_{av} - (30)$$

where  $d$  is a constant.

144(3) : Plane Wave Solution for Spin Velocity and Acceleration in the Limit of Vanishing Convection.

In the limit of vanishing spin convection the spin velocity is :

$$\frac{\underline{w}}{c} = \underline{\nabla} \times \underline{s} \quad \dots \quad (1)$$

and the spin acceleration is :

$$\frac{\underline{\alpha}_{\text{spin}}}{c^2} = \frac{\underline{\nabla} \times (\underline{\nabla} \times \underline{s})}{c^2} \quad \dots \quad (2)$$

The plane wave solution to these equations is :

$$\underline{s} = \frac{r}{\sqrt{2}} (i - i\underline{j}) \exp(i(at - kr)) \quad \dots \quad (3)$$

where

$$r = |\underline{s}| \quad \dots \quad (4)$$

Eq (3) is the function describing a helix about

$\underline{z}$  :

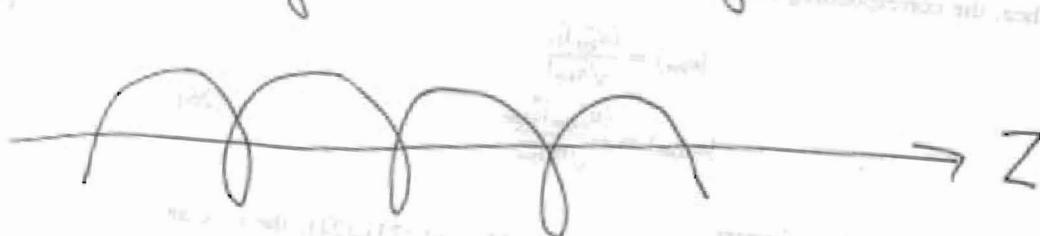


Fig (1)

Therefore:

$$\frac{\underline{w}}{c} = ik\underline{r}, \quad \dots \quad (5)$$

$$\boxed{\underline{w} = i\omega \underline{r}} \quad \dots \quad (6)$$

$$\omega = kc \quad \dots \quad (7)$$

is the angular frequency of a particle travelling along the helix.

2)

Similarly:

$$\frac{d}{dt} \underline{\text{spin}} = \omega^2 \underline{r}$$

-(8)

Therefore  $\underline{\omega}$  and  $\frac{d}{dt} \underline{\text{spin}}$  are the velocity and acceleration of the photon along the helical path in Fig (7), because for the photon, eq. (7) is true.

Therefore eqs. (1) and (2) are equations of electrodynamics as well as equations of dynamics. They are kinematic equations of the photon. This is clearly a limit of special relativity because it is the limit in which the spin current goes to zero but is not identically zero, because for the latter case, the  $\underline{\omega}$  and  $\frac{d}{dt} \underline{\text{spin}}$  also identically vanish, and the basic hypothesis of paper 143 is a  $\sqrt{\mu_0} = (0 \wedge r)^{\mu_0}$  - (9)

and this leads to the kinematic equation of the photon. In this example, the spin velocity

3) is the velocity of the photon along a helix, and the spin acceleration  $\frac{d^2 \vec{s}^{(i)}}{dt^2}$  is the acceleration of the photon along the helical trajectory.

Interestingly, a finite spin correction will produce generally covariant corrections to the trajectory of a photon in free space. The may occur in cosmological situation for example. One well known example is the light bending by gravitation. In this development it is seen from eq. (3) that the frame itself propagates along  $\Sigma$ . This is the philosophy of general relativity, the spin velocity  $\vec{w}$  and the spin acceleration  $\frac{d^2 \vec{s}^{(i)}}{dt^2}$  are defined in terms of a propagating frame of reference.

In deriving eq. (8) we have used:

$$\underline{\Sigma} \times (\underline{\Sigma} \times \underline{\Sigma}) = \underline{\Sigma} (\underline{\Sigma} \cdot \underline{\Sigma}) - \nabla^2 \underline{\Sigma} \quad -(10)$$

For eq. (3):

$$\underline{\Sigma} \cdot \underline{\Sigma} = 0 \quad -(11)$$

so  $\underline{\Sigma} \times (\underline{\Sigma} \times \underline{\Sigma}) = -\nabla^2 \underline{\Sigma} \quad -(12)$

144(4) : Place Wave Solution for Velocity and Acceleration  
In the Limit of Vanishing Spin Condition

In this case:

$$\underline{\underline{v}} = \frac{1}{c} \frac{d\underline{r}}{dt} + c \nabla \underline{r}_0 \quad (1)$$

By anti-symmetry:

$$\frac{1}{c} \frac{d\underline{r}}{dt} = c \nabla \underline{r}_0 \quad (2)$$

The dir function is:  $i(c t - k z)$

$$\underline{\underline{v}} = \frac{c}{\sqrt{2}} (i - i \hat{j}) e^{i(c t - k z)} \quad (3)$$

which is a plane wave. The position vector is:

$$\underline{\underline{r}} = (r_0, \underline{\underline{v}}) \quad (4)$$

Therefore:

$$\underline{\underline{v}} = \frac{d\underline{r}}{dt} = c \nabla \underline{r}_0 = i \omega \underline{\underline{r}} \quad (5)$$

and

$$\nabla \underline{r}_0 = i \frac{\omega}{c} \underline{\underline{r}} = i k \underline{\underline{r}} \quad (6)$$

Therefore

$$\underline{\underline{v}} = 2 i \omega \underline{\underline{r}} \quad (7)$$

The acceleration is:

$$\underline{\underline{a}}_{\text{orbital}} = \frac{d\underline{\underline{v}}}{dt} + c \nabla \underline{v}_0 \quad (8)$$

