

# Special Relativity (The Big Idea)

①

Physics is the same in all inertial reference frames.

Inertial reference frame = no external forces = constant velocity

⇒ No local experiment can tell if you're moving at a constant velocity or not.

⇒ If you wake up locked in a train car with no windows, there is nothing you can do to tell if you're already moving.

②

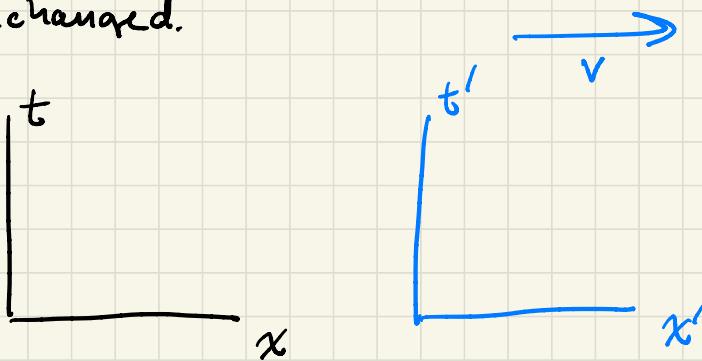
The speed of light is the same in all inertial frames.

↳ Why? Because experimentally, it works.

Einstein saw that Maxwell's equations only make sense if the speed of light is the same in all inertial reference frames.

① can also be rephrased as a statement about coordinates.

If you apply a coordinate transformation between two inertial frames, physics is unchanged.



Such coordinate transformations are known as Lorentz transformations.

matrix notation:  $\begin{array}{c} \leftrightarrow \\ \Lambda \end{array} \vec{x} = \vec{x}'$

component notation:  $\Lambda^{\mu'}_{\nu} x^\nu = x^{\mu'}$

code notation: [Sum( $\Lambda[\text{muprime}, i] x[i]$ )  
for muprime in 1..4]

$$\Lambda^{\mu'}_{\nu} = \begin{pmatrix} \gamma & -\gamma v & 0 \\ -\gamma v & \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

also for "boost" in  $x$ -direction by  $v$   
 3 rotations  $\gamma \equiv (\sqrt{1-v^2})^{-1}$

We can also write this as

$$\Lambda^{\mu'}_{\nu} = \frac{dx^{\mu'}}{dx^{\nu}} \quad \begin{matrix} \text{new coordinates} \\ \text{for } x^{\mu}, x^{\nu} \text{ inertial frames} \\ \text{old coordinates} \end{matrix}$$

Lorentz transformations ensure  
 $c$  is the same before and after  
 coordinate change.

# General Relativity (The Big Idea)

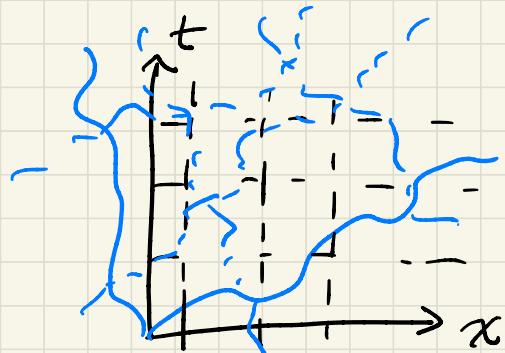
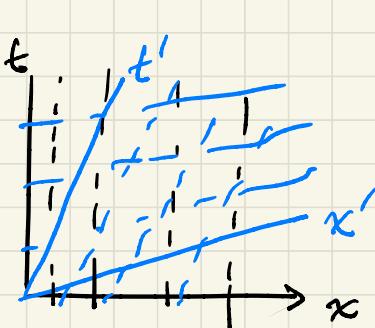
We saw that in SR, the laws of physics should not change under Lorentz transformations. (Rotations + boosts)

Let's elevate this statement. . .

## ① Principle of General Covariance.

The laws of physics should not change under ANY coordinate transformation,

⇒ \*COORDINATES ARE NOT PHYSICAL\*



## ② Equivalence Principle

Many ways to state it...

