

Exercise 3:

Let  $V_0 = L_x L_y L_z$  and  $\phi_{\vec{k}}(\vec{x}) = \frac{1}{\sqrt{V_0}} e^{i\vec{k} \cdot \vec{x}}$ .

1.: Boundary conditions:  $\phi_{\vec{k}}(x, y, z) = \phi_{\vec{k}}(x + L_x, y, z) = \phi_{\vec{k}}(x, y + L_y, z) = \phi_{\vec{k}}(x, y, z + L_z)$

Thus for  $x$ :  $\phi_{\vec{k}}(x, y, z) = \phi_{\vec{k}}(x + L_x, y, z)$

$$\Rightarrow e^{i(L_k x + L_y y + L_z z)} = e^{i(L_k(x + L_x) + L_y y + L_z z)}$$

$$\Rightarrow \exp[i L_k x] = \exp[i L_k (x + L_x)] \quad | \text{multiplying by } 1 = e^{i 2\pi n_x}$$

$$\Rightarrow 2\pi \cdot n_x + L_k \cdot x = L_k (x + L_x)$$

$$\Rightarrow 2\pi \cdot n_x = L_k \cdot L_x$$

$$\Rightarrow L_k = 2\pi \frac{n_x}{L_x}$$

Analogously this is done for  $y$  and  $z$ . Thus:

$$\vec{k} = 2\pi \left( \frac{n_x}{L_x}, \frac{n_y}{L_y}, \frac{n_z}{L_z} \right). \quad \square$$

2.:

$$\langle \phi_{\vec{k}} | \phi_{\vec{k}'} \rangle = \int_{V_0} d\vec{x} \phi_{\vec{k}}^*(\vec{x}) \phi_{\vec{k}'}(\vec{x}) \quad | \text{boundary conditions: } \mathbb{R} = \mathbb{Z} = V_0$$

$$= \int_{V_0} d\vec{x} \frac{1}{\sqrt{V_0}} e^{-i\vec{k} \cdot \vec{x}} \cdot \frac{1}{\sqrt{V_0}} e^{i\vec{k}' \cdot \vec{x}}$$

$$\text{I:} \quad = \frac{1}{V_0} \int_{V_0} d\vec{x} e^{i\vec{x} \cdot (\vec{k}' - \vec{k})}$$

$$= \frac{1}{V_0} \int_{V_0} d\vec{x} \exp\left(i\left[\frac{x}{L_x}(n'_x - n_x) + \frac{y}{L_y}(n'_y - n_y) + \frac{z}{L_z}(n'_z - n_z)\right]\right)$$

$$= \frac{1}{V_0} \int_{L_x} e^{i \frac{x}{L_x}(n'_x - n_x)} dx \int_{L_y} e^{i \frac{y}{L_y}(n'_y - n_y)} dy \int_{L_z} e^{i \frac{z}{L_z}(n'_z - n_z)} dz \quad | \int_a^a \delta(x) = \frac{1}{a} \delta(x)$$

$$= \frac{1}{V_0} L_x \delta(n'_x - n_x) L_y \delta(n'_y - n_y) L_z \delta(n'_z - n_z) \quad | n_i \in \mathbb{N} \Rightarrow \delta(n'_i - n_i) = \delta_{n'_i, n_i}$$

$$= \delta_{n'_x, n_x} \delta_{n'_y, n_y} \delta_{n'_z, n_z} \quad \square$$

From (I) we could already see, that  $\langle \phi_{\vec{k}} | \phi_{\vec{k}} \rangle = \delta_{\vec{k}, \vec{k}}$ .

