
**PHYSICS OF ELEMENTARY PARTICLES
AND ATOMIC NUCLEI. EXPERIMENT**

Inelastic Scattering and Clusters Transfer in $^{3,4}\text{He} + ^9\text{Be}$ Reactions¹

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Abstract—A study of inelastic scattering and multi-particle transfer reactions was performed by alpha and ^3He beams on a ^9Be target at energy about 50 MeV. Angular distributions of the differential cross sections for the $^9\text{Be}(\alpha, \alpha')^9\text{Be}^*$, $^9\text{Be}(\alpha, ^3\text{He})^{10}\text{Be}$, $^9\text{Be}(\alpha, t)^{10}\text{B}$, $^9\text{Be}(^3\text{He}, ^6\text{Li})^6\text{Li}$ and $^9\text{Be}(^3\text{He}, ^6\text{Be})^6\text{He}$ reactions were measured. Experimental angular distributions of the differential cross sections for the ground state and a few low-lying states were analyzed in the framework of the optical model, coupled channels and distorted-wave Born approximation. The information on the cluster structure of the reaction products are obtained. An analysis of the obtained spectroscopic factors was performed.

DOI: 10.1134/S1547477115050052

1. INTRODUCTION

In last years the study of light weakly-bound nuclei [3] was intensified due to the significant progress made with radioactive beam facilities. It has been shown that in light nuclei the nucleons tend to group into clusters, whose relative motion mainly defines the properties of these nuclei. Along with unstable nuclei the interest in the study of light stable nuclei has been renewed. The cluster structures of the ground as well as low-lying excited states of these nuclei are in the focus of studies. As examples, nuclei ^6Li and ^7Li are both well described by two-body cluster models ($\alpha + d$ and $\alpha + t$, respectively). Another interesting nuclide is ^9Be , which is usually treated as an $\alpha + n + \alpha$ three-body configuration, at the same time one may also suppose different two-body configuration $^8\text{Be} + n$ or $^5\text{He} + \alpha$ for the excited states of ^9Be . The addition of a second valence neutron to ^9Be leads to another interesting nucleus, ^{10}Be where the $^6\text{He} + \alpha$ and the $^5\text{He} + ^5\text{He}$ configurations together with the $\alpha + n + n + \alpha$ one may play important role.

Recently, special attention has been focused on the role of the extra “valence” nucleons, and their influ-

ence on the cluster structure of the excited states. This subject is discussed for example in Ref. [4], where the two-center molecular states in ^9B , ^9Be , ^{10}Be , and ^{10}B nuclei were considered in the framework of a molecular-type model.

The standard tool to study nuclear structure is the scattering of the projectiles like the protons or $^{3,4}\text{He}$ by a target nucleus, the structure of which is going to be studied. This method is based on the angular-distribution measurements of the projectile-like products with fixed excitation energy in the (in)elastic and transfer reaction channels.

This article is an attempt to shed light on the internal structure of $^{9,10}\text{Be}$ and ^{10}B nuclei as well as on the mechanism of the nucleon-cluster transfers by the study of the $^{3,4}\text{He}$ induced reactions with the ^9Be target. We expected that the sensitivity of the high precision scattering data to the cluster nuclear structure of light nuclei could be demonstrated. In this work we supplement the data measured before [5–7] with the new ones extending the ranges of the exciting energy and scattering angle of the products. The performed data analysis is based on the optical model, distorted-wave Born approximation (DWBA) and coupled-channel (CC) approach.

¹ The article is published in the original.

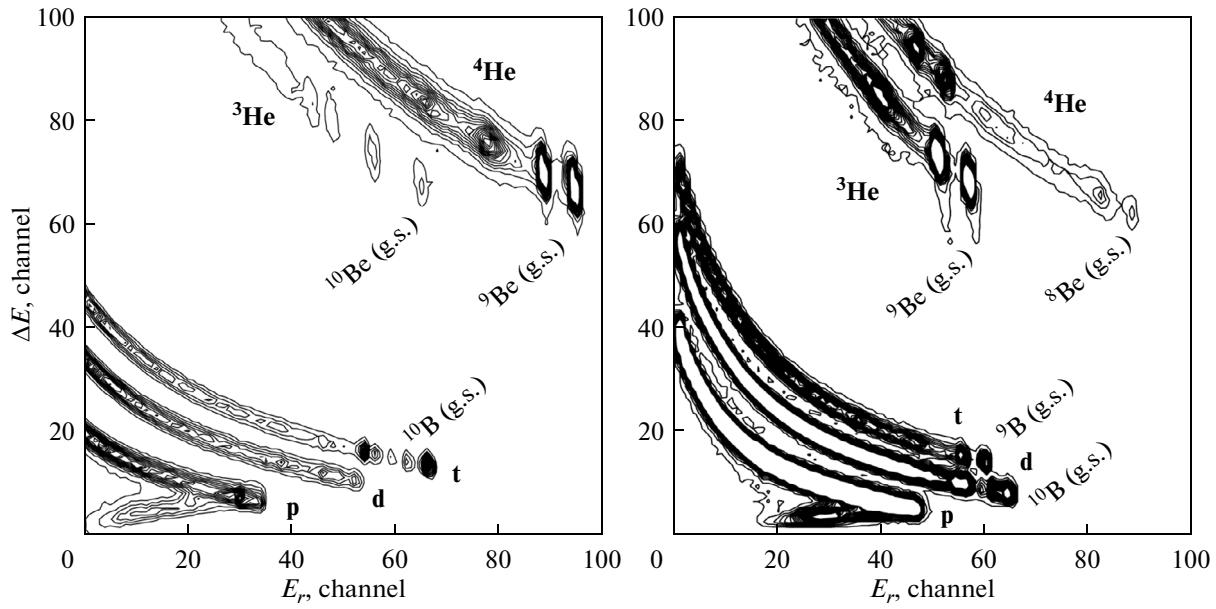


Fig. 1. ${}^4\text{He} + {}^9\text{Be}$ (left panel) ${}^3\text{He} + {}^9\text{Be}$ (right panel) reactions product yields versus energy loss ΔE and residual energy E_r , measured by the Si-Si(Li) telescope. The loci for ${}^3, {}^4\text{He}$, p, d and t are indicated.

