assignment-4

January 10, 2024

0.1 Assignment 4: Time-dependent problems

0.1.1 SN: 20121702

```
import numpy as np
import matplotlib.pyplot as plt
import numba
from numba import cuda, njit, prange
from mpl_toolkits.mplot3d import Axes3D
import scipy
import time
from math import sqrt
from scipy.sparse import coo_matrix, identity
import pandas as pd
import timeit
from mpl_toolkits.axes_grid1.inset_locator import inset_axes
from scipy.sparse.linalg import spsolve
```

0.2 1. Introduction

In this assignment, we will solve a mathematical problem involving the temperature evolution of a square plate. The plate with some boundary conditions: - Square plate with sides = [-1,1] X [-1,1] - BC = at t = 0 T is u = 5 on one side and u = 0 on the other sides.

To solve this problem, we will use finite difference schemes, both explicit and implicit, analyze their stability, and assess their accuracy as we increase discretization points.

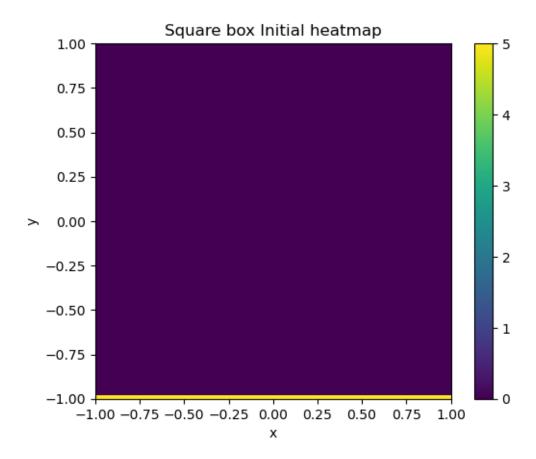
Our goal is to compute the time t^* at which the target temperature is reached, with a precision of 12 digits ($t^* = 0.42401138033$).

Initial Setup of plate

```
[6]: N = 100
u = np.zeros((N,N))
u[-1,:] = 5.0
T_fin = 0.424011387033
plt.imshow(u, extent=(-1,1,-1,1))
plt.xlabel('x')
plt.ylabel('y')
plt.title('Square box Initial heatmap')
```

plt.colorbar()

[6]: <matplotlib.colorbar.Colorbar at 0x7ff3b6312590>



0.2.1 Finite Difference Method

The finite difference method serves as a fundamental numerical technique in solving partial differential equations (PDE's). Among these PDE's, heat conduction problems involve second-order linear equations. One approach to solving these PDEs involves representing the derivative as a field function and then solving for this field function. The Taylor expansion offers an optimal approach for calculating derivatives within this context.

• Heat Equation

$$\frac{\partial u(t)}{\partial t} = \left(\frac{\partial^2 u(x)}{\partial x^2} + \frac{\partial^2 u(y)}{\partial y^2}\right)$$

• Taylor Expand a function u(x) this allows us to approximate a function u(x) around a point, expressing it as a sum of an infinite series of its derivatives:

$$u(x + \Delta x) = u(x) + \frac{1}{1!}u'(x)\Delta x + \frac{1}{2!}u''(x)(\Delta x)^{2} + \dots$$
$$u(x - \Delta x) = u(x) - \frac{1}{1!}u'(x)\Delta x + \frac{1}{2!}u''(x)(\Delta x)^{2} + \dots$$

• Manipulating the system above to find an expression for u''(x) and u''(y)

$$u''(x) = \frac{u(x + \Delta x) + u(x - \Delta x) - 2u(x)}{\Delta x^2}$$

$$u''(y) = \frac{u(y + \Delta y) + u(y - \Delta y) - 2u(y)}{\Delta y^2}$$

• Ignore higher order terms than $\partial u/\partial t$ this simplification results in a more manageable equation:

$$u'(t) = \frac{u(t + \Delta t) - u(t)}{\Delta t}$$

• The heat conduction problem is isotropic in space, this means, that the equations are uniform and Δx and Δy are equal and u(x) and u(y) can be aggregated in u(x,y). therefore we can obtain a propagation formula:

$$\frac{u(x+\Delta x)+u(x-\Delta x)+u(y+\Delta y)+u(y-\Delta y)-4u(x,y)}{\partial \Delta x^2}=\frac{u(t+\Delta t)-u(t)}{\Delta t}$$

To simplify the Courant number will be:

$$C = \frac{\Delta t}{\Delta x^2}$$

Now we can express it in a slightly different form:

$$(1-4C)u_t(i,j) + C[u_t(i+1,j) + u_t(i-1,j) + u_t(i,j+1) + u_t(i,j-1)] = u_{t+1}(i,j)$$

Now, we can determine the field's distribution at the next time step, based on the distribution at the moment before (Explicit method). At the start it will be driven by the initial conditions, but with this equation the distribution will be updated iteratively.

