

1) 128(1) : Metric Development of E(E) and Derivation of the Engineering Model

The metric is conventionally defined as :

$$x_\mu = g_{\mu\nu} x^\nu \quad - (1)$$

and so :  $x^\mu = g^{\mu\nu} x_\nu \quad - (2)$

where :  $g^{\mu\nu} = g^{\mu a} g_{a\nu} \quad - (3)$

In general :  $\nabla^\mu = g^{\mu\nu} \nabla_\nu \quad - (4)$

where  $g^{\mu\nu}$  is a tetrad. These equations are defined in the general spacetime of any dimension. If the tetrad is defined by :

$$\nabla^a = g^a_\mu \nabla^\mu \quad - (5)$$

The first Cartan structure equation is :

$$T^\alpha_{\mu\nu} = \partial_\mu g^\alpha_\nu - \partial_\nu g^\alpha_\mu + \omega^\alpha_{\mu b} g^\beta_\nu - \omega^\alpha_{\nu b} g^\beta_\mu \quad - (6)$$

E.g. (2) implies that it is possible to define a torsion tensor through the metric  $g^{\mu\nu}$ :

$$T^\lambda_{\mu\nu} = \partial_\mu g^\lambda_\nu - \partial_\nu g^\lambda_\mu + \Gamma^\lambda_{\mu k} g^\kappa_\nu - \Gamma^\lambda_{\nu k} g^\kappa_\mu \quad - (7)$$

and so the field density is :

$$2) \boxed{F_{\mu\nu}^{\lambda} = \partial_{\mu}A_{\nu}^{\lambda} - \partial_{\nu}A_{\mu}^{\lambda} + \Gamma_{\mu k}^{\lambda}A_{\nu}^k - \Gamma_{\nu k}^{\lambda}A_{\mu}^k} \quad -(8)$$

This means that the potential density of the electric field is directly proportional to the metric:

$$A_{\mu}^{\lambda} = A^{(0)} g_{\mu}^{\lambda} \quad -(9)$$

$$\boxed{A_{\mu}^{\lambda} = A^{(0)} g^{\lambda \nu} g_{\nu \mu}} \quad -(10)$$

In field theory (Ryder, 1996) the angular momentum is obtained from the angular energy / momentum density by an integration over:

$$\lambda = 0 \quad -(11)$$

$$J_{\mu\nu} = \int J_{\mu\nu} dV \quad -(12)$$

The eln field density is directly proportional to the angular momentum density as shown in paper 127, so:

$$F_{\mu\nu} = \int F_{\mu\nu} dV \quad -(13)$$

Take:

$$F_{\mu\nu}^{\circ} = \partial_{\mu}A_{\nu}^{\circ} - \partial_{\nu}A_{\mu}^{\circ} + \Gamma_{\mu k}^{\circ}A_{\nu}^k - \Gamma_{\nu k}^{\circ}A_{\mu}^k \quad -(14)$$

3) Therefore:

$$F_{\mu\nu} = \int (\partial_\mu A_\nu - \partial_\nu A_\mu + \Gamma_{\mu\nu}^\sigma A_\sigma - \Gamma_{\nu\mu}^\sigma A_\sigma) dV \quad - (15)$$

expresses the electromagnetic field in terms of the metric. In eq. (15)

$$\partial_\mu A_\nu = \int \partial_\mu A_\nu dV \quad - (16)$$

$$\Gamma_{\mu\nu}^\sigma A_\sigma = \int \Gamma_{\mu\nu}^\sigma A_\sigma dV \quad - (17)$$

and so on. Therefore:

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + \Gamma_{\mu\nu}^\sigma A_\sigma - \Gamma_{\nu\mu}^\sigma A_\sigma$$

Diagonal Metric

If the metric is diagonal, then:

$$A^\circ = A^{(0)} g^{00} = A^{(0)} \quad - (19)$$

$$A^1 = A^{(0)} g^{11} = A^{(0)} \quad - (20)$$

Similarly:  $A^2 = A^{(0)} g^{22} = A^{(0)}$

few other terms such as:  $(g^{00} g_{01} + \dots + g^{03} g_{31})$  - (21).

$$A^0_1 = A^{(0)} (g^{00} g_{01} + \dots + g^{03} g_{31})$$

which exist if there are off diagonal elements of the metric.

The engineering model is derived from eq. (18).