2. EXPERIMENTAL PROCEDURE

2.1. $\alpha(63 \text{ MeV}) + {}^9\text{Be}$ Measurements

First run of the $\alpha + {}^9\text{Be}$ experiments was performed at the K130 Cyclotron facility of the Accelerator Laboratory of the Physics Department of Jyväskylä University and in the Nuclear Physics Institute (NPI), Řež, Czech Republic. The beam energy of ${}^4\text{He}$ ions was 63 MeV. The average beam current during the experiment was maintained at 3 nA. The self-supporting Be target was prepared from a 99% pure thin foil of beryllium. The target thickness was 12 μm . Peaks due to carbon and oxygen contaminations were not observed in the energy spectra.

To measure (in)elastically scattered ions, two telescopes each consisting of Si-Si(Li) detectors with thicknesses of 100 μm and 3 mm, respectively, were used. Each pair of detectors was mounted at a distance of about 45 cm from the target. Particle identification was performed based on the energy-loss measurements of ΔE and residual energy E_r , i.e. the so-called $\Delta E - E$ method. The Si-telescopes were mounted on rotating supports, which allowed to obtain data from $\theta_{lab} = 20^\circ$ to $\theta_{lab} = 107^\circ$ in steps of $1^\circ - 2^\circ$.

The overall energy resolution of the telescopes was nearly 200 keV. An example of two-dimensional plot (yield versus energy loss ΔE and residual energy E_r , measured by Si-Si(Li) detectors) is shown in the left panel of Fig. 1. Excellent energy resolutions of both ΔE and E detectors allowed identifying ${}^3, {}^4\text{He}$, t, d and p unambiguously.

The channel leading to the production of ${}^7\text{Be} + {}^6\text{He}$ has minimal probability due to the low Q-value. Other

reaction channels take places at higher Q-values and consequently have larger cross sections, as it is shown in Fig. 1. The production yield of ${}^6\text{He}$ starts to be visible only when plotting the z-axis (yield) in logarithmic scale; it is not shown in Fig. 1.

Comparing with the experimental technique of Ref. [6] we have the advantage to distinguish the particles p, d, t, ${}^3\text{He}$ and ${}^4\text{He}$ and determine their total deposited energies. The total energies were obtained after energy calibration of all Si-detectors and summing of energy deposits in the ΔE and E_r detectors (see Fig. 2). All the observed peaks were identified and found to belong to the ground and excited states of ${}^9\text{Be}$, ${}^{10}\text{Be}$ and ${}^{10}\text{B}$, as the complementary products to detected particles ${}^4\text{He}$, ${}^3\text{He}$, and t, respectively.

We found excellent agreement in the identification of the excited states observed in our experiment with those previously measured for ${}^9\text{Be}$ [5, 6], ${}^{10}\text{Be}$ [8, 9], and ${}^{10}\text{Be}$ [8, 10]. Because the incident beam energy was rather high (about 15 MeV/u), the observed states are most likely populated in one-step direct transfer reactions.

2.2. ${}^3, {}^4\text{He}(30, 40 \text{ MeV}) + {}^9\text{Be}$ Measurements

Next two runs of the experiment were performed at the K130 Cyclotron facility (Jyväskylä University) and later using the isochronous cyclotron U-120M Nuclear Physics Institute (NPI, Řež, Czech Republic). The beam energy of ${}^4\text{He}$ ions was 39 MeV (Rez) and 30 MeV and 40 MeV in the case of ${}^3\text{He}$ ion beams (Rez and Jyväskylä, respectively). The experiments were performed at sufficiently high energy to ensure sup-

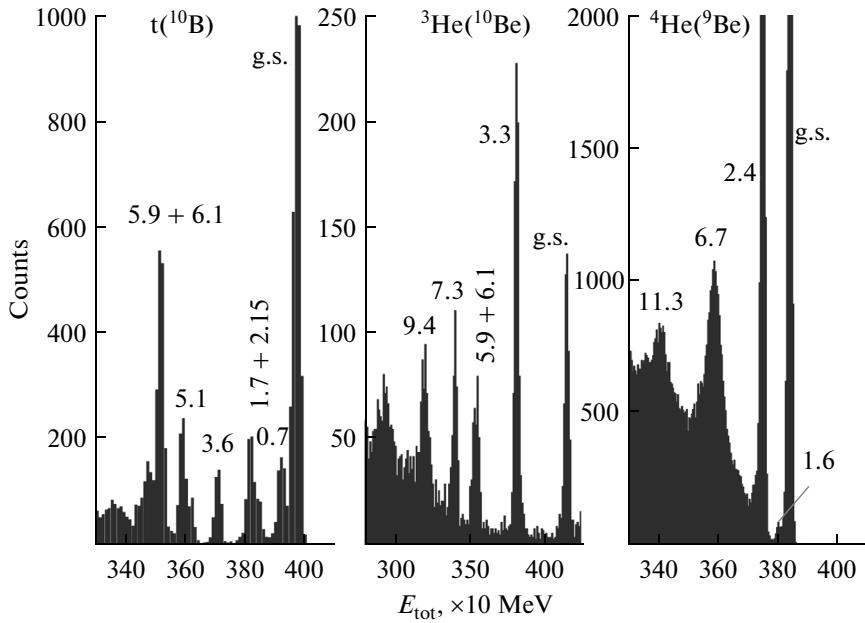


Fig. 2. Measured spectra of total energies for ${}^9\text{Be}(\alpha, t){}^{10}\text{B}$ (left), ${}^9\text{Be}(\alpha, {}^3\text{He}){}^{10}\text{Be}$ (middle) and ${}^9\text{Be}(\alpha, \alpha'){}^9\text{Be}$ (right) reaction channels. The ground and most of the populated excited states of ${}^{10}\text{B}$, ${}^{10}\text{Be}$ and ${}^9\text{Be}$ are unambiguously identified.

pression of compound-nucleus contributions. The average beam current during the experiment was maintained at 10 nA. The self-supporting enriched (99%) Be target with 12 μm thickness was used.

The experimental conditions, detector system and particle identification method were the same or similar to that ones realized in first run. The $\Delta E - E$ distribution for the experiment performed with the ${}^3\text{He}$ beams is shown in right panel of Fig. 1. This plot demonstrates that the energy spectrums of the detected ejectiles are the fingerprint of the complementary nuclei, which are the object of our interest.

3. RESULTS

3.1. Elastic and Inelastic Scattering

Measured angular distributions of differential cross sections of the ground and low-lying excited states for ${}^9\text{Be}$ are presented in Fig. 3. Due to low statistics we were not able to get the angular distribution for the first-excited $1/2^+$ state of ${}^9\text{Be}$ at 1.6 MeV.