0.3 2.Explicit Implementation

NUMBA edition Numba implementation with explicit time-step

```
[22]: @njit(parallel=True)
def explicit_step_cpu(u0, N, T):
    C = (N - 1) * (N - 1) / (4 * (T - 1))
    u1 = np.copy(u0)
    for i in prange(1, N - 1):
        for j in range(1, N - 1):
```

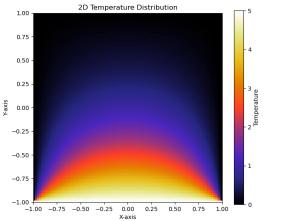
```
u1[i, j] = (1 - 4 * C) * u0[i, j] + C * (u0[i - 1, j] + u0[i + 1, u])
 \rightarrowj] + u0[i, j - 1] + u0[i, j + 1])
    return u1
def explicit method cpu(N,T):
    i = 0
    u = np.zeros((N, N), dtype=np.float64)
    u[0, :] = 5.0
    mid_point_T=[]
    while u[(N-1)//2, (N-1)//2] < 1.0:
        u = explicit_step_cpu(u,N,T)
        mid_point = u[(N-1)//2, (N-1)//2]
        mid_point_T.append(mid_point)
    mid_point_T.append(u[(N-1)//2, (N-1)//2])
    timer = (i / (T - 1))
    print("T(e-cpu): {:.16f}".format(timer))
    return u,timer, mid_point_T
u_ex_cpu, timer_CPU, center_temps = explicit_method_cpu(1011, 80000)
```

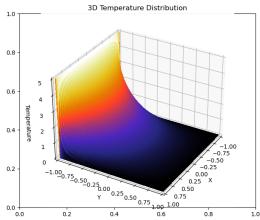
T(e-cpu): 0.0063125789072363

```
[8]: def plot_temperature_distributions(u,azimuthal):
         fig, ax = plt.subplots(1, 2, figsize=(16, 6))
         # 2D Temperature Distribution
         im = ax[0].imshow(u, origin='lower', extent=(-1, 1, -1, 1), cmap='CMRmap')
         fig.colorbar(im, ax=ax[0], label='Temperature')
         ax[0].set_xlabel('X-axis')
         ax[0].set_ylabel('Y-axis')
         ax[0].set_title('2D Temperature Distribution')
         # 3D Temperature Distribution
         x = np.linspace(-1.0, 1.0, u.shape[0])
         y = np.linspace(-1.0, 1.0, u.shape[1])
         X, Y = np.meshgrid(x, y)
         ax[1] = fig.add_subplot(122, projection='3d')
         ax[1].set_xlim(-1, 1)
         ax[1].set_ylim(-1, 1)
         ax[1].set_zlim(0, 5)
         ax[1].set xlabel('X')
         ax[1].set_ylabel('Y')
         ax[1].set_zlabel('Temperature')
         ax[1].view_init(elev=30, azim=azimuthal)
```

```
surf = ax[1].plot_surface(X, Y, u, rstride=1, cstride=1, cmap='CMRmap',
linewidth=0, antialiased=False)
ax[1].set_title('3D Temperature Distribution')
plt.show()

plot_temperature_distributions(u_ex_cpu, 30)
```





Examining the image and output from above, it is noticeable that when the matrix dimension is set to N=101 and the total time steps are M=80000, the time calculated by this numerical scheme is approximately 0.4240, which compared to t*=0.424011387033. This result demonstrates a precision of up to 4 decimal places. Increasing the discretization in both time and space, we suggest that the accuracy will be improve in our solution.

CUDA edition Finite difference method based on CUDA using an explicit time-step.

Employing a CUDA kernel function (explicit_step_gpu) based on the formula derived before. Performing na iterative update calculation of the matrix.

Another function explicit_method_gpu, that takes N and T as inputs. It returns the temperature matrix and computation time to reach a center temperature of 1.0. To optimize performance, we use global spaces U1_device and U2_device for matrix data and U_Central_device to track center temperature.

```
[9]: gpu = cuda.get_current_device()
max_threads_per_block = gpu.MAX_THREADS_PER_BLOCK
print("Maximum threads per block:", max_threads_per_block)
```

