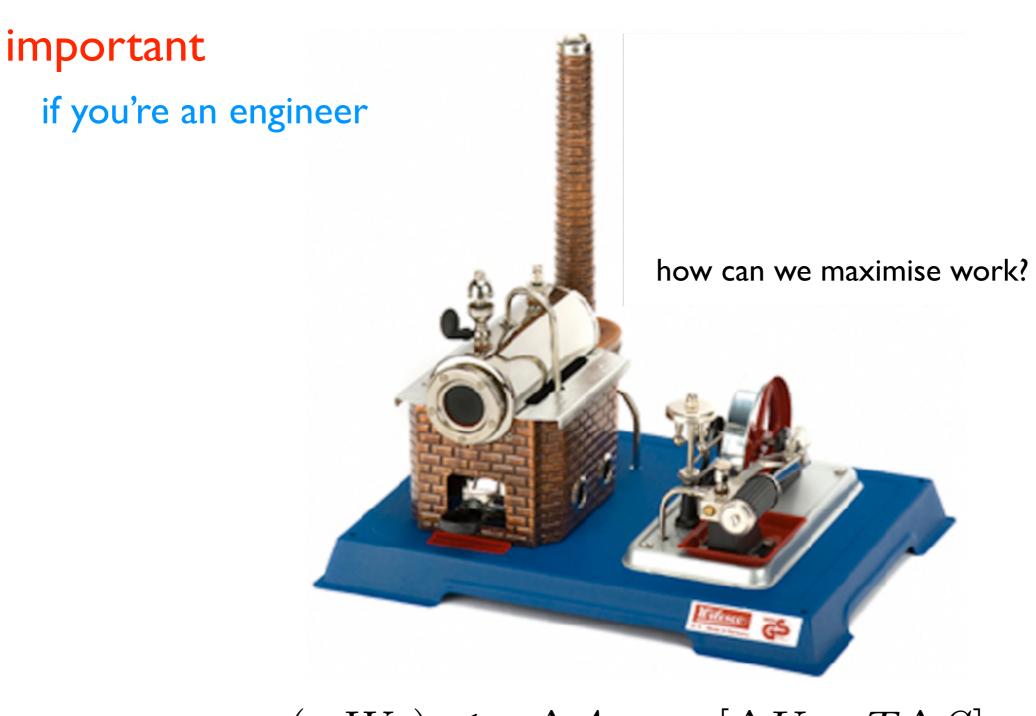
# Lecture 3

first law: conservation of energy second law: direction of spontaneous change



$$(-W_{\rm u}) \le -\Delta A = -\left[\Delta U - T\Delta S\right]$$

# Helmholtz free energy

internal energy

entropy

$$A = U - TS$$

extensive properties

# Helmholtz free energy

internal energy

entropy

$$A = U - TS$$

extensive properties

chemical/physical change at constant temperature

$$\Delta A = \Delta U - T\Delta S$$

$$\Delta A = -T\Delta S^{\text{tot}}$$

spontaneous

$$\Delta A < 0$$

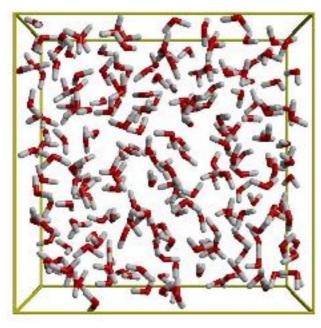
equilibrium

$$\Delta A = 0$$

entropy of isolated system

micro-states (realisations)

box with 216 waters:

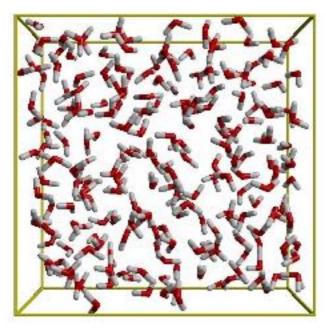


microstate 1

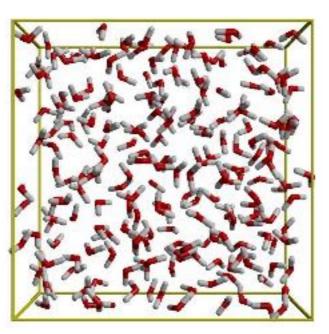
entropy of isolated system

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box with 216 waters:



