

Spin of Slater determinant

Spin operators

electron spin

$$\hat{s} = \hat{s}_x \mathbf{x} + \hat{s}_y \mathbf{y} + \hat{s}_z \mathbf{z} \quad \hat{s}^2 = \hat{s}_x^2 + \hat{s}_y^2 + \hat{s}_z^2$$

$$\hat{s}_x |\alpha\rangle = \frac{1}{2} |\beta\rangle \quad \hat{s}_y |\alpha\rangle = \frac{i}{2} |\beta\rangle \quad \hat{s}_z |\alpha\rangle = \frac{1}{2} |\alpha\rangle$$

$$\hat{s}_x |\beta\rangle = \frac{1}{2} |\alpha\rangle \quad \hat{s}_y |\beta\rangle = -\frac{i}{2} |\alpha\rangle \quad \hat{s}_z |\beta\rangle = -\frac{1}{2} |\beta\rangle$$

commutation relations

$$[\hat{s}_x, \hat{s}_y] = \hat{s}_x \hat{s}_y - \hat{s}_y \hat{s}_x = i \hat{s}_z$$

$$[\hat{s}_y, \hat{s}_z] = i \hat{s}_x$$

$$[\hat{s}_z, \hat{s}_x] = i \hat{s}_y$$

Spin of Slater determinant

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electron spin

$$\hat{s} = \hat{s}_x \mathbf{x} + \hat{s}_y \mathbf{y} + \hat{s}_z \mathbf{z} \quad \hat{s}^2 = \hat{s}_x^2 + \hat{s}_y^2 + \hat{s}_z^2$$

$$\hat{s}_x |\alpha\rangle = \frac{1}{2} |\beta\rangle \quad \hat{s}_y |\alpha\rangle = \frac{i}{2} |\beta\rangle \quad \hat{s}_z |\alpha\rangle = \frac{1}{2} |\alpha\rangle$$

$$\hat{s}_x |\beta\rangle = \frac{1}{2} |\alpha\rangle \quad \hat{s}_y |\beta\rangle = -\frac{i}{2} |\alpha\rangle \quad \hat{s}_z |\beta\rangle = -\frac{1}{2} |\beta\rangle$$

step operators

$$\hat{s}_+ = \hat{s}_x + i\hat{s}_y \quad \hat{s}_+ |\alpha\rangle = 0 \quad \hat{s}_+ |\beta\rangle = |\alpha\rangle$$

$$\hat{s}_- = \hat{s}_x - i\hat{s}_y \quad \hat{s}_- |\alpha\rangle = |\beta\rangle \quad \hat{s}_- |\beta\rangle = |\alpha\rangle$$

$$\hat{s}^2 = \hat{s}_+ \hat{s}_- - \hat{s}_z + \hat{s}_z^2$$

$$\hat{s}^2 = \hat{s}_- \hat{s}_+ + \hat{s}_z + \hat{s}_z^2$$

Spin of Slater determinant

many electron spin operators

$$\hat{S}^2 = \hat{S} \cdot \hat{S} = \sum_k^N \sum_l^N \hat{s}(k) \cdot \hat{s}(l) \quad \hat{S}_z = \sum_k^N \hat{s}_z(k)$$

$$\hat{S}_+ = \sum_k^N \hat{s}_+(k) \quad \hat{S}_- = \sum_k^N \hat{s}_-(k)$$

Spin eigenfunctions

$$\hat{S}_z |\phi_1 \phi_2 \dots \phi_3\rangle = \frac{1}{2}(N^\alpha - N^\beta) |\phi_1 \phi_2 \dots \phi_3\rangle$$

closed shell

$$\hat{S}^2 |\phi_1 \phi_2 \dots \phi_3\rangle = S(S+1) |\phi_1 \phi_2 \dots \phi_3\rangle$$

open shell?

configuration state functions

Spin of Slater determinant

minimal basis hydrogen molecule

restricted Hartree-Fock wave function

$$\psi_0(\mathbf{x}_1, \mathbf{x}_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{x}_1)\phi_2(\mathbf{x}_2) - \phi_2(\mathbf{x}_1)\phi_1(\mathbf{x}_2)]$$

$$\phi_1(\mathbf{x}) = \varphi_1(\mathbf{r})\alpha(s)$$

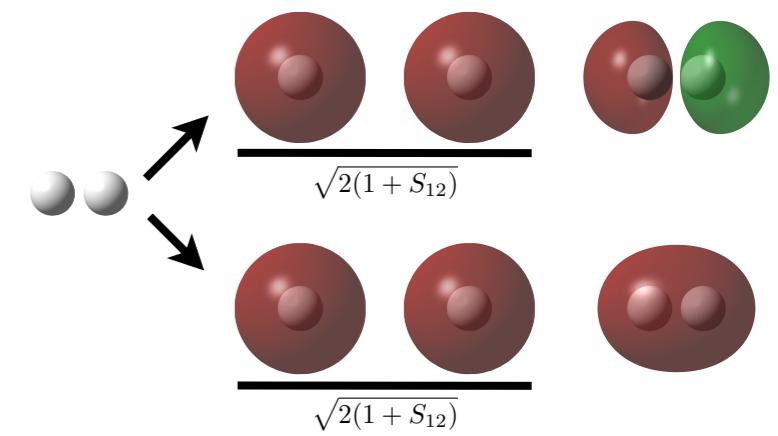
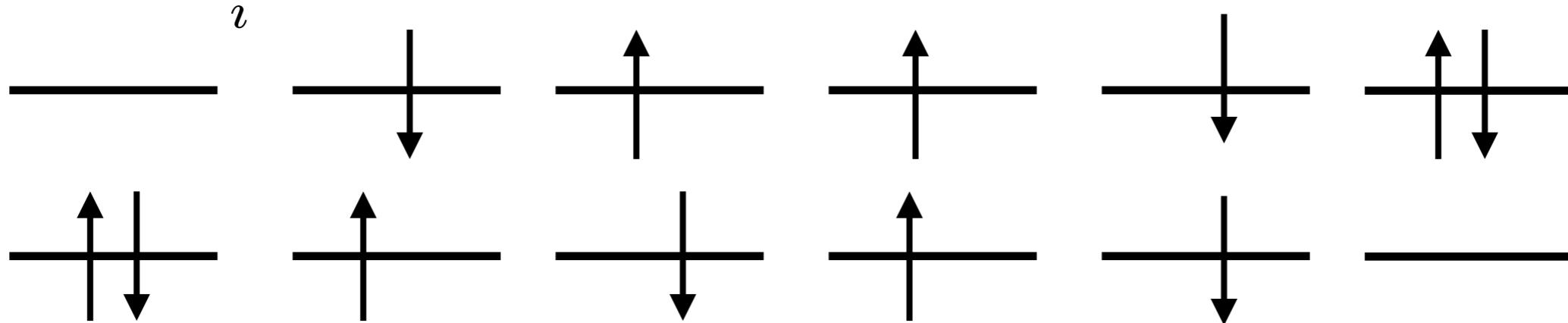
$$\phi_3(\mathbf{x}) = \varphi_2(\mathbf{r})\alpha(s)$$

$$\phi_2(\mathbf{x}) = \varphi_1(\mathbf{r})\beta(s)$$

$$\phi_4(\mathbf{x}) = \varphi_2(\mathbf{r})\beta(s)$$

configuration interaction

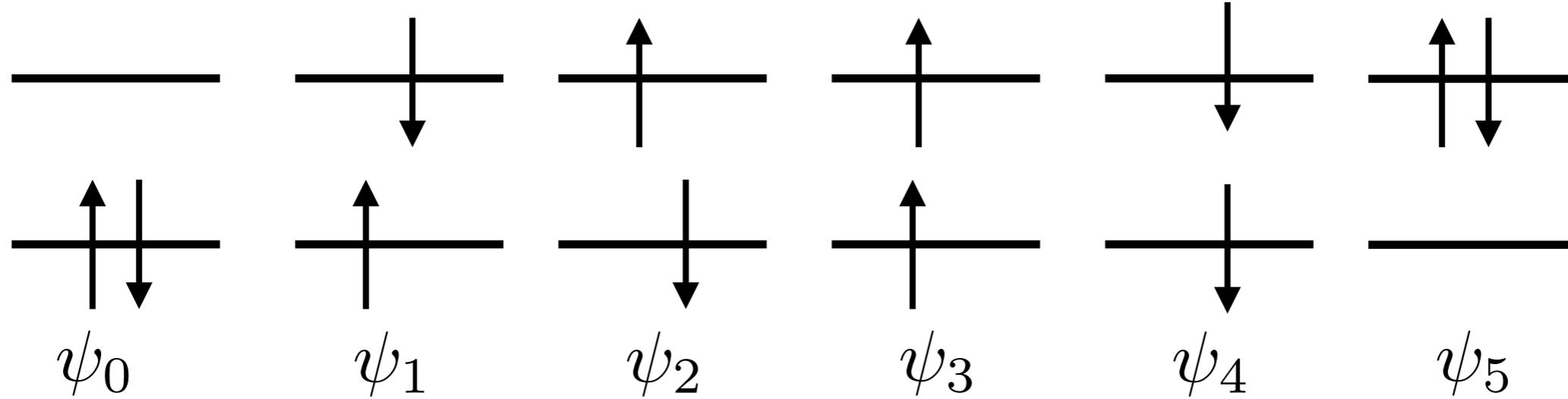
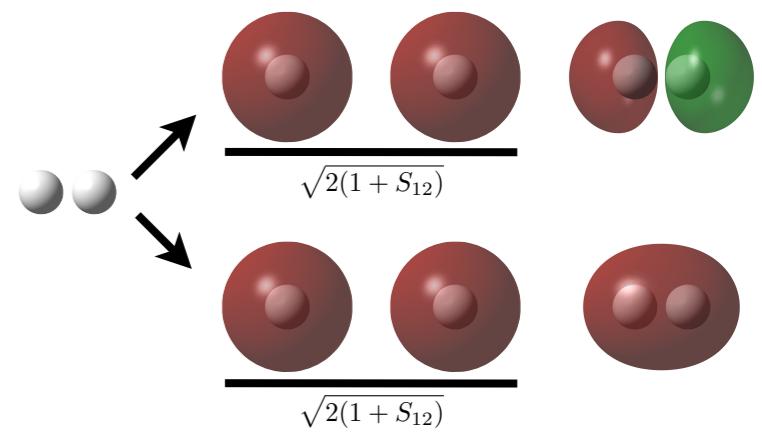
$$\Psi = \sum_i \psi_i \quad \binom{4}{2} = 6$$



