Lecture 5

Hyperbolic Equations of first order - Part 2

Finite Volume Method for Conservation Laws

SF2521

General conservation law - Model problem

For

- ▶ a smooth flux $f: \mathbb{R}^m \to \mathbb{R}^m$
- ▶ and initial value $\mathbf{v}: [\mathbf{0}, \mathbf{1}] \to \mathbb{R}^m$

we seek

$$\mathbf{u} = \mathbf{u}(x,t) : [0,1] \times \mathbb{R}^+ \to \mathbb{R}^m$$

with

$$\partial_t \mathbf{u} + \partial_x \mathbf{f}(\mathbf{u}) = \mathbf{o}$$
 and $\mathbf{u}(x, \mathbf{o}) = \mathbf{v}(x)$.

for
$$0 < x < 1$$
 and $t \in (0, \infty)$.

General conservation law - Model problem

For 0 < x < 1 and $t \in (0, \infty)$ find $\mathbf{u}(x, t)$ with

$$\partial_t \mathbf{u} + \partial_x f(\mathbf{u}) = \mathbf{o}$$
 and $\mathbf{u}(x, \mathbf{o}) = \mathbf{v}(x)$.

Note: the problem can be

- ▶ Scalar; m = 1, e.g.
 - $f(\mathbf{u}) = \mathbf{a} \mathbf{u}$, linear wave propagation;
 - $f(\mathbf{u}) = \frac{1}{2}\mathbf{u}^2$, Burgers' equation;
- ► System; *m* > 1, e.g.
 - linear case: for $\mathbf{u} = \begin{pmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \end{pmatrix}$ we have $f(\mathbf{u}) = \mathbf{A} \begin{pmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \end{pmatrix}$
 - nonlinear case: for $\mathbf{u} = \begin{pmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \end{pmatrix}$ we have $f(\mathbf{u}) = \begin{pmatrix} f_1(\mathbf{u}_1, \mathbf{u}_2) \\ f_2(\mathbf{u}_1, \mathbf{u}_2) \end{pmatrix}$

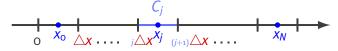
Finite Volume Discretization

For 0 < x < 1 and $t \in (0, \infty)$ find $\mathbf{u}(x, t)$ with

$$\partial_t \mathbf{u} + \partial_x f(\mathbf{u}) = \mathbf{o}$$
 and $\mathbf{u}(x, \mathbf{o}) = \mathbf{v}(x)$.

Finite Volume Discretization: more or less as before.

1. Discretize in space into *N* cells of size $\triangle x = 1/N$.



Here
$$x_j = \frac{\Delta x}{2} + j\Delta x$$
 for $j = 0, 1, 2, \dots, N-1$.

2. Discretize in time. For time step size $\triangle t$ we have $t^n = n \triangle t$ where $n = 0, 1, 2, \cdots$.

Finite Volume Discretization

3. Derive exact update formula as follows.

Conservation law on $C_j = [x_{j-\frac{1}{2}}, x_{j+\frac{1}{2}}]$ scaled with $1/\Delta x$:

$$o = \frac{1}{\Delta x} \int_{C_j} \partial_t \mathbf{u}(x,t) \, dx + \frac{1}{\Delta x} \int_{C_j} \partial_x f(\mathbf{u}(x,t)) \, dx$$
$$= \frac{1}{\Delta x} \int_{C_j} \partial_t \mathbf{u}(x,t) \, dx + \frac{f(\mathbf{u}(x_{j+\frac{1}{2}},t)) - f(\mathbf{u}(x_{j-\frac{1}{2}},t))}{\Delta x}.$$