microstate 1

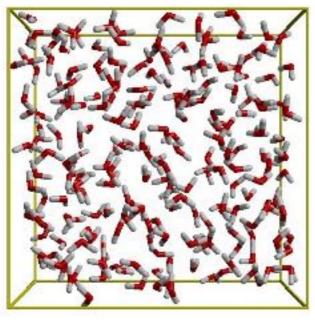


microstate 2

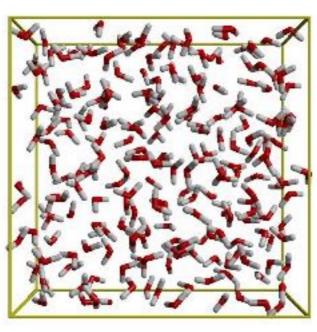
# entropy of isolated system

micro-states (realisations)

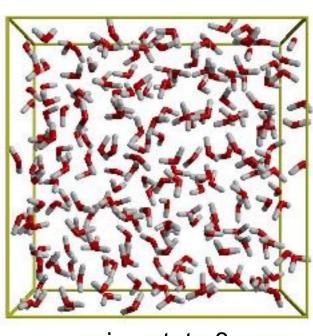
box with 216 waters:



microstate 1



microstate 2

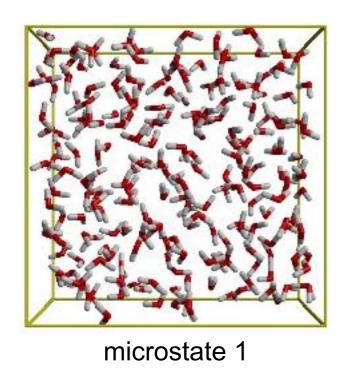


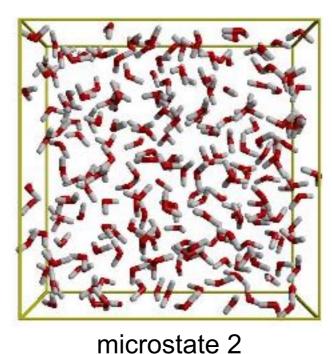
microstate 3

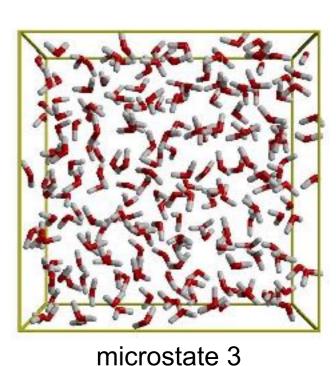
entropy of isolated system

micro-states (realisations)

box with 216 waters:







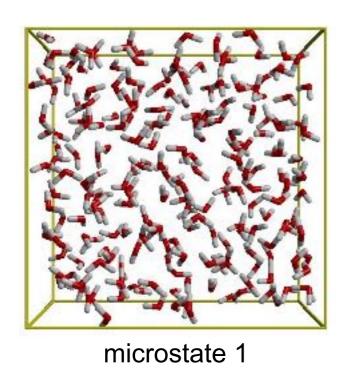
total number of micro-states (realisations)



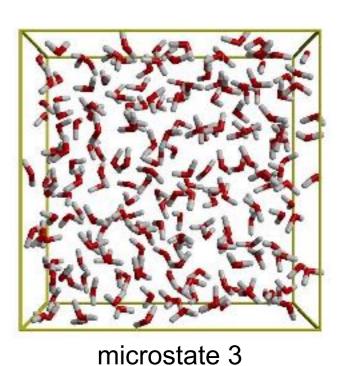
## entropy of isolated system

micro-states (realisations)

box with 216 waters:



microstate 2



total number of micro-states (realisations)

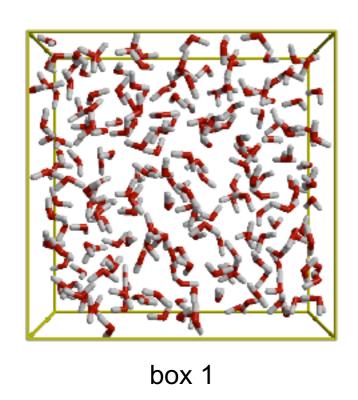
 $\Omega$ 

if you know how many, you pass the course today!