Comparison with the results for the ground state of the previous measurements [6] (open symbols) demonstrates a good agreement at small scattering angles. The disagreement is observed at angles larger than 70° where our data are smaller than those of Ref. [6]. From the technical point of view, this difference could be explained by absence of particle identification in Ref. [6] where Si(Li) detectors were used to measure the total energy only, without a ΔE measurement that would allow Z and A identification of the detected particles. Another reason could also be due to a different method used for subtraction of the continuum

under the peak. The same reasons are responsible for the difference between our data and those of Ref. [6] for the level at 6.76 MeV in the angular range $30-60$ degrees (see Fig. 3).

Figure 3 shows measured differential cross sections for the elastic and inelastic scattering (symbols) together with the results of theoretical calculations (curves) performed within optical model and coupled-channel approach. Theoretical curves were obtained with the aid of NRV server optical model code [11] and the ECIS06 coupled-channel code [12, 13].

Firstly, let us consider the analysis of the elastic scattering cross section. The optical potential was chosen in the usual Woods-Saxon form

$$V(r) = -V_0 f(r, R_v, a_v) - i W_0 f(r, R_w, a_w),$$

where the function $f(r, R, a) = (1 + e^{(r-R)/a})^{-1}$ and radii $R_i = r_i A^{1/3}$ depend on the mass of the heavier fragment A and corresponding reduced radius r_i . Potential parameters fitted within the optical model to the measured experimental data are listed in Table 2. Corresponding curve is shown in Fig. 3 as a dashed line and demonstrates good agreement with obtained data. Note that the optical potentials recommended in previous studies [5, 6] do not provide a proper description of the elastic scattering data, since they were obtained by the fitting the data in narrower angular range. Thus in subsequent calculations in order to describe the entrance channel wave function we used the optical potential with parameters listed in Table 1.

The simplest view of the ${}^9\text{Be}$ nucleus is that it is the Borromean system, i.e. strongly deformed three-body

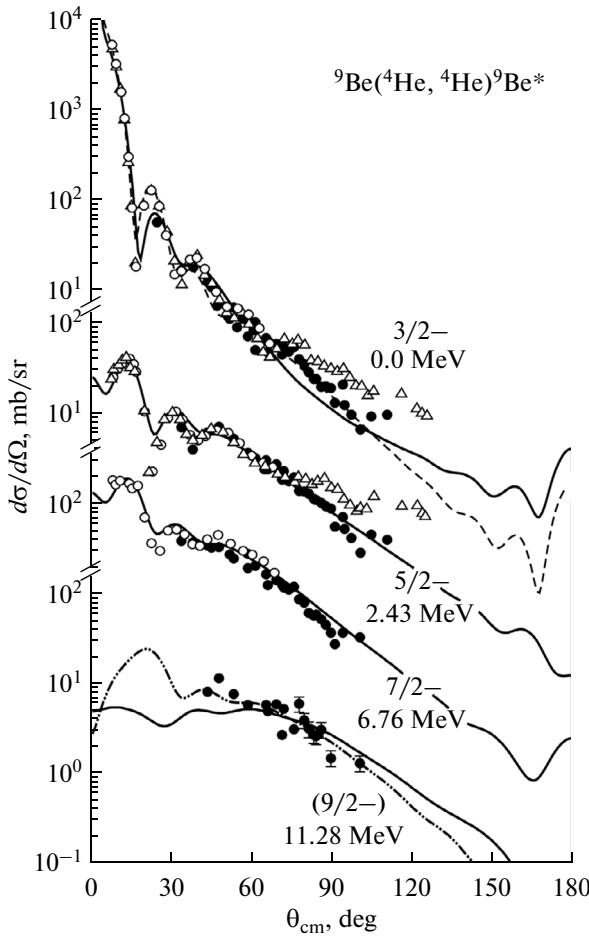


Fig. 3. Angular distributions of the differential cross sections for the ${}^4\text{He}(63 \text{ MeV}) + {}^9\text{Be}$ elastic and inelastic scattering. Data obtained in the present work are shown by ● symbols, data from Refs. [4, 5] are denoted by ○ and △ symbols, respectively. Elastic scattering cross section was considered within optical model and shown by the dashed curve. Theoretical analysis (solid and dash-dotted curves) was performed within the coupled-channel approach assuming the rotational character of the excitations. Details of the calculation see the text.

system consisting of two α particles held together by a weakly bound neutron. It is very natural that different molecule-like states may appear in the excited states. The aim of the present experiment was to study the peculiarity of the angular distributions of elastic and inelastic scattering, mainly for $5/2^-$, $7/2^-$ and $9/2^-$ states, trying to learn more about their cluster structure.

Analysis of inelastic scattering data within the DWBA or CC approach allows to extract the information on the deformation of an excited nucleus treating these states as collective rotational excitations. Corresponding coupling matrix elements in addition to the radial form-factor includes the deformation length $\beta_\lambda R_V$, where quantity β_λ is a deformation parameter, λ is a multipolarity of the transition defined by the transferred angular momentum and $R_V = r_V A^{1/3}$ is an inter-

action radius depending on the mass A of excited nucleus.

It is known [15–17] that ${}^9\text{Be}$ has a rotational band ($K^\pi = 3/2^-$) built on its ground state. In previous studies only ground ($3/2^-$) and ($5/2^-$) excited states of the band were analyzed together in the CC framework. One may expect [18, 19] that all the angular distributions shown in Fig. 3 are related to the same rotational band. Note that the values of spin and parity of the 11.28 MeV state are uncertain. This level was listed either $7/2^-$ or $9/2^-$ state in the literature and databases. Following Ref. [16], we consider this state to belong to the rotational band and therefore to have spin-parity $9/2^-$.

The solid lines in Fig. 3 represent the results of a CC calculation within the symmetric rotational model taking also into account Coulomb excitation and reorientation terms. The ECIS06 code was employed. The parameters of the optical potential used in the CC calculations are given in Table 1 (marked as CC). They were fitted to the data shown in Fig. 3, using the optical model parameters as an initial set. It was found that inelastic scattering data for the first three states of the rotational band may be well described if one assumes $\beta_\lambda R_V = 1.574 \text{ fm}$ and $\beta_2 = 0.64$.