128(2) : New General Condition for any Metric  
 Start with the definition of the metric + tetrad:

$$V^a = \varphi^a_\mu V^\mu - (1)$$

A particular case of this is :

$$x^k = g^k_\mu x^\mu - (2)$$

where  $g^k_\mu$  is the metric. Consider:

$$D_\mu V^\nu = \partial_\mu V^\nu + \Gamma^\nu_{\mu\lambda} V^\lambda - (3)$$

and:

$$D_\mu V^a = \partial_\mu V^a + \omega^a_{\mu b} V^b - (4)$$

Eqs. (3) and (4) imply the tetrad postulate:

$$\partial_\mu \varphi^a_\nu + \omega^a_{\mu b} \varphi^b_\nu - \Gamma^\lambda_{\mu\nu} \varphi^a_\lambda = 0, - (5)$$

i.e.  $\Gamma^\lambda_{\mu\nu} = \varphi^{\lambda}_a (\partial_\mu \varphi^a_\nu + \omega^a_{\mu b} \varphi^b_\nu) - (6)$

using  $\varphi^a_\lambda \varphi^d_a = \delta^d_\lambda. - (7)$

The special case of eq. (5) implies that  
 eq. (4) becomes:

$$\partial_\mu g^k_\nu + \Gamma^k_{\mu\lambda} g^\lambda_\nu - \Gamma^\lambda_{\mu\nu} g^k_\lambda = 0 - (8)$$

i.e. a is replaced by k, b by λ, and ω by

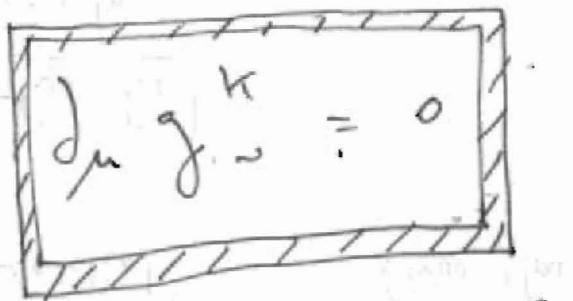
Γ. In eq. (8) :

$$\Gamma^k_{\mu\lambda} g^\lambda_\nu = \Gamma^k_{\mu\nu} - (9)$$

and

$$\Gamma^\lambda_{\mu\nu} g^k_\lambda = \Gamma^k_{\mu\nu} - (10)$$

2) Therefore:



$$\partial_\mu g_{\alpha\beta}^K = 0 \quad - (11)$$

This is an important new fundamental equation for the metric:

$$g_{\alpha\beta}^K = g^{\alpha\mu} g_{\mu\beta}^K \quad - (12)$$

Any Riemannian metric obeys eq. (11), and in general any metric in any spacetime of any dimension, in general a spacetime with torsion and curvature.

### Diagonal Metric

In this case off-diagonals are zero, so eq. (11) produces:

$$\partial_\mu (g^{00} g_{00}) = 0 \quad - (13)$$

$$\partial_\mu (g^{11} g_{11}) = 0 \quad - (14)$$

$$\partial_\mu (g^{22} g_{22}) = 0 \quad - (15)$$

$$\partial_\mu (g^{33} g_{33}) = 0 \quad - (16)$$

$$\partial_\mu (g^{00} g_{00}) = 0 \quad - (17)$$

Eqs. (13) - (16) are true

for any  $\mu$ .

because  $g^{00} g_{00} = g^{11} g_{11} = g^{22} g_{22} = g^{33} g_{33} = 1$  - (17)

3) In general, in four dimensions:

$$\partial_\mu (g^{Kd} g_{d\mu}) = 0 \quad (18)$$

where:

$$g^{Kd} g_{d\mu} = g^{K0} g_{0\mu} + g^{K1} g_{1\mu} \quad (19)$$

$$+ g^{K2} g_{2\mu} + g^{K3} g_{3\mu}$$

and where the metric has diagonal and off-diagonal elements.

The review of the abitual theorem of paper III says eq. (11) because it is a diagonal metric.

### Computer Test

It is possible now to test metrics with off-diagonal elements by using computer algebra with eq. (18).

1) Tetrad Postulate is the Base Manifold (Note 128(3)) .

