

The background features a large, light blue watermark of the KTH logo. It consists of a crown at the top, a circular wreath of oak leaves in the middle, and the text 'KTH VETENSKAP OCH KONST' at the bottom.

# Numerical solutions of differential equations

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Course **SF2521**, 7.5 ECTS, VT18



## Lecture 10

# Entropy solutions



# Applications of the Entropy Condition

# Entropy - Consequence for numerical schemes

**Recall:** von Neumann analysis for linear problem with periodic BC revealed:

Example ( $\partial_t u + \partial_x u = 0$  - Advection equation)

Consider **forward Euler** and **central differences** for the advection equation:

$$Q_j^{n+1} = Q_j^n + \frac{1}{2} \lambda_{\text{CFL}} (Q_{j+1}^n - Q_{j-1}^n), \quad \lambda_{\text{CFL}} = \frac{\Delta t}{\Delta x}.$$

Again, here  $m = M = 1$ , but  $b_{-1} = -\frac{\lambda_{\text{CFL}}}{2}$ ,  $b_0 = 1$ ,  $b_1 = \frac{\lambda_{\text{CFL}}}{2}$ . Then

$$g_k(\Delta t, \Delta x) = -\frac{\lambda_{\text{CFL}}}{2} e^{-ik\Delta x} + 1 + \frac{\lambda_{\text{CFL}}}{2} e^{ik\Delta x} = 1 + \lambda_{\text{CFL}} i \sin(k\Delta x).$$

Hence,

$$\max_k |g_k(\Delta t, \Delta x)| = \max_k \sqrt{1 + \lambda_{\text{CFL}}^2 \sin^2(k\Delta x)} > 1,$$

and the method is unstable for all fixed  $\lambda_{\text{CFL}}$ .

# Entropy - Consequence for numerical schemes

General (nonlinear) case with convex flux  $f$ :

**central differences are inappropriate** to discretize the problem.

## Example

Let  $f'' > 0$  with  $f'(-1) < 0 < f'(1)$  and  $f(1) = f(-1)$ , e.g.  
 $f(u) = u^2$ .

Consider initial value  $v_0(x) = \begin{cases} 1, & x \geq 0 \\ -1, & x < 0 \end{cases}$ .

Central difference discretization:

$$Q_j^{n+1} := Q_j^n - \frac{\Delta t}{2\Delta x} (f(Q_{j+1}^n) - f(Q_{j-1}^n)),$$

where  $Q_j^n \approx u(j\Delta x, n\Delta t)$  on an equidistant mesh.

# Entropy - Consequence for numerical schemes

## Example (- Part 2)

Let  $f'' > 0$  with  $f'(-1) < 0 < f'(1)$  and  $f(1) = f(-1)$ .

Consider initial value  $v_0(x) = \begin{cases} 1, & x \geq 0 \\ -1, & x < 0 \end{cases}$ .

Hence:  $f(Q_j^0) = f(\pm 1) = f(1) \Rightarrow f(Q_{j+1}^0) - f(Q_{j-1}^0) = 0$  for all  $j$ .

Since

$$Q_j^{n+1} := Q_j^n - \frac{\Delta t}{2\Delta x} (f(Q_{j+1}^n) - f(Q_{j-1}^n)),$$

we conclude

$$Q_j^n = \begin{cases} 1, & j > 0 \\ -1, & j < 0 \end{cases} \quad \text{for all } n \in \mathbb{N}.$$

## Entropy - Consequence for numerical schemes

### Example (- Part 3)

Let  $f'' > 0$  with  $f'(-1) < 0 < f'(1)$  and  $f(1) = f(-1)$  and

$$\text{initial value } v_0(x) = \begin{cases} 1, & x \geq 0 \\ -1, & x < 0 \end{cases}.$$

Central difference approximation:

$$Q_j^n = \begin{cases} 1, & j > 0 \\ -1, & j < 0 \end{cases} \quad \text{for all } n \in \mathbb{N}.$$

**Question:**

$$\text{Is } u(x, t) = \begin{cases} 1, & x > 0 \\ -1, & x \leq 0 \end{cases} \quad \text{also Lax-Entropy solution?}$$

# Entropy - Consequence for numerical schemes

## Example (- Part 4)

Is  $u(x, t) = \begin{cases} 1, & x > 0 \\ -1, & x \leq 0 \end{cases}$  also Lax-Entropy solution?

