How to create superposition of quantum states?

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Qubit

 The superposition of states of classical bits is used to perform computations in quantum informatics

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

Examples of systems that can be used as a qubit: spin ½; photon polarization; two-level atom.

- If we have a system of two qubits and the first one is described by the space \mathcal{H}_A and the second one is described by the space \mathcal{H}_B then the whole system is described by the space $\mathcal{H}_A \otimes \mathcal{H}_B$
- A qubit $|\psi\rangle_{AB} \in \mathcal{H}_A \otimes \mathcal{H}_B$ that describes a system of two qubits $|\psi_a\rangle_A \in \mathcal{H}_A$ and $|\psi_b\rangle_B \in \mathcal{H}_B$ can be describe in such way $|\psi\rangle_{AB} = |\psi_a\rangle_A \otimes |\psi_b\rangle_B = |\psi_a\psi_b\rangle$

Density matrix

■ Set of states $\{|\psi_i\rangle\}$ and set of probabilities $\{p_i\}$ is introduced a density matrix

$$\rho = \sum p_i |\psi_i\rangle \langle \psi_i|$$

- The density matrix for a particular state $|\psi\rangle$ looks as follows $\rho_{\psi}=|\psi\rangle\,\langle\psi|$
- Density matrix ρ for the state of n several systems $\{\rho_i\}$ is equal $\rho = \rho_1 \otimes \rho_2 \otimes ... \otimes \rho_n$

Quantum Gates

Quantum gates are used to implement quantum algorithms

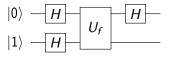


Figure: Deutsch's algorithm

lacksquare The quantum gate operator is a unitary operator: $H^\dagger H = \mathbb{I}$

Quantum gates

$$\begin{array}{c} \blacksquare \ X,Y,Z \\ X \ |\psi\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ Y \ |\psi\rangle = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ Z \ |\psi\rangle = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \\ \end{array}$$

■ H - Hadamard operator:

$$H|0\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}} H|1\rangle = \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$



Quantum gates

■ Control Not Gate: $|\psi\rangle \rightarrow (|0\rangle \langle 0| \otimes \mathbb{I} + |1\rangle \langle 1| \otimes X) |\psi\rangle$ $|00\rangle \rightarrow |00\rangle$ $|01\rangle \rightarrow |01\rangle$ $|10\rangle \rightarrow |11\rangle$ $|11\rangle \rightarrow |10\rangle$

$$|\psi_1\rangle$$
 $\psi_2\rangle$

Figure: CNOT

Universal Hadamard Gate

■ Creation of superposition from the general state:

$$|\psi
angle
ightarrowrac{|\psi
angle+|\psi^{\perp}
angle}{\sqrt{2}}=|\Psi
angle$$
, in general $reve{H}|\psi
angle_{a}|Q
angle_{b...}=rac{|\psi
angle+|\psi^{\perp}
angle}{\sqrt{2}}_{a}|Q_{\psi}
angle_{...}$

$$\ket{\psi} - oxed{reve{H}} - oxed{\ket{\psi}+\ket{\psi^\perp}}{\sqrt{2}} \ \ket{Q_\psi}$$

Figure: U-Hadamard

Universal Hadamard Gate

Nonlinearity of the U-Hadamard operator

$$\begin{split} \hat{G}|\psi\rangle_{a}|Q\rangle_{b} &= \hat{G}(\alpha|0\rangle_{a} + \beta|1\rangle_{a})|Q\rangle_{b} = \\ \alpha \hat{G}|0\rangle_{a}|Q\rangle_{b} + \beta \hat{G}|1\rangle_{a}|Q\rangle_{b} &= \alpha \hat{G}|1\rangle_{a}|Q_{0}\rangle_{b} + \beta \hat{G}|0\rangle_{a}|Q_{1}\rangle_{b} \end{split}$$

This operator is not linear, which means it is not unitary and cannot be realized as a quantum gate



Bures Fidelity

If we cannot create an exact state σ , we can create a state ρ as close to the required state as possible to implement the algorithm.