Quadrupole moment Q_{20} of the ${}^9\text{Be}$ nucleus is known to be equal $+53 \text{ mb}$ [20, 21] indicating a prolate deformation for the ground state. Previous studies (e.g. Refs. [5, 6]) have shown a quite large deformation parameter β_2 lying in the range 0.5 to 0.7. It provided rather good agreement with our data on elastic and inelastic (2.43 MeV and 6.76 MeV states) scattering. The obtained large β_2 value may be considered as the confirmation of the cluster structure of the low-lying states of ${}^9\text{Be}$. However it doesn't allow to give unambiguous preference to the one of the possible configurations, for example, $(\alpha + \alpha + n)$ or $(\alpha + {}^5\text{He})$.

In Fig. 3, one may see quite poor agreement between CC calculation and the experimental data (see solid line in the bottom part of Fig. 3) in the case of 11.28 MeV state attributed as $9/2^-$ rotational state. It may testify either the fallacy of the assumption on the $9/2^-$ state nature or may indicate the different structure of this state. In order to improve the fits for this state, an additional hexadecapole term β_4 in the definition of the ${}^9\text{Be}$ radius was added. The dash-double-dotted line in Fig. 3 demonstrates the result obtained with the same β_2 value and $\beta_4 = 0.27$. It agrees much better with the data. There is insignificant influence of the β_4 parameter on the cross sections for the $3/2^-$, $5/2^-$ and $7/2^-$ states. This may be evidence of the different structure of the $9/2^-$ state of the ${}^9\text{Be}$ nucleus. It should be noted that data on inelastic scattering to the 11.28 MeV state were measured in the middle range of the angles, where two theoretical predictions are quite close. Thus, in order to draw final conclusion, additional measurements are required in a broader region of the scattering angles.

Table 1. Potential parameters used within optical model and CC approaches

	V_0 , MeV	r_v , fm	a_v , fm	W_0 , MeV	r_w , fm	a_w , fm	r_C , fm
$\alpha + {}^9\text{Be}$	101.0	1.40	0.75	32.70	1.50	0.75	1.3
$\alpha + {}^9\text{Be(CC)}$	96.8	1.19	0.75	11.84	1.61	0.75	1.3
${}^3\text{He} + {}^{10}\text{Be}^a$	95.0	0.95	0.82	8.00	1.60	0.73	1.1
$t + {}^{10}\text{B}^b$	95.0	1.04	0.82	3.00	1.87	0.47	1.1
${}^3\text{He} + {}^9\text{Be}$	106.2	1.08	0.88	20.63	1.78	0.89	1.3
${}^6\text{Li} + {}^6\text{Li}$	81.9	1.08	0.96	11.72	2.38	0.89	1.2

^a Parameters were taken from Ref. [14]. Real and imaginary depths and radii were modified within 10–15% of magnitude in order to fit experimental data on transfer to the ${}^{10}\text{Be}$ ground state.

^b The ${}^3\text{He} + {}^{10}\text{Be}$ parameters (were taken from Ref. [14]) were used as initial set and then parameters were fitted to reproduce experimental data on transfer to the ${}^{10}\text{B}$ ground state.

3.2. ${}^9\text{Be}({}^4\text{He}, {}^3\text{He}){}^{10}\text{Be}$ Reactions

If ${}^9\text{Be}$ shows molecular cluster structure [4], then ${}^{10}\text{Be}$ might be expected to show more sophisticated internal structure. Molecular structure of the ${}^{10}\text{Be}$ nucleus is formed by two alpha particles and two neutrons. Such constitution attracts even more interest, since one neutron added to ${}^9\text{Be}$ makes the ${}^{10}\text{Be}$ nucleus as more compact configuration [15].

In this work we performed measurements of the angular distributions for the ${}^9\text{Be}({}^4\text{He}, {}^3\text{He}){}^{10}\text{Be}$ reaction, leading to different ${}^{10}\text{Be}$ excited states. Angular distributions of the differential cross sections for the ground and low-lying excited states for ${}^{10}\text{Be}$ are plotted in the left panel of Fig. 4. Results of the present experiment are shown by solid symbols; data from Ref. [6] are presented by the open symbols. Solid lines are the result of the finite-range DWBA calculations with help of the DWUCK5 code [22] available via the Internet on the site of the NRV project [23].

In order to perform the DWBA calculations the potential parameters listed in Table 1 were chosen to calculate distorted waves in the entrance and exit channels. The potential parameters for the exit channel ${}^3\text{He} + {}^{10}\text{Be}$ were chosen close to the potential recommended in Ref. [14]. The Q-values for the ${}^9\text{Be}({}^4\text{He}, {}^3\text{He}){}^{10}\text{Be}$ reaction channel is -13.8 MeV, that reduce the relative energy in the exit channel significantly. It legitimizes a slight variation of the optical model parameters for the exit channel (within 10%) for better agreement of the calculations with the data. In the analysis reported below we varied only the depths of the real and imaginary parts within indicated limits.

The single-particle wave functions in the entrance and exit channels were defined within standard potential model [24, 25]. The interaction for $n + {}^3\text{He}$ system was chosen of the Gaussian form

$$V(r) = -V_G \exp\left(-\frac{r^2}{R_G^2}\right),$$

where the radius $R_G = 2.452$ fm [25], while the potential depth V_G is fitted to reproduce the correct value of neutron binding energy $E_n = -20.58$ MeV in the ${}^4\text{He}$ nucleus. The $n + {}^9\text{Be}$ potential in the final state was defined as a real Woods–Saxon potential with radius $R_V = 1.26 A_{\text{Be}}^{1/3}$ fm and diffuseness $a_V = 0.6$ fm. Potential depth V_0 was defined in the same manner as V_G .