entropy of isolated system

micro-states (realisations)

two boxes of water:



box 2

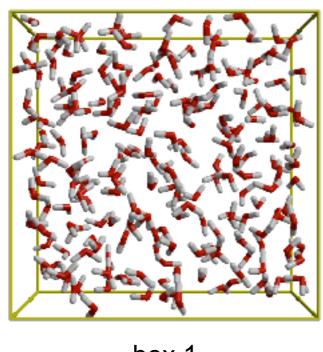
total number of micro states (realisations)

$$\Omega^{\mathrm{tot}} =$$

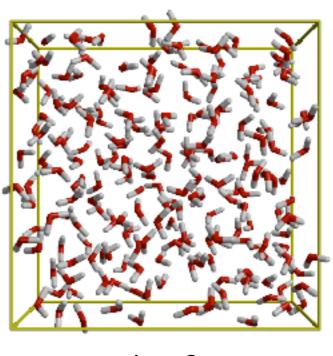
## entropy of isolated system

micro-states (realisations)

two boxes of water:



box 1



box 2

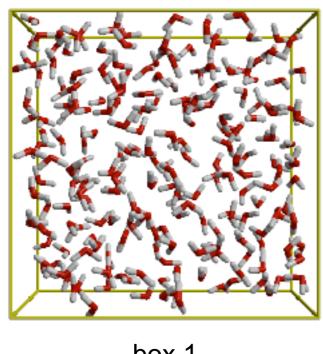
total number of micro states (realisations)

$$\Omega^{\rm tot} = \Omega_1 \cdot \Omega_2$$

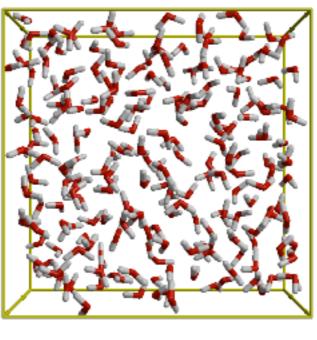
entropy of isolated system

micro-states (realisations)

two boxes of water:



box 1



box 2

total number of micro states (realisations)

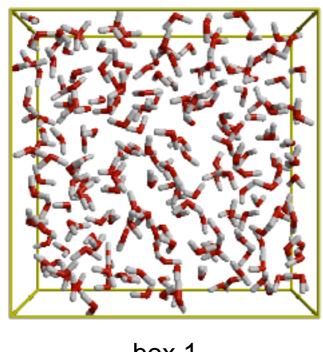
$$\Omega^{\mathrm{tot}} = \Omega_1 \cdot \Omega_2$$

entropy is extensive, so should be sum:

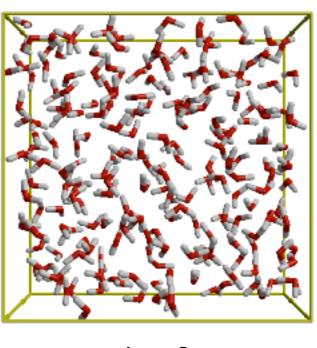
## entropy of isolated system

micro-states (realisations)

two boxes of water:



box 1



box 2

total number of micro states (realisations)

$$\Omega^{\mathrm{tot}} = \Omega_1 \cdot \Omega_2$$

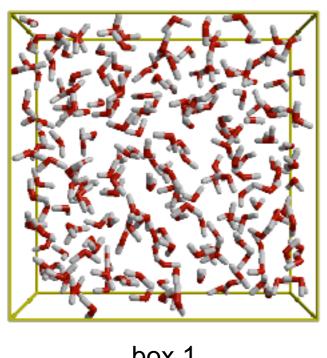
entropy is extensive, so should be sum:

$$S = k \ln \Omega$$

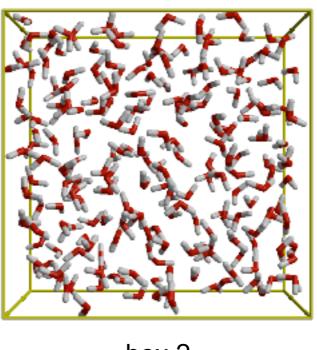
## entropy of isolated system

micro-states (realisations)

two boxes of water:



box 1



box 2

total number of micro states (realisations)

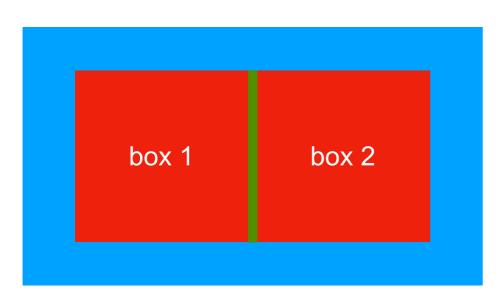
$$\Omega^{\mathrm{tot}} = \Omega_1 \cdot \Omega_2$$

entropy is extensive, so should be sum:

$$S^{\text{tot}} = k \ln[\Omega_1 \cdot \Omega_2] = k \ln \Omega_1 + k \ln \Omega_2 = S_1 + S_2$$