- All things fall at the same rate
- Gravitational and inertial mass are equal (We saw there are two eq.s with mass,  $F_g = \frac{GMm}{r^2}$  and  $F=ma$ )
- No local experiment can tell you if you're in a gravitational field or are being accelerated.
  - ↳ no reason those "m's need to be the same!"
- If you wake up in an elevator with no windows, there is no way to tell if it's accelerating upwards or if you're at rest in a gravitational field.

(Much like in SR, Einstein had one pretty obvious and believable postulate and one that was... inspired.)

We like SR and want it to stay true.  
Non-accelerating observers should just see SR.

- ⇒ There is always a local frame that looks inertial in the SR sense.  
(though it will change between  $x$  and  $x+dx$ )
- ⇒ On small enough scales, things look like a Minkowski metric.
- ⇒ There is no way to detect a gravitational field by a local experiment.
- ⇒ If you wake up floating in an elevator with no windows, you can't tell if you're about to hit the ground.

## Gedankenexperiments

- ① You are in an elevator in free-fall and drop a ping-pong ball.
- ② You are standing on Earth and drop a ping-pong ball.
- ③ You are in an elevator accelerating up and drop a ping-pong ball.
- ④ You are in an elevator accelerating up and hold a ping-pong ball in water, then release.
- ⑤ You are in a falling elevator and hold a ping-pong ball in water, then release.
- ⑥ You are orbiting the Earth on the ISS and hold a ping-pong ball in water, then release.
- ⑦ You are on a spinning amusement park ride and hold a ping-pong ball in water, then release.

# Manifolds

N.B. We're playing fast and loose here with the maths.

How do we build a mathematical structure that

- ① Doesn't depend on coordinates
- ② Looks locally flat
- ③ We can do physics on (Aka take derivatives)

Such a structure is called a manifold.

We're going to make space-time a manifold.

Manifold:  $n$ -dimensional manifold looks like  $\mathbb{R}^n$  at any individual point

Permits a notion of differentiation.

Dimensionality  $n$  is the same at all points.

Easiest to understand by example.

Is it a manifold?

$\mathbb{R}^n$  ? Yes

Sphere ? Yes

Torus? Yes

Cone? Sort of (not smooth)

Two cones? No!

Line segment? Sort of (has boundary)

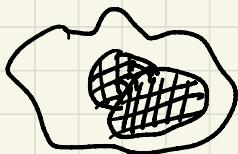
(Ex.) Which alphabet letters are 1-manifolds?

C, D, G, L, M, N, O, S, U, V, W, Z

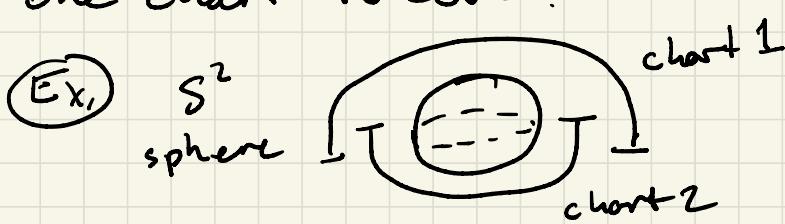
COLUMNS is the longest word I  
can find

# Coordinates on Manifolds

Covered by atlas of charts  $\rightarrow$  coordinate systems in various patches  
↓ coordinate charts must overlap smoothly,



Many manifolds require more than one chart to cover.



The existence of world maps  $\Rightarrow$  the Earth is flat???

Ex. Which letters require more than one chart?

D, O

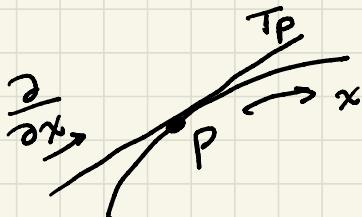
DODO is the longest word I can find

## Vectors

At any point  $P$  in the manifold, there exists a tangent space  $T_p$  spanned by the partial derivative operators.

i.e. if you choose some coordinates

$x^m$ , then you can express any vector in  $T_p$  as  $\nabla = V^m \partial_m$  where  $\partial_m \equiv \frac{\partial}{\partial x^m}$



Often, people just write  $V^m$  as a vector but that is sloppy.