Table 2. Spectroscopic information for the ${}^9\text{Be}({}^4\text{He}, {}^3\text{He}){}^{10}\text{Be}$ and ${}^9\text{Be}({}^4\text{He}, t){}^{10}\text{B}$ reactions as obtained from the DWBA analysis

E_x , (MeV)	$J\pi$	l	S_f , Ref. [5]	S_f , present
${}^9\text{Be}(\alpha, {}^3\text{He}){}^{10}\text{Be}$				
g.s.	0^+	1	1.58	1.65
3.368	2^+	1	0.38	1.00
5.958	2^+	1	≤ 0.73	≤ 1.40
5.960	1^-	2	≤ 0.14	≤ 0.43
6.179	0^+	1	—	—
6.263	2^-	2	0.08	≤ 0.26
7.371	3^-	2	0.26	0.28
7.542	2^+	1	—	—
9.27	(4^-)	2	≤ 0.18	0.10
9.56	2^+	1	—	0.23
${}^9\text{Be}(\alpha, t){}^{10}\text{B}$				
g.s.	3^+	1	0.89	0.59
0.781	1^+	1	1	1.0
1.76	0^+	1	1.58	1.38
2.1	1^+	1	0.52	0.30
3.6	2^+	1	0.28	0.23
5.11	2^-	2	≤ 0.27	≤ 0.16
5.16	2^+	1	≤ 1.85	≤ 0.75
5.18	1^+	1	≤ 3.14	≤ 1.0
5.93	2^+	1	0.48	≤ 0.95
6.13	3^-	2	0.24	≤ 0.19

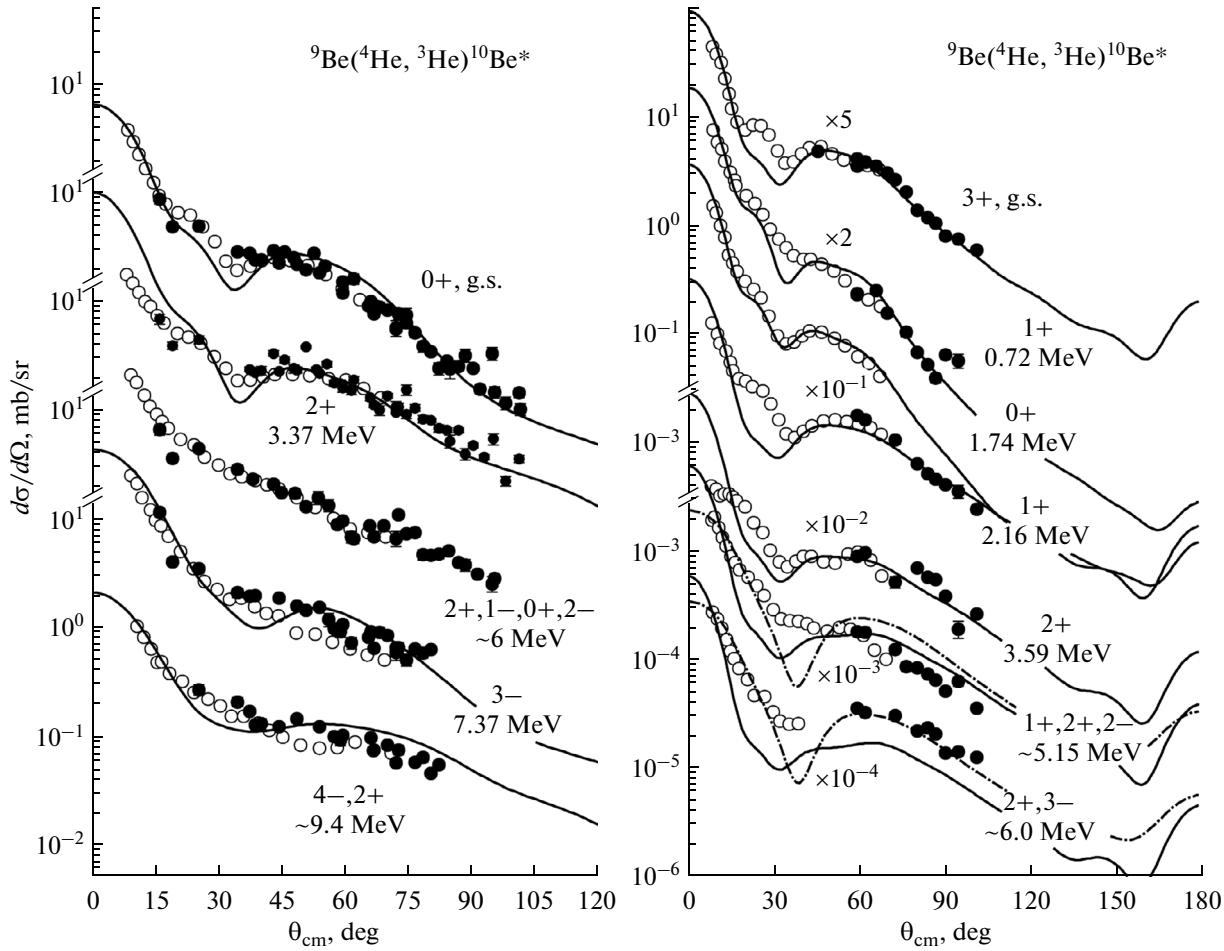


Fig. 4. Angular distributions of the differential cross sections for the ground and low-lying excited states of the ^{10}Be (left panel) and ^{10}B (right panel) nuclei formed in the $^4\text{He} + ^9\text{Be}$ reaction. Results of the present experiment are shown by solid symbols. Data from Ref.[5] are shown by open symbols. The curves demonstrate the calculation results explained in the text.

parameter. For states unbound to the neutron emission in ^{10}Be , the neutron was assumed to be bound by 0.1 MeV, as it is usually recommended [5].

Relative angular momentum of neutron state in the projectile or target-like fragment was fixed by the total momentum J and parity π conservation laws. In particular, the ground state of the $^{10}\text{Be}(0^+) = n(1/2^+) + ^9\text{Be}(3/2^-)$ nucleus was considered as $1p_{3/2}$ neutron state, while the excited states of ^{10}Be with negative parity were treated as $1d_{5/2}$ neutron states. All spectroscopic properties of the ^{10}Be excited states are listed in Table 2.

The DWBA differential cross section for the considered stripping reactions can be compared with experimental data in the following way [22]

$$\frac{d\sigma_{\text{exp}}}{d\Omega} = S_f \frac{(2J_f + 1)}{(2J_i + 1)} \sigma_{DW}(q),$$

where $J_i = 3/2$ and J_f are the angular momenta of the ^9Be target and the final state populated in ^{10}Be , respectively, $\sigma_{DW}(\theta)$ is the output from DWUCK5, S_f is the

targetlike fragment spectroscopic factors. Figure 4 demonstrates how good are the obtained absolute values of the spectroscopic factors S_f , which were obtained from the comparison of the measured angular distributions and DWUCK5 calculations for the different ^{10}Be final states. The values of the spectroscopic factors are listed in Table 2 together with S_f reported in Ref. [5].