## two systems in thermal equilibrium, isolated from world

diathermic walls (only energy can transfer)



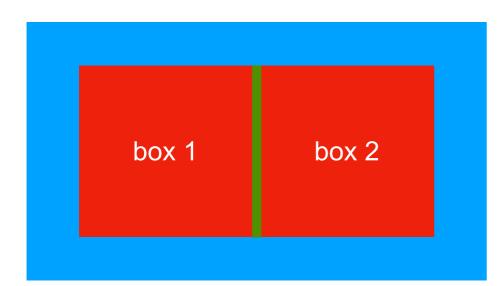
$$N_1 + N_2 = N$$

$$E_1 + E_2 = E$$

$$\frac{dE_2}{dE_1} = -1$$

## two systems in thermal equilibrium, isolated from world

### diathermic walls (only energy can transfer)



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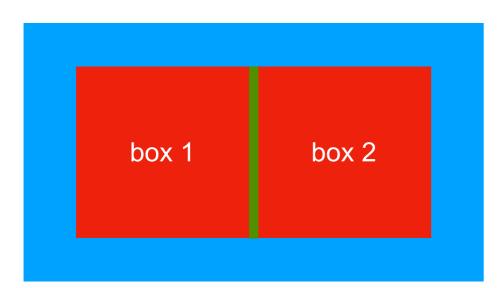
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### equilibrium: no net changes in total entropy

$$\frac{\partial S^{\text{tot}}}{\partial E_1} = \frac{\partial S_1}{\partial E_1} + \frac{\partial S_2}{\partial E_2} \frac{dE_2}{dE_1} = 0$$

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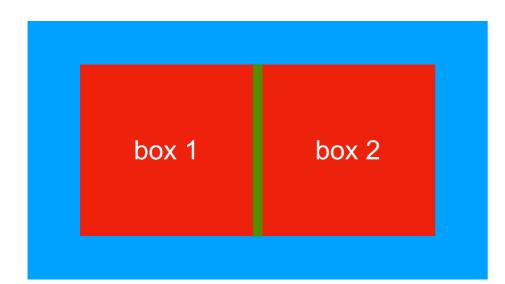
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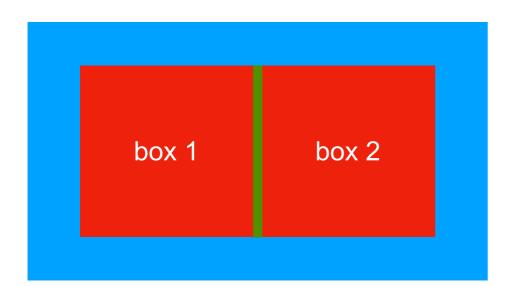
$$\frac{\partial S_1}{\partial E_1} = \frac{\partial S_2}{\partial E_2}$$

### definition of temperature:

$$\frac{\partial S_1}{\partial E_1} = \frac{1}{T_1}$$

## two systems in thermal equilibrium, isolated from world

diathermic walls (only energy can transfer)



$$N_1 + N_2 = N$$

$$E_1 + E_2 = E$$

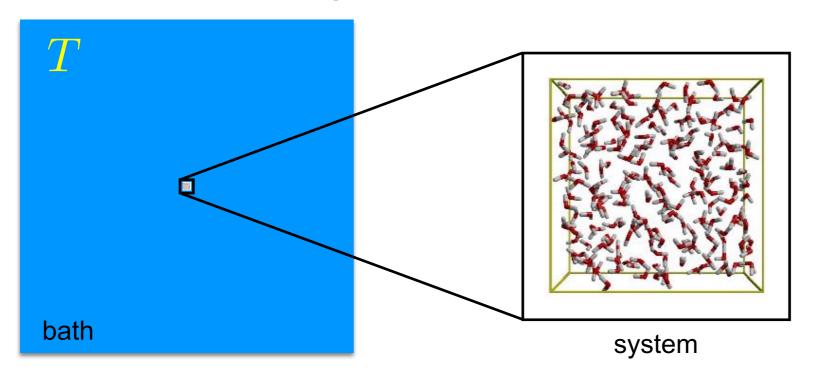
no equilibrium: total entropy must increase