2) where:  $\frac{d\vec{v}}{dt} = c \nabla v_0 - (9)$

By antisymmetry: So:

$$\underline{a}_{\text{orbital}} = -4\omega^2 r - (10)$$

and

$$\nabla v_0 = -2\frac{\omega^2 r}{c} - (11)$$

From eq. (6) there are equations for  $v_0$  as follows:

$$\frac{d v_0}{dx} = i \frac{\kappa r}{\sqrt{2}} \exp(i(\omega t - \kappa z)) - (12)$$

$$\frac{d v_0}{dy} = \frac{\kappa r}{\sqrt{2}} \exp(i(\omega t - \kappa z)) - (13)$$

From eq. (11)

$$\frac{d v_0}{dx} = -2 \frac{\omega^2}{\sqrt{2} c} \exp(i(\omega t - \kappa z)) - (14)$$

$$\frac{d v_0}{dy} = 2 \frac{i \omega^2}{\sqrt{2} c} \exp(i(\omega t - \kappa z)) - (15)$$

Here

$$C^{\mu} = (v_0, \underline{v}) - (16)$$

$$\underline{v}^{\mu} = (v_0, \nabla) - (17)$$

Here  $v_0$  and  $\underline{v}$  are interpreted using the Stokes Theorem as follows:

$$3) r_0 = \frac{1}{c} \int \underline{v} \cdot d\underline{s} = \frac{1}{c} \int_S \underline{\Delta} \times \underline{v} \cdot \frac{dA}{ds} \quad -(18)$$

$$\text{and } v_0 = \frac{1}{c} \int \underline{a} \cdot d\underline{s} = \frac{1}{c} \int_S \underline{\Delta} \times \underline{a} \cdot dA \quad -(19)$$

where

$$\underline{a} = 2 \frac{d\underline{v}}{dt} \quad -(20)$$

so

$$v_0 = 2 \frac{dr_0}{dt} \quad -(21)$$

If the four momentum is defined as

$$\underline{p}^\mu = m \underline{v}^\mu = \left( \frac{E}{c}, \underline{p} \right) \quad -(22)$$

then

$$E = mc v_0 \quad -(23)$$

Further

$$E = 2mc \frac{dr_0}{dt} \quad -(24)$$

The rest energy for a photon is defined by de Broglie's theorem:

$$E_0 = h v c = mc \quad -(25)$$

" which case "

$$c = 2 \frac{dr_0}{dt} \quad -(26)$$

so  $r_0 \approx 1.5 \times 10^{-8}$  metres  
and is a rest length.

144(5) : New

Fields.

These are found from the minimal prescription:

$$p_\mu^a = m v_\mu^a = e A_\mu^a \quad (1)$$

$$\text{so } A_\mu^a = \frac{m}{e} v_\mu^a \quad (2)$$

$$\text{so } A_\mu^a = \frac{m}{e} v_\mu^a$$

Vector  
Orbital Potential

The space-like part of this is defined by

$$A_{\text{orb}}^a = \frac{m}{e} \left( \frac{dr^a}{dt} + c \nabla^a r^b - (\omega^a_b - c r^b \omega^a_b) \right) \quad (3)$$

Vector  
Spin Potential

The space-like part of this is defined by

$$A_{\text{spin}}^a = \frac{m}{e} \left( \vec{\omega} \times \vec{r}^a - \frac{e}{c} \vec{b} \times \vec{r}^b \right) \quad (4)$$

The spin potential is time smaller than the orbital vector potential.

The usual electric and magnetic fields are derived from the orbital part of the potential, by convention, so:

$$E_{\text{orb}}^a = - \frac{d A_{\text{orb}}^a}{dt} - c \nabla A_{\text{orb}}^a - c \omega^a_b A_{\text{orb}}^b + c A^b_{\text{orb}} \omega^a_b \quad (5)$$

and

$$2). \underline{\underline{B}}_{\text{orb}}^a = \nabla \times \underline{\underline{A}}_{\text{orb}}^a - \underline{\omega}^a b \times \underline{\underline{A}}_{\text{orb}}^b - (6)$$

New Types of Electric and Magnetic Field Derived From the Spin Potential.

These are  $c$  times smaller in magnitude than the well known electric and magnetic fields, and are defined by:

$$\underline{\underline{E}}_{\text{spin}}^a = -\frac{1}{c} \left( \frac{\partial \underline{\underline{A}}_{\text{spin}}^a}{\partial t} + c \nabla \underline{\underline{A}}_{\text{spin}}^a + c \omega_{\text{orb}}^a \underline{\underline{A}}_{\text{spin}}^b - c \underline{\underline{A}}_{\text{spin}}^b \underline{\omega}^a_b \right) \quad (7)$$

$$\underline{\underline{B}}_{\text{spin}}^a = \frac{1}{c} \left( \nabla \times \underline{\underline{A}}_{\text{spin}}^a - \underline{\omega}^a_b \times \underline{\underline{A}}_{\text{spin}}^b \right) \quad (8)$$

It can be seen that the orbital and spin vector potentials each have an internal structure defined by eqns. (3) and (4). The new types of electric and magnetic fields are  $c$  times smaller than the well known type of electric and magnetic fields. This means they are smaller in magnitude.

3) The internal structure of the usual vector potential (axial vector potential) is given in Eq. (3), leads to the possibility of many new types of spin conservation.

### Dynamics and gravitational Theory

In this theory there exist new types of spin vector potentials in dynamics and theory of gravitation, and new types of spin conservation.

### New Field Equation

There exist new field equation for the spin electric and magnetic fields.