Integrating in time over $[t^n, t^{n+1}]$:

$$o = \frac{1}{\triangle x} \int_{C_j} \int_{t^n}^{t^{n+1}} \partial_t \mathbf{u}(x, t) dt dx + \int_{t^n}^{t^{n+1}} \frac{f(\mathbf{u}(x_{j+\frac{1}{2}}, t)) - f(\mathbf{u}(x_{j-\frac{1}{2}}, t))}{\triangle x} dt$$

$$= \frac{1}{\triangle x} \left(\int_{C_j} \mathbf{u}(x, t^{n+1}) dx - \int_{C_j} \mathbf{u}(x, t^n) dx \right) + \int_{t^n}^{t^{n+1}} \frac{f(\mathbf{u}(x_{j+\frac{1}{2}}, t)) - f(\mathbf{u}(x_{j-\frac{1}{2}}, t))}{\triangle x} dt$$

Finite Volume Discretization

4. Approximation of equation. We have

$$\underbrace{\frac{1}{\triangle \mathbf{x}} \int_{C_j} \mathbf{u}(x, t^{n+1}) dx}_{\text{cell average} \approx Q_i^{n+1}} = \underbrace{\frac{1}{\triangle \mathbf{x}} \int_{C_j} \mathbf{u}(x, t^n) dx}_{\text{cell average} \approx Q_i^n} + \underbrace{\frac{1}{\triangle \mathbf{x}} \int_{t^n}^{t^{n+1}} f(\mathbf{u}(x_{j+\frac{1}{2}}, t)) - f(\mathbf{u}(x_{j-\frac{1}{2}}, t)) dt}_{\text{cell average} \approx Q_i^n}$$

Introduce suitable approximations:

► for the cell average of the wave

$$Q_j^n \approx \frac{1}{\triangle x} \int_{C_i} \mathbf{u}(x, t^n) dx.$$

▶ for the time average of the flux (depending on Q_j^n and Q_{j+1}^n)

$$F(Q_j^n,Q_{j+1}^n) pprox rac{1}{\triangle t} \int_{t^n}^{t^{n+1}} f(\mathbf{u}(x_{j+rac{1}{2}},t)) dt.$$

► Hence
$$Q_j^{n+1} = Q_j^n - \frac{\Delta t}{\Delta x} \left(F(Q_j^n, Q_{j+1}^n) - F(Q_{j-1}^n, Q_j^n) \right)$$

Finite Volume Discretization

Finite Volume Scheme in conservation form:

$$Q_{j}^{n+1} = Q_{j}^{n} - \frac{\Delta t}{\Delta x} \left(F(Q_{j}^{n}, Q_{j+1}^{n}) - F(Q_{j-1}^{n}, Q_{j}^{n}) \right)$$

Possible choices for $F(Q_i^n, Q_{i+1}^n)$? Examples:

► Lax-Friedrich scheme

$$F(U,V) = \frac{1}{2} \left(f(U) + f(V) - \frac{\triangle x}{\triangle t} (U - V) \right).$$

Lax-Wendroff scheme

$$F(U,V) = \frac{1}{2} \left(f(U) + f(V) - \frac{\triangle t}{\triangle x} \frac{(f(V) - f(U))^2}{V - U} \right).$$

Finite Volume Discretization - Lax-Friedrich

Lax-Friedrich scheme Motivation:

in $\partial_t \mathbf{u} + \partial_x f(\mathbf{u}) = \mathbf{o}$ we approximate

$$\partial_t \mathbf{u}(x_j, t^n) \approx \frac{\mathbf{u}(x_j, t^{n+1}) - \frac{\mathbf{u}(x_{j+1}, t^n) + \mathbf{u}(x_{j-1}, t^n)}{2}}{\triangle t} \approx \frac{Q_j^{n+1} - \frac{Q_{j+1}^n + Q_{j-1}^n}{2}}{\triangle t}$$

and

$$\partial_{\mathbf{x}} f(\mathbf{u}(\mathbf{x}_{j}, t^{n})) \approx \frac{f(\mathbf{u}(\mathbf{x}_{j+1}, t^{n})) - f(\mathbf{u}(\mathbf{x}_{j}, t^{n}))}{\triangle \mathbf{x}} \approx \frac{f(Q_{j+1}^{n}) - f(Q_{j-1}^{n})}{2 \triangle \mathbf{x}}$$