Spin of Slater determinant

minimal basis hydrogen molecule

configurations



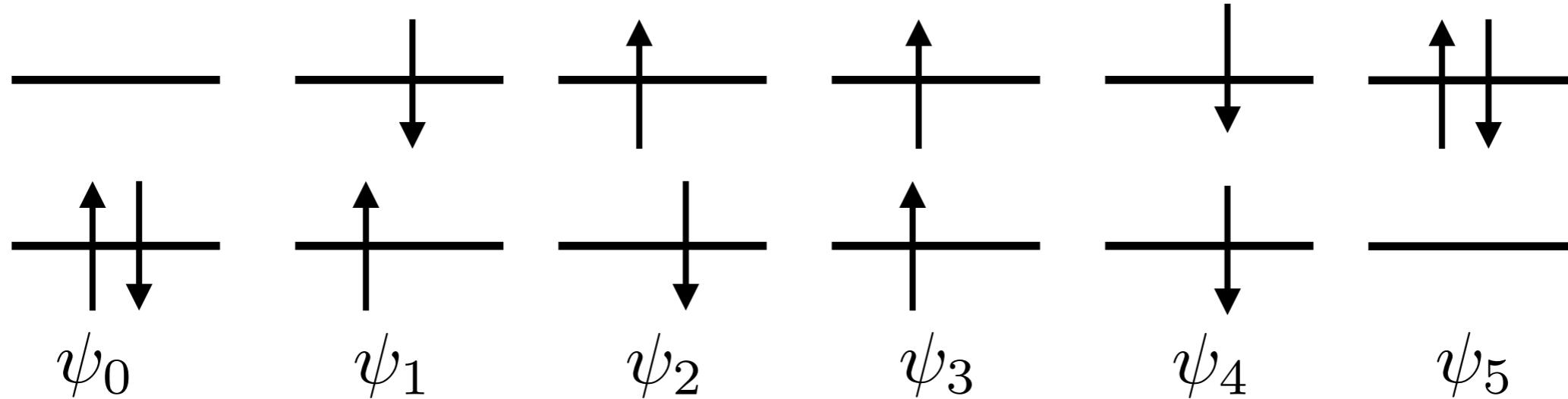
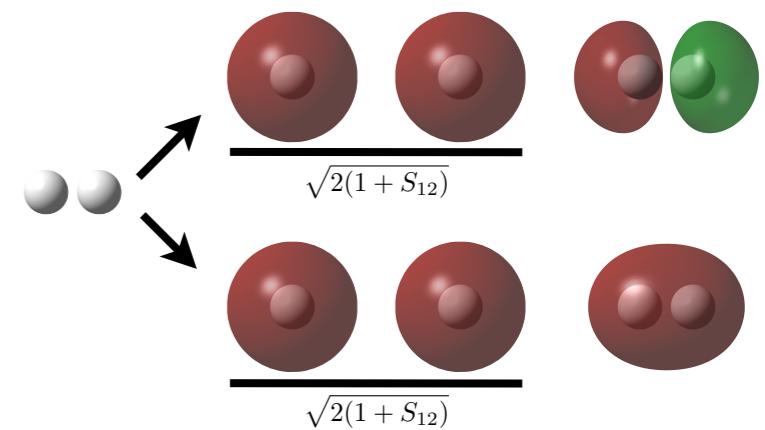
total spin

$$\begin{aligned} S^2 \psi_0 &= (S_- S_+ + S_z + S_z^2) \psi_0 = S(S+1) \psi_0 \\ &= \frac{1}{\sqrt{2}} [(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2)) \\ &\quad + s_z(1)(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2)) + s_z(2)(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2))] \\ &\quad + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2))] \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



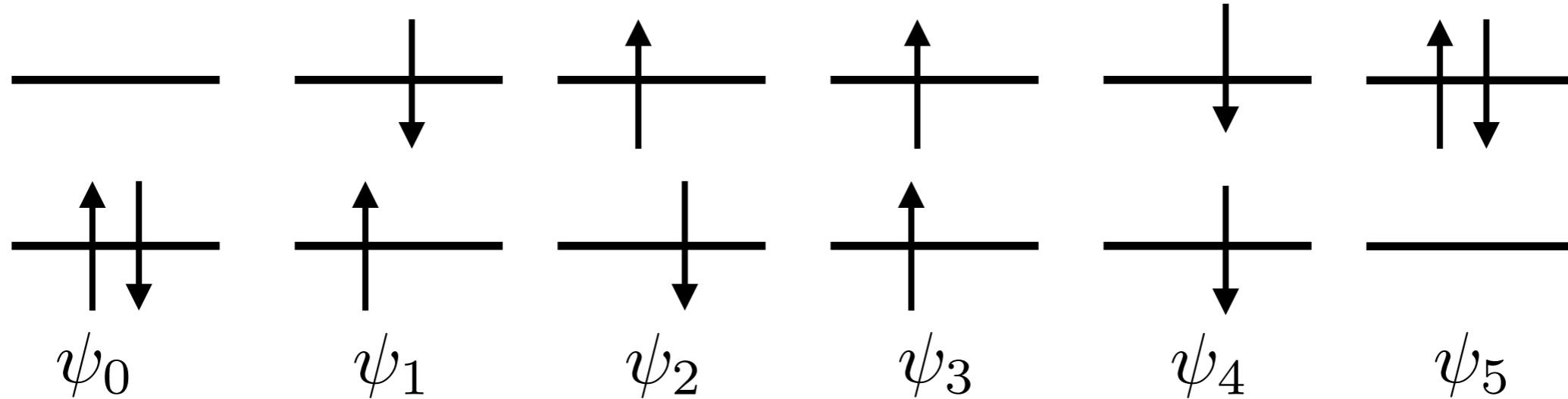
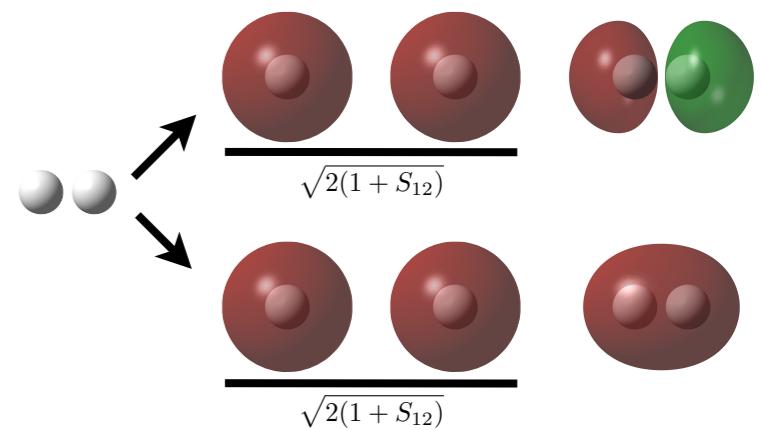
total spin

$$\begin{aligned} S^2 \psi_0 = & \frac{1}{\sqrt{2}} [(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2)) \varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\ & + s_z(1)\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\ & + s_z(2)\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\ & + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2)) \varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2))] \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



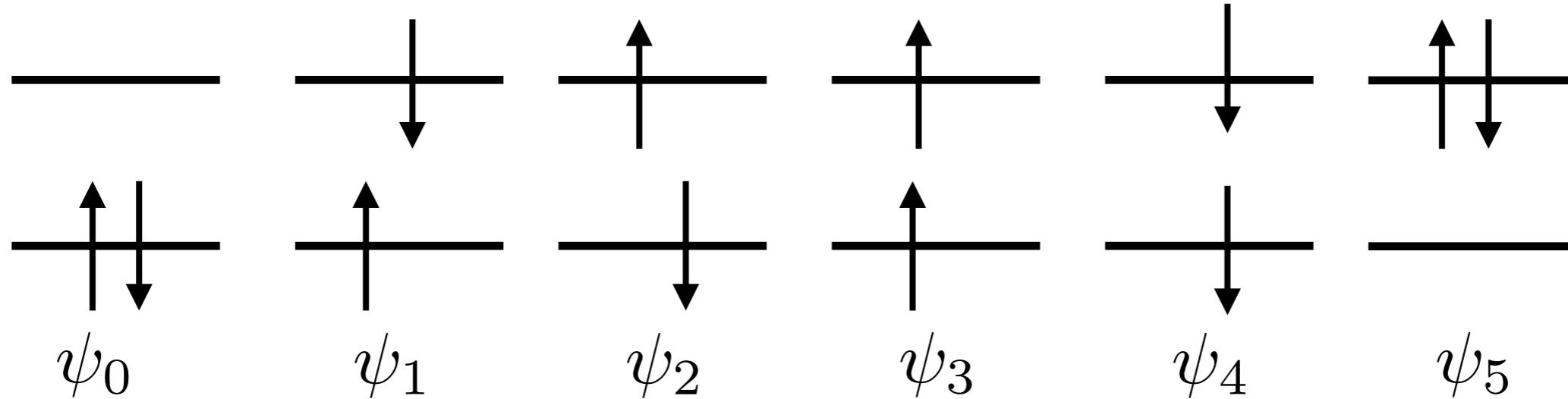
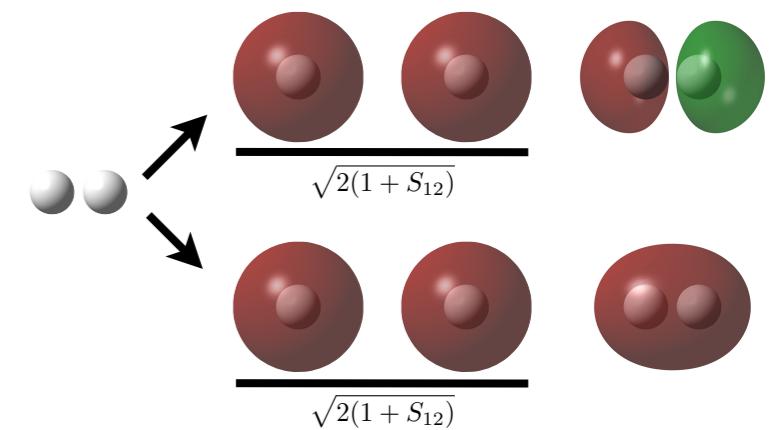
total spin