It is seen that obtained cross sections agree well with data. The spectroscopic factors extracted from our analysis are very close to the ones listed in Ref. [5] (see Table 2), except for the ^{10}Be state at 3.368 MeV, where spectroscopic values differ by more than a factor two. The reason for this discrepancy is the following. The spectroscopic factor in our work was defined by adjusting the theoretical curve to the data measured in middle angle domain, while in Ref. [5] it was fitted to the forward experimental points near $\theta_{cm} \approx 10^\circ$.

In spite of the high energy resolution, we were not able to separate the excited states nearby 6 MeV. Low statistics did not allow to observe the 2^+ state at

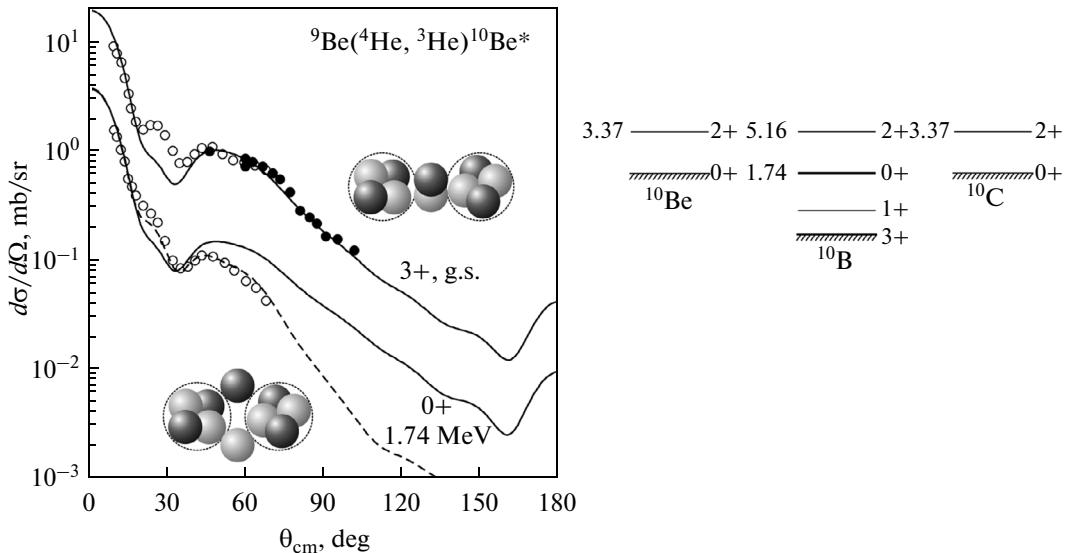


Fig. 5. On the left side the differential cross sections for reaction ${}^9\text{Be}({}^4\text{He}, t){}^{10}\text{B}$ leading to the ground and 1.74 MeV excited states of ${}^{10}\text{B}$ nucleus. Solid symbols show the results of the present experiment, open ones is the data from Ref. [5]. Curves show the results of DWBA calculations (see text for details). On the right side the level diagrams of low-lying states of members of the $A = 10$ isobaric multiplet are shown.

7.54 MeV. The experimental cross sections corresponding to two overlapping states near 9.5 MeV were described as a sum of DWUCK5 outputs multiplied by the corresponding spectroscopic factors. Table 2 contains the S_f values providing the best fit.

Note that in order to describe the data one needs much smaller radius of the real part of the optical potential for the exit channel ($r_V = 0.95$ fm) in comparison to the radius in the entrance channel ($r_V = 1.40$ fm) that indicates the compactness of the ${}^{10}\text{Be}$ nucleus.

Nuclear charge radii of ${}^{7,9,10}\text{Be}$ have been measured by high precision laser spectroscopy [18, 19]: the charge radius decreases from ${}^7\text{Be}$ to ${}^{10}\text{Be}$. Comparing the Coulomb parameter r_C with that of ${}^9\text{Be}$, we obtained a smaller value of r_C for ${}^{10}\text{Be}$. In Ref. [18, 19], the decrease was explained as probably caused by the clusterization of ${}^7\text{Be}$ into an α and triton clusters, whereas ${}^{9,10}\text{Be}$ were considered to be $\alpha + \alpha + n$ and $\alpha + \alpha + n + n$ systems, respectively, and were more compact. The experimental trend was shown [15], to change beyond ${}^{10}\text{Be}$ with an increase of the charge radius with atomic mass. Furthermore, the large experimental value of the charge radius for ${}^{12}\text{Be}$ is consistent with a breakdown of the $N = 8$ shell closure.

3.3. ${}^9\text{Be}({}^4\text{He}, t){}^{10}\text{B}$ Reactions

Differential cross sections versus angles in cm-system for the ground and low-lying excited states of ${}^{10}\text{B}$ are plotted in the right panel of Fig. 4. Results of the present experiment are shown by solid symbols, and data from Ref. [5] are presented as open symbols. DWBA calculations [23] for the ${}^9\text{Be}(\alpha, t){}^{10}\text{B}$ reaction

were performed with the DWUCK5 code [22] fitting the differential cross sections for the ground and low-lying states. Results are shown by the solid lines in Fig. 4. One may notice a very good agreement between data obtained in Ref. [5] and our measurements. Theoretical results (solid and dash-dotted curves) fairly reproduce the data in case of well-defined final states. For the unresolved mixture of states at excitation energies of about 5.15 MeV and 6 MeV one may conclude that negative parity states corresponding to $l = 2$ (shown by dash-dotted curves) provide better agreement with data than positive parity ones ($l = 1$, solid curves).

Spectroscopic factors S_f for the different states populated in the reaction ${}^9\text{Be}({}^4\text{He}, t){}^{10}\text{B}$ are listed on the right side in Table 2. For the data corresponding to the mixture of a few levels an upper limit of spectroscopic factor was obtained, describing the data by one component only. S_f values are in good agreement with those reported in the literature.