1. Verify Rankine-Hugoniot jump condition:

$$s = 0, \quad f(u_l) - f(u_r) = 0 \quad \Rightarrow \quad (u_l - u_r)s = f(u_l) - f(u_r)$$

$\Rightarrow u$  is weak solution.

2. Verify Lax entropy condition:

We have  $f'(u_l) = f'(-1)$  and  $f'(u_r) = f'(1)$ . Hence:

$$f'(-1) = f'(u_l) < \underbrace{s}_{=0} < f'(u_r) = f'(1) \quad \text{"Contradiction"}.$$

$\Rightarrow u$  does not fulfill Lax entropy condition.

The central difference scheme does not produce an entropy solution!



# Entropy - Consequence for numerical schemes

Observe:

If  $f'' > 0$  with  $f'(-1) < 0 < f'(1)$  and  $f(1) = f(-1)$ , and for initial value  $v_0(x) = \begin{cases} 1, & x \geq 0 \\ -1, & x < 0 \end{cases}$ ,

the upwind scheme

$$Q_j^{n+1} := Q_j^n - \frac{\Delta t}{\Delta x} (f(Q_j^n) - f(Q_{j-1}^n)),$$

suffers from the same problem.

# Entropy solutions to the Riemann-problem

## Riemann-problem.

Let  $f \in C^2(\mathbb{R})$  be a convex flux, i.e.  $f'' > 0$ , and let

$$v_0(x) = \begin{cases} u_l, & x < 0 \\ u_r, & x > 0 \end{cases}$$

We seek the **entropy solution**  $u \in L^\infty(\mathbb{R} \times \mathbb{R}^+)$  to

$$\partial_t u + \partial_x f(u) = 0 \quad \text{and} \quad u(\cdot, 0) = v_0.$$

The equation is understood in the weak sense!

Can we explicitly state the entropy solution to this problem?

# Entropy solutions to the Riemann-problem

## Observation:

Let  $u$  be a weak solution to  $\partial_t u + \partial_x f(u) = 0$ .

Then,  $u_\lambda(x, t) := u(\lambda x, \lambda t)$  is also a weak solution for all  $\lambda > 0$ . In particular:

$$u_\lambda(x, 0) = u(\lambda x, 0) = \begin{cases} u_l, & \lambda x < 0 \\ u_r, & \lambda x > 0 \end{cases} = v_0(x).$$

Since the entropy solution is unique, it must hold  $u(\lambda x, \lambda t) = u(x, t)$  for all  $\lambda > 0$ .

We therefore consider solutions of the form

$$u(x, t) = v\left(\frac{x}{t}\right).$$

## Entropy solutions to the Riemann-problem

Let  $u(x, t) = v\left(\frac{x}{t}\right)$ .

In regions in which  $v$  is smooth we have:

$$\begin{aligned} 0 &= \partial_t u(x, t) + \partial_x f(u(x, t)) \\ &= -\frac{x}{t^2} v'\left(\frac{x}{t}\right) + f'\left(v\left(\frac{x}{t}\right)\right) v'\left(\frac{x}{t}\right) \frac{1}{t} \\ &= v'\left(\frac{x}{t}\right) \frac{1}{t} \left(f'\left(v\left(\frac{x}{t}\right)\right) - \frac{x}{t}\right). \end{aligned}$$

Hence for all  $\xi = \frac{x}{t} \in \mathbb{R}$  we have either

$$f'(v(\xi)) - \xi = 0$$

or

$$v'(\xi) = 0.$$

# Entropy solutions to the Riemann-problem

From the conservation we have

$$f'(v(\xi)) - \xi = 0 \quad \text{or} \quad \underline{v'(\xi) = 0} \quad \text{for all } \xi = \frac{x}{t} \in \mathbb{R}.$$

We can distinguish 3 cases.