$V^m$  are JUST THE COMPONENTS OF

$\nabla$  for some choice of coordinates  $x^m$

Same as the difference between

$$\vec{v} = (1, 3, -4) \quad \text{and} \quad \vec{v} = 1\hat{i} + 3\hat{j} + (-4)\hat{k}$$

Vector components transform under coordinate change like

$$\mathbf{v}^m' = \frac{\partial x^m'}{\partial x^m} \mathbf{v}^m$$

for  $x^m \rightarrow x'^m$

Follows from chain rule since  $\partial_m = \frac{\partial}{\partial x^m}$

$$\rightarrow \partial_{m'} = \frac{\partial x^m}{\partial x^{m'}} \frac{\partial}{\partial x^m}$$

✓ Doesn't change - it is a coordinate independent object.

$$\mathbf{v} = v^m \partial_m \rightarrow v^{m'} \partial_{m'} = \cancel{\frac{\partial x^{m'}}{\partial x^m}} v^m \cancel{\frac{\partial x^m}{\partial x^{m'}}} \cancel{\partial_m} \\ = v^m \partial_m = \mathbf{v}$$

→ You are going to become very familiar with this transformation rule.

## Dual vectors

You might be wondering what's up with the indices up top versus down below.

When the index is up top, it is called contravariant and transforms

like a vector, i.e.  $v^m' = \frac{\partial x^m}{\partial x^{m'}} v^m$ .

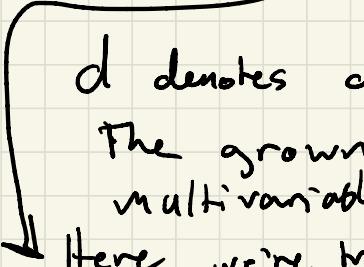
When the index is down below, it is called covariant and transforms

like a "dual vector", i.e.  $v_m' = \frac{\partial x^m}{\partial x^{m'}} v_m$

Dual vectors live in dual tangent space,

hence we can define a basis for it

$$\text{via } dx^m \left( \frac{\partial}{\partial x^\nu} \right) = \frac{\partial x^m}{\partial x^\nu} = \delta^m_\nu$$

d denotes gradient operator:  $df(\frac{d}{dx}) = \frac{df}{dx}$

The grown-up version of  $\vec{\nabla}$  from multivariable.

Here, we're treating  $x^m$  like a function  $x^m(x^\sigma) = x^m$  kind of weird.  
 $x(\vec{x}) = x$ ,  $y(\vec{x}) = y \dots$

Dual vectors map vectors to real numbers.

$$W_\mu V^\mu \in \mathbb{R}.$$

This number doesn't change under coordinate transformations.

Proof:

$$\frac{\partial x^{\mu}}{\partial x^{\mu}}, W_\mu \frac{\partial x^{\mu'}}{\partial x^{\mu}} V^{\mu'} = W_\mu V^{\mu'}$$
$$= W_\mu V^\mu$$

Quantities that don't change under coordinate transformations are called SCALARS. Just a  $\mathbb{R}$  number at a point  $p$ .

Scalar field:  $\mathbb{R}$  at all points  $p \in M$

Vector field:  $V \in T_p$  for all  $p \in M$

# Tensors

Map  $k$  vectors and  $\ell$  dual vectors to  $\mathbb{R}$ .

$$T^{u_1 u_2 u_3 \dots u_k} \\ v_1 v_2 \dots v_\ell$$

(again, these are the \*COMPONENTS\*.)

Tensors, like vectors, are invariant, but their components and bases change)

e.g.  $T = T^{\mu\nu} \sigma_\mu (\partial_\mu \otimes \partial_\nu \otimes dx^\alpha \otimes dx^\beta)$

Contraction: sum over matching co and contra  $T^{\mu}_{\mu}$

Tensor components transform exactly as

You'd expect.. -

$$T^{\mu'}_{\nu'} = \frac{\partial x^{\mu'}}{\partial x^\mu} \frac{\partial x^\nu}{\partial x^{\nu'}} T^\mu_{\nu}$$

(Ex) Do the transformation of  $S_{\mu\nu} = \begin{pmatrix} 1 & x^2 \\ 0 & 1 \end{pmatrix}$

$$x' = \frac{2x}{y}, \quad y' = \frac{y}{2} \text{ two ways.}$$

1. Use above formula

$$2. S = 1 dx^2 + x^2 dy^2, \quad dx = y' dx' + x' dy' \\ dy = 2 dy'$$