

The background features a large, light blue watermark of the KTH logo. It consists of a crown at the top, a circular wreath of oak leaves in the middle, and the text 'KTH VETENSKAP OCH KONST' at the bottom.

# Numerical solutions of differential equations

Patrick Henning

[pathe@kth.se](mailto:pathe@kth.se)

Division of Numerical Analysis, KTH, Stockholm

Course **SF2521**, 7.5 ECTS, VT18

# Finite Volumes Schemes of Higher Order

## The Lax-Wendroff Scheme

## Finite Differences of Higher Order

### Goal:

Construct higher order methods using Taylor expansion:

As usual: let  $x_j = \frac{\Delta x}{2} + j\Delta x$  for  $j \in \mathbb{Z}$  and  $t_n = n\Delta t$  for  $n \in \mathbb{N}_0$ .

Let  $u$  be a smooth solution to  $\partial_t u + \partial_x f(u) = 0$ . Set

$$u_j^n := u(x_j, t_n).$$

Then Taylor expansion yields:

$$u_j^{n+1} \stackrel{\text{Taylor}}{=} u_j^n + \Delta t \partial_t u_j^n + \frac{\Delta t^2}{2} \partial_{tt} u_j^n + \mathcal{O}(\Delta t^3)$$

**Idea:** Replace time derivatives by space derivatives using the conservation law.

## Taylor Expansion

We consider

$$u_j^{n+1} = u_j^n + \Delta t \partial_t u_j^n + \frac{\Delta t^2}{2} \partial_{tt} u_j^n + \mathcal{O}(\Delta t^3).$$

It holds

$$\partial_t u_j^n = -\partial_x f(u_j^n) = -\frac{f(u_{j+1}^n) - f(u_{j-1}^n)}{2\Delta x} + \mathcal{O}(\Delta x^2).$$

and it also holds

$$\begin{aligned} \partial_t \partial_t u_j^n &= -\partial_t \partial_x f(u_j^n) = -\partial_x (f'(u_j^n) \partial_t u_j^n) \\ &= -\partial_x (f'(u_j^n) (-\partial_x f(u_j^n))) = \partial_x (f'(u_j^n)^2 \partial_x u_j^n) \\ &= \frac{f'(u_{j+\frac{1}{2}}^n)^2 \partial_x u_{j+\frac{1}{2}}^n - f'(u_{j-\frac{1}{2}}^n)^2 \partial_x u_{j-\frac{1}{2}}^n}{\Delta x} + \mathcal{O}(\Delta x) \\ &= f'(u_{j+\frac{1}{2}}^n)^2 \frac{u_{j+1}^n - u_j^n}{\Delta x^2} - f'(u_{j-\frac{1}{2}}^n)^2 \frac{u_j^n - u_{j-1}^n}{\Delta x^2} + \mathcal{O}(\Delta x) \end{aligned}$$

## Taylor Expansion

We have

$$u_j^{n+1} = u_j^n + \Delta t \partial_t u_j^n + \frac{\Delta t^2}{2} \partial_{tt} u_j^n + \mathcal{O}(\Delta t^3),$$

where we approximate

$$\partial_t u_j^n = -\frac{f(u_{j+1}^n) - f(u_{j-1}^n)}{2\Delta x} + \mathcal{O}(\Delta x^2).$$

and

$$\partial_{tt} u_j^n = f'(u_{j+\frac{1}{2}}^n)^2 \frac{u_{j+1}^n - u_j^n}{\Delta x^2} - f'(u_{j-\frac{1}{2}}^n)^2 \frac{u_j^n - u_{j-1}^n}{\Delta x^2} + \mathcal{O}(\Delta x).$$

Question: How do we choose  $f'(u_{j\pm\frac{1}{2}}^n)$ ?

## Taylor Expansion

Question: How do we choose  $f'(u_{j\pm\frac{1}{2}}^n)$ ? It holds

$$f'(u_{j+\frac{1}{2}}^n) = \frac{f(u_{j+1}^n) - f(u_j^n)}{u_{j+1}^n - u_j^n} + \underbrace{\mathcal{O}(u_{j+1}^n - u_j^n)}_{\mathcal{O}(\Delta x)}$$

Hence

$$f'(u_{j+\frac{1}{2}}^n)^2 (u_{j+1}^n - u_j^n) = \frac{(f(u_{j+1}^n) - f(u_j^n))^2}{u_{j+1}^n - u_j^n} + \mathcal{O}(\Delta x^3).$$

We obtain

$$\begin{aligned} \partial_{tt} u_j^n &= f'(u_{j+\frac{1}{2}}^n)^2 \frac{u_{j+1}^n - u_j^n}{\Delta x^2} - f'(u_{j-\frac{1}{2}}^n)^2 \frac{u_j^n - u_{j-1}^n}{\Delta x^2} + \mathcal{O}(\Delta x) \\ &= \frac{1}{\Delta x^2} \frac{(f(u_{j+1}^n) - f(u_j^n))^2}{u_{j+1}^n - u_j^n} - \frac{1}{\Delta x^2} \frac{(f(u_j^n) - f(u_{j-1}^n))^2}{u_j^n - u_{j-1}^n} + \mathcal{O}(\Delta x) \end{aligned}$$