The derivative of Cartan geometry's tetrad postulate is based on the use of a base manifold labelled  $\mu$  and a tangent spacetime labelled  $a$ . The derivation of the postulate is given in proof tree and is reviewed here before specializing to the case of one manifold. The covariant derivative of Riemann geometry is:

$$D_\mu V^\lambda = \partial_\mu V^\lambda + \Gamma_{\mu\lambda}^\lambda V^\lambda \quad (1)$$

and is defined in the base manifold. The tetrad is defined by:

$$V^a = g_{\mu}^a V^\mu \quad (2)$$

In the tangent spacetime the covariant derivative is defined by:

$$D_\mu V^a = \partial_\mu V^a + \omega_{\mu b}^a V^b \quad (3)$$

The basis elements for this tangent space (a Minkowski space) are  $\{e_a\}$ , and the basis elements of the base manifold are  $\{q_\mu\}$ . The advantage of using the tangent space is that  $e_a$  can be Pauli matrices, for example, so the Dirac equation and spinors can be considered. Therefore,  $\omega_{\mu b}^a$  is known as the spin connection.

The complete vector field is the same for eqn. (1) &

(3):

$$DV = D_\mu V^\lambda dx^\mu \otimes e_\lambda = D_\mu V^a dx^\mu \otimes \hat{V}_a \quad (4)$$

The vector components and basis elements are related by

equations similar to eqn. (2):

$$\hat{V}_a = q_a^\sigma d\sigma, \quad (5)$$

$$V^a = q_a^\sigma V^\sigma. \quad (6)$$

2. Therefore in eq. (4) :

$$DV = \left( \partial_\mu (\sqrt{a} V^\sigma) + \omega_{\mu b}^a \sqrt{a} V^\lambda \right) dx^\mu \otimes (\sqrt{a} d\sigma) - (7)$$

$$= \sqrt{a} \left( \partial_\mu (\sqrt{a} V^\sigma) + \omega_{\mu b}^a \sqrt{a} V^\lambda \right) dx^\mu \otimes d\sigma. - (8)$$

The Cartan convention is :

$$\sqrt{a} \sqrt{a} = \delta_a^a - (9)$$

where:  $\delta_a^a = 1 \text{ if } a = a, - (10)$

$$\delta_a^a = 0 \text{ if } a \neq a. - (11)$$

Eq. (8) is :

$$DV = \sqrt{a} \sqrt{a} \partial_\mu V^\sigma dx^\mu \otimes d\sigma + \dots - (12)$$

$$= \delta_a^a \partial_\mu V^\sigma dx^\mu \otimes d\sigma + \dots - (13)$$

$$= \delta_0^0 \partial_\mu V^0 dx^\mu \otimes d_0 + \delta_1^1 \partial_\mu V^1 dx^\mu \otimes d_1$$

$$+ \delta_2^2 \partial_\mu V^2 dx^\mu \otimes d_2 + \delta_3^3 \partial_\mu V^3 dx^\mu \otimes d_3 + \dots - (14)$$

$$= \delta_0^0 \partial_\mu V^0 dx^\mu \otimes d_0 + \delta_1^1 \partial_\mu V^1 dx^\mu \otimes d_1$$

$$+ \delta_2^2 \partial_\mu V^2 dx^\mu \otimes d_2 + \delta_3^3 \partial_\mu V^3 dx^\mu \otimes d_3 + \dots - (15)$$

$$= \partial_\mu V^0 dx^\mu \otimes d_0 + \partial_\mu V^1 dx^\mu \otimes d_1$$

$$+ \partial_\mu V^2 dx^\mu \otimes d_2 + \partial_\mu V^3 dx^\mu \otimes d_3 + \dots - (16)$$

$$= \partial_\mu V^\sigma dx^\mu \otimes d_\sigma + \dots - (17)$$

So:

$$DV = \left( \partial_\mu V^\sigma + \sqrt{a} \left( \partial_\mu \sqrt{a} + \omega_{\mu b}^a \sqrt{a} \right) V^\lambda \right) dx^\mu \otimes d_\sigma - (18)$$

From eqs. (1), (4) and (18):

$$\tilde{\Gamma}_{\mu\lambda}^a = \tilde{g}^a_b (\partial_\mu \tilde{g}^\lambda_b + \omega_{\mu b}^a \tilde{g}^\lambda_b) - (19)$$

The covetia can therefore be expanded in terms of the tetrad and the spin connection. As can be seen from eq. (4) this is a basic property of the complete vector field DV, and so is a very fundamental result.

