

Exercise 1.

1:

Lorentz trans. ($SO(3, 1)$): ($\Lambda: \Lambda^T g \Lambda = g$) with $g = \text{diag}(1, -1, -1, -1)$

- Associativity: is trivial since associativity is provided by matrix mult..
- Identity element: $\mathbb{1} = \text{diag}(1, 1, 1, 1) \Rightarrow A \cdot \mathbb{1} = A$ with $A \in \mathbb{R}^4$
(for Λ := Lorentz boost: $\Lambda(u=0) = \mathbb{1}$ since $\gamma(u=0)=1$ and $\beta(u=0)=0$.)

- Inverse element: $g = \Lambda g \Lambda^T = (\Lambda g) \Lambda^T = (\Lambda (\Lambda g)^T)^T$
 \Rightarrow for $\Lambda g \rightarrow \Lambda \cdot g$ the inverse is
 $\Lambda^{-1} g \rightarrow (\Lambda \cdot g^T)^T = g \cdot \Lambda^T$, since
 $\Lambda \circ (\Lambda^{-1} g) \rightarrow (\Lambda (\Lambda g)^T)^T = \Lambda g \Lambda^T = g$

$$\text{Rotation } (SO(3, 1)): O_z(\theta) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & \sin \theta & 0 \\ 0 & -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

- Associativity: is trivial since associativity is provided by matrix mult..
- Identity element: $O_z(\theta = 2\pi n) = \mathbb{1}$ with $n \in \mathbb{N}$
 $\mathbb{1} \cdot A = A$ with $A \in \mathbb{R}^4$

- Inverse element: $O_z^{-1}(\theta) = O_z(-\theta) \Rightarrow O_z(-\theta) \cdot O_z(\theta) = \mathbb{1}$

$$\Rightarrow \mathbb{1} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & \sin \theta & 0 \\ 0 & -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos(-\theta) & \sin(-\theta) & 0 \\ 0 & -\sin(-\theta) & \cos(-\theta) & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad \left| \begin{array}{l} \cos(-\theta) = \cos \theta \\ \sin(-\theta) = -\sin \theta \end{array} \right.$$

$$= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & \sin \theta & 0 \\ 0 & -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta & -\sin \theta & 0 \\ 0 & \sin \theta & \cos \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos^2 \theta + \sin^2 \theta & \sin^2 \theta + \cos^2 \theta & 0 \\ 0 & \cos^2 \theta - (-)\sin^2 \theta & -(-)\sin^2 \theta + \cos^2 \theta & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{cases} -(-)1 = 1 \\ \sin^2 \theta + \cos^2 \theta = 1 \end{cases}$$

$$= \text{diag}(1, 1, 1, 1)$$

$$= \mathbb{1}$$

□

For $O_x(\theta)$ and $O_y(\theta)$ the prove is analogously.

Thus with $\mathbb{1}$, $O_z(\theta)$, $O_x(\theta)$ and

$O_y(\theta)$ we can form a group: $SO(3, 1)$

2.: As shown in part 1, $SO(3, 1)$ contains the rotations around the x , y and z axis. Thus this set of all rotations is a subgroup:

$$SO(3) \in SO(3, 1) \text{ with } SO(3) = (O_i(\theta))_i \text{ and } SO(3, 1) = (O_i(\theta))_i \cup \mathbb{1}$$

That the set $SO(3)$ forms indeed a group was also shown.

3. In part 1 we already showed that the rotation at $O_z(\theta)$ around the (x, y) -plane forms a group. This is analogously for $O_x(\theta)$ around the (y, z) -plane. So since $O_x(\theta) \in SO(3)$, $O_x(\theta)$ is a subgroup of $SO(3)$. Here the set of all rotations for one axis means: $\{O_i(\theta) : \theta \in \mathbb{R}\}$.

Successive rotations α_1 or α_2 in one rotation: $O_x(\alpha_1) \cdot O_x(\alpha_2) = O_x(\alpha_1 + \alpha_2)$.

For that relation the trig. identity $\sin(\alpha_1 + \alpha_2) = \sin \alpha_1 \cos \alpha_2 + \cos \alpha_1 \sin \alpha_2$

and $\cos(\alpha_1 + \alpha_2) = \cos \alpha_1 \sin \alpha_2 - \sin \alpha_1 \cos \alpha_2$ is crucial. The remaining part of the prove is matrix multiplication.