Maximum threads per block: 1024

```
[10]: @cuda.jit
      def explicit_step_gpu(u0, u1, uc, N, T, Iter):
          '''CUDA Kernel function'''
          i, j = cuda.grid(2)
          C = ((N - 1) * (N - 1)) / (4 * (T - 1))
          if i > 0 and i < N - 1 and j > 0 and j < N - 1:
              if Iter % 2 == 0:
                   u1[i, j] = (1 - 4 * C) * u0[i, j] + C * (u0[i - 1, j] + u0[i + 1, u])
       uc[Iter] = u1[int((N - 1) / 2), int((N - 1) / 2)] # Center_{l}
       \hookrightarrow temperature
              else:
                  u0[i, j] = (1 - 4 * C) * u1[i, j] + C * (u1[i - 1, j] + u1[i + 1, u])
       \rightarrowj] + u1[i, j - 1] + u1[i, j + 1])
                   uc[Iter] = u0[int((N - 1) / 2), int((N - 1) / 2)] # Center_{l}
       \hookrightarrow temperature
      def explicit_method_gpu(N, T, max_threads=1024):
          # maximum threads--> 1024
          block size = int(sqrt(max threads))
          block_dim = (block_size, block_size)
          n_blocks_x = (N+(block_dim[0]-1))//block_dim[0] # Blocks along the X_
       \hookrightarrow dimension
          n_blocks_y = (N+(block_dim[1]-1))//block_dim[1] # Blocks along the Y_
       \hookrightarrow dimension
          grid_dim = (n_blocks_x, n_blocks_y)
          Iter = 0
          # Intialize arrays on the host
          U0_host = np.zeros((N, N), dtype=np.float64)
          U1_host = np.zeros((N, N), dtype=np.float64)
          U0_host[-1, :] = 5.0
          U1_host[-1, :] = 5.0
          # Device transfer
          U0 device = cuda.to device(U0 host)
          U1_device = cuda.to_device(U1_host)
          UC device = cuda.device array(int(T), dtype=np.float64)
          central_point_temps = []
          for i in range(int(T*0.425)):
              explicit_step_gpu[grid_dim, block_dim](UO_device, U1_device, UC_device, u
       \rightarrowN, T, i)
```

/home/juan/anaconda3/lib/python3.11/site-packages/numba/cuda/dispatcher.py:536: NumbaPerformanceWarning: Grid size 49 will likely result in GPU under-

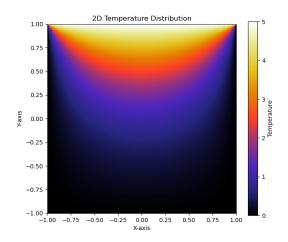
utilization due to low occupancy.

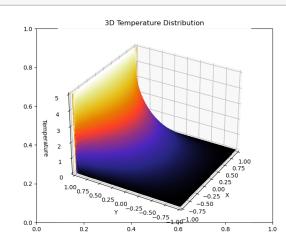
warn(NumbaPerformanceWarning(msg))

T(e-gpu): 0.4240125000000000

Since we found that a bigger discretization value might involve a more precise solution. We use N = 201 (odd number). Indeed its closer to the solution ahrieving a 6 digit precision.

[11]: plot_temperature_distributions(u_ex_gpu, 210)





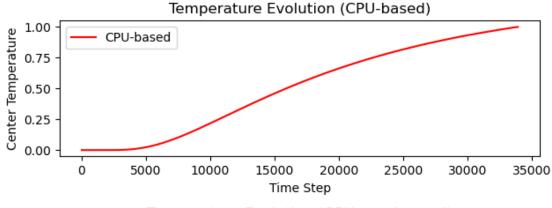
/home/juan/anaconda3/lib/python3.11/site-packages/numba/cuda/dispatcher.py:536: NumbaPerformanceWarning: Grid size 16 will likely result in GPU under-

```
utilization due to low occupancy.
warn(NumbaPerformanceWarning(msg))
```

T(e-gpu): 0.4240125000000000

We can conclude the **GPU** Implementation using CUDA, result in a **more precise solution** achieving up to 4 decimal places precision

```
[12]: fig, ax = plt.subplots(2, 1)
      # Plot center temperatures using the CPU-based method in red
      ax[0].plot(center_temps, color='red', label='CPU-based')
      # Plot center temperatures using the GPU-accelerated method in pink
      ax[1].plot(center_temps_gpu, color='pink', label='GPU-accelerated')
      # Add labels and titles
      ax[0].set_xlabel('Time Step')
      ax[0].set_ylabel('Center Temperature')
      ax[0].set_title('Temperature Evolution (CPU-based)')
      ax[0].legend()
      ax[1].set_xlabel('Time Step')
      ax[1].set ylabel('Center Temperature')
      ax[1].set_title('Temperature Evolution (GPU-accelerated)')
      ax[1].legend()
      # Show the plot
      plt.tight_layout()
      plt.show()
```



Temperature Evolution (GPU-accelerated) 1.00 Center Temperature GPU-accelerated 0.75 0.50 0.25 0.00 Ó 5000 10000 15000 20000 25000 30000 35000 Time Step