## 14.4(6): New Spin Conservation Resource Structures

Defire:  $R_{\mu}^a = \frac{m}{e} r_{\mu}^a - (1)$

Then:  $A_{orb}^a = \frac{dR^a}{dt} + c \nabla R^a \cdot \underline{R}^b - c R^b \underline{\omega}^a{}_b$   $- (2)$

$$\frac{1}{c} A_{spin}^a = \nabla \times \underline{R}^a - \underline{\omega}^a{}_b \times \underline{R}^b - (3)$$

From the Cartan and Evans identifier:

$$\nabla \cdot A_{spin}^a = 0 - (4)$$

$$\nabla \times A_{orb}^a + \frac{1}{c} \frac{dA_{spin}^a}{dt} = 0 - (5)$$

$$\nabla \cdot A_{orb}^a = \underline{g}^a - (6)$$

$$\nabla \times A_{spin}^a - \frac{1}{c} \frac{dA_{orb}^a}{dt} = \underline{g}^a - (7)$$

eqs. (4) to

Eqs (2) and (3), summed in a new resource structure

(7), produce various spin conservation

$$R_{\mu}^a = (R_{\mu}^a, - \underline{R}^a) - (8)$$

$$= \frac{m}{e} r_{\mu}^a$$

This means that there is a "resonance" property of  $\underline{A}^a_{spin}$  and  $\underline{A}^a_{orb}$ .

144(7): Interaction of Spin and Orbital E and B

The basic structure of the theory is:

$$A = \frac{m}{e} D \wedge r \quad - (1)$$

$$H = F \left[ \sum_i [D_i \wedge A_i] \right] \quad - (2)$$

$$H = F \left[ \sum_i [D_i \wedge A_i] \right] \quad - (2)$$

The two-form  $\underline{A}$  on the LHS of eq. (1) is re-arranged to give the one-form potential  $A_{orb}^a$  we get in eq. (2). This gives:

$$\underline{A}_{orb}^a = \frac{m}{e} \left( \frac{dr^a}{dt} + c \nabla r^a + c \omega_{ab} \underline{\omega}^b - cr^a \underline{\omega}^b \right) \quad - (3)$$

$$\underline{A}_{spin}^a = \frac{m}{e} \left( \nabla \times \underline{\omega}^a - \underline{\omega}^a b \times \underline{\omega}^b \right) \quad - (4)$$

It is seen that  $\underline{A}_{spin}^a$  is  $c$  times smaller than  $\underline{A}_{orb}^a$ .

Analogously, the magnetic flux density  $\underline{B}$  is  $c$  times smaller in magnitude than the electric field strength  $\underline{E}$ . Similarly, the gravitational field  $\underline{g}$  is  $c$  times smaller than the gravitational quantity of Newtonian theory. Therefore  $\underline{A}_{spin}^a$  is going to be a small correction to the usual (Heaviside)  $\underline{A}_{orbital}$ . We can write:

$$|\underline{A}_{\text{orb}}^a| = c |\underline{A}_{\text{spur}}^a|. \quad (5)$$

Analogously:  $|\underline{E}| = c |\underline{B}| \quad (6)$

in S.I. units, and:  $|\underline{g}| = c |\underline{\Omega}|. \quad (7)$

Now write:  $\underline{B}^a(t_{\text{tot}}) = \underline{B}_{\text{orb}}^a + c \underline{B}_{\text{spur}}^a \quad (8)$

$$\underline{E}^a(t_{\text{tot}}) = \underline{E}_{\text{orb}}^a + c \underline{E}_{\text{spur}}^a \quad (9)$$

Note that:  $|\underline{E}_{\text{orb}}^a| = c |\underline{E}_{\text{spur}}^a| \quad (10)$

$$|\underline{B}_{\text{orb}}^a| = c |\underline{B}_{\text{spur}}^a| \quad (11)$$

Thus:  $\nabla \cdot \underline{B}^a(t_{\text{tot}}) = 0 \quad (12)$

$$\nabla \times \underline{E}^a(t_{\text{tot}}) + \frac{\partial \underline{B}^a(t_{\text{tot}})}{\partial t} = 0 \quad (13)$$

$$\nabla \cdot \underline{E}^a(t_{\text{tot}}) = \rho^a(t_{\text{tot}}) / \epsilon_0 \quad (14)$$

$$\nabla \times \underline{B}^a(t_{\text{tot}}) - \frac{1}{c^2} \frac{\partial \underline{E}^a(t_{\text{tot}})}{\partial t} = \mu_0 \underline{J}^a(t_{\text{tot}}) \quad (15)$$

The total  $\underline{E}^a$  and  $\underline{B}^a$  fields in eqs (8) and (9), are dominated by the orbital

fields, which have a magnitude & kind greater  
than the spin fields. However, a possible solution of  
eq. (13) is :

$$\nabla \times \underline{E}^a(\text{orbital}) + \frac{d\underline{B}^a}{dt}(\text{spin}) = 0 \quad -(16)$$

and another possible solution is :

$$\nabla \times \underline{E}^a(\text{spin}) + \frac{d\underline{B}^a}{dt}(\text{orbital}) = 0 \quad -(17)$$

In certain circumstances the orbital electric  
field may induce a spin magnetic field, and  
vice versa. In other two possible Faraday  
laws of induction are :

$$\nabla \times \underline{E}^a(\text{orbital}) + \frac{d\underline{B}^a}{dt}(\text{orbital}) = 0 \quad -(18)$$

$$\nabla \times \underline{E}^a(\text{spin}) + \frac{d\underline{B}^a}{dt}(\text{spin}) = 0 \quad -(19)$$

The Coulomb Law

This is :

$$\nabla \cdot \underline{E}(\text{tot.}) = \rho^{(\text{total})}/\epsilon_0 \quad -(20)$$

where

$$\underline{E}(\text{tot.}) = \underline{E}(\text{orb}) + \underline{E}(\text{spin}) \quad -(21)$$

4)

Here:

$$\underline{E}(\text{orb}) = - \frac{\partial \underline{A}}{\partial t} - c \nabla A \cdot (\underline{\omega}_b) - c \omega_{ob} \underline{A}_{\text{orb}}^b + c A^b \underline{\omega}_b \quad -(22)$$

$$\underline{E}(\text{spin}) = - \frac{\partial \underline{A}}{\partial t} - c \nabla A \cdot (\underline{\omega}_{\text{spin}}) - c \omega_{ob} \underline{A}_{\text{spin}}^b + c A^b \underline{\omega}_b \quad -(23)$$