4.: We already showed in part 1, that the any rank two tensor Λ satisfying $\Lambda^T g \Lambda = g$ and $\Lambda = \mathbb{1}$ forms a group and together with $SO(3)$ form $SO(3, 1)$.

Thus we need to show that Λ being the Lorentz boost:

$$\Lambda = \begin{pmatrix} \gamma & \beta\gamma & 0 & 0 \\ \beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \text{ satisfies } \Lambda^T g \Lambda = g \text{ and } \Lambda = \Lambda^{-1}.$$

At this point we should mention, that we use Λ very loosely and talk about it being a set. But Λ is an element of the set we talk about.

More correct would be calling

$$\Lambda \in \{ \Lambda | (\Lambda^T g \Lambda = g : g = \text{diag}(1, -1, -1, -1)) \wedge (\exists \Lambda : \Lambda = \Lambda^{-1}) \}.$$

Back to the prove:

$$\begin{aligned} \Lambda^T g \Lambda &= \begin{pmatrix} \gamma & \beta\gamma & 0 & 0 \\ \beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot g \cdot \begin{pmatrix} \gamma & \beta\gamma & 0 & 0 \\ \beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \gamma - \beta^2\gamma & 0 & 0 & 0 \\ -\beta\gamma - \gamma\beta & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \gamma & \beta\gamma & 0 & 0 \\ \beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \gamma^2 - \beta^2\gamma^2 & 0 & 0 & 0 \\ \beta^2\gamma^2 - \gamma^2 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad | \quad \gamma^2 - \beta^2\gamma^2 = \gamma^2(1 - \beta^2) = \frac{1 - \beta^2}{1 - \beta^2} = 1 \\ &= g \end{aligned}$$

As hinted in part 1: $\Lambda(u=0) = \mathbb{1}$, thus we get our identity element.

Successive boosts β_1 and β_2 :

$$\begin{aligned} \Lambda(\beta_2) \cdot \Lambda(\beta_1) &= \begin{pmatrix} \gamma_2 & \beta_2\gamma_2 & 0 & 0 \\ \beta_2\gamma_2 & \gamma_2 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \gamma_1 & \beta_1\gamma_1 & 0 & 0 \\ \beta_1\gamma_1 & \gamma_1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \gamma_1\gamma_2 + \beta_1\gamma_1\beta_2\gamma_2 & \gamma_1\beta_1\gamma_2 + \gamma_1\beta_2\gamma_2 & 0 & 0 \\ \gamma_1\beta_1\gamma_2 + \gamma_1\beta_2\gamma_2 & \gamma_1\gamma_2 + \beta_1\gamma_1\beta_2\gamma_2 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \end{aligned}$$

$$\Rightarrow x_\mu = \Lambda_\mu^\kappa(\beta_1) \cdot \Lambda_\kappa^\nu(\beta_2) \cdot x_\nu$$

$$= \begin{pmatrix} \gamma_1 \gamma_2 + \beta_1 \gamma_1 \beta_2 \gamma_2 & \gamma_1 \beta_1 \gamma_2 + \gamma_1 \beta_2 \gamma_2 & 0 & 0 \\ \gamma_1 \beta_1 \gamma_2 + \gamma_1 \beta_2 \gamma_2 & \gamma_1 \gamma_2 + \beta_1 \gamma_1 \beta_2 \gamma_2 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$$

$$= \begin{pmatrix} ct(\gamma_1 \gamma_2 + \beta_1 \gamma_1 \beta_2 \gamma_2) + x(\gamma_1 \beta_1 \gamma_2 + \gamma_1 \beta_2 \gamma_2) \\ ct(\gamma_1 \beta_1 \gamma_2 + \gamma_1 \beta_2 \gamma_2) + x(\gamma_1 \gamma_2 + \beta_1 \gamma_1 \beta_2 \gamma_2) \\ y \\ z \end{pmatrix}$$

Or for simply: $x_t = \gamma_2(x - v_2 t)$ with $x = \gamma_1(x_i - v_1 t)$

$$\Rightarrow x_t = \frac{\frac{x_i - v_1 t}{\sqrt{1 - \beta_1^2}} - v_2 t}{\sqrt{1 - \beta_2^2}} = \frac{\frac{x_i - \beta_1 ct}{\sqrt{1 - \beta_1^2}} - \beta_2 ct}{\sqrt{1 - \beta_2^2}}$$

5:

Let β_x and β_y be two successive boosts in x and y direction. The order does not matter.