Hence

$$Q_{j}^{n+1} = \frac{1}{2} \left(Q_{j+1}^{n} + Q_{j-1}^{n} \right) - \frac{\Delta t}{2 \Delta x} \left(f(Q_{j+1}^{n}) - f(Q_{j-1}^{n}) \right)$$

Finite Volume Discretization - Lax-Friedrich

This coincides with the Lax-Friedrichs scheme, because

$$\begin{split} Q_{j}^{n+1} &= \frac{1}{2} \left(Q_{j+1}^{n} + Q_{j-1}^{n} \right) - \frac{\Delta t}{2\Delta x} \left(f(Q_{j+1}^{n}) - f(Q_{j-1}^{n}) \right) \\ &= Q_{j}^{n} - \frac{\Delta t}{2\Delta x} \left(f(Q_{j+1}^{n}) - f(Q_{j-1}^{n}) \right) + \frac{1}{2} \left(Q_{j+1}^{n} - 2Q_{j}^{n} + Q_{j-1}^{n} \right) \\ &= Q_{j}^{n} - \frac{\Delta t}{\Delta x} \left\{ \frac{1}{2} \left(f(Q_{j+1}^{n}) + f(Q_{j}^{n}) - f(Q_{j}^{n}) - f(Q_{j-1}^{n}) \right) \right. \\ &\left. + \frac{\Delta x}{2\Delta t} \left((Q_{j+1}^{n} - Q_{j}^{n}) - (Q_{j}^{n} - Q_{j-1}^{n}) \right) \right\} \\ &= Q_{j}^{n} - \frac{\Delta t}{\Delta x} \left(F(Q_{j+1}^{n}, Q_{j}^{n}) - F(Q_{j}^{n}, Q_{j-1}^{n}) \right) \end{split}$$

with the Lax-Friedrich flux

$$F(U,V) := \frac{1}{2} \left(f(U) + f(V) - \frac{\triangle X}{\triangle t} (U - V) \right).$$

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Finite Volume Discretizations - Short comparison

Example: linear, scalar Riemann problem

$$\partial_t \mathbf{u} + \mathbf{a} \partial_x \mathbf{u} = \mathbf{0}$$
 and

$$\mathbf{u}(x, 0) = \mathbf{v}(x).$$

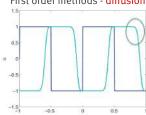
Lax-Friedrichs scheme

$$Q_{j}^{n+1} = \frac{1}{2} \left(Q_{j+1}^{n} + Q_{j-1}^{n} \right) - \frac{\triangle ta}{2 \triangle x} \left(Q_{j+1}^{n} - Q_{j-1}^{n} \right).$$

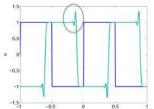
Lax-Wendroff scheme

$$\begin{split} Q_j^{n+1} &= Q_j^n - \frac{\Delta ta}{2\Delta x} \left(Q_{j+1}^n - Q_{j-1}^n\right) \\ &+ \frac{1}{2} \left(\frac{\Delta ta}{\Delta x}\right)^2 \left(Q_{j+1}^n - 2Q_j^n + Q_{j-1}^n\right). \end{split}$$

First order methods - diffusion

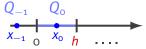


Second order methods - dispersion



Finite Volume Discretizations - Boundary conditions

Note on boundary conditions (cf. HW2):



If boundary does not coincide with nodes \Rightarrow use ghost points.

- Reflecting boundary condition at x = 0, e.g. $\mathbf{u}(0,t) = 0$.

 Approximated by average $\frac{Q_{-1}+Q_0}{2} = 0 \Rightarrow Q_0 = -Q_{-1}$.
- Non-reflecting boundary condition at x = 0: wave behaves as if there is no boundary \Rightarrow Approximated by extrapolation.