Here:

$$\nabla \cdot \underline{A}(\underline{\omega}_{\text{spin}}) = 0 \quad -(24)$$

$$\nabla \times \underline{A}(\underline{\omega}_{\text{spin}}) + \frac{\partial \underline{A}}{\partial t} = 0 \quad -(25)$$

$$\nabla \cdot \underline{A}(\underline{\omega}_{\text{orb}}) = 0 \quad -(26)$$

$$\nabla \times \underline{A}(\underline{\omega}_{\text{orb}}) - \frac{1}{c^2} \frac{\partial \underline{A}}{\partial t} = 0 \quad -(27)$$

$$\nabla \times \underline{A}(\underline{\omega}_{\text{orb}}) - \frac{1}{c^2} \frac{\partial \underline{A}}{\partial t} = 0 \quad -(28)$$

Also:

$$\underline{A}_{\text{orb}} = \frac{m}{e} \left( \frac{\partial \underline{r}}{\partial t} + c \nabla r_0 + c \omega_{ob} \underline{r}^b - c \underline{r} \cdot \underline{\omega}_b \right)$$

$$\underline{A}_{\text{spin}} = \frac{m}{e} \left( \nabla \times \underline{r} - \underline{\omega}_b \times \underline{r}^b \right) \quad -(29)$$

144(10) : Experiment to Detect the E and B Spin Fields

Consider an electron trajectory:

$$\underline{r} = \frac{r}{\sqrt{2}} (\underline{i} - \underline{j}) \sin(\omega(at - nz)) \quad (1)$$

where  $r = |\underline{r}| \quad (2)$

As it notes 144(3) this produces the spin velocity:

$$\frac{\underline{v}}{c} = \underline{\nabla} \times \underline{r} \quad (3)$$

and spin acceleration:

$$\frac{d\underline{\text{spin}}}{c^2} = \underline{\nabla} \times (\underline{\nabla} \times \underline{r}) \quad (4)$$

Thus:

$$\underline{w} = i\omega \underline{r}, \quad \frac{d\underline{\text{spin}}}{c^2} = \omega^2 \underline{r}. \quad (5)$$

In electrodynamics, eq. (3) translates into:

$$\frac{\underline{A}}{c}(\text{spin}) = \frac{m}{e} \underline{\nabla} \times \underline{r} \quad (6)$$

i.e.  $\boxed{\underline{A}(\text{spin}) = i \frac{m}{e} \omega \underline{r}} \quad (7)$

In the limit of zero spin correction, and using antisymmetry:

$$\underline{E}(\text{spin}) = -2 \frac{d\underline{A}(\text{spin})}{dt} \quad (8)$$

$$= -2i \frac{m}{e} \omega \frac{dr}{dt}$$

$$= 2 \frac{m}{e} \omega^2 r \quad (9)$$

2)

So:

$$\boxed{\underline{E}(\text{spii}) = \frac{2m}{e} \omega^2 \underline{r}} \quad - (10)$$

This has an  $\propto \omega^2$  dependence which could be tested experimentally.

However, in the limit of zero spin correction

$$\underline{E}(\text{orb}) = -2 \frac{\partial \underline{A}(\text{orb})}{\partial t} \quad - (11)$$

and  $\underline{A}(\text{orb}) = 2 \frac{m}{e} \frac{d\underline{r}}{dt}$  - (12)

So

$$\underline{E}(\text{orb}) = -4 \frac{m}{e} \frac{d^2 \underline{r}}{dt^2} \quad - (13)$$

For a plane wave such as eq. (1):

$$\boxed{\underline{E}(\text{orb}) = 4 \frac{m}{e} \omega^2 \underline{r}} \quad - (14)$$

and

$$|\underline{E}(\text{orb})| = c \left| \frac{\underline{E}(\text{spii})}{c} \right| \quad - (15)$$

The force on a test electron or charge for eq. (14) is:

$$\boxed{\underline{F}_i = e \underline{E}(\text{orb}) = 4m \omega^2 \underline{r}} \quad - (16)$$

144 (ii) : Resonant Coulomb Law, Orbital Field

In this case:

$$\underline{E} = -\frac{\partial \underline{A}}{\partial t} - c \nabla A_0 - c \omega_{ob} \underline{A}^b + c \underline{A}_0^b \underline{\omega}_b \quad (1)$$

$$\text{where: } \underline{A} = \frac{m}{e} \left( \frac{\partial \underline{r}}{\partial t} + c \nabla r_0 + c \omega_{ob} \underline{r}^b - c r_0^b \underline{\omega}_b \right) \quad (2)$$

and

$$\nabla \cdot \underline{E} = \rho / \epsilon_0 \quad (3)$$

$$\text{In eq. (1)}: c \omega_{ob} \underline{A}^b = c \omega_{ob} \underline{A}^b = c \omega_0 \underline{A} \quad (4)$$

$$\text{and} \quad c \underline{A}_0^b \underline{\omega}_b = \cancel{\phi}^b \cancel{\omega}_b = \phi \underline{v}^b \underline{\omega}_b \\ = \phi \underline{\omega} \quad (5)$$

So

$$\boxed{\underline{E} = -\frac{\partial \underline{A}}{\partial t} - \nabla \phi - \omega_0 c \underline{A} + \phi \underline{\omega}} \quad (6)$$

i.e.

$$\underline{E} = -\left( \frac{\partial \underline{A}}{\partial t} + \omega_0 c \right) \underline{A} - (\nabla - \underline{\omega}) \phi \quad (7)$$

In the limit of special relativity:

$$\omega^\mu = (\omega_0, \underline{\omega}) \rightarrow 0 \quad (8)$$

so

$$\underline{E} \rightarrow -\frac{\partial \underline{A}}{\partial t} - \nabla \phi \quad (9)$$

which is the usual textbook result.