$$\begin{aligned} S^2 \psi_0 = & \frac{1}{\sqrt{2}} \varphi_1(1) \varphi_1(2) [(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2))] \\ & + s_z(1)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\ & + s_z(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\ & + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2))] \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

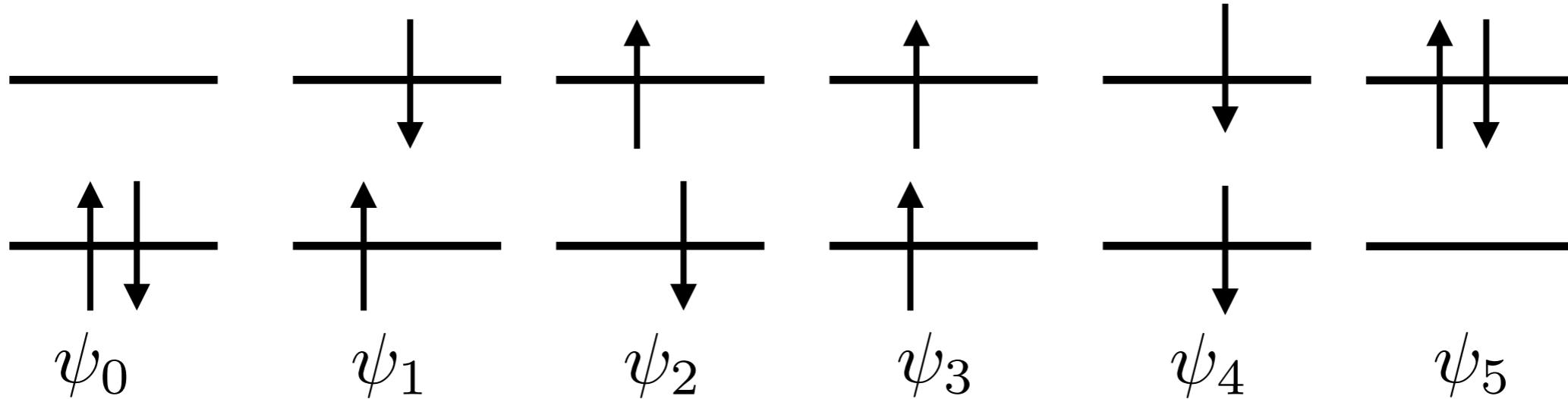
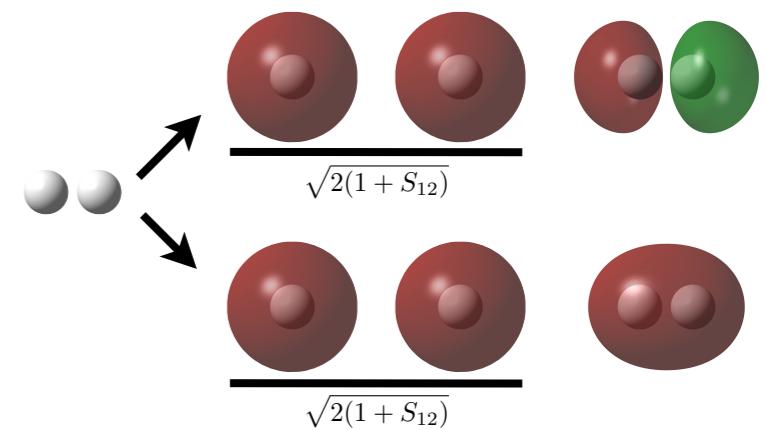
$$\begin{aligned} S^2 \psi_0 = & \frac{1}{\sqrt{2}} \varphi_1(1) \varphi_1(2) [-\beta(1)\alpha(2) + \beta(1)\alpha(2) - \alpha(1)\beta(2) + \alpha(1)\beta(2) \\ & + \frac{1}{2}\alpha(1)\beta(2) + \frac{1}{2}\beta(1)\alpha(2) \\ & - \frac{1}{2}\alpha(1)\beta(2) - \frac{1}{2}\beta(1)\alpha(2) \\ & + (\frac{1}{4} - \frac{1}{4} - \frac{1}{4} + \frac{1}{4} - \frac{1}{4} + \frac{1}{4} + \frac{1}{4} - \frac{1}{4})(\alpha(1)\beta(2) - \beta(1)\alpha(2))] \\ = & 0 \end{aligned}$$

singlet

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

$$\psi_3 = (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\alpha(1)\alpha(2)$$

$$S^2\psi_3 = (S_-S_+ + S_z + S_z^2)\psi_3 = S(S+1)\psi_3$$

$$= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\alpha(1)\alpha(2)$$

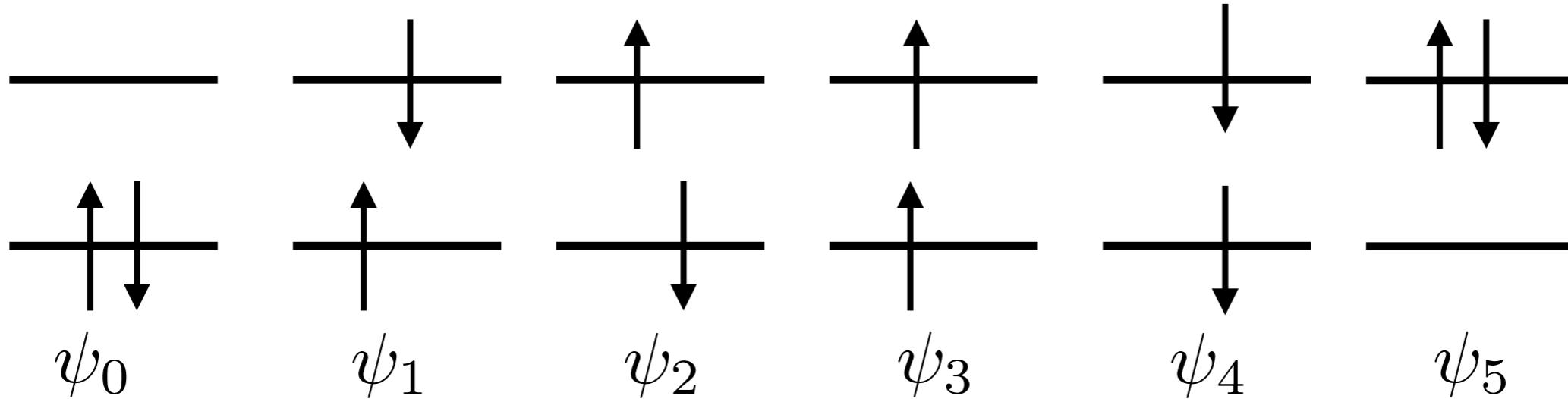
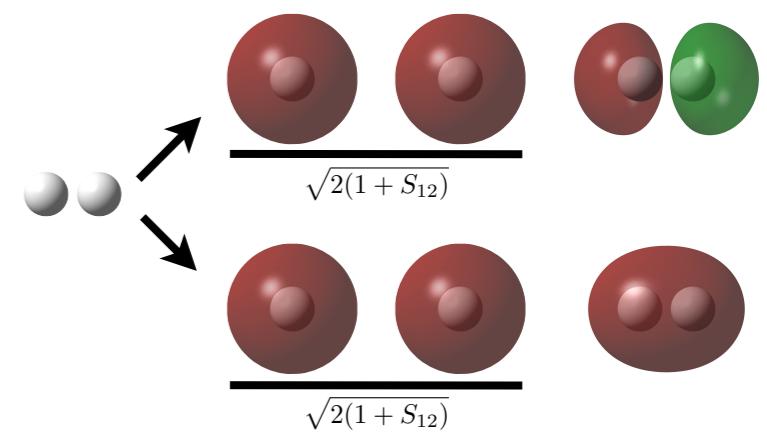
$$+ s_z(1)\alpha(1)\alpha(2) + s_z(2)\alpha(1)\alpha(2)]$$

$$+ (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\alpha(1)\alpha(2)]$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