Multiply both sides of eq. (19) by  $\tilde{g}^\lambda_a$  and we:

$$\tilde{g}^\lambda_a \tilde{g}^\mu_b = 1 - (20)$$

To state the tetrad postulate:

$$\begin{aligned} \partial_\mu \tilde{g}^\lambda_a &= \partial_\mu \tilde{g}^\lambda_a + \omega_{\mu b}^a \tilde{g}^\lambda_b - \tilde{\Gamma}_{\mu\lambda}^a \tilde{g}^\lambda_a - (21) \\ &= 0 \end{aligned}$$

Eq. (21) was the rule for the covariant derivative of the mixed index rank two tensor  $\tilde{g}^\lambda_a$ .

Tetrad Postulate in the base manifold.

Consider the basic definition (2) in the base manifold, when:

$$a = \nu. - (22)$$

Then:  $\nabla^\nu = \tilde{g}_\mu^\nu \nabla^\mu - (23)$

where  $\nabla^\nu$  and  $\nabla^\mu$  are any vectors. In the special

case :

$$\nabla^{\mu} = x^{\mu}, \quad \nabla^{\sim} = x^{\sim} - (24)$$

then:

$$x^{\sim} = \sqrt{g_{\mu\nu}} x^{\mu} - (25)$$

However, the metric is defined by:

$$x_{\mu} = g_{\mu\nu} x^{\nu}, - (26)$$

so:

$$\boxed{x^{\sim} = g^{\sim}_{\mu} x^{\mu}}, - (27)$$

where:

$$g^{\sim}_{\mu} = g^{\sim\alpha} \gamma_{\alpha\mu} - (28)$$

$$= g^{\sim\alpha} g_{\alpha\mu} + \dots + g^{\sim 3} \gamma_{3\mu} - (29)$$

In the case of eq. (22), eq. (3) becomes:

$$\partial_{\mu} \nabla^{\sim} = \partial_{\mu} \nabla^{\sim} + \omega^{\sim}_{\mu b} \nabla^b. - (30)$$

However, it is known that:

$$\partial_{\mu} \nabla^{\sim} = \partial_{\mu} \nabla^{\sim} + \Gamma^{\sim}_{\mu\lambda} \nabla^{\lambda} - (31)$$

so:  $\omega^{\sim}_{\mu b} \nabla^b = \Gamma^{\sim}_{\mu\lambda} \nabla^{\lambda} - (32)$

Therefore:  $b = \lambda. - (33)$

The tetrad postulate (21) becomes:

$$\partial_{\mu} g^{\sim}_{\nu} = \partial_{\mu} g^{\sim}_{\nu} + \Gamma^{\sim}_{\mu\lambda} g^{\sim}_{\nu} - \Gamma^{\sim}_{\nu\lambda} g^{\sim}_{\mu} = 0$$

- (34)

5.) It is seen that:

$$\nabla_\mu \tilde{g}_{\nu} = \tilde{g}_{\mu\nu} - (35)$$

In Cartesian geometry there are results such as:

$$T_{\mu\nu}^a = g_{\lambda}^a T_{\mu\nu}^\lambda - (36)$$

and  $\omega_{\mu b}^a = \omega_{\mu b}^\lambda g_{\lambda}^b - (37)$

Therefore using eqn. (35) :

$$\Gamma_{\mu\nu}^K = g_{\lambda}^\lambda \Gamma_{\mu\nu}^\lambda, - (38)$$

$$\Gamma_{\mu\nu}^K = g_{\lambda}^\lambda \Gamma_{\mu\nu}^\lambda. - (39)$$

So we obtain:

$$\boxed{\partial_\mu g_{\nu}^K = \partial_\nu g_{\mu}^K = 0} - (40)$$

This is a very fundamental result, valid  
for any Riemannian spacetime in any dimension.

Computer Tests

Test whether solution of the Einstein field  
equations obey:

$$\boxed{\partial_\mu (g^{\mu\lambda} g_{\lambda\nu}) = 0} - (41)$$

# Note 128(k) : Basic Hypotheses.

There are two fundamental hypotheses of ECE theory.

$$A_{\mu}^a = A^{(0)} g_{\mu}^a \quad - (1)$$

which implies:  $F_{\mu\nu}^a = A^{(0)} T_{\mu\nu}^a \quad - (2)$

and:  $J_{\mu\nu}^a = \frac{c}{k} T_{\mu\nu}^a \quad - (3)$

Here  $A_{\mu}^a$  is the electromagnetic potential density and  $J_{\mu\nu}^a$  the angular energy/momentum density.

and  $J_{\mu\nu}^a$  the angular energy/momentum density.