2) By antisymmetry:

$$-\left(\frac{\partial}{\partial t} + \omega_0 c\right) \underline{A} = -\left(\underline{\nabla} - \underline{\omega}\right) \phi \quad -(10)$$

So:

$$\underline{E} = -2\left(\frac{\partial}{\partial t} + \omega_0 c\right) \underline{A} = -2\left(\underline{\nabla} - \underline{\omega}\right) \phi \quad -(11)$$

Similarly:

$$\underline{A} = \frac{q m}{e} \left( \left( \frac{\partial}{\partial t} + (c\omega_0) \right) \underline{\zeta} + c(\underline{\nabla} - \underline{\omega}) \underline{r}_0 \right). \quad -(12)$$

By antisymmetry:

$$\underline{A} = \frac{q m}{e} \left( \frac{\partial}{\partial t} + (c\omega_0) \right) \underline{\zeta} = \frac{q m c}{e} (\underline{\nabla} - \underline{\omega}) \underline{r}_0 \quad -(13)$$

The sign change between eqs. (11) and (13)

is a convention.  
Therefore:

$$\begin{aligned} \underline{E} &= -\frac{q m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \left( \frac{\partial}{\partial t} + \omega_0 c \right) \underline{\zeta} \\ &= -2(\underline{\nabla} - \underline{\omega}) \phi \end{aligned} \quad -(14)$$

$$3) \text{ and } \underline{\Delta} \cdot \underline{E} = \rho / \epsilon_0 - (15)$$

$$\text{In the limit } \omega^m \rightarrow 0 - (16)$$

$$\underline{E} = -2\underline{\Delta} \phi = -\frac{8m}{e} \frac{d^2 r}{dt^2} = -2 \frac{\partial A}{\partial t} - (17)$$

$$\text{In this limit } \underline{A} = \frac{4m}{e} \frac{dr}{dt} = \frac{4m}{e} v - (18)$$

Eq. (18) is the minimal prescription:

$$\underline{P} = mv = \frac{e}{2} \underline{A} - (19)$$

From eqs. (14) and (15):

$$\boxed{\frac{d^2}{dt^2} (\underline{\Delta} \cdot \underline{\Sigma}) + 2\omega_0 c \frac{d}{dt} (\underline{\Delta} \cdot \underline{\Sigma}) + c^2 \omega_0^2 (\underline{\Delta} \cdot \underline{\Sigma}) = -\frac{\rho e}{4\pi m \epsilon_0}} - (20)$$

$$\text{Denote } \underline{\Delta} \cdot \underline{\Sigma} = f(t) - (21)$$

$$\text{Then: } \boxed{\frac{d^2 f}{dt^2} + 2\omega_0 c \frac{df}{dt} + c^2 \omega_0^2 f = -\frac{\rho e}{4\pi m \epsilon_0}} - (22)$$

4) This is an Euler resonance equation if

$$\rho = \bar{\rho}_0 \cos(\omega t) - (23)$$

At resonance the time dependent  $f(t)$  become

very large so  $\underline{A} \cdot \underline{s}(t)$  becomes very large and the electric mvs very quickly. From eq. (18),  $\underline{A}$  becomes very large and  $\underline{E}$  becomes very large.

Here is a resonant amplification.

The electric field strength  $E$  (volts per metre) and a sharp max. min in voltage.

### C. Resonant frequency

Now the resonance is said to be complete, i.e., maximum power is added to the circuit due to resonance and the current is zero.

At resonance ( $\omega_0$ ) the inductor current is zero and the capacitor voltage is zero.

$$i_L = \sum_{n=1}^N V_n e^{j\omega_0 t}, \quad V_0 = 1/V \quad (24)$$

Since the voltage across the capacitor is zero the current through the capacitor

144(12) : SCR is the Coulomb Law of the Electric Spin Field.

For the orbital electric field, and if by convention we choose constant signs,

$$\underline{E} = - \left( \frac{d}{dt} + \omega_0 c \right) \underline{A} - (\nabla - \underline{\omega}) \phi \quad (1)$$

$$\underline{E} = - \left( \frac{d}{dt} + \omega_0 c \right) \underline{A} - c(\nabla - \underline{\omega}) \underline{r}_0 \quad (2)$$

$$\underline{A} = - \frac{m}{e} \left( \frac{d}{dt} + \omega_0 c \right) \underline{r}$$

giving us SCR equation of the type:

$$j \frac{d^2 \underline{A}}{dt^2} + 2\omega_0 c \frac{d \underline{A}}{dt} + c^2 \omega_0^2 \underline{A} = \rho e \frac{e}{4\pi m \epsilon_0} \quad (3)$$

The Spin Electric Field  
This is defined as  $c$  times smaller than the orbital electric field, so

$$\underline{E} := \underline{E} / c \quad (4)$$

for the units of magnetic flux density (Tesla) SW is an electric type of field. So:

$$\underline{E} := \frac{1}{c} \underline{E} = - \frac{1}{c} \left( \frac{d}{dt} + \omega_0 c \right) \underline{A} - \frac{1}{c} (\nabla - \underline{\omega}) \phi \quad (5)$$

$$\text{where } \frac{1}{c} \underline{A} = - \frac{m}{e} (\nabla \times \underline{r} - \underline{\omega}_b \times \underline{r}^b) \quad (6)$$

Now eliminate the 5 indices:

$$\begin{aligned} \underline{\omega}_b \times \underline{r}^b &= r \underline{\omega}_b \times \underline{r}^b \\ &= \underline{\omega}_b \times (r \underline{r}^b) = \underline{\omega} \times \underline{r} \end{aligned} \quad (7)$$

$$2) \text{ So: } \frac{1}{c} \underline{A} = -\frac{m}{e} (\underline{\nabla} \times \underline{c} - \underline{\omega} \times \underline{c}) \quad -(8)$$