In 2-d space we have: $\Lambda = \begin{pmatrix} \gamma & \beta \gamma & 0 \\ \beta \gamma & \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix}$

Let $\vec{u} = (\cos \theta, \sin \theta)$.

If $\vec{u}_T \perp \vec{u} \Rightarrow \theta_T = \theta + \frac{\pi}{2} \Rightarrow \vec{u}_T = (-\sin \theta, \cos \theta)$.

$$\vec{x} = a \vec{u} + b \vec{u}_T$$

$$\Rightarrow \vec{x} = \Lambda_{\vec{u}_T} \Lambda_{\vec{u}} \vec{x} \\ = \gamma_x(a - v_x t) \vec{u} + \gamma_y(b - v_y t) \vec{u}_T$$

$$\begin{aligned}\vec{x}' &= Q_2(\theta' = \theta + \tan^{-1}(\frac{v_y}{v_x})) \cdot \vec{x} \quad | \quad v_i = c\beta_i, \quad v = \sqrt{v_x^2 + v_y^2} \Rightarrow \beta = \sqrt{\beta_x^2 + \beta_y^2} \\ &= Q_2(\theta' = \theta + \tan^{-1}(\frac{v_y}{v_x})) \cdot \begin{pmatrix} \gamma(a - vt) \\ \gamma(b - vt) \end{pmatrix} \\ &= \begin{pmatrix} \cos(\theta + \tan^{-1}(\frac{v_y}{v_x}))\gamma(a - vt) + \sin(\theta + \tan^{-1}(\frac{v_y}{v_x}))\gamma(b - vt) \\ -\sin(\theta + \tan^{-1}(\frac{v_y}{v_x}))\gamma(a - vt) + \cos(\theta + \tan^{-1}(\frac{v_y}{v_x}))\gamma(b - vt) \end{pmatrix}\end{aligned}$$

= ...

Essentially we see, that $\gamma \neq \sqrt{\gamma_x^2 + \gamma_y^2}$, thus we can't just do a boost with $v = |\vec{v}|$. Tbh I don't get why that's not possible.

But I don't know what other thing the exercise wants from me.

I think there should be a boost that is equivalent to to several boosts, since one could always just use the absolute value of the total boost velocity and just rotate to the new direction $\theta = \tan^{-1}(\frac{v_y}{v_x})$.

2) 4-Vectors and Relativistic Kinematics

1) Show: $a' \cdot b' = a \cdot b$ where $a'^\mu = \Lambda^\mu_\nu a^\nu$, $b'^\mu = \Lambda^\mu_\nu b^\nu$ and $\Lambda = \begin{pmatrix} \gamma & 0 & \beta\gamma & 0 \\ 0 & 1 & 0 & 0 \\ \beta\gamma & 0 & \gamma & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$

• $a' \cdot b' = a \cdot b$

$\Leftrightarrow a'_\mu b'^\mu = a_\mu b^\mu$

$\Leftrightarrow g_{\mu\nu} a'^\mu b'^\nu = g_{\mu\nu} a^\mu b^\nu \quad | \quad x'^\mu = \Lambda^\mu_\nu x^\nu$

$\Leftrightarrow g_{\mu\nu} \Lambda^\mu_\alpha \Lambda^\nu_\beta a^\alpha b^\beta = g_{\mu\nu} a^\mu b^\nu$

$\Leftrightarrow g_{\mu\nu} \Lambda^\mu_\alpha \Lambda^\nu_\beta a^\alpha b^\beta = g_{\alpha\beta} a^\alpha b^\beta \quad \checkmark \rightarrow \text{correct, because } g_{\mu\nu} \Lambda^\mu_\alpha \Lambda^\nu_\beta = g_{\alpha\beta} (*) \quad \square$

(*) Show: $g_{\mu\nu} \Lambda^\mu_\alpha \Lambda^\nu_\beta = g_{\alpha\beta}$

• for $(\alpha, \beta) = (0, 0)$

$\Rightarrow g_{\mu\nu} \Lambda^\mu_0 \Lambda^\nu_0 = g_{00} \Lambda^0_0 \Lambda^0_0 + g_{11} \Lambda^1_0 \Lambda^1_0 + g_{22} \Lambda^2_0 \Lambda^2_0 + g_{33} \Lambda^3_0 \Lambda^3_0 = \gamma^2 - (\beta\gamma)^2 = \gamma^2(1 - \beta^2) = 1 = g_{00} \quad \checkmark$