3.: Let  $\hat{T} + \hat{V} = \hat{H} = -\frac{\hbar^2}{2m} \Delta + \hat{V}(x)$ , thus  $\hat{T} = -\frac{\hbar^2}{2m} \Delta$ .

$$\begin{aligned} \Rightarrow \langle \psi' | \hat{T} | \psi \rangle &= \langle \psi' | (-\frac{\hbar^2}{2m}) \Delta | \psi \rangle \quad | \Delta(x) = \Delta(x) \rangle \\ &= -\frac{\hbar^2}{2m} \langle \psi' | x \rangle \Delta \langle x | \psi \rangle \quad | \langle x | \psi \rangle = \phi_{\vec{k}}(\vec{x}) \\ &= -\frac{\hbar^2}{2m} \langle \psi' | x \rangle \Delta \frac{1}{\sqrt{V_0}} e^{i\vec{k}\vec{x}} \\ &= \frac{\hbar^2}{2m} \langle \psi' | x \rangle \vec{k}^2 \frac{1}{\sqrt{V_0}} e^{i\vec{k}\vec{x}} \\ &= \frac{\hbar^2}{2m} \vec{k}^2 \langle \psi' | \psi \rangle \\ &= \frac{\hbar^2}{2m} \vec{k}^2 \delta_{\vec{k}'\vec{k}} \end{aligned}$$

□

4:

$$\begin{aligned} \frac{1}{V_0} a_{\vec{k}-\vec{k}'} &= \langle \vec{k}' | U | \vec{k} \rangle \\ &= \int_{\mathbb{R}} d\vec{x} \phi_{\vec{k}'}^*(\vec{x}) U(x) \phi_{\vec{k}}(\vec{x}) \quad | \text{boundary conditions} \\ &= \int_{V_0} d\vec{x} \frac{1}{V_0} U(x) e^{i\vec{x}(\vec{k}-\vec{k}')} \\ &= \int_{\mathbb{R}} d\vec{x} U(x) e^{i\vec{x}(\vec{k}-\vec{k}')} \\ &= \mathcal{F}[U(x)](\vec{k}-\vec{k}') \end{aligned}$$

□

5:

With  $\langle x | \vec{p}, \vec{k} \rangle = \phi_{\vec{p}, \vec{k}}(x) = \frac{1}{\sqrt{V_0}} e^{i\vec{x}\vec{p}} e^{i\vec{x}\vec{k}}$  and  $V(\vec{x}-\vec{x}') = \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} follows:$

$$\begin{aligned} \langle \vec{p}', \vec{k}' | V(\vec{x}-\vec{x}') | \vec{p}, \vec{k} \rangle &= \int_{V_0} d\vec{x} \frac{1}{V_0} e^{i\vec{x}(\vec{p}'-\vec{p})} e^{i\vec{x}(\vec{k}'-\vec{k})} \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} \\ &= \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} \int_{\mathbb{R}} d\vec{x} e^{i\vec{x}(\vec{p}'-\vec{p})} e^{i\vec{x}(\vec{k}'-\vec{k})} e^{i\vec{q}\cdot(\vec{x}-\vec{x}')} \\ &= \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} \delta_{\vec{q}, \vec{p}-\vec{p}'} \delta_{\vec{q}, \vec{k}-\vec{k}'} \end{aligned}$$

□

6.: Let  $\hat{H} = \hat{T} + \hat{U} + \hat{V}$ .

Find  $\hat{T}$ :  $\langle \vec{h}' | \hat{T} | \vec{h} \rangle = \frac{\hbar^2}{2m} \vec{h}^2 \delta_{\vec{h}', \vec{h}}$

$$\begin{aligned} \Rightarrow \hat{T} &= \sum_{\vec{h}, \vec{h}'} \langle \vec{h}' | \hat{T} | \vec{h} \rangle |\vec{h}'\rangle \langle \vec{h}| = \sum_{\alpha} |i_{\alpha}\rangle \langle j_{\alpha}| = \hat{a}_i^{\dagger} \hat{a}_j \\ &= \sum_{\vec{h}, \vec{h}'} \frac{\hbar^2}{2m} \vec{h}^2 \delta_{\vec{h}', \vec{h}} \hat{a}_{\vec{h}'}^{\dagger} \hat{a}_{\vec{h}} \\ &= \sum_{\vec{h}} \frac{(\hbar \vec{h})^2}{2m} \hat{a}_{\vec{h}}^{\dagger} \hat{a}_{\vec{h}} \end{aligned}$$