The Coulomb law for  $\underline{E}$  is

$$\underline{\nabla} \cdot \underline{E} = \frac{f_{E_0 c}}{c \mu_0} = \frac{e c \mu_0 \rho}{2m} \quad -(9)$$

where:  $f_{E_0 c} = 1/c^2 \quad -(10)$

using differentiation:

$$\underline{E} = -\frac{2}{c} \left( \frac{d}{dt} + \omega_{0c} \right) \underline{A} \quad -(11)$$

$$\underline{E} = \frac{2m}{e \mu_0} \left( \frac{d}{dt} + \omega_{0c} \right) \left( \underline{\nabla} \times \underline{c} - \underline{\omega} \times \underline{c} \right) \quad -(12)$$

From eqs. (9) and (12):

$$-\left( \frac{d}{dt} + \omega_{0c} \right) \underline{\nabla} \cdot \underline{\omega} \times \underline{c} = \frac{e c \mu_0 \rho}{2m}$$

because  $\underline{\nabla} \cdot \underline{\nabla} \times \underline{c} = 0 \quad -(14)$

$$\underline{\nabla} \cdot \underline{\omega} \times \underline{c} = \underline{\omega} \cdot \underline{\nabla} \times \underline{c} - \underline{c} \cdot \underline{\nabla} \times \underline{\omega} \quad -(15)$$

Now use:

$$\underline{\nabla} \cdot \underline{\omega} \times \underline{c} = \underline{\omega} \cdot \underline{\nabla} \times \underline{c} - \underline{c} \cdot \underline{\nabla} \times \underline{\omega}$$

to get:  $\frac{df}{dt} + \omega_{0c} f = \frac{e c \mu_0 \rho}{2m} \quad -(16)$

where:

where:

$$f = \underline{e} \cdot \nabla \times \underline{\zeta} = \underline{e} \cdot \nabla \times \underline{\omega} \quad - (17)$$

Differentiate eq. (16) & find:

$$\frac{d^2 f}{dt^2} + \omega_0 c \frac{df}{dt} + c \left( \frac{d\omega_0}{dt} \right) f = \frac{e \mu_0}{2m} \frac{dp}{dt}$$

- (18)

This is a Euler resonance equation in the variable  $f$ . If it is assumed that  $\underline{\omega}$  is a fixed property of space, then at resonance  $\nabla \times \underline{\zeta}$  becomes very large.

(Conclusion)

The spin electric field  $\underline{E}$  is much smaller than the orbital (usual) electric field  $\underline{E}$  and is defined by eq. (15) in terms of the position  $\underline{\zeta}$  of the electron or ion. The SCL equation of the Coulomb law is eq. (18).

# 144(13) : Orbital and Spin Magnetic Fields

## Orbital Magnetic Field

This is defined by:

$$\underline{B} = \nabla \times \underline{A} - \omega \times \underline{A} \quad - (1)$$

$$\underline{A} = -\frac{2m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \underline{r} \quad - (2)$$

$$\underline{A} = -\frac{2m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \underline{r} \quad - (2)$$

$$\underline{B}(\text{orbital}) = -\frac{2m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \nabla \times \underline{r} \quad - (3)$$

$$+ \frac{2m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \omega \times \underline{r}$$

$$\boxed{\underline{B}(\text{orbital}) = \frac{2m}{e} \left( \frac{\partial}{\partial t} + \omega_0 c \right) (\omega \times \underline{r} - \nabla \times \underline{r})} \quad - (4)$$

Assuming the existence of a magnetic monopole  $\rho_m$

$$\text{then } \nabla \cdot \underline{B}(\text{orbital}) = \rho_m \quad - (5)$$

From eqns. (4) and (5)

$$\boxed{\left( \frac{\partial}{\partial t} + \omega_0 c \right) \nabla \cdot (\omega \times \underline{r}) = \frac{e}{2m} \rho_m} \quad - (6)$$

because

$$\nabla \cdot \nabla \times \underline{r} = 0 \quad - (7)$$

If there is no magnetic monopole :

$$\left( \frac{d}{dt} + \omega_0 c \right) \underline{\Sigma} \cdot (\underline{\omega} \times \underline{\Sigma}) = 0 \quad -(8)$$

Eqs. (6) and (8) can be solved by computer algebra for various models.

### Spiral Magnetic Field

$$\left( \frac{\underline{B}}{c} \right) (\text{spin}) = \frac{1}{c} \left( \underline{\Sigma} \times \underline{A} - \underline{\omega} \times \underline{A} \right) \quad -(9)$$

$$\frac{\underline{A}}{c} = \frac{m}{e} \left( \underline{\Sigma} \times \underline{\Sigma} - \underline{\omega} \times \underline{\omega} \right) \quad -(10)$$

In eqs. (9) and (10), positive signs are used. The constant positive signs are used in eqs. (9) and (10). The spin field is  $c$  times smaller than the orbital field.

So:

$$\left( \frac{\underline{B}}{c} \right) (\text{spin}) = \frac{m}{e} \left( \underline{\Sigma} \times (\underline{\Sigma} \times \underline{\Sigma}) - \underline{\omega} \times \underline{\Sigma} \right) - \underline{\omega} \times (\underline{\Sigma} \times \underline{\Sigma}) \quad -(11)$$

If it is assumed that there is a magnetic monopole then:

$$\underline{\Sigma} \cdot \left( \frac{\underline{B}}{c} (\text{spin}) \right) = \rho_m / c \quad -(12)$$

3)

Now use :

$$\nabla \cdot \nabla \times (\nabla \times \underline{r}) = 0 \quad - (13)$$

$$\nabla \cdot \nabla \times (\underline{\omega} \times \underline{r}) = 0 \quad - (14)$$

So:

$$\boxed{\nabla \cdot (\underline{\omega} \times (\underline{\omega} \times \underline{r}) - \underline{\omega} \times (\nabla \times \underline{r})) = \frac{\rho_m e}{c m}} \quad - (15)$$

Eq. (15) can be solved with various models.