0.4 3. Implicit Implementation

This method consists of using the implicit finite difference approach to solve differential equations. Deriving the two-point backwards formula for the first derivative, we find:

$$u'(x) = \frac{u(t + \Delta t) - u(t)}{\Delta t}$$

Additionally, employing backward difference, the implicit time step formula generally becomes:

$$\frac{u_{t+\Delta t}(x+\Delta x)+u_{t+\Delta t}(x-\Delta x)+u_{t+\Delta t}(y+\Delta y)+u_{t+\Delta t}(y-\Delta y)-4u_{t+\Delta t}(x,y)}{\partial \Delta x^2}=\frac{u_{t+\Delta t}-u_{t+\Delta t}(x+\Delta y)+u_{t+\Delta t}(x+\Delta y)+u_{t+\Delta t}(x+\Delta y)+u_{t+\Delta t}(x+\Delta y)}{\Delta t}$$

Rearanged as:

$$u_t(i,j) = -C[u_{t+1}(i+1,j) + u_{t+1}(i-1,j) + u_{t+1}(i,j+1) + u_{t+1}(i,j-1)] + (1+4C)u_{t+1}(i,j)$$

This can be respresented in matrix form:

$$(I - \Delta t A)u_{t+1} = u_t$$

Here, I represents the identity matrix and A is a sparse matrix representation of the laplacian operator in the finite difference scheme (NxN grid size).

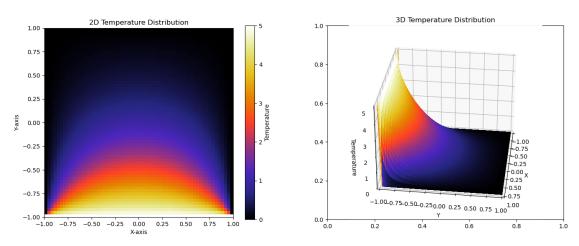
This will iteratively solve the equation Ax = b. We could expect that this will make this mehtod more inefficient since it will be computing a sparrspea matrix vector problem repeatively, compared to the explicit method.

```
[25]: def laplace_coo_matrix(N):
          n_{elements} = 5 * N**2 - 16 * N + 16
          rows = np.empty(n_elements, dtype=np.float64)
          cols = np.empty(n_elements, dtype=np.float64)
          data = np.empty(n_elements, dtype=np.float64)
          h = 2/(N-1)
          iteration = 0
          for j in range(N):
              for i in range(N):
                  # Boundary Conditions
                  if i == 0 or i == N-1 or j == 0 or j == N-1:
                      rows[iteration] = cols[iteration] = j*N + i
                      data[iteration]=1
                      iteration += 1
                  else:
                      rows[iteration:iteration+5] = j*N + i
                      # Diagonal
                      cols[iteration] = j*N + i
                      cols[iteration] = j*N + i
                      cols[iteration+1] = j*N + i+1
                      cols[iteration+2] = j*N + i-1
                      cols[iteration+3] = (j+1)*N+i
                      cols[iteration+4] = (j-1)*N+i
                      data[iteration] = -4/(h**2)
                      data[iteration+1 : iteration+5] = 1/(h**2)
                      iteration += 5
          return coo_matrix((data,(rows,cols)),shape=(N**2,N**2)).tocsr()
      def implicit_method(N,T):
          A = laplace_coo_matrix(N)
          dt = 1/(T-1)
          I = identity(N*N)
          u = np.zeros((N,N), dtype=np.float64)
```

```
iteration = 0
while u[int((N-1)/2), int((N-1)/2)] < 1.0:
    u = u.reshape((N*N))
    u_1 = spsolve((I-dt*A), u)
    u_1 = u_1.reshape((N,N))
    u_1[0,:] = 5.0
    u =u_1.copy()
    iteration = iteration +1
    time = iteration *dt
    print("T(i): {:.17f}".format(time))
    return u, time

# smaller N and T values to avoid time consuming tasks
u_implicit, T_implicit = implicit_method(61,10001)
plot_temperature_distributions(u_implicit,7)</pre>
```

T(i): 0.42410000000000003



This solution only achieves 3 decimal places precision, with N=51 and T=80001

0.4.1 CUDA version

```
[26]: @cuda.jit
def laplace_coo_matrix_kernel(rows, cols, data, N):
    i, j = cuda.grid(2)  # Global position in the grid

if i >= N or j >= N:
    return # Check boundary

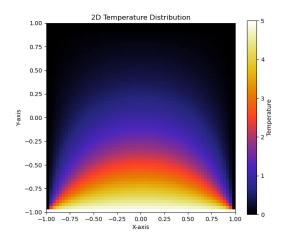
index = i * N + j # Convert 2D position to 1D index in the matrix
```