3.4. Multiplet $A = 10$

The structure of ${}^{10}\text{Be}$, ${}^{10}\text{B}$ and ${}^{10}\text{C}$ nuclei was usually considered as two α -clusters in the presence of two extra nucleons. Level diagrams for the low-energy excited states for these nuclei are shown on the right panel of Fig. 5. One may see that the ${}^{10}\text{B}$ ground state is shifted down by 2 MeV approximately. It may be treated as a three cluster configuration ${}^{10}\text{B} = \alpha + d + \alpha$ where the pairing of proton and neutron results in formation of a deuteron cluster inside. The 3^+ spin of this state also supports this assumption. 1.74 MeV excited state might be considered as a state, where the deuteron cluster

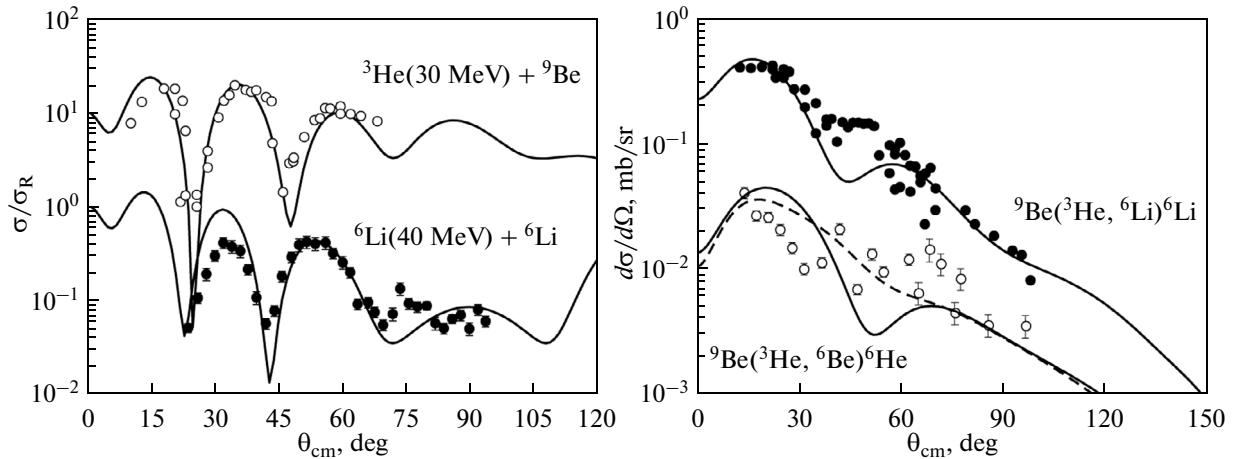


Fig. 6. Left panel: the elastic scattering angular distribution for the ${}^3\text{He} + {}^9\text{Be}$ and ${}^6\text{Li} + {}^6\text{Li}$ reactions. The curves correspond to the optical model fit with parameters from the Table 1. Right panel: differential cross section for the ${}^3\text{H}$ and ${}^3\text{He}$ pickup channels of the ${}^3\text{He} + {}^9\text{Be}$ reaction at the energy 40 MeV. Curves show the calculations within DWBA (see text for details). Experimental data from this work except the ${}^6\text{Li} + {}^6\text{Li}$ data which are from Ref. [26].

becomes unbound. Thus the four clusters configuration ${}^{10}\text{B}(0^+, 1.74 \text{ MeV}) = \alpha + p + n + \alpha$ with uncorrelated proton and neutron manifests itself. Two mirror ground states in ${}^{10}\text{Be}$ and ${}^{10}\text{C}$ in this case have to be of similar structure $\alpha + N + N + \alpha$. One of consequence of such an internal organization is the absence of the di-neutron component in the ${}^{10}\text{Be}$ ground state wave function.

Difference in the structure of the ground state and the 1.74 MeV state in ${}^{10}\text{B}$ may also reveal itself in the difference of optical potentials for these exit channels. In Fig. 5 the corresponding experimental data are compared with the results of DWBA calculations performed in the same manner as for the ${}^3\text{He} + {}^{10}\text{Be}$ exit channel. Solid curves show theoretical cross sections obtained with the exit channel optical potential from Table 1. This potential was chosen on the basis of OM potential compilation from Ref. [14] with additional adjustment of parameters to the present data since Ref. [14] contains recommended optical potential for the lower collision energies. One may see quite good agreement between calculation and data on the case of proton transfer, leading to the ${}^{10}\text{B}$ ground-state. Applying the same potential for the transfer to the 1.74 MeV state one gets the noticeable overestimation in the cross section at large angles. We found that in order to improve the agreement in the last case it is necessary to use the following parameters: $V_0 = 85 \text{ MeV}$, $r_v = 1.14 \text{ fm}$ and $W_0 = 8 \text{ MeV}$. Corresponding result is shown in Fig. 5 by the dashed curve and demonstrates excellent fit of the data. The obtained parameters turn out to be close to the OM potential for ${}^3\text{He} + {}^{10}\text{Be}(\text{g.s.})$ channel.

3.5. Cluster Transfer in ${}^3\text{He} + {}^9\text{Be}$ Reaction

Here we report the data on the ${}^3\text{He} + {}^9\text{Be}$ interactions followed with the triton and ${}^3\text{He}$ pickup at the

incident energy 40 MeV. These reaction channels are of specific interest since provide the direct information on the cluster structure of initial and final states.

First of all we fit the optical model parameters in order to describe the elastic scattering data. The data for the entrance channel were measured during the reported experiment, while the data for the ${}^6\text{Li} + {}^6\text{Li}$ elastic scattering were found in [26]. Obtained potential parameters are listed in Table 1. In the left panel of Fig. 6 the elastic scattering cross sections are shown. One may see the acceptable agreement between data and optical model calculations. It is necessary to note one peculiarities of the ${}^6\text{Li} + {}^6\text{Li}$ optical potential which has a very expanded imaginary part. Indeed the radius R_w for this combination is more than 4 fm, while the real part has radius less than 2 fm. The optical potentials fitting the data at different energies [26] demonstrate the same properties. There is no clear explanation of this fact.

In order to perform the DWBA calculation of the transfer reactions one defined the bound states parameters in the entrance and exit channels. The ${}^9\text{Be}(3/2^-)$ ground state is treated now as a two-body system

$${}^3\text{H}\left(\frac{1}{2}^+\right) + {}^6\text{Li}(1^+)$$

with binding energy $\varepsilon_b = -17.7 \text{ MeV}$. According to the selection rules on the momentum coupling and the parity conservation we can define the relative state as $2p_{3/2}$ one. The Woods-Saxon potential with parameters $V_0 = -142.6 \text{ MeV}$, $r_v = 1.05 \text{ fm}$, $a_v = 0.5 \text{ fm}$ was applied to describe the ${}^3\text{H} + {}^6\text{Li}$ bound state. In the same manner the bound state ${}^6\text{Li} = {}^3\text{H} + {}^3\text{He}$ in the exit channel is treated. The Gaussian potential with depth $V_0 = -148.6 \text{ MeV}$ and radius $r_v = 1.25 \text{ fm}$ describes the $2s_{1/2}$ state with binding energy $\varepsilon_b = -15.8 \text{ MeV}$.