These can be expressed in the base manifold  $\mathcal{K}$ :

$$A_{\mu}^K = A^{(0)} g_{\mu}^K \quad - (4)$$

$$J_{\mu\nu}^K = \frac{c}{k} T_{\mu\nu}^K \quad - (5)$$

and  $J_{\mu\nu}^K = \frac{c}{k} T_{\mu\nu}^K \quad - (5)$

In the base manifold:

$$T_{\mu\nu}^K = \partial_{\mu} g_{\nu}^K - \partial_{\nu} g_{\mu}^K + \Gamma_{\mu\lambda}^K g_{\nu}^{\lambda} - \Gamma_{\nu\lambda}^K g_{\mu}^{\lambda} \quad - (6)$$

$$= \Gamma_{\mu\nu}^K - \Gamma_{\nu\mu}^K$$

because:  $\partial_{\mu} g_{\nu}^K = \partial_{\nu} g_{\mu}^K = 0 \quad - (7)$

$$\Gamma_{\mu\nu}^K = \Gamma_{\mu\lambda}^K g_{\nu}^{\lambda} \quad - (8)$$

and  $\Gamma_{\nu\mu}^K = \Gamma_{\nu\lambda}^K g_{\mu}^{\lambda} \quad - (9)$

Therefore:

$$J_{\mu\nu}^K = \frac{c}{k} \left( \Gamma_{\mu\nu}^K - \Gamma_{\nu\mu}^K \right) \quad (10)$$

However:

$$J_{\mu\nu}^K = x_\mu p_\nu - x_\nu p_\mu \quad (11)$$

which is a generalization of:

$$J = c \times p \quad (12)$$

Here:

$$x_\mu = (ct, -\vec{x}) \quad (13)$$

and

$$p_\sim^K = p^{(0)} g_\sim^K \quad (14)$$

so:

$$J_{\mu\nu}^K = p^{(0)} \left( x_\mu g_\sim^K - x_\nu g_\mu^K \right) \quad (15)$$

and

$$\Gamma_{\mu\nu}^K = \frac{k p^{(0)}}{c} x_\mu g_\sim^K \quad (16)$$

The units of  $k p^{(0)} / c$  are  $1 / \text{Ar}$ , so:

$$\Gamma_{\mu\nu}^K = \frac{1}{\text{Ar}} x_\mu g_\sim^K \quad (17)$$

Eq. (14) defines energy momentum density  $p_\sim^K$  in terms of the metric  $g_\sim^K$ .

3) The electromagnetic field density is:

$$F_{\mu\nu}^K = \Gamma_{\mu\nu}^K A^\lambda - \Gamma_{\nu\lambda}^K A_\mu^\lambda - (18)$$
$$= A^{(0)} (\Gamma_{\mu\nu}^K - \Gamma_{\nu\mu}^K)$$

Using:

$$\Gamma_{\mu\nu}^K = -\Gamma_{\nu\mu}^K - (19)$$

then

$$\boxed{F_{\mu\nu}^K = 2A^{(0)}\Gamma_{\mu\nu}^K} - (20)$$

This result is consistent with:

$$F_{\mu\nu}^K = \sqrt{a} \Gamma_{\mu\nu}^K F^a$$
$$= A \sqrt{a} \left( \partial_\mu \sqrt{a} - \partial_\nu \sqrt{a} + \omega_{\mu b}^a \sqrt{a} - \omega_{\nu b}^a \sqrt{a} \right)$$
$$= A^{(0)} (\Gamma_{\mu\nu}^K - \Gamma_{\nu\mu}^K) - (21)$$

Therefore the analysis is self-consistent.

128(5): General Interpretation of the Tetrad in the Base Manifold.