$$\frac{dS^{\text{tot}}}{dt} = \left(\frac{1}{T_1} - \frac{1}{T_2}\right) \frac{dE_1}{dt} > 0$$

energy flows from higher to lower temperature

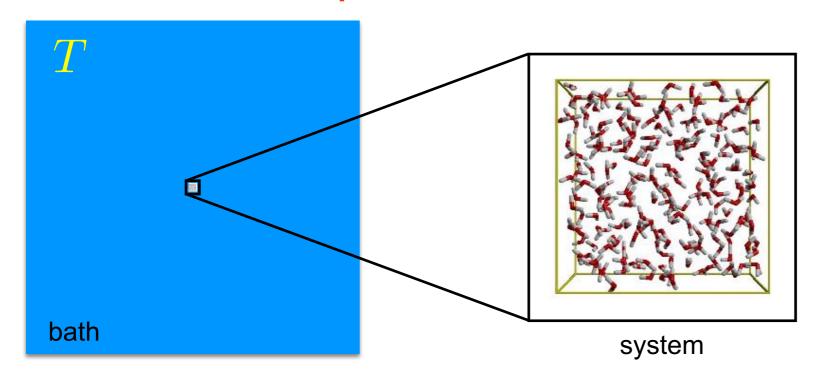
### canonical ensemble

system in thermal equilibrium with bath



#### canonical ensemble

system in thermal equilibrium with bath



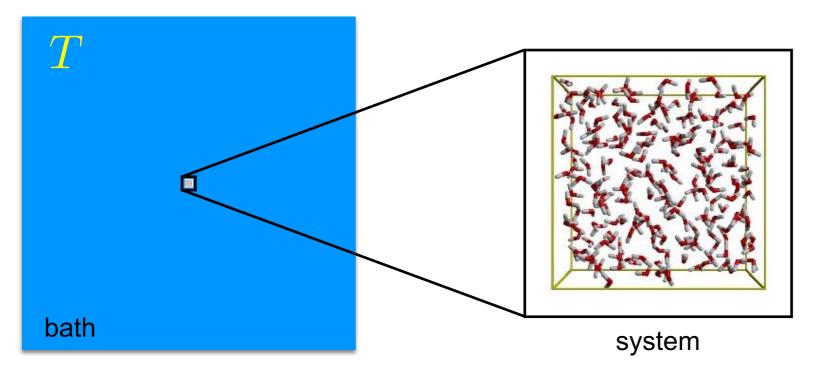
#### micro-states of system

each with different energy

$$E_1 < E_2 < E_3 < \dots < E_i < \dots$$

#### canonical ensemble

system in thermal equilibrium with bath



#### micro-states of system

each with different energy

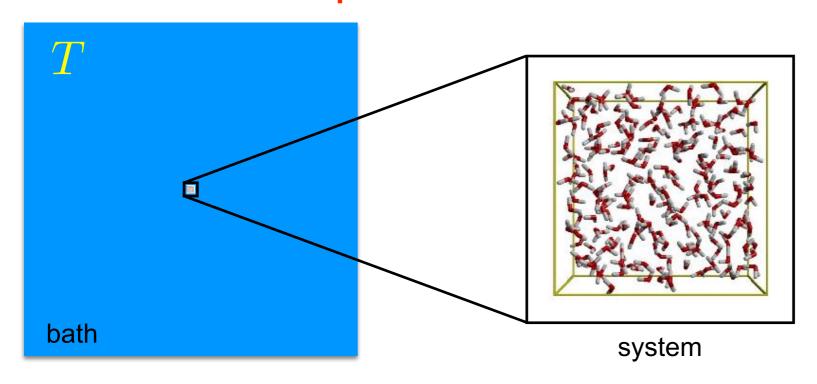
$$E_1 < E_2 < E_3 < \dots < E_i < \dots$$