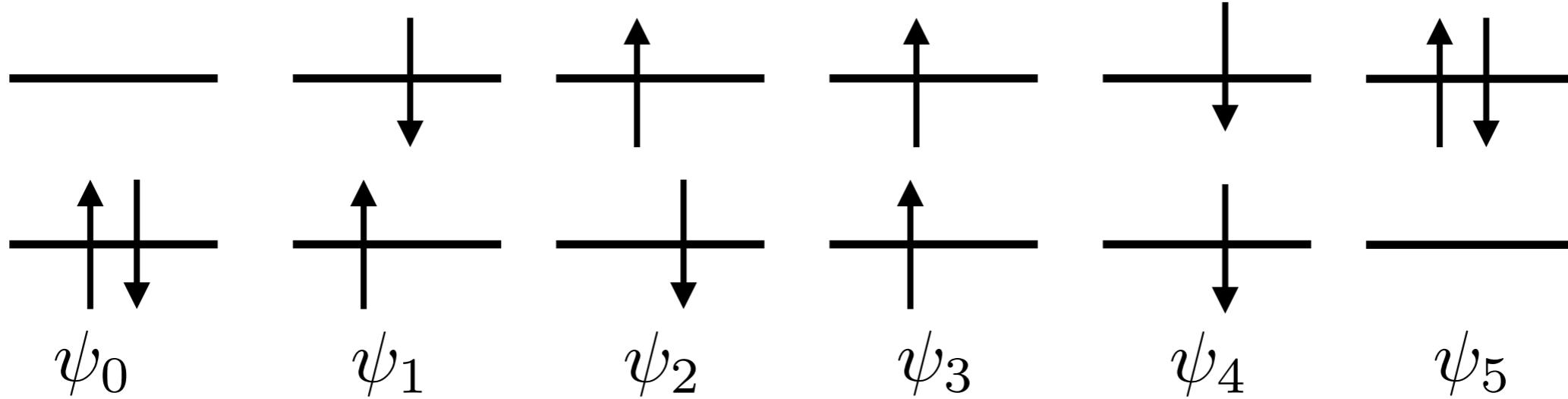
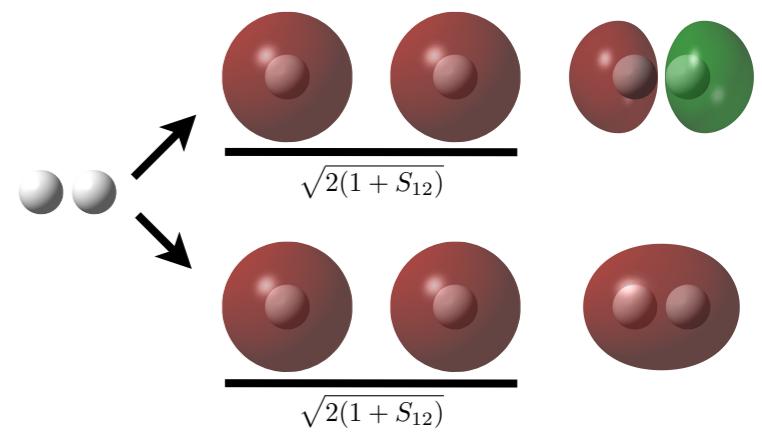
$$\begin{aligned} S^2 \psi_3 &= \frac{1}{\sqrt{2}} (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2)) [\\ &\quad (0 + 0 + 0 + 0)\alpha(1)\alpha(2) \\ &\quad + (\frac{1}{2} + \frac{1}{2})\alpha(1)\alpha(2) \\ &\quad + (\frac{1}{4} + \frac{1}{4} + \frac{1}{4} + \frac{1}{4})\alpha(1)\alpha(2)] \\ &= 2\frac{1}{\sqrt{2}} (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\alpha(1)\alpha(2) \end{aligned}$$

triplet

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

$$\psi_4 = (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\beta(1)\beta(2)$$

$$S^2\psi_4 = (S_-S_+ + S_z + S_z^2)\psi_4 = S(S+1)\psi_4$$

$$= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\beta(1)\beta(2)]$$

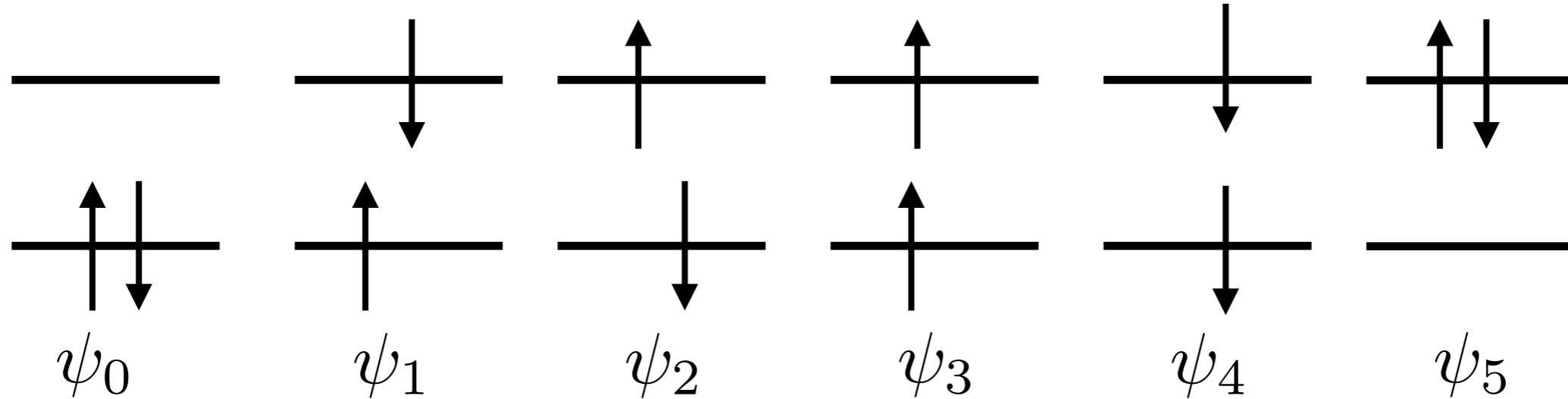
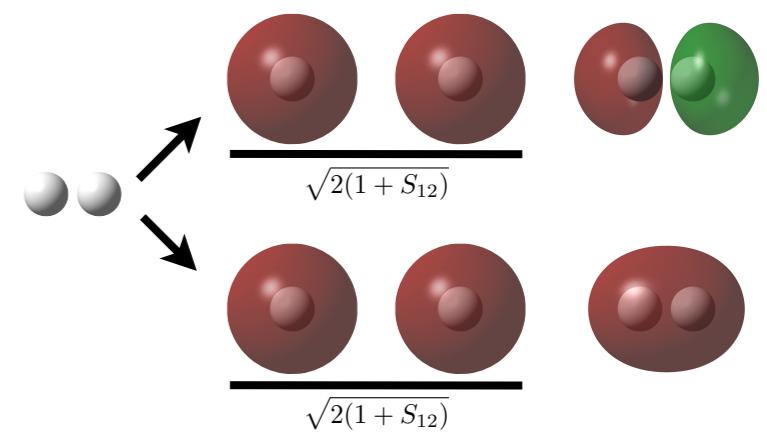
$$+ s_z(1)\beta(1)\beta(2) + s_z(2)\beta(1)\beta(2)$$

$$+ (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\beta(1)\beta(2)]$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

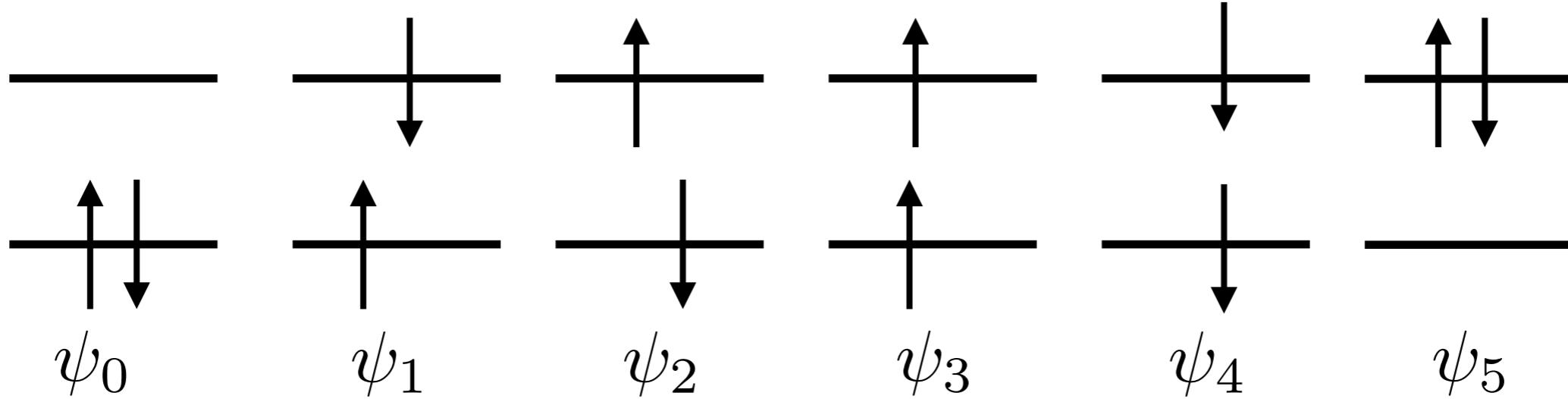
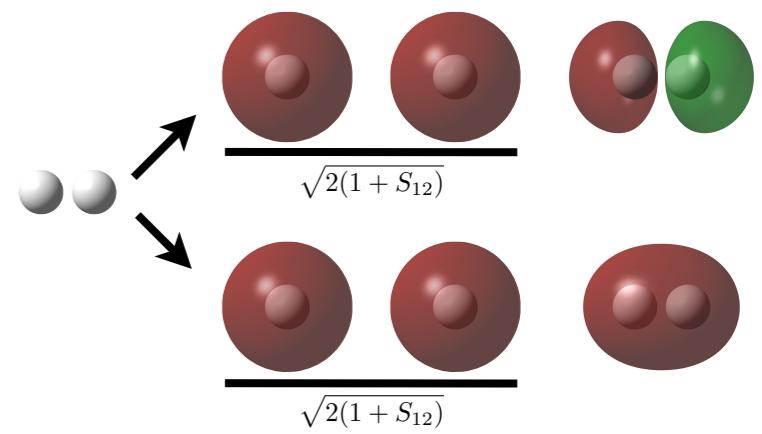
$$\begin{aligned} S^2 \psi_4 &= \frac{1}{\sqrt{2}} (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2)) [\\ &\quad (1 + 0 + 0 + 1)\beta(1)\beta(2) \\ &\quad + (-\frac{1}{2} - \frac{1}{2})\beta(1)\beta(2) \\ &\quad + (\frac{1}{4} + \frac{1}{4} + \frac{1}{4} + \frac{1}{4})\beta(1)\beta(2)] \\ &= 2 \frac{1}{\sqrt{2}} (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\beta(1)\beta(2) \end{aligned}$$