In general:  $\tilde{\nabla}^\mu = g^\mu_\nu \nabla^\nu - (1)$

where all quantities are expressed in the base manifold. Here  $\nabla^\mu$  and  $\tilde{\nabla}^\mu$  are vectors and  $g^\mu_\nu$  is the tetrad. It is proven here that under all conditions:

$$\partial_\lambda g^\mu_\nu = 0. - (2)$$

Proof

The torsion is:  $- (3)$

$$\begin{aligned} T^\kappa_{\mu\nu} &= \Gamma^\kappa_{\mu\nu} - \Gamma^\kappa_{\nu\mu} \\ &= \partial_\mu g^\kappa_\nu - \partial_\nu g^\kappa_\mu + \Gamma^\kappa_{\mu\lambda} g^\lambda_\nu - \Gamma^\kappa_{\nu\lambda} g^\lambda_\mu. \end{aligned}$$

However:  $\Gamma^\kappa_{\mu\nu} = \Gamma^\kappa_{\mu\lambda} g^\lambda_\nu - (4)$

$$\Gamma^\kappa_{\nu\mu} = \Gamma^\kappa_{\nu\lambda} g^\lambda_\mu, - (5)$$

so:  $\partial_\mu g^\kappa_\nu - \partial_\nu g^\kappa_\mu = 0. - (6)$

It is known from the tetrad postulate that:

$$\partial_\mu g^\kappa_\nu = 0, - (7)$$

$$\partial_\nu g^\kappa_\mu = 0. - (8)$$

In the special case:

$$\tilde{\nabla}^\mu = x^\mu - (9)$$

then:

$$2) \quad \nabla^{\mu} = g^{\mu}_{\nu} - (10)$$

where the metric  $g^{\mu}_{\nu}$  is defined:

$$x^{\mu} = g^{\mu}_{\nu} x^{\nu} - (11)$$

### Cartesian Interpretation

An example of the meaning of eqs. (7) and (8) may be given in a Cartesian space, "which

for example:  $x^1 = x, x^2 = y - (12)$

In this case there is no connection:

$$\Gamma^{\mu}_{\mu\nu} = 0 - (13)$$

because the space is a flat space. So there is no

Cartesian. Therefore:

$$\partial_{\mu} \nabla^1 = 0 - (14)$$

$$\partial_{\mu} \nabla^2 = \nabla^1 \nabla^2 - (15)$$

then:

$$\nabla^1 = \nabla^1 \nabla^2 - (16)$$

Similarly:  $\partial_{\mu} g^{12} = 0 - (17)$

where  $x^1 = g^{12} x^2 - (18)$

with

$$g^{12} = g^{1d} g_{d2} - (19)$$

However:

$$g_{\mu\nu} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} - (20)$$

so

$$g^{12} = 0 - (21)$$

### 3) The Existence of Spin Conservation Rule (SCR)

SCR can exist if and only if:

$$\nabla^a = \nabla_\mu \nabla^\mu - (21)$$

and:

$$\nabla_\mu^a = \partial_\mu q^a - \partial_a q_\mu + \omega_{\mu b}^a q^b - \omega_{ab}^a q_\mu^b - (22)$$

This means that the representation space of  $\nabla^a$  must be distinct from that of  $\nabla^\mu$ . In this case,

for example:

$$\nabla^a = (ct, (1), (2), (3)), - (23)$$

$$\nabla^\mu = (ct, x, y, z) - (24)$$

where respectively, the complex circular and Cartesian representations are used for the space-like parts of  $\nabla^a$  and  $\nabla^\mu$ . In this case:

$$\partial_\mu q^a = 0, - (25)$$

$$\partial_\mu q^a \neq 0. - (26)$$

Therefore the tetrad may be interpreted as a Circular spinor. Carter is forced both to tetards and spinors. The latter are of course wed in the Dirac equation, where

4)

In basis elements are the Pauli matrices ( $SU(2)$ ).  
 In strong field theory the basis elements are the Gell-Mann matrices, ( $SU(3)$ ).

Three dimensional space may be represented by a basis set consisting either of the Cartesian unit vectors,  $\hat{i}$ ,  $\hat{j}$  and  $\hat{k}$ , or of the complex circular unit vectors,  $\hat{e}^{(1)}$ ,  $\hat{e}^{(2)}$  and  $\hat{e}^{(3)}$ . Both are  $O(3)$  representation spaces. The electromagnetic potential density may be represented by  $A_\mu^a$  and the electromagnetic field density by  $F_{\mu\nu}^a$ . It is this representation that gives rise to SCR. The latter does not arise from eqn. (1) because the latter leads to eqns. (7) and (8).

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