probability of micro state i proportional to number of micro states of bath

$$p_i = \text{const} \cdot \Omega_{\text{bath}}(E^{\text{tot}} - E_i)$$

#### canonical ensemble

system in thermal equilibrium with bath



#### micro-states of system

each with different energy

$$E_1 < E_2 < E_3 < \dots < E_i < \dots$$

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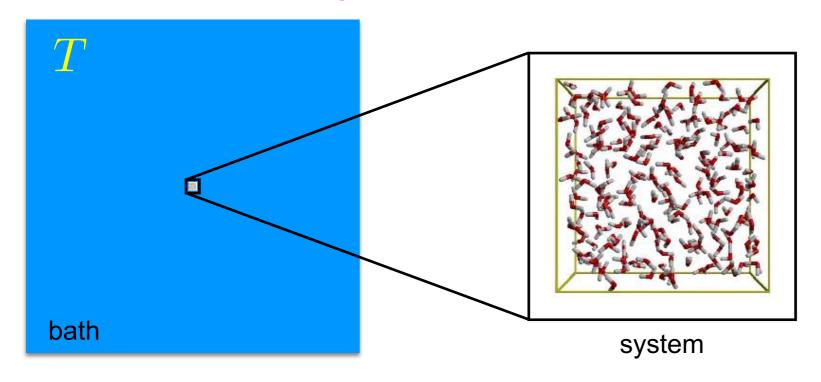
$$p_i = \text{const} \cdot \Omega_{\text{bath}}(E^{\text{tot}} - E_i)$$

normalisation (const):

$$p_i = \frac{\Omega_{\text{bath}}(E^{\text{tot}} - E_i)}{\sum_i \Omega_{\text{bath}}(E^{\text{tot}} - E_i)}$$

#### canonical ensemble

system in thermal equilibrium with bath



#### micro-states of system

probability of micro-state i proportional to number of micro states of bath

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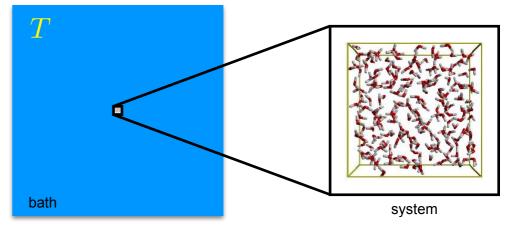
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#### canonical ensemble

system in thermal equilibrium with bath

micro-states of system



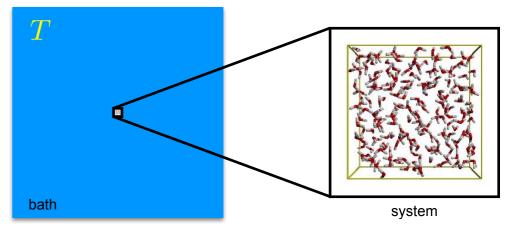
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system in thermal equilibrium with bath

## micro-states of system



probability of micro state i proportional to number of micro-states of bath

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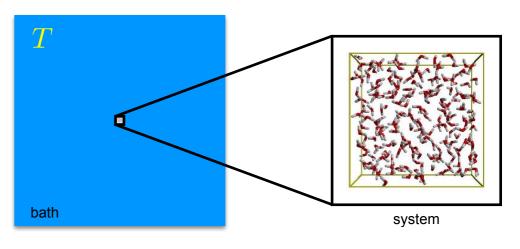
with definition of entropy:  $S = k \ln \Omega$ 

$$p_i = \text{const} \cdot \exp[S_{\text{bath}}(E^{\text{tot}} - E_i)/k]$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



probability of micro state i proportional to number of micro-states of bath

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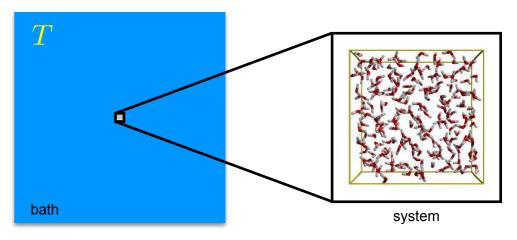
bath much larger than system:

$$E^{\rm tot} \gg E_i$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



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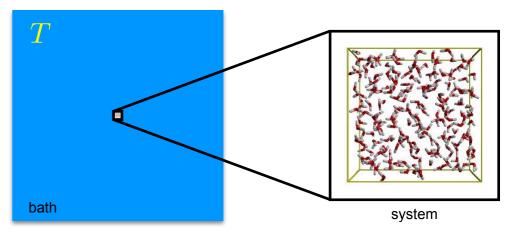
Taylor expansion of  $S_{\text{bath}}$  around  $E^{\text{tot}}$ 

$$\frac{1}{k}S_{\text{bath}}(E^{\text{tot}} - E_i) = \frac{1}{k}S_{\text{bath}}(E^{\text{tot}}) - \frac{E_i}{k}\frac{\partial S_{\text{bath}}(E)}{\partial E}\bigg|_{E = E^{\text{tot}}} + \dots$$