Some Vecto Ident. Vis

$$\nabla \times (\nabla \times \underline{r}) = -\nabla^2 \underline{r} + \nabla (\nabla \cdot \underline{r}) \quad - (16)$$

$$\nabla \times (\underline{\omega} \times \underline{r}) = \underline{\omega} \nabla \cdot \underline{r} - \underline{r} \nabla \cdot \underline{\omega} + (\underline{r} \cdot \nabla) \underline{\omega} - (\underline{\omega} \cdot \nabla) \underline{r} \quad - (17)$$

$$\underline{\omega} \times (\nabla \times \underline{r}) = (\underline{\omega} \times \nabla) \times \underline{r} + \underline{\omega} \nabla \cdot \underline{r} - (\underline{\omega} \cdot \nabla) \underline{r} \quad - (18)$$

$$\underline{\omega} \times (\underline{\omega} \times \underline{r}) = (\underline{r} \cdot \underline{\omega}) \underline{\omega} - (\underline{\omega} \cdot \underline{\omega}) \underline{r} \quad - (19)$$

In the limit of vanishing spin connection,

44(14) : Summary of Results t. Date  
The results of paper 144 t. date are as follows:-

$$\underline{E}(\text{orbital}) = -\frac{4m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) \left( \frac{d}{dt} + \omega_{oc} \right) \Sigma - (1)$$

$$\underline{B}(\text{orbital}) = -\frac{2m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) (\nabla \times \underline{\Sigma} - \underline{\Omega} \times \underline{\Sigma}) - (2)$$

$$\underline{E}(\text{spin}) = \frac{2m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) (\nabla \times \underline{\Sigma} - \underline{\Omega} \times \underline{\Sigma}) - (3)$$

$$\underline{B}(\text{spin}) = \frac{m}{e} (\underline{\Sigma} - \underline{\Omega}) \times (\nabla \times \underline{\Sigma} - \underline{\Omega} \times \underline{\Sigma}) - (4)$$

It is seen that the electric spin field and the magnetic orbital field have the same characteristics.

The magnetic orbital field is conventionally therefore the magnetic field.

observed, is a spin electric field.  
The completely new entity, the spin magnetic field, which is  $c^2$  times smaller than the orbital electric field, i.e.  $9 \times 10^{16}$  times smaller.

The analysis may now be extended using the Faraday law of induction and Ampere's Maxwell laws.

144(15) : Faraday Law of Induction for Orbital and Spin Fields.

The Faraday Law of induction from ECE theory  
in the assumed absence of a magnetic monopole is:

$$\underline{\nabla} \times \underline{E} + \frac{d\underline{B}}{dt} = \underline{0} \quad - (1)$$

Orbital Law

$$\underline{E} = \frac{4m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) \left( \frac{d}{dt} + \omega_{oc} \right) \underline{\Sigma} \quad - (2)$$

$$\underline{B} = \frac{2m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) \left( \underline{\omega} \times \underline{r} - \underline{\nabla} \times \underline{\Sigma} \right) \quad - (3)$$

So the law is:

$$\left( \frac{d}{dt} + \omega_{oc} \right) \left( \frac{d}{dt} \left( \underline{\nabla} \times \underline{\Sigma} + \underline{\omega} \times \underline{r} \right) + 2\omega_{oc} \right) = \underline{0} \quad - (4)$$

For a model of the spin conversion

$$\omega_m = (\omega_o, \underline{\omega}) \quad - (5)$$

eq. (4) can be solved for the electric trajectory  $\underline{\Sigma}$ .

Spin Law

$$\frac{\underline{E}}{c} = \frac{2m}{e} \left( \frac{d}{dt} + \omega_{oc} \right) \left( \underline{\omega} \times \underline{r} - \underline{\nabla} \times \underline{\Sigma} \right) \quad - (6)$$

$$\frac{\underline{B}}{c} = \frac{m}{e} \left( \underline{\nabla} - \underline{\omega} \right) \times \left( \underline{\nabla} \times \underline{\Sigma} - \underline{\omega} \times \underline{\Sigma} \right) \quad - (7)$$

2) The law "therefore":

$$\frac{d}{dt} (\underline{\Sigma} + \underline{\Omega}) \times (\underline{\Sigma} \times \underline{\Sigma} - \underline{\Omega} \times \underline{\Sigma}) = \underline{0} \\ + 2\omega_0 c (\underline{\Sigma} \times \underline{\Sigma} - \underline{\Omega} \times \underline{\Sigma}) \quad - (8)$$

It is seen that if the same rule of the spin connection is used in eqs. (4) and (8), the trajectory of the electron is different. This means that the orbital law and spin law have a different physical meaning.

Limit of zero Spin Connection

The orbital law becomes:

$$\frac{d^2}{dt^2} (\underline{\Sigma} \times \underline{\Sigma}) = \underline{0} \quad - (9)$$

and the spin law becomes:

$$\frac{d}{dt} (\underline{\Sigma} \times (\underline{\Sigma} \times \underline{\Sigma})) = \underline{0} \quad - (10)$$

This means that in the limit of zero spin connection all fields vanish, because in this limit the torsion vanishes.

# H44(16): Ampere Maxwell Law for O.S.Tal and S.p.i Fields

## O.S.Tal Law

$$\left( \frac{\partial}{\partial t} + \omega_0 c \right) \left( \nabla \times (\underline{E} \times \underline{C} - \underline{E} \times \underline{C}) - \frac{2}{c^2} \frac{\partial}{\partial t} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \underline{C} \right) = \frac{\epsilon \mu_0}{2m} \underline{I} \quad - (1)$$

$\underline{E} = \underline{C}/c$  वाले द्वारा दर्शाया गया है कि इसकी विशेषता यह है कि यह एक स्थिर तापमान के बराबर है। इसकी विशेषता यह है कि यह एक स्थिर तापमान के बराबर है। इसकी विशेषता यह है कि यह एक स्थिर तापमान के बराबर है।