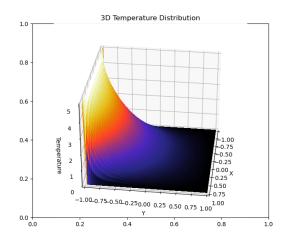
```
h = 2.0 / (N - 1)
# Calculate the number of elements before the current row
base\_index = 5 * i * N if i > 0 else 0
if i == 0 or i == N - 1 or j == 0 or j == N - 1:
    # Boundary conditions
    rows[base_index + j] = index
    cols[base_index + j] = index
    data[base_index + j] = 1.0
else:
        rows[base_index:base_index+5] = index
        # Diagonal
        cols[base_index] = index
        cols[base_index] = index
        cols[base_index+1] = index+1
        cols[base_index+2] = index-1
        cols[base_index+3] = (j+1)*N+i
        cols[base_index+4] = (j-1)*N+i
        data[base\_index] = -4/(h**2)
        data[iteration+1 : iteration+5] = 1/(h**2)
        pass
```

```
[28]: def implicit_method(N,T):
          A = laplace_coo_matrix(N)
          dt = 1/(T-1)
          I = identity(N*N)
          u = np.zeros((N,N), dtype=np.float64)
          u[0,:] = 5.0
          iteration = 0
          while u[int((N-1)/2), int((N-1)/2)] < 1.0:
              u = u.reshape((N*N))
              u_1 = spsolve((I-dt*A), u)
              u_1 = u_1.reshape((N,N))
              u_1[0,:] = 5.0
              u = u_1 \cdot copy()
              iteration = iteration +1
          time = iteration *dt
          return u, time
      u_implicit, T_implicit = implicit_method(61,10001)
      print(T_implicit)
```

```
plot_temperature_distributions(u_implicit,7)
```

0.42410000000000003





As expected the cuda version also lacks of precision compared to the explicit method, achieving again only 3 decimal precision.

0.5 4. Analysis of results

4.1 Number of discretisation points As previously theorized, larger matrix dimensions (N) could lead to more precise solutions. In this experiment, we aim to identify the optimal value of N that achieves the highest accuracy for a fixed time (T) using the explicit method.

```
[29]: target_time = 0.42011387033
      best_diff_cpu = float('inf')
      best_diff_gpu = float('inf')
      best_N_cpu = None
      best_N_gpu = None
      constant_T = 2555001 # Fixed value for T
      # Assuming min_N is an odd number
      min_N = 151 # for example, start from 151
      max_N = 301 # for example, up to 301
      step_N = 30 # increment by 30 to ensure N is always odd
      for N in range(min_N, max_N + 1, step_N):
          try:
              # Run the explicit methods on CPU
              _, exp_time_cpu, _ = explicit_method_cpu(N, constant_T)
              diff_cpu = abs(exp_time_cpu - target_time)
              if diff_cpu < best_diff_cpu:</pre>
                  best_diff_cpu = diff_cpu
```

```
best_N_cpu = N
         # Run the explicit methods on GPU
         _, exp_time_gpu, _ = explicit_method_gpu(N, constant_T)
        diff_gpu = abs(exp_time_gpu - target_time)
        if diff_gpu < best_diff_gpu:</pre>
            best_diff_gpu = diff_gpu
            best_N_gpu = N
    except Exception as e:
        print(f"Failed for N={N}, T={constant_T}: {e}")
print(f"Best N for CPU: {best_N_cpu}, Difference: {best_diff_cpu}")
print(f"Best N for GPU: {best_N_gpu}, Difference: {best_diff_gpu}")
T(e-cpu): 0.4240117416829746
T(e-gpu): 0.4240121330724070
T(e-cpu): 0.4240117416829746
T(e-gpu): 0.4240121330724070
T(e-cpu): 0.4240113502935421
T(e-gpu): 0.4240117416829746
T(e-cpu): 0.4240113502935421
/home/juan/anaconda3/lib/python3.11/site-packages/numba/cuda/dispatcher.py:536:
NumbaPerformanceWarning: Grid size 64 will likely result in GPU under-
utilization due to low occupancy.
  warn(NumbaPerformanceWarning(msg))
T(e-gpu): 0.4240117416829746
T(e-cpu): 0.4240113502935421
/home/juan/anaconda3/lib/python3.11/site-packages/numba/cuda/dispatcher.py:536:
NumbaPerformanceWarning: Grid size 81 will likely result in GPU under-
utilization due to low occupancy.
  warn(NumbaPerformanceWarning(msg))
T(e-gpu): 0.4240117416829746
T(e-cpu): 0.4240113502935421
/home/juan/anaconda3/lib/python3.11/site-packages/numba/cuda/dispatcher.py:536:
NumbaPerformanceWarning: Grid size 100 will likely result in GPU under-
utilization due to low occupancy.
  warn(NumbaPerformanceWarning(msg))
T(e-gpu): 0.4240117416829746
Best N for CPU: 211, Difference: 0.003897479963542061
Best N for GPU: 211, Difference: 0.0038978713529745357
```

From this code we find out that the best N for both for fixed time 2555001 is an N value on the

```
[32]: exp_distribution_cpu, exp_time_cpu, exp_central_temp_history_cpu = explicit_method_cpu(211, 2560001)
exp_temp_distribution_gpu, exp_time_gpu, exp_central_temp_history_gpu = explicit_method_gpu(211, 2560001)
```

T(e-cpu): 0.4240113281250000 T(e-gpu): 0.4240117187500000

Best results, achieving a 7 decimal places precision.