• for $(\alpha, \beta) = (1, 1)$

$\Rightarrow g_{\mu\nu} \Lambda^\mu_1 \Lambda^\nu_1 = g_{00} \Lambda^0_1 \Lambda^0_1 + g_{11} \Lambda^1_1 \Lambda^1_1 + g_{22} \Lambda^2_1 \Lambda^2_1 + g_{33} \Lambda^3_1 \Lambda^3_1 = -1 = g_{11} \quad \checkmark$

• for $(\alpha, \beta) = (2, 2)$

$\Rightarrow g_{\mu\nu} \Lambda^\mu_2 \Lambda^\nu_2 = g_{00} \Lambda^0_2 \Lambda^0_2 + g_{11} \Lambda^1_2 \Lambda^1_2 + g_{22} \Lambda^2_2 \Lambda^2_2 + g_{33} \Lambda^3_2 \Lambda^3_2 = (\beta\gamma)^2 - \gamma^2 = -\gamma^2(1 - \beta^2) = -1 = g_{22} \quad \checkmark$

• for $(\alpha, \beta) = (3, 3)$

$\Rightarrow g_{\mu\nu} \Lambda^\mu_3 \Lambda^\nu_3 = g_{00} \Lambda^0_3 \Lambda^0_3 + g_{11} \Lambda^1_3 \Lambda^1_3 + g_{22} \Lambda^2_3 \Lambda^2_3 + g_{33} \Lambda^3_3 \Lambda^3_3 = -1 = g_{33} \quad \checkmark$

• for (α, β) where $\alpha \neq \beta$

$\Rightarrow g_{\mu\nu} \Lambda^\mu_\alpha \Lambda^\nu_\beta = g_{00} \Lambda^0_\alpha \Lambda^0_\beta + g_{11} \Lambda^1_\alpha \Lambda^1_\beta + g_{22} \Lambda^2_\alpha \Lambda^2_\beta + g_{33} \Lambda^3_\alpha \Lambda^3_\beta = -1 \cdot \Lambda^0_\alpha \Lambda^0_\beta + 1 \cdot \Lambda^1_\alpha \Lambda^1_\beta = 0 = g_{\alpha\beta} (\alpha \neq \beta) \quad \checkmark \quad \square$

2) Show: $f(\vec{x}, t) = e^{\frac{i}{\hbar} p \cdot x}$ where $k^\mu = \frac{p^\mu}{\hbar}$ behaves like a plane wave

• $f(\vec{x}, t) = e^{\frac{i}{\hbar} p \cdot x} = e^{i k^\mu x_\mu} \quad | \quad k^\mu = \frac{p^\mu}{\hbar} = \left(\frac{E}{\hbar}, \frac{\vec{p}}{\hbar} \right), \quad x^\mu = (ct, \vec{x}) \Rightarrow k^\mu x_\mu = \frac{E}{\hbar} - \frac{\vec{p} \cdot \vec{x}}{\hbar} = \omega t - \vec{k} \cdot \vec{x} \quad \text{where } \omega = \frac{E}{\hbar} \wedge \vec{k} = \frac{\vec{p}}{\hbar}$

$\Rightarrow f(\vec{x}, t) = e^{i(\omega t - \vec{k} \cdot \vec{x})} \rightarrow \text{plane wave} \quad \square$

3) Show: $m_1 + m_2 \leq M$

• $p^\mu = k_1^\mu + k_2^\mu$

$\Leftrightarrow M u^\mu = m_1 u_1^\mu + m_2 u_2^\mu$

$\Rightarrow M u^0 = m_1 u_1^0 + m_2 u_2^0$

$\Leftrightarrow \gamma M c = \gamma_1 m_1 c + \gamma_2 m_2 c \quad | \quad \text{rest frame of } M: \gamma = 1$

$\Leftrightarrow M = \gamma_1 m_1 + \gamma_2 m_2 \quad | \quad \gamma \geq 1$

$\Rightarrow M \geq m_1 + m_2 \quad \square$

4) $\pi^0(p) \rightarrow \gamma(k_1) \gamma(k_2)$ (pion $\pi \rightarrow$ two photons γ)

• $\frac{d\Gamma}{dE_1^*} = \frac{d\Gamma}{d\Omega^*} \frac{d\Omega^*}{dE_1^*} \quad | \quad d\Omega^* = \sin\theta d\varphi d\theta$