triplet

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



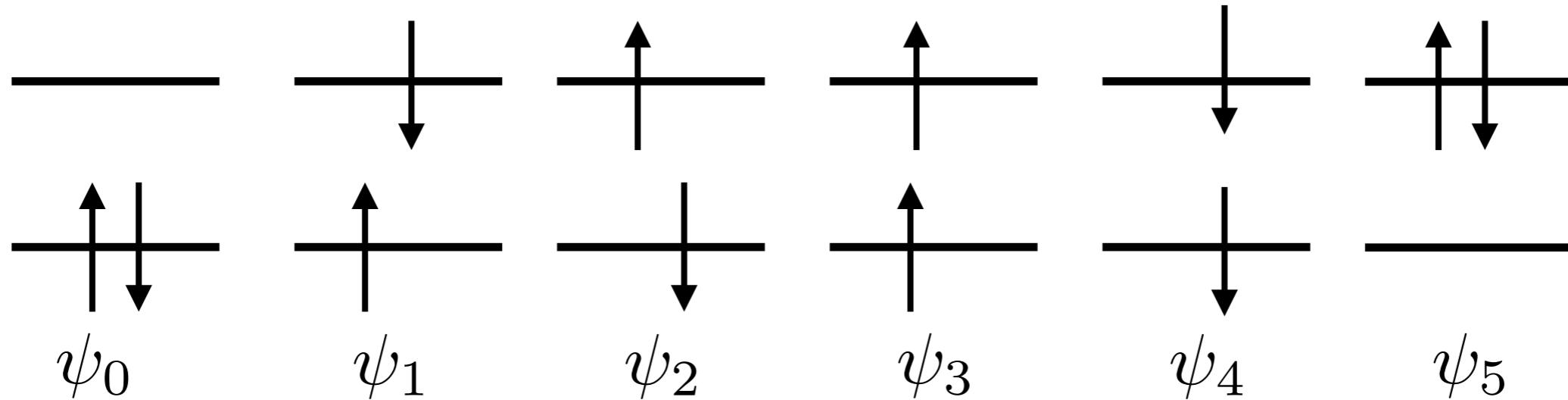
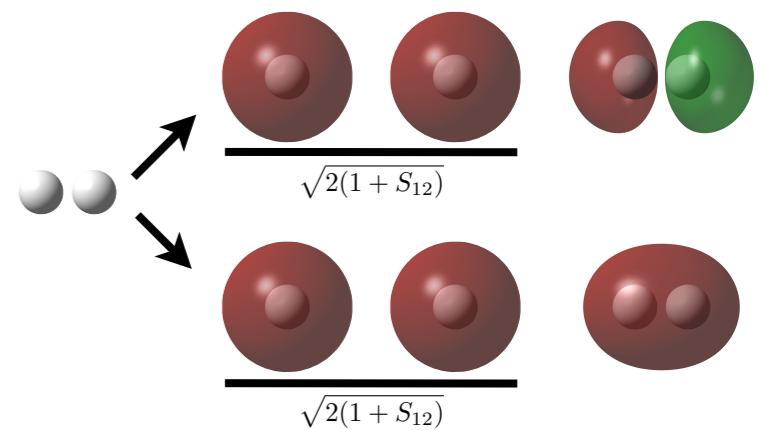
total spin

$$\psi_1 = \frac{1}{\sqrt{2}} [\varphi_1(1)\alpha(1)\varphi_2\beta(2) - \varphi_2(1)\beta(2)\varphi_1(2)\alpha(2)]$$

Spin of Slater determinant

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configurations



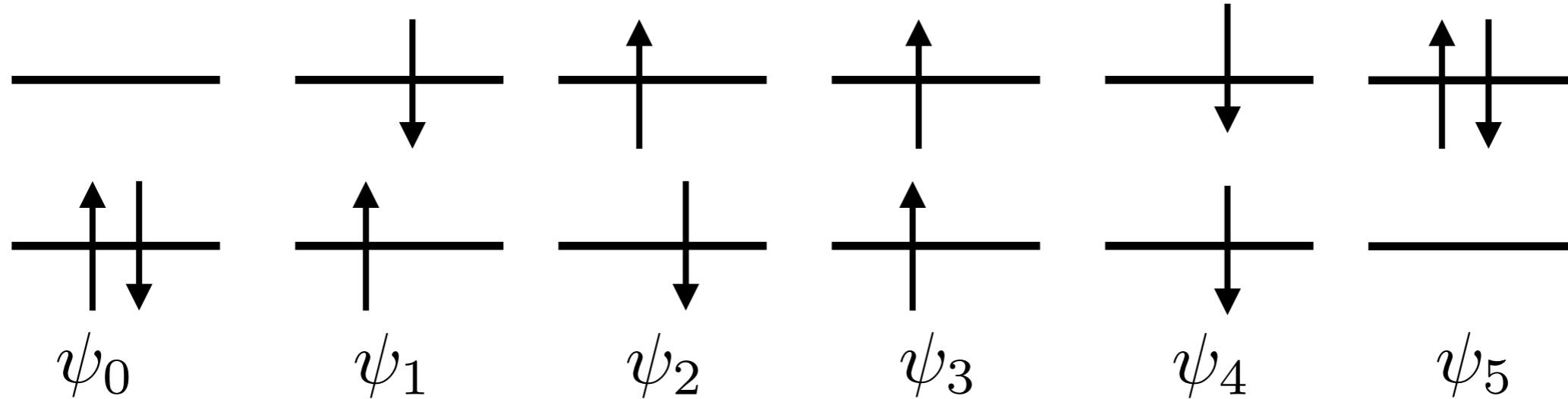
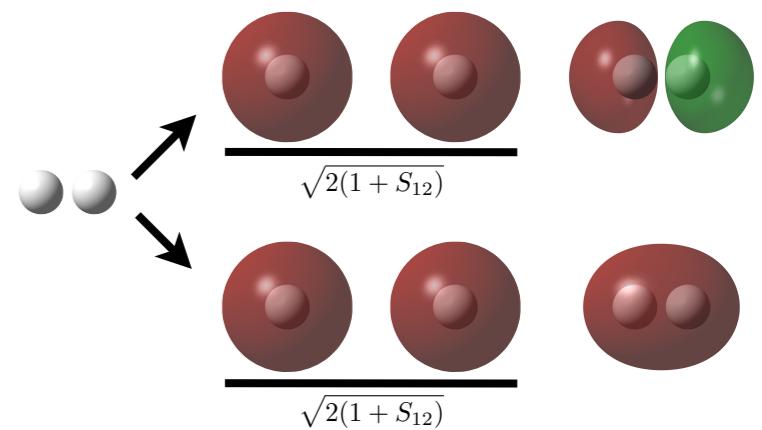
total spin

$$\begin{aligned} S^2 \psi_1 = & \frac{1}{\sqrt{2}} [(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2)) \varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) \\ & - (s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2)) \varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\ & + s_z(1)\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) - s_z(1)\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\ & + s_z(2)\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) - s_z(2)\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\ & + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2)) \varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) \\ & - (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2)) \varphi_2(1)\varphi_1(2)\beta(1)\alpha(2)] \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

$$S^2 \psi_1 = \frac{1}{\sqrt{2}} [\varphi_1(1)\alpha(1)\varphi_2(2)\beta(2) - \varphi_2(1)\beta(1)\varphi_1(2)\alpha(2)]$$

$$+ [\varphi_1(1)\beta(1)\varphi_2(2)\alpha(2) - \varphi_2(1)\alpha(1)\varphi_1(2)\beta(2)]$$