## S.p.i Law

$$\left( \nabla \times (\underline{E} - \underline{C}) + \frac{2}{c^2} \frac{\partial}{\partial t} \left( \frac{\partial}{\partial t} + \omega_0 c \right) \right) \left( \underline{E} \times \underline{C} - \underline{E} \times \underline{C} \right) = \frac{\epsilon \mu_0}{2m c} \underline{I} \quad - (2)$$

$$E = E^x - E^y \quad I^x = 0$$

$\omega_m \rightarrow 0 \quad - (3)$

$$M = 1/(k^2 + q^2)^{1/2}$$

eq. (1) becomes:

$$\frac{\partial}{\partial t} \left( \nabla \times (\underline{E} \times \underline{C}) + \frac{2}{c^2} \frac{\partial^2 \underline{C}}{\partial t^2} \right) = - \frac{\epsilon \mu_0}{2m} \underline{I} \quad - (4)$$

and eq. (2) becomes:

$$\nabla \times \left( \nabla \times (\underline{E} \times \underline{C}) + \frac{2}{c^2} \frac{\partial^2 \underline{C}}{\partial t^2} \right) \underline{E} \times \underline{C} = \frac{\epsilon \mu_0}{2m c} \underline{I} \quad - (5)$$

using the identity

$$2) \quad \nabla \times (\nabla \times \underline{\Sigma}) = -\nabla^2 \underline{\Sigma} + \nabla (\nabla \cdot \underline{\Sigma}) - (6)$$

and assuming:

$$\nabla \cdot \underline{\Sigma} = 0 \quad - (7)$$

then eq. (4) becomes:

$$\frac{\partial}{\partial t} \left( -\nabla^2 \underline{\Sigma} + \frac{2}{c^2} \frac{\partial^2 \underline{\Sigma}}{\partial t^2} \right) = -\frac{e \mu_0}{2 \pi n} \underline{J}$$

but we will ignore the right hand side which is considered to be negligible.

consequently we can ignore the term  $\frac{\partial}{\partial t} \left( -\nabla^2 \underline{\Sigma} \right)$ .

so we have  $\frac{\partial}{\partial t} \left( \frac{2}{c^2} \frac{\partial^2 \underline{\Sigma}}{\partial t^2} \right) = -\frac{e \mu_0}{2 \pi n} \underline{J}$

which is the same as eq. (8).

Integrating both sides with respect to time we get  $\frac{2}{c^2} \frac{\partial \underline{\Sigma}}{\partial t} = -\frac{e \mu_0}{2 \pi n} \underline{J} t + C$

where  $C$  is a constant of integration. This is the same as eq. (5).

Integrating again with respect to time we get  $\underline{\Sigma} = -\frac{e \mu_0}{4 \pi n c^2} \underline{J} t^2 + C'$

where  $C'$  is a constant of integration. This is the same as eq. (6).

Note that the current appearing in (5) is  $\underline{J} / c$ , which is smaller than the current  $\underline{J}$  of eq. (1).

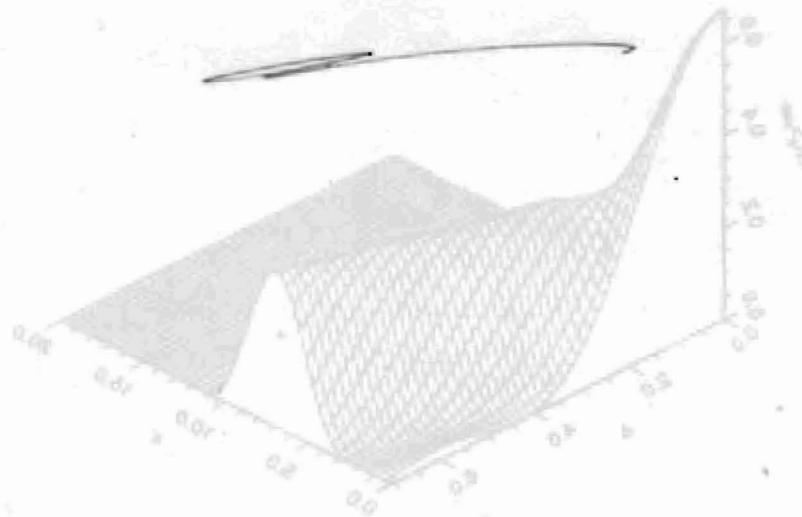


Figure 1: A periodic wave function  $\sin(\omega t)$  shown over a unit square domain. All except one point in a diagonal band around  $t=0.5 \pm 0.1$  lie in regions where the field is constant ( $\omega = 0$ ).

$$4) \quad \underline{B}(\text{orbital}) \rightarrow -\frac{2m}{e} \underline{\nabla} \times \frac{\underline{dr}}{dt} \quad -(20)$$

$$\frac{\underline{B}}{c}(\text{Spin}) \rightarrow \frac{m}{e} \underline{\nabla} \times (\underline{\nabla} \times \underline{\zeta}) \quad -(21)$$

$$= \frac{m}{e} \left( -\nabla^2 \underline{\zeta} + \underline{\nabla} (\underline{\zeta} \cdot \underline{\zeta}) \right) \quad -(22)$$

### Computer Modelling

The characteristics of the orbital magnetic field (4) and the spin magnetic field (11) can be worked out by computer algebra and graphed for values of  $\underline{r}$ ,  $\omega_0$  and  $\underline{\omega}$ . The spin magnetic field can be written as:

$$\boxed{\left( \frac{\underline{B}}{c} \right)(\text{Spin}) = \frac{m}{e} \left( (\underline{\nabla} - \underline{\omega}) \times (\underline{\nabla} \times \underline{\zeta} - \underline{\omega} \times \underline{\zeta}) \right)} \quad -(23)$$

and the orbital as:

$$\boxed{\underline{B}(\text{orbital}) = -\frac{2m}{e} \left( \frac{d}{dt} + \omega_0 c \right) \left( \underline{\nabla} \times \underline{\zeta} - \underline{\omega} \times \underline{\zeta} \right)} \quad -(24)$$