0.5.1 4.2 Convergence Analysis

To evaluate convergence, we employ the method of controlled variables. The size of the matrix N is increased with a constant time steps T. And the other way around with a constant N size matrix, we increase time steps T.

```
[46]: # Testing values
      dimension_sizes = np.array([9,17,37,57,77,97])
      time_steps = np.array([2001,4001,5001, 6001, 8001, 9001])
      grid_spacings = np.zeros(len(dimension_sizes))
      # Relative error results for each combination of dimension size and time step
      error explicit cuda variable dim = np.zeros(len(dimension sizes))
      error_implicit_variable_dim = np.zeros(len(dimension_sizes))
      error_explicit_cuda_fixed_dim = np.zeros(len(time_steps))
      error_implicit_fixed_dim = np.zeros(len(time_steps))
      T_actual = target_time
      for i, dim_size in enumerate(dimension_sizes):
          # For explicit_method_gpu, assuming it returns three values
          _, time_explicit_cuda_var_dim, _ = explicit_method_gpu(dim_size, 10001)
          # For implicit method
          _, time_implicit_var_dim = implicit_method(dim_size, 10001)
          # For explicit_method_gpu with fixed dimension
          _, time_explicit_cuda_fixed_dim, _ = explicit_method_gpu(51, time_steps[i])
          # For implicit_method with fixed dimension
          _, time_implicit_fixed_dim = implicit_method(51, time_steps[i])
          error_explicit_cuda_variable_dim[i] = np.abs((time_explicit_cuda_var_dim -u

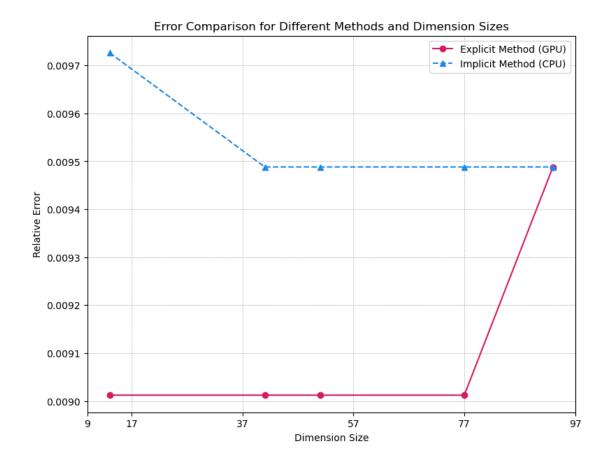
¬T_actual) / T_actual)
          error_implicit_variable_dim[i] = np.abs((time_implicit_var_dim - T_actual) /
       → T actual)
```