Find  $\hat{U}$ :  $\langle \vec{h}' | \hat{U} | \vec{h} \rangle = \frac{1}{V_0} U_{\vec{h}' - \vec{h}}$

$$\begin{aligned} \Rightarrow \hat{U} &= \sum_{\vec{h}', \vec{h}} \langle \vec{h}' | \hat{U} | \vec{h} \rangle |\vec{h}'\rangle \langle \vec{h}| \\ &= \sum_{\vec{h}', \vec{h}} \frac{1}{V_0} U_{\vec{h}' - \vec{h}} \hat{a}_{\vec{h}'}^{\dagger} \hat{a}_{\vec{h}} \\ &= \frac{1}{V_0} \sum_{\vec{h}', \vec{h}} U_{\vec{h}' - \vec{h}} \hat{a}_{\vec{h}'}^{\dagger} \hat{a}_{\vec{h}} \end{aligned}$$

Find  $\hat{V}$ :  $\langle \vec{p}', \vec{h}' | V(\vec{x} - \vec{x}') | \vec{p}, \vec{h} \rangle = \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} \delta_{\vec{q}, \vec{p}' - \vec{p}} \delta_{\vec{q}, \vec{h} - \vec{h}'}$

From exercise 2 we know:  $\hat{F} = \frac{\epsilon}{2} \sum_{\alpha \neq \beta} \sum_{i,j,k,l} \langle i,j | \hat{f} | k,l \rangle |i\rangle_{\alpha} |j\rangle_{\beta} \langle k|_{\alpha} \langle l|_{\beta}$

$$\begin{aligned} \text{Thus: } \hat{V} &= \frac{\epsilon}{2} \sum_{\alpha \neq \beta} \sum_{\vec{p}', \vec{h}', \vec{p}, \vec{h}} \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} \delta_{\vec{q}, \vec{p}' - \vec{p}} \delta_{\vec{q}, \vec{h} - \vec{h}'} |\vec{p}'\rangle_{\alpha} |\vec{h}'\rangle_{\beta} \langle \vec{h}|_{\alpha} \langle \vec{p}|_{\beta} \\ &= \frac{1}{2} \sum_{\vec{p}', \vec{h}', \vec{p}, \vec{h}} \frac{1}{V_0} \sum_{\vec{q}} V_{\vec{q}} \delta_{\vec{q}, \vec{p}' - \vec{p}} \delta_{\vec{q}, \vec{h} - \vec{h}'} \hat{a}_{\vec{p}'}^{\dagger} \hat{a}_{\vec{h}'}^{\dagger} \hat{a}_{\vec{h}} \hat{a}_{\vec{p}} \\ &= \frac{1}{2V_0} \sum_{\vec{h}, \vec{p}, \vec{q}} V_{\vec{q}} \hat{a}_{\vec{q} + \vec{p}}^{\dagger} \hat{a}_{\vec{h} - \vec{q}}^{\dagger} \hat{a}_{\vec{h}} \hat{a}_{\vec{p}} \end{aligned}$$

Finally we construct  $\hat{H}$ :

$$\hat{H} = \sum_{\vec{h}} \frac{(\hbar \vec{h})^2}{2m} \hat{a}_{\vec{h}}^{\dagger} \hat{a}_{\vec{h}} + \frac{1}{V_0} \sum_{\vec{h}', \vec{h}} U_{\vec{h}' - \vec{h}} \hat{a}_{\vec{h}'}^{\dagger} \hat{a}_{\vec{h}} + \frac{1}{2V_0} \sum_{\vec{h}, \vec{p}, \vec{q}} V_{\vec{q}} \hat{a}_{\vec{q} + \vec{p}}^{\dagger} \hat{a}_{\vec{h} - \vec{q}}^{\dagger} \hat{a}_{\vec{h}} \hat{a}_{\vec{p}}$$

$\hat{T}$ , the kinetic energy, is independent of other particles, thus only dependant on the own wave vector  $\vec{h}$ .  $\hat{U}$ , an arbitrary external potential, can be dependant of other particles, thus being also dependant on the other wave vectors  $\vec{h}'$ . In particular it is dependant of its own wave vector. This is not the case for  $\hat{V}$ , the two particle-potential.  $\hat{V}$  describes the forces created by the other particle, thus there can't be any if the other particle would be identical in every way.