Case 1:  $u_l = u_r$ .

We have the classical solution  $u(x, t) \equiv u_l$  for all  $(x, t) \in \mathbb{R} \times \mathbb{R}^+$ .

Since it is a classical solution, it must be the unique entropy solution.

# Entropy solutions to the Riemann-problem

From the conservation we have

$$f'(v(\xi)) - \xi = 0 \quad \text{or} \quad \underline{v'(\xi) = 0} \quad \text{for all } \xi = \frac{x}{t} \in \mathbb{R}.$$

Case 2:  $u_l > u_r$ . Then

$$u(x, t) = \begin{cases} u_l & \text{for } x < st \\ u_r & \text{for } x > st \end{cases}$$

with  $s = \frac{f(u_l) - f(u_r)}{u_l - u_r}$  is the unique **entropy solution**, because

$$f'(u_r) < \frac{f(u_l) - f(u_r)}{u_l - u_r} < f'(u_l).$$

This is the *Lax shock*.

## Entropy solutions to the Riemann-problem

Ansatz:  $u(x, t) = v(\frac{x}{t})$ . From the conservation we have

$$\underline{f'(v(\xi)) - \xi = 0} \quad \text{or} \quad v'(\xi) = 0 \quad \text{for all } \xi = \frac{x}{t} \in \mathbb{R}.$$

Case 3:  $u_l < u_r$ . Recall: **discontinuous solutions are only possible if the Lax entropy condition is violated**. Since  $v'(\xi) = 0$  is therefore not everywhere possible, it must hold  $f'(v(\xi)) - \xi = 0$ . Hence:

$$f'(v(\xi)) = \xi \quad \Rightarrow \quad (f')^{-1}(f'(v(\xi))) = (f')^{-1}(\xi)$$

$$\Rightarrow \quad v(\xi) := (f')^{-1}(\xi).$$

We obtain

$$u(x, t) := \begin{cases} u_l & \text{for } \frac{x}{t} \leq f'(u_l) \\ v(\frac{x}{t}) & \text{for } f'(u_l) < \frac{x}{t} < f'(u_r) \\ u_r & \text{for } \frac{x}{t} > f'(u_r) \end{cases}$$

This entropy solution is called **rarefaction-wave**.

# Entropy solutions to the Riemann-problem

**Summary.** Riemann-problem:  $\partial_t u + \partial_x f(u) = 0$ .

Let  $f \in C^2(\mathbb{R})$  be a convex flux, i.e.  $f'' > 0$ , and let

$$u(x, 0) = v_0(x) = \begin{cases} u_l, & x < 0 \\ u_r, & x > 0 \end{cases}$$

The entropy solution  $u \in L^\infty(\mathbb{R} \times \mathbb{R}^+)$  is given by:

If  $u_l = u_r$ :

$$u(x, t) \equiv u_l.$$

If  $u_l > u_r$ :

$$u(x, t) = \begin{cases} u_l & \text{for } x < st \\ u_r & \text{for } x > st \end{cases} \quad \text{where } s = \frac{f(u_l) - f(u_r)}{u_l - u_r}.$$

If  $u_l < u_r$ :

$$u(x, t) := \begin{cases} u_l & \text{for } \frac{x}{t} \leq f'(u_l) \\ (f')^{-1}\left(\frac{x}{t}\right) & \text{for } f'(u_l) < \frac{x}{t} < f'(u_r) \\ u_r & \text{for } \frac{x}{t} > f'(u_r) \end{cases}$$



# Entropy solutions to the Riemann-problem

## Remarks:

- ▶ It is possible to show that these **unique entropy solutions** are obtained by the viscosity limit.
- ▶ For the **Riemann-problem** we have now an **explicit formula** to state the solutions for quite general nonlinearities.
- ▶ Unfortunately, it is **not always that easy** and we mostly **need numerical methods**.
- ▶ Can we use the Lax entropy condition for numerical solutions?