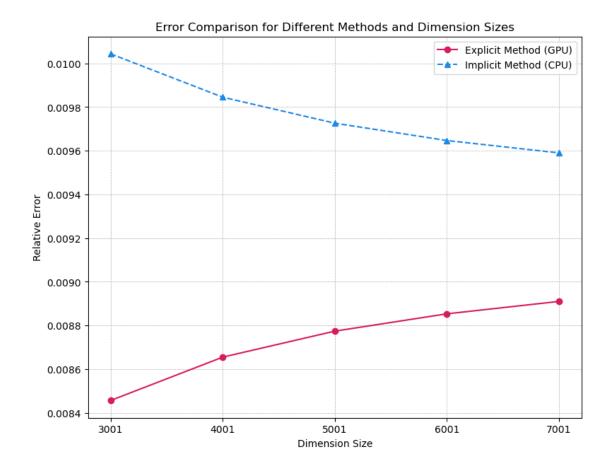
```
error_explicit_cuda_fixed_dim[i] = np.abs((time_explicit_cuda_fixed_dim -_u
       →T_actual) / T_actual)
          error_implicit_fixed_dim[i] = np.abs((time_implicit_fixed_dim - T_actual) / ___
       →T actual)
          grid_spacings[i] = 2 / (dim_size - 1)
     T(e-gpu): 0.4240000000000000
     T(e-gpu): 0.0350000000000000
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.4237500000000000
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.4238000000000000
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.42383333333333333
     T(e-gpu): 0.424100000000000
     T(e-gpu): 0.4238750000000000
     T(e-gpu): 0.424100000000000
     T(e-gpu): 0.4238888888888889
[47]: # Testing values
      grid_sizes = np.array([13, 41, 51, 77, 93])
      time_steps = np.array([3001, 4001, 5001, 6001, 7001])
      # Arrays to store grid spacings and relative errors
      grid_spacings = np.zeros(len(grid_sizes))
      error_explicit_gpu_variable_grid = np.zeros(len(grid_sizes))
      error_implicit_variable_grid = np.zeros(len(grid_sizes))
      error_explicit_gpu_fixed_grid = np.zeros(len(time_steps))
      error_implicit_fixed_grid = np.zeros(len(time_steps))
      T_actual = target_time
      for i in range(len(grid_sizes)):
          # Run explicit and implicit methods for variable grid sizes
          _, time_explicit_gpu_var_grid, _ = explicit_method_gpu(grid_sizes[i], 10001)
          _, time_implicit_var_grid= implicit_method(grid_sizes[i], 10001)
          # Calculate relative errors for variable grid sizes
          error_explicit_gpu_variable_grid[i] = np.abs((time_explicit_gpu_var_grid -u
       →T_actual) / T_actual)
          error_implicit_variable_grid[i] = np.abs((time_implicit_var_grid -_

¬T_actual) / T_actual)

          # Run explicit and implicit methods for fixed grid size with variable time_{\sqcup}
       \hookrightarrowsteps
          _, time_explicit_gpu_fixed_grid, _ = explicit_method_gpu(51, time_steps[i])
```

```
_, time_implicit_fixed_grid = implicit_method(51, time_steps[i])
          # Calculate relative errors for fixed grid size with variable time steps
          error_explicit_gpu_fixed_grid[i] = np.abs((time_explicit_gpu_fixed_grid -__
       →T_actual) / T_actual)
          error implicit fixed grid[i] = np.abs((time implicit fixed grid - T actual)
       →/ T_actual)
          # Calculate grid spacing for each grid size
          grid_spacings[i] = 2 / (grid_sizes[i] - 1)
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.4236666666666667
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.4237500000000000
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.4238000000000000
     T(e-gpu): 0.4239000000000000
     T(e-gpu): 0.42383333333333333
     T(e-gpu): 0.424100000000000
     T(e-gpu): 0.4238571428571429
[48]: # Style
      plt.style.use('seaborn-v0_8-pastel')
      color_explicit = '#D81B60'
      color_implicit = '#1E88E5'
      plt.figure(figsize=(9, 7))
      plt.plot(grid_sizes, error_explicit_gpu_variable_grid, color=color_explicit,_u
       ⇔label='Explicit Method (GPU)', marker='o', linestyle='-')
      plt.plot(grid sizes, error implicit variable grid, color=color implicit,
       ⇔label='Implicit Method (CPU)', marker='^', linestyle='--')
      plt.grid(True, which='both', linestyle='--', linewidth=0.5)
      plt.legend()
      plt.title('Error Comparison for Different Methods and Dimension Sizes')
      plt.xlabel('Dimension Size')
      plt.ylabel('Relative Error')
      plt.xticks(dimension_sizes)
      plt.show()
```





What we can get from these results is that no matter the version of the method. As N (the matrix dimension) and T (total time steps) increases, the relative errors decrease.