#### canonical ensemble

system in thermal equilibrium with bath

micro-states of system



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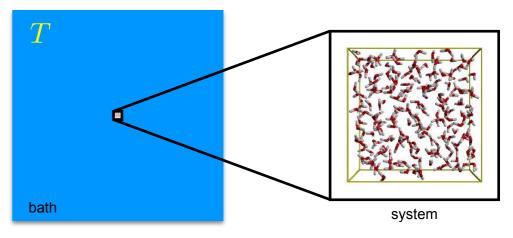
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system in thermal equilibrium with bath

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with definition of temperature

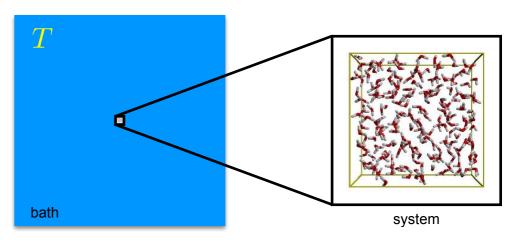
$$\frac{\partial S}{\partial E} = \frac{1}{T}$$

$$\frac{1}{k}S_{\text{bath}}(E^{\text{tot}} - E_i) = \frac{1}{k}S_{\text{bath}}(E^{\text{tot}}) - \frac{1}{kT}E_i$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



probability of micro-state i proportional to number of micro-states of bath

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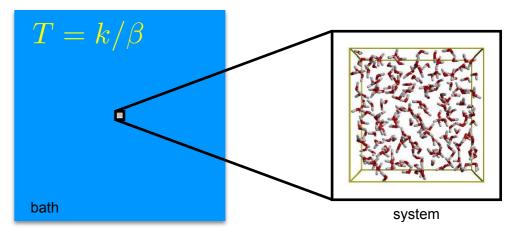
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#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



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Taylor expansion of  $S_{\text{bath}}$  around  $E^{\text{tot}}$ 

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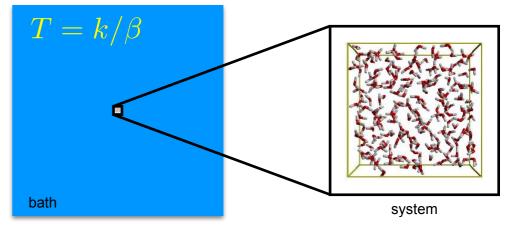
define

$$\beta \equiv \frac{1}{kT}$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



probability of micro-state i proportional to number of micro-states of bath

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define

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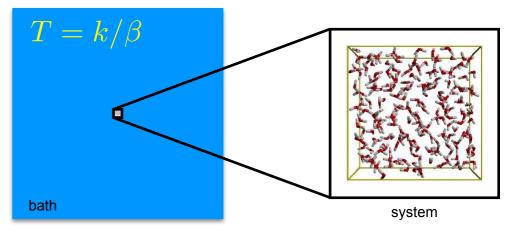
so that

$$\frac{1}{k}S_{\text{bath}}(E^{\text{tot}} - E_i) = \frac{1}{k}S_{\text{bath}}(E^{\text{tot}}) - \beta E_i$$

#### canonical ensemble

system in thermal equilibrium with bath

micro-states of system



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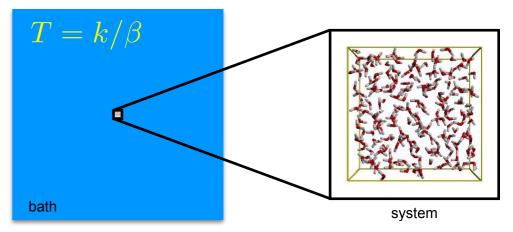
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system in thermal equilibrium with bath

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probability of micro-state i

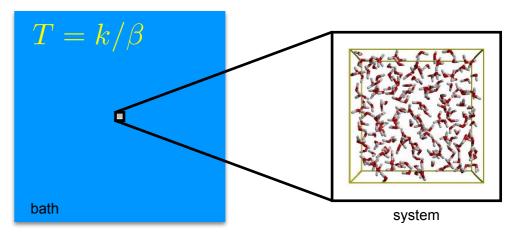
Boltzmann distribution

$$p_i = \frac{1}{Z}e^{-\beta E_i}$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



probability of micro state i proportional to number of micro-states of bath

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probability of micro-state i