$$= \psi_1 + \psi_2$$

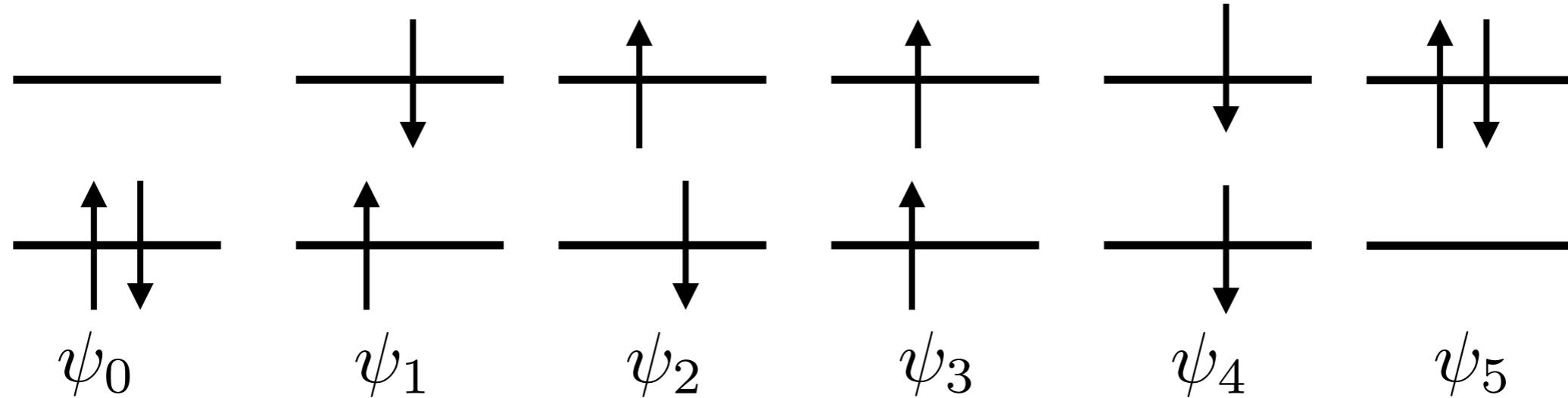
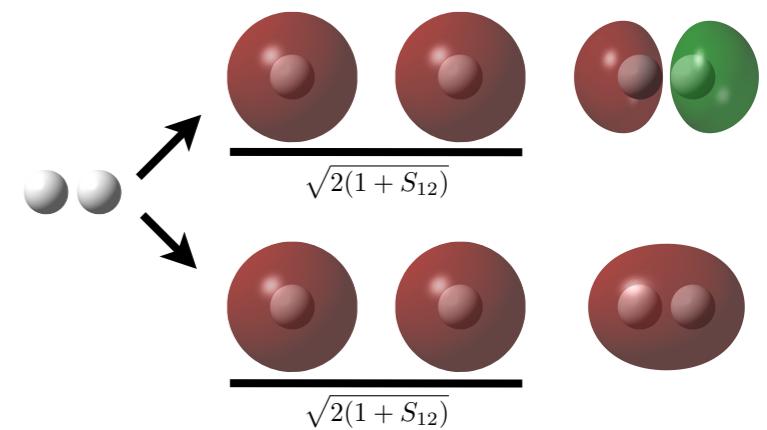
$$\neq S(S+1)\psi_1$$

no eigenfunction of total spin operator!

Spin of Slater determinant

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spin-adapted linear combinations:

$$\psi_A = \psi_1 + \psi_2$$

$$= (\varphi_1(1)\varphi_2(1) - \varphi_2(1)\varphi_1(2))(\alpha(1)\beta(2) + \beta(1)\alpha(2))$$

$$\psi_B = \psi_1 - \psi_2$$

$$= (\varphi_1(1)\varphi_2(1) + \varphi_2(1)\varphi_1(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2))$$

homework: show that both are eigenfunction of total spin operator.

Hartree-Fock Theory for n electrons

Expectations values for one and two electron operators

Hartree product (no spin)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{r}_1) d\mathbf{r}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{r}_1) \phi_b^*(\mathbf{r}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{r}_1) \phi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Slater determinant (spin, Pauli principle)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{x}_1) d\mathbf{x}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$- \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_b(\mathbf{x}_1) \phi_a(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) (1 - \hat{p}_{12}) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

Hartree-Fock Theory for n electrons

Expectations values for one and two electron operators

Hartree product (no spin)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{r}_1) d\mathbf{r}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{r}_1) \phi_b^*(\mathbf{r}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{r}_1) \phi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Slater determinant (spin, Pauli principle)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{x}_1) d\mathbf{x}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$-\frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_b(\mathbf{x}_1) \phi_a(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) (1 - \hat{p}_{12}) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Variational principle

$$E \geq E^0$$

minimum

$$\partial E = 0$$

when for all orbitals

$$\varphi_i(\mathbf{x}) = \varphi_i(\mathbf{x}) + \partial\varphi_i(\mathbf{x})$$

constraint

$$\langle \varphi_i(\mathbf{x}) | \varphi_k(\mathbf{x}) \rangle = \delta_{ik}$$

Langrange multipliers

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero, or actually n^2 smart zeros)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & & [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle \\ &&& - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \\ &+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & & [\langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \partial \varphi_k \rangle \\ &&& - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k | \varphi_i \rangle - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \partial \varphi_i \rangle] \end{aligned}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E = & \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle \\ & & - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \end{aligned}$$

$$\begin{aligned} & + \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \partial \varphi_k \rangle \\ & & - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k | \varphi_i \rangle - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \partial \varphi_i \rangle] \end{aligned}$$

complex conjugate

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\partial E = \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k$$

changing variables

$$[\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle]$$

$$+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k [\langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i \varphi_k \rangle + \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \partial \varphi_k \rangle]$$

$$- \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k \varphi_i \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \partial \varphi_i \rangle]$$

complex conjugate

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \\ &\quad + \text{c.c} \end{aligned}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

rearranging

$$\begin{aligned}\partial E = & \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ & - \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle \\ & + \text{c.c.}\end{aligned}$$

coulomb

exchange

for each orbital:

$$\langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial \varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial \varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

rearranging a bit more

$$\langle \partial\varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial\varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial\varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

local coulomb operator

$$\hat{J}_k |\varphi_i(1)\rangle = \frac{|\varphi_k|^2}{r_{12}} |\varphi_i(1)\rangle$$

non-local exchange operator

$$\hat{K}_k |\varphi_i(1)\rangle = \langle \varphi_k(2) | \frac{1}{r_{12}} | \varphi_i(2) \rangle | \varphi_k(1) \rangle$$

Fock operator

$$\hat{f} = \hat{h} + \sum_k [\hat{J}_k - \hat{K}_k]$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

going back the the lagrangian

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

thus

$$\sum_i \left[\langle \partial \varphi_i | \hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ik} \langle \partial \varphi_i | \varphi_k \rangle \right] = 0$$

$$\langle \partial \varphi_i | \left[\hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ik} | \varphi_k \rangle \right] = 0$$

system of n coupled non-linear equations

$$\hat{f} | \varphi_i \rangle = \sum_k \epsilon_{ik} | \varphi_k \rangle$$

let's try to uncouple them!

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

system of n coupled non-linear equations

$$\hat{f}|\varphi_i\rangle = \sum_k \epsilon_{ik} |\varphi_k\rangle$$

matrix notation

$$\mathbf{F}|\varphi\rangle = \boldsymbol{\epsilon}|\varphi\rangle$$

unitary transformation

$$|\varphi'_a\rangle = \hat{U}|\phi_a\rangle = \sum_k |\varphi_k\rangle U_{ka} \quad \hat{U}^\dagger \hat{U} = 1 \quad \hat{U}^\dagger = \hat{U}^{-1}$$

Fock equation

$$\mathbf{U}\mathbf{F}\mathbf{U}^\dagger\mathbf{U}|\varphi\rangle = \mathbf{U}\boldsymbol{\epsilon}\mathbf{U}^\dagger\mathbf{U}|\varphi\rangle$$

choose unitary transformation matrix that diagonalizes $\boldsymbol{\epsilon}$

$$\mathbf{F}'|\varphi'\rangle = \boldsymbol{\epsilon}'|\varphi'\rangle \quad \epsilon'_{ik} = \delta_{ik} \epsilon_k$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

$$\hat{f}|\varphi'_i\rangle = \epsilon'_i |\varphi'_i\rangle$$

orbital energies (ionization energies and electron affinities)

$$\epsilon_i = \langle \varphi_i | \hat{h} | \varphi \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x})_1|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

$$- \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

$$E = \sum_i h_{ii} + \sum_i \sum_k J_{ik} - \sum_i \sum_k K_{ik}$$

total energy (correct double counting, also first order energy in PT)

$$E = \sum_i h_{ii} + \frac{1}{2} \sum_i \sum_k J_{ik} - \frac{1}{2} \sum_i \sum_k K_{ik}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

$$\hat{f}|\varphi_i\rangle = \epsilon_i |\varphi_i\rangle$$

orbital energies (ionization energies and electron affinities)

$$\epsilon_i = \langle \varphi_i | \hat{h} | \varphi \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x})_1|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

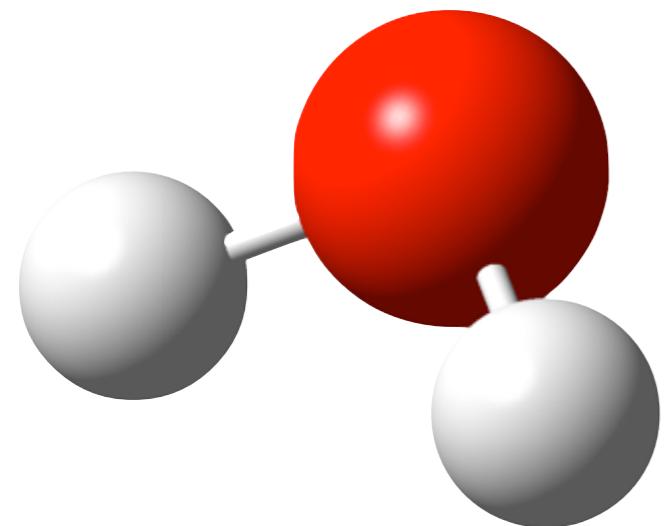
$$- \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

$$E = \sum_i h_{ii} + \sum_i \sum_k J_{ik} - \sum_i \sum_k K_{ik}$$

total energy (correct double counting)

$$E = \sum_i h_{ii} + \frac{1}{2} \sum_i \sum_k J_{ik} - \frac{1}{2} \sum_i \sum_k K_{ik}$$