Bures Fidelity

$$F = Tr(\sqrt{\sqrt{
ho}\sigma\sqrt{
ho}}) \in (0, 1)$$
 or for pure state $\sigma = |\Psi\rangle \langle \Psi| \Rightarrow F = \langle \Psi| \, \rho \, |\Psi\rangle$

Universal Cloner and Universal NOT Gates

- $|\psi\rangle \to \frac{|\psi\rangle + |\psi^{\perp}\rangle}{\sqrt{2}}$
- Universal Cloner $|\psi\phi\rangle \rightarrow |\psi\psi\rangle = U_c |\psi\phi\rangle$

Nonlinearity of U-Cloner

$$|\psi\phi\rangle = \alpha |0\phi\rangle + \beta |1\phi\rangle \rightarrow \alpha |00\rangle + \beta |11\rangle \neq |\psi\psi\rangle$$

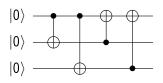
■ Universal NOT Gate $\alpha |0\rangle + \beta |1\rangle = |\psi\rangle \rightarrow |\psi^{\perp}\rangle = \beta^* |0\rangle - \alpha^* |1\rangle$

Nonlinearity of U-NOT

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle = \alpha |1\rangle + \beta |0\rangle \neq |\psi^{\perp}\rangle$$



U-Cloner and U-NOT Machine



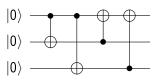
$$|\psi\rangle_a |\Xi_{00}\rangle_{bc} = |\psi\rangle \left(\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)\right) \rightarrow |\psi\rangle_a |\Xi\rangle_{bc}$$

$$|\psi\rangle_{a} |\Xi_{0+}\rangle_{bc} = |\psi\rangle \left(\frac{1}{\sqrt{2}}(|00\rangle + |01\rangle)\right) \rightarrow |\psi\rangle_{b} |\Xi\rangle_{ac}$$

$$|\Xi\rangle_{bc} = c_0 |\Xi_{00}\rangle_{bc} + c_1 |\Xi_{0x}\rangle_{bc}$$



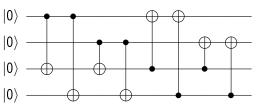
U-Cloner and U-NOT Machine



- $\begin{array}{l} \bullet \ \, \rho_1^{out} = \left(c_0^2 + c_0 \, c_1\right) |\psi\rangle \, \langle\psi| + \frac{1}{2} \, c_1^2 \mathbb{I} \\ \rho_2^{out} = \left(c_0^1 + c_0 \, c_1\right) |\psi\rangle \, \langle\psi| + \frac{1}{2} \, c_0^2 \mathbb{I} \\ \rho_3^{out} = c_0 \, c_1 \, |\psi^*\rangle \, \langle\psi^*| + \frac{1}{2} (1 c_0 \, c_1) \mathbb{I} \end{array}$
- $c_0 = c_1 = \frac{1}{\sqrt{3}} \Rightarrow \rho_1^{out} = \rho_2^{out} = \frac{5}{6} |\psi\rangle \langle\psi| + \frac{1}{6} |\psi^{\perp}\rangle \langle\psi^{\perp}|$ $F_{|\psi\rangle} = \frac{5}{6}$
- $(iY)\rho_3^{out}(iY)^{\dagger} = \frac{1}{3} |\psi\rangle \langle \psi| + \frac{2}{3} |\psi^{\perp}\rangle \langle \psi^{\perp}|$ $F_{|\psi^{\perp}\rangle} = \frac{2}{3}$



Idea of U-Hadamard

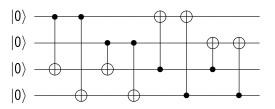


$$\rho_2^{out}{}_a \otimes (iY) \rho_3^{out} (iY)_b^\dagger \otimes \check{\rho}_{cd} \rightarrow \rho_b^{goal} \tilde{\rho}_{acd}$$



Idea of U-Hadamard

By varying the constants C_i , $i = \{1, 2, 3, 4\}$, the maximum value of F can be obtained



$$\begin{split} & \check{\rho} = C_1 \left| 00 \right\rangle \left\langle 00 \right| + C_2 \left| 01 \right\rangle \left\langle 00 \right| + C_3 \left| 00 \right\rangle \left\langle 01 \right| + C_4 \left| 11 \right\rangle \left\langle 11 \right| \\ & \rho^{goal} = A \left| \psi \right\rangle \left\langle \psi \right| + B \left| \psi \right\rangle \left\langle \psi^{\perp} \right| + C \left| \psi^{\perp} \right\rangle \left\langle \psi \right| + D \left| \psi^{\perp} \right\rangle \left\langle \psi^{\perp} \right| \\ & F = \frac{\left\langle \psi \right| + \left\langle \psi^{\perp} \right|}{\sqrt{2}} \rho^{goal} \frac{\left| \psi \right\rangle + \left| \psi^{\perp} \right\rangle}{\sqrt{2}} = \frac{1}{2} (1 + C + D) \ \textit{To be continued...} \end{split}$$