## Taylor Expansion

Combining  $u_j^{n+1} = u_j^n + \Delta t \partial_t u_j^n + \frac{\Delta t^2}{2} \partial_{tt} u_j^n + \mathcal{O}(\Delta t^3)$  with

$$\partial_t u_j^n = -\frac{f(u_{j+1}^n) - f(u_{j-1}^n)}{2\Delta x} + \mathcal{O}(\Delta x^2).$$

and

$$\partial_{tt} u_j^n = \frac{1}{\Delta x^2} \frac{\left(f(u_{j+1}^n) - f(u_j^n)\right)^2}{u_{j+1}^n - u_j^n} - \frac{1}{\Delta x^2} \frac{\left(f(u_j^n) - f(u_{j-1}^n)\right)^2}{u_j^n - u_{j-1}^n} + \mathcal{O}(\Delta x)$$

we obtain

$$\begin{aligned} u_j^{n+1} = & u_j^n - \frac{\Delta t}{2\Delta x} [f(u_{j+1}^n) - f(u_{j-1}^n)] \\ & + \frac{\Delta t^2}{2\Delta x^2} \left[ \frac{\left(f(u_{j+1}^n) - f(u_j^n)\right)^2}{u_{j+1}^n - u_j^n} - \frac{\left(f(u_j^n) - f(u_{j-1}^n)\right)^2}{u_j^n - u_{j-1}^n} \right] + \Delta t \mathcal{O}(\Delta x^2) \end{aligned}$$

## Taylor Expansion

Based on

$$u_j^{n+1} = u_j^n - \frac{\Delta t}{2\Delta x} [f(u_{j+1}^n) - f(u_{j-1}^n)] \\ + \frac{\Delta t^2}{2\Delta x^2} \left[ \frac{(f(u_{j+1}^n) - f(u_j^n))^2}{u_{j+1}^n - u_j^n} - \frac{(f(u_j^n) - f(u_{j-1}^n))^2}{u_j^n - u_{j-1}^n} \right] + \Delta t \mathcal{O}(\Delta x^2)$$

we can formulate a scheme with **consistency order 2**. Note that the term  $\mathcal{O}(\Delta x^2)$  is just the local truncation error.

The resulting scheme is called **Lax-Wendroff Scheme** with **numerical flux**

$$g(v, w) := \frac{1}{2} \left[ f(v) + f(w) - \lambda \frac{(f(w) - f(v))^2}{w - v} \right], \quad \lambda = \frac{\Delta t}{\Delta x}.$$



# Lax-Wendroff Scheme

## Summary:

With Lax-Wendroff flux

$$g(v, w) := \frac{1}{2} \left[ f(v) + f(w) - \lambda \frac{(f(w) - f(v))^2}{w - v} \right], \quad \lambda = \frac{\Delta t}{\Delta x},$$

the Lax-Wendroff Scheme is given by

$$Q_j^{n+1} = Q_j^n - \frac{\Delta t}{\Delta x} (g(Q_j^n, Q_{j+1}^n) - g(Q_{j-1}^n, Q_j^n)).$$

# Lax-Wendroff Scheme

## The Lax-Wendroff Scheme

- ▶ is a scheme in conservation form,
- ▶ is consistent of order 2,
- ▶ is not monotone,
- ▶ yields typically no converge to the entropy solution,
- ▶ produces typically strong oscillations close to the discontinuities.
- ▶ Idea by Sweby: “Erase” the higher order contributions close to the shocks.

## Modified equations - a comparison

Consider the **linear problem**

$$\partial_t u + a \partial_x u = 0, \quad \text{for } a > 0.$$

We saw for **monotone schemes** that they better approximate **modified equations** with an **artificial diffusion term**. In the **Lax-Wendroff scheme**, we “killed” this term artificially to obtain a **higher order**.

### Comparison:

- **Lax-Friedrichs Scheme**. Modified equation:

$$\partial_t u + a \partial_x u = \frac{\Delta x}{2\lambda} \partial_{xx} u \quad \text{diffusive} - \underline{\text{everything smeared out}}$$

- **Lax-Wendroff Scheme**. Modified equation:

$$\partial_t u + a \partial_x u = \frac{\Delta x^2}{6} a(1 - (a\lambda)^2) \partial_{xxx} u \quad \text{dispersion} - \underline{\text{artificial oscillations}}$$

## Remark: Godunov Theorem

**Question:** Can we find a “better” second order scheme without oscillations?

**Answer:** No. The so called **Godunov Theorem** says:

Any linear scheme for solving conservation laws with the property that it does not create new extrema is at most of order 1.

Hence: there is no scheme of consistency order 2 that is free from the artificial oscillations.

## Lax-Wendroff Scheme - Example

We consider the **linear problem**

$$\partial_t u + a \partial_x u = 0, \quad \text{for } a > 0.$$

In this case, the **Lax-Wendroff flux** reads

$$g(v, w) = \frac{1}{2} a [v + w - \lambda a (v - w)]$$

and the scheme becomes

$$\begin{aligned} Q_j^{n+1} &= Q_j^n - \lambda a (Q_j^n - Q_{j-1}^n) - \frac{\lambda a}{2} (1 - \lambda a) (Q_{j+1}^n - 2Q_j^n + Q_{j-1}^n) \\ &= Q_j^n - \underbrace{\lambda a \Delta_- Q_j^n}_{\text{"upwind"-part 1. order}} - \underbrace{\frac{\lambda a}{2} (1 - \lambda a) \Delta_- \Delta_+ Q_j^n}_{\text{higher order corrector}} \end{aligned}$$