Wassermolekül



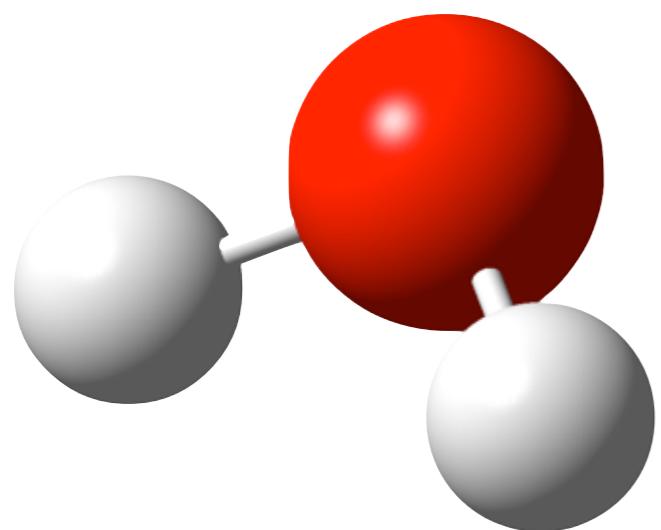
2 H⁺

1 O⁸⁺

10 Elektronen

10 Molekülorbitale

Wassermolekül

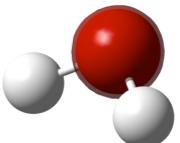
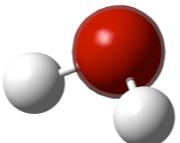
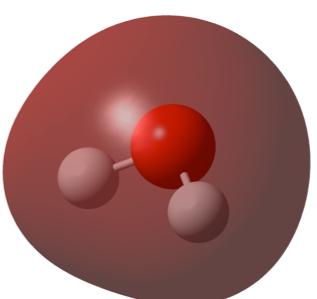
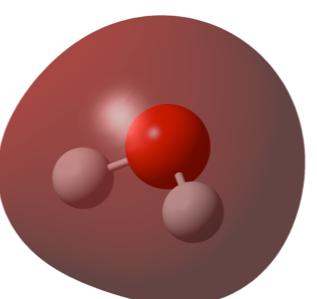
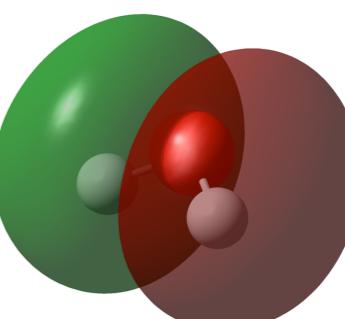
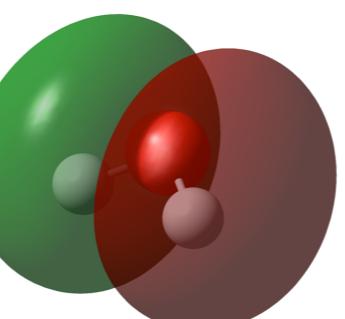
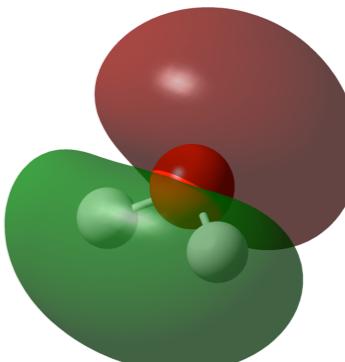
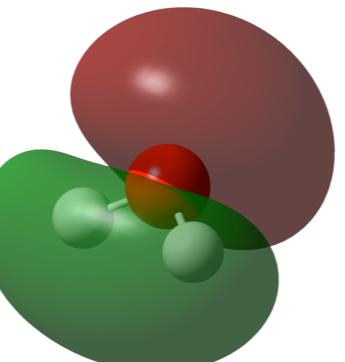
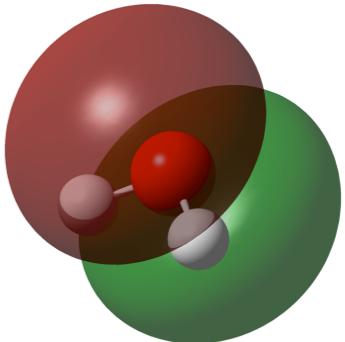
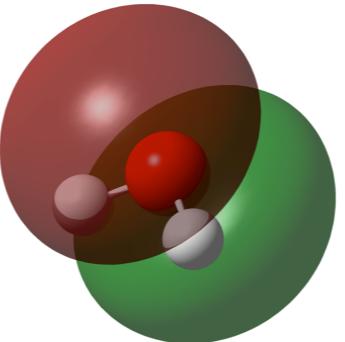


2 H⁺

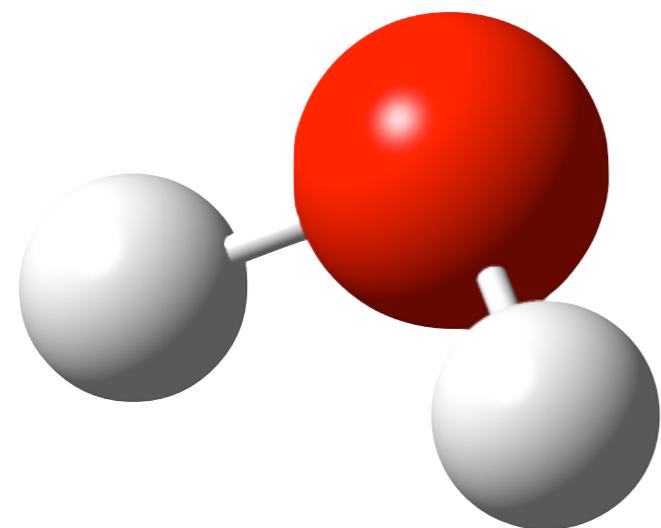
1 O⁸⁺

10 Elektronen

10 Molekülorbitale



Wassermolekül

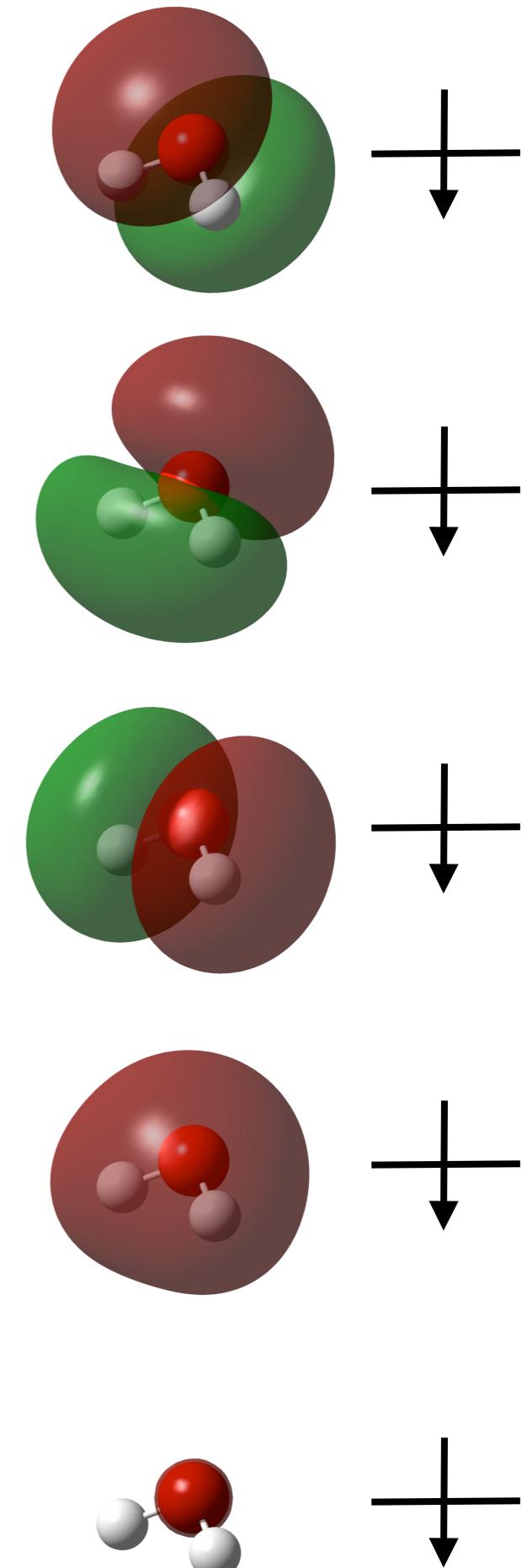
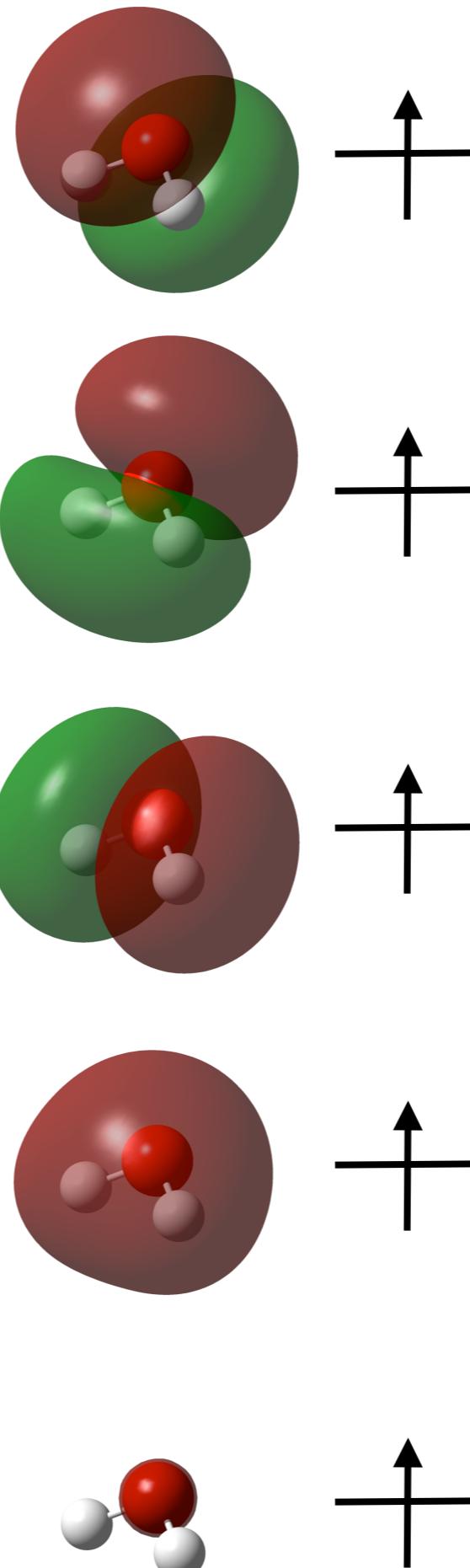