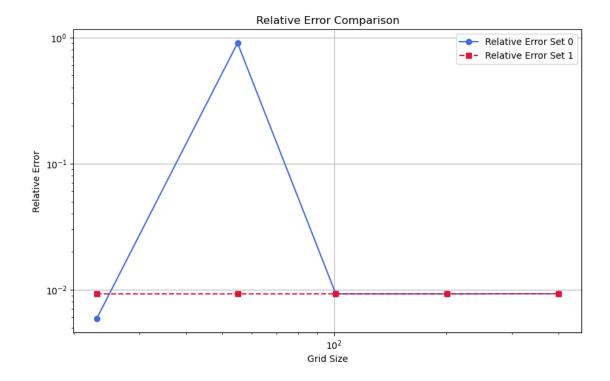
Furthermore, to test the hypothesis of second-order convergence with spatial second-order and temporal first-order finite differences, matrix dimensions N are set to values of 23, 55, 101, 201, and 401. The total number of time steps T is varied accordingly to preserve Courant numbers C of 0.25 and 0.125, in each case.

```
[54]: grid_sizes = np.array([23, 55, 101, 201, 401])
    time_steps_set_0 = np.array([621, 2601, 10001, 44001, 160001])
    time_steps_set_1 = np.array([1251, 5001, 20001, 80001, 320001])

# Arrays to store grid spacings and time step intervals
    grid_spacing_values = np.zeros(5)
    time_step_intervals_0 = np.zeros(5)
    time_step_intervals_1 = np.zeros(5)

# Arrays to store relative error
    relative_errors_0 = np.zeros(5)
    relative_errors_1 = np.zeros(5)
```

```
target_time = 0.42011387033
      for index in range(len(grid_sizes)):
          # Run explicit_numba method for different time steps
          _, result_time_0, _ = explicit_method_cpu(grid_sizes[index],_
       stime_steps_set_0[index])
          _, result_time_1, _ = explicit_method_cpu(grid_sizes[index],__
       →time_steps_set_1[index])
          # Calculate relative errors
          relative errors 0[index] = np.abs((result time 0 - target time) / ____
       →target_time)
          relative errors_1[index] = np.abs((result_time 1 - target_time) /__
       →target_time)
          # Calculate grid spacings and time step intervals
          grid_spacing_values[index] = 2 / (grid_sizes[index] - 1)
          time_step_intervals_0[index] = 1 / (time_steps_set_0[index] - 1)
          time_step_intervals_1[index] = 1 / (time_steps_set_1[index] - 1)
     T(e-cpu): 0.4225806451612903
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.0403846153846154
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.4240000000000000
     T(e-cpu): 0.4240062500000000
     T(e-cpu): 0.4240093750000000
[55]: # Plot for relative errors
      plt.figure(figsize=(10, 6))
      plt.loglog(grid_sizes, relative_errors_0, label='Relative Error Set 0', u
       ⇔marker='o', linestyle='-', color='royalblue')
      plt.loglog(grid_sizes, relative_errors_1, label='Relative Error Set 1', u
       →marker='s', linestyle='--', color='crimson')
      plt.title('Relative Error Comparison')
      plt.xlabel('Grid Size')
      plt.ylabel('Relative Error')
      plt.legend()
      plt.grid(True)
      plt.show()
```



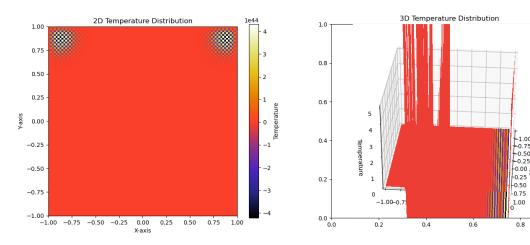
Both converge with second-order accuracy

0.5.2 4.3 Stability

Previous results suggest that the finit difference method with explicit time step can lead to instability. This is the case when the Courant number (C) exceeds 0.25 Therefore a case of N=101 and T=5001, could lead to unstable solutions.

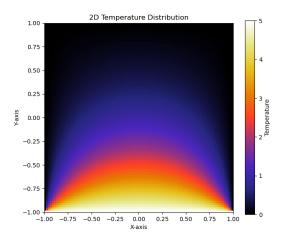
[43]: u_exp_stability, T_exp_stability, timer = explicit_method_gpu(101,5001) plot_temperature_distributions(u_exp_stability,5)

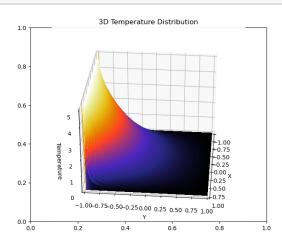
T(e-gpu): 0.0204000000000000



This solution does not converge and is unstable. However the implicit method does not have this disadvantage. While the explicit scheme is straightforward, its stability and convergence are contingent upon the Courant number. The Courant number C must not exceed 0.25, restricting the choice of matrix dimension N and the number of steps T. This limitation bounds the range of the numerical solution that can be accurately described. In contrast, the implicit scheme with backward differences is unconditionally stable, affording greater flexibility in parameter selection.

[44]: u_imp_stability, T_imp_stability = implicit_method(101,5001) plot_temperature_distributions(u_imp_stability,5)





[45]: print(T_imp_stability)

0.4242

This proves the flexibility of the implicit method compared to the explicit method

0.5.3 5. Conclusions

In this study, we have applied the finite difference method to address the two-dimensional heat conduction problem. Our efforts included implementing an explicit scheme utilizing CUDA and Numba for computational acceleration, as well as developing a finite difference approach with implicit time stepping.

We have enhanced the number of discretization points to obtain up to seven decimal places of precision in part 4.1 and have examined the convergence in terms of actual versus ideal time, along with the stability of various schemes. We can state that while the explicit solution offers quicker computation times and in many cases a higher precision, the implicit solution demonstrates superior convergence and stability.

However due to limitations, some aspects have not been fully investigated, such as the impact of the courant number in stability and other aspects. A more comprehensive investigation could incorporate these considerations for an in-depth analysis.

[]:[