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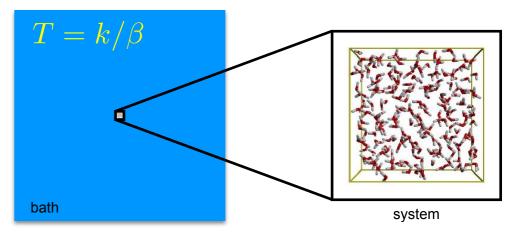
partition function

$$Z = \sum_{i} e^{-\beta E_i}$$

#### canonical ensemble

system in thermal equilibrium with bath

## micro-states of system



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probability of micro state i

Boltzmann distribution

$$p_i = \frac{1}{Z}e^{-\beta E_i}$$

partition function

$$Z = \sum_{i} e^{-\beta E_i}$$

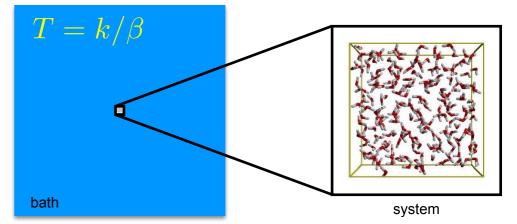
from microscopic to macroscopic with partition function

canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro-state i

$$p_i = \frac{1}{Z}e^{-\beta E_i} \qquad Z = \sum_i e^{-\beta E_i}$$



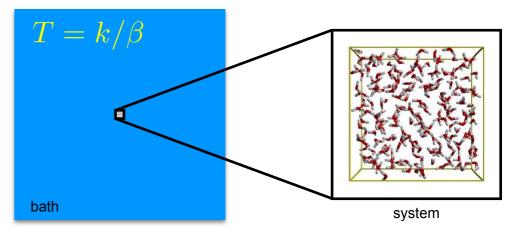
#### canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro-state i

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$$\langle E \rangle =$$



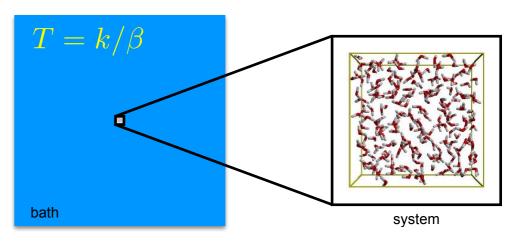
#### canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro-state i

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$$\langle E \rangle = \sum_{i} p_i E_i$$



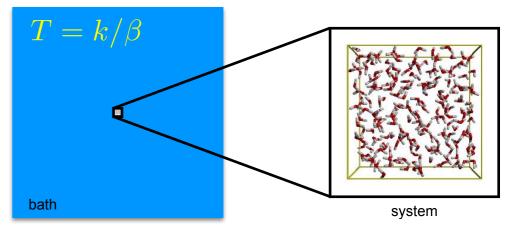
#### canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro-state i

$$p_i = \frac{1}{Z}e^{-\beta E_i} \qquad Z = \sum_i e^{-\beta E_i}$$

$$\langle E \rangle = \sum_{i} p_{i} E_{i} = \frac{\sum_{i} E_{i} e^{-\beta E_{i}}}{\sum_{i} e^{-\beta E_{i}}}$$



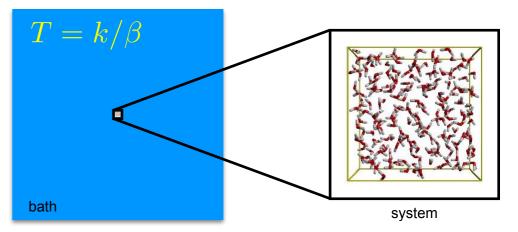
#### canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro-state i

$$p_i = \frac{1}{Z}e^{-\beta E_i} \qquad Z = \sum_i e^{-\beta E_i}$$

$$\langle E \rangle = \sum_{i} p_{i} E_{i} = \frac{\sum_{i} E_{i} e^{-\beta E_{i}}}{\sum_{i} e^{-\beta E_{i}}} = -\frac{\partial \ln Z}{\partial \beta}$$



#### canonical ensemble

system in thermal equilibrium with bath entropy of system

probability of micro state i

$$p_i = \frac{1}{Z}e^{-\beta E_i} \qquad Z = \sum_i e^{-\beta E_i}$$

average energy of system

$$\langle E \rangle = \sum_{i} p_{i} E_{i} = \frac{\sum_{i} E_{i} e^{-\beta E_{i}}}{\sum_{i} e^{-\beta E_{i}}} = -\frac{\partial \ln Z}{\partial \beta}$$

what about entropy?