2 H⁺

1 O⁸⁺

10 Elektronen

10 Molekülorbitale



Hartree-Fock Theory for n electrons

Hartree-Fock eigenvalue equations

$$\hat{f}(\mathbf{r})\varphi_i(\mathbf{x}) = \epsilon_i\varphi(\mathbf{x})$$

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

step 2: expand spatial orbitals in basis functions

restricted Hartree Fock

electron pair with opposite spin in same spatial orbital

$$\varphi_i(\mathbf{x}) = \phi_j(\mathbf{r})\alpha(s)$$

$$\varphi_{i+1}(\mathbf{x}) = \phi_j(\mathbf{r})\beta(s)$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) \\ &\quad + \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\alpha(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad + \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\beta(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad - \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad - \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\beta(s_1) d\mathbf{r}_2 ds_2\end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) = \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1)$$

$$+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \alpha^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \alpha(s_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \beta^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \beta(s_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \alpha^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \alpha(s_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \beta^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \alpha(s_2) \phi_k(\mathbf{r}_1) \beta(s_1) d\mathbf{r}_2 ds_2$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) = \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1)$$

$$+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

$$+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

$$- \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\begin{aligned} \int \alpha^*(s_1) \hat{f}(\mathbf{x}_1) \alpha(s_1) ds_1 \phi_i(\mathbf{r}_1) &= \int \alpha^*(s_1) \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \alpha(s_1) ds_1 \\ &\quad + \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &\quad + \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &\quad - \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \end{aligned}$$

Hartree-Fock eigenvalue equation for spatial orbitals

$$\begin{aligned} \hat{f}(\mathbf{r}_1) \phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \\ &\quad + 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &\quad - \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r}_2 \\ &= \epsilon_i \phi_i(\mathbf{r}_1) \end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1) \\ &\quad + 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &\quad - \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r}_2 \\ &= \epsilon_i \phi_i(\mathbf{r}_1)\end{aligned}$$

step 2: expand spatial orbitals in basis functions (basisset)

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

Hartree-Fock Theory for n electrons

linear combination of atomic orbitals

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

hydrogen-like orbitals (one possibility out of many....)

$$\gamma_1 = \psi^{1s}(\zeta_1)$$

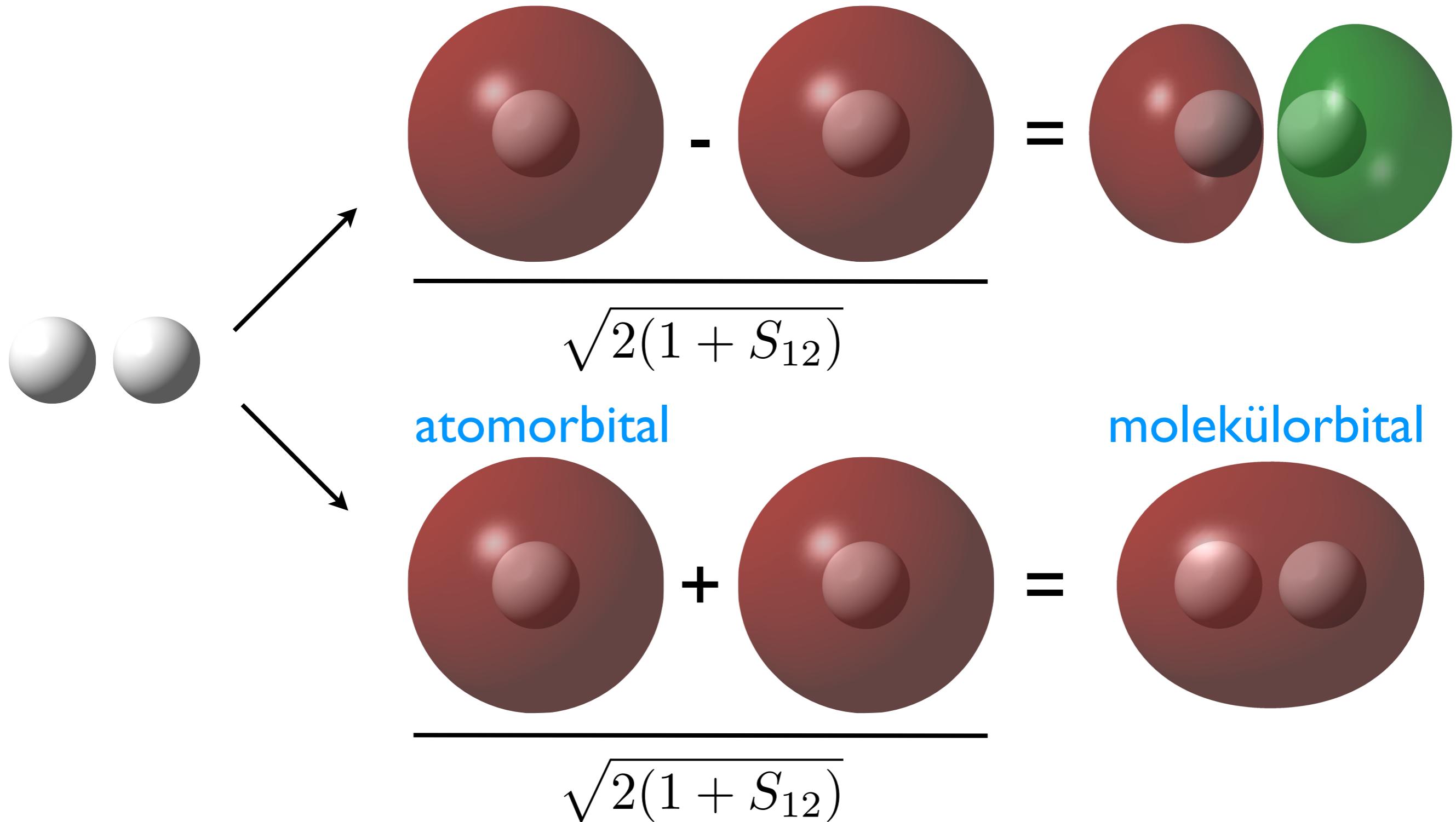
$$\gamma_2 = \psi^{2s}(\zeta_2)$$

$$\gamma_3 = \psi^{2p}(\zeta_3)$$

$$\gamma_4 = \dots$$

Wasserstoffmolekül

Lineare Kombination von einzelne Wasserstoff-Orbitale



Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions

$$\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) = \epsilon_i\phi_i(\mathbf{r}_1)$$

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij}\gamma_j(\mathbf{r} - \mathbf{R}_j)$$

$$\hat{f}(\mathbf{r}_1) \sum_\nu c_{\nu i} \gamma_\nu(\mathbf{r}_1) = \epsilon_i \sum_\nu c_{\nu i} \gamma_\nu(\mathbf{r}_1)$$

$$\sum_\nu c_{\nu i} \int \gamma_\mu^*(\mathbf{r}_1) \hat{f}(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 = \epsilon_i \sum_\nu c_{\nu i} \int \gamma_\mu(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions (basisset)

$$\sum_{\nu} c_{\nu i} \int \gamma_{\mu}^{*}(\mathbf{r}_1) \hat{f}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 = \epsilon_i \sum_{\nu} c_{\nu i} \int \gamma_{\mu}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1$$

express in terms of matrices

$$\sum_{\nu} F_{\mu\nu} c_{\nu i} = \epsilon_i \sum_{\nu} S_{\mu\nu} c_{\nu i}$$

$$\mathbf{F}\mathbf{C} = \mathbf{S}\mathbf{C}\boldsymbol{\epsilon}$$

solution if, and only if

$$|\mathbf{F} - \epsilon_i \mathbf{S}| = 0$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

non-linear: \mathbf{F} depends on \mathbf{C}

$$F_{\mu\nu} = \int \gamma_\mu^*(\mathbf{r}_1) \hat{h}^0(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1$$

$$+ 2 \sum_a \int \int \gamma_\mu^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_a(\mathbf{r}_2) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$- \sum_a \int \int \gamma_\mu^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\nu(\mathbf{r}_2) \phi_a(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$F_{\mu\nu} = h_{\mu\nu}^0$$

$$+ 2 \sum_a \sum_\kappa \sum_\lambda c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_\mu^*(\mathbf{r}_1) \gamma_\kappa^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\lambda(\mathbf{r}_2) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$- \sum_a \sum_\kappa \sum_\lambda c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_\mu^*(\mathbf{r}_1) \gamma_\kappa^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\nu(\mathbf{r}_2) \gamma_a(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

Hartree-Fock

Roothaan-Hall equations

non-linear eigenvalue problem

$$Fc = \epsilon_i Sc$$

practical algorithm

iterate until self-consistency

pre-compute integrals of basisset

$$S_{\mu\nu} = \langle \gamma_\mu | \gamma_\nu \rangle$$