In the right panel of Fig. 6 the solid dots and corresponding curve show the experimental data and DWBA calculation of the angular distribution for the triton pickup reaction. The figure demonstrate fairly good agreement between data and calculations with spectroscopic factors $S_p S_f = 0.9$, confirming the large contribution of the (${}^3\text{H} + \text{X}$) configuration into the ground state wave functions of the ${}^9\text{Be}$ and ${}^6\text{Li}$ nuclei.

Following the same procedure we calculated the ${}^3\text{He}$ pickup reaction. The bound state parameters for both channels were chosen the same except the binding energies and depth of the potentials. The energy of the ${}^3\text{He} + {}^6\text{He}$ system is -21.2 MeV, that results in the potential depth $V_0 = -151.4$ MeV. The separation energy of the ${}^3\text{He}$ from the ${}^6\text{Be}$ nucleus equal to 11.5 MeV. The solid curve together with open symbol in the right panel of Fig. 6 demonstrates the differential cross section calculated using the same distortion potentials. The obtained spectroscopic factor for this reaction is 0.6 .

One may see from Fig. 6 that in spite of comparable spectroscopic factors the cross sections for the considered reactions are approximately one order of magnitude different. The reason is mainly the Q-values which are -1.9 MeV for the ${}^3\text{H}$ pickup and -9.7 MeV in the case of ${}^3\text{He}$ transfer.

One may see that the data for the ${}^6\text{He} + {}^6\text{Be}$ exit channels are in poor agreement with calculations. Since both nuclei in this exit channel are exotic weakly bound and even unstable it is appropriate to analyze the dependance of the cross section on the model parameters. Since the ${}^6\text{Be}$ nucleus is unstable to the double-proton emission one may expect, in particular, the extended spatial distribution for this nucleus. To simulate it we increase the potential radius for the bound state in the exit channel to the value $r_V = 1.9$ fm (more than 50%), that leads to the depth $V_0 = -79.8$ MeV. Resulting cross section is shown by the dashed curve in Fig. 6. It is very promising to see that the difference in the calculated cross sections are mainly at small scattering angles where the data are measured. It stimulates further investigation of the reactions of such kind with better statistics.

4. CONCLUSIONS

Angular distributions of the differential cross sections for the ${}^9\text{Be}(\alpha, \alpha){}^9\text{Be}^*$, ${}^9\text{Be}({}^4\text{He}, {}^3\text{He}){}^{10}\text{Be}$ and ${}^9\text{Be}({}^4\text{He}, t){}^{10}\text{B}$ reactions were measured. Experimental angular distributions were described within the optical model, coupled channel approach and distorted-wave Born approximation. Spectroscopic factors for the ground and excited states of ${}^{10}\text{B}$ and ${}^{10}\text{Be}$ were deduced. We found pretty good agreement between our results and the previous data.

The performed analysis of the experimental data shows that the potential parameters are quite sensitive to the exit channel and hence to the cluster structures

of the excited states. It allows to make the general conclusions and assumptions on the peculiarities of the internal structure of tested nuclei. To study their cluster structure, a complicated experiment is planned in which decay of excited states by cluster emission will be investigated. However, to distinguish break-up into ${}^4\text{He}$ and ${}^5\text{He}$ will be not be a trivial kinematical problem.

The values $9/2^-$ were assigned to the spin and parity of the 11.28 MeV state in ${}^9\text{Be}$. The obtained large deformation confirms the cluster structure of the low-lying states of ${}^9\text{Be}$. However, it doesn't allow to give unambiguous preference to one of the possible configurations $\alpha + \alpha + n$ or $\alpha + {}^5\text{He}$. In order to improve the agreement between the theoretical prediction and the experimental data, related to the $9/2^-$ state, an additional hexadecapole term is suggested in the definition of the ${}^9\text{Be}$ radius.

Our analysis supports an evidence of the compactness of the ${}^{10}\text{Be}$ ground state. It was found that in order to describe the data one needs a much smaller radius of the real part of the optical potential for the exit channel ($r_V = 0.95$ fm) in comparison to the radius in the entrance channel ($r_V = 1.40$ fm).

The comparison of the angular distributions of the differential cross sections for the isobaric analog states of ${}^{10}\text{Be}$ and ${}^{10}\text{B}$ was done. The structure of ${}^{10}\text{Be}$, ${}^{10}\text{B}$ and ${}^{10}\text{C}$ nuclei was usually considered as two α -clusters in the presence of two extra nucleons. It is shown that the ${}^{10}\text{B}$ ground state could be treated as a three cluster configuration ${}^{10}\text{B} = \alpha + d + \alpha$, where the pairing of proton and neutron results in formation of a deuteron cluster inside ${}^{10}\text{B}$. In the case of the 1.74 MeV excited state the deuteron cluster becomes unbound and system is considered as the four-body configuration, where two α -clusters coexist with an uncorrelated proton and neutron pair.

The cross section of the ${}^3\text{H}$ and ${}^3\text{He}$ transfer were measured in the ${}^3\text{He} + {}^9\text{Be}$ collision at energy of 40 MeV. Proposed optical potential provides good fit of the elastic scattering in the entrance and exit channels. The DWBA calculations are well agreed with the transfer reaction data. The spectroscopic factors of both reactions are close to unit that confirm significant contribution of the considered cluster configurations into the structure of ground states. We show also the possibility to extract the structural information from the comparative analysis of the ${}^9\text{Be}({}^3\text{He}, {}^6\text{Li}){}^6\text{Li}$ and ${}^9\text{Be}({}^3\text{He}, {}^6\text{Be}){}^6\text{He}$ reactions.

ACKNOWLEDGMENTS

We would like to thank the JYFL Accelerator Laboratory and NPI (Řež) for giving us the opportunity to perform this study and the cyclotron staff for the excellent beam quality. This work was supported in part by the Russian Foundation for Basic Research (project numbers: 13-02-00533 and 14-02-91053), the CANAM (IPN ASCR) and by grants to JINR

(Dubna) of the Czech and Polish Republics, and by mobility grant from the Academy of Finland.

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