

ON LEARNING PSEUDO-SENSORS TO IMPROVE EGOMOTION ESTIMATION FOR
MOBILE AUTONOMY

by

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Abstract

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The ability to estimate *egomotion*, that is, to track one’s own pose through an unknown environment, is at the heart of safe and reliable mobile autonomy. By inferring pose changes from sequential sensor measurements, egomotion estimation forms the basis of mapping and navigation pipelines, and permits mobile robots to self-localize within environments where external localization sources are intermittent or unavailable. Visual and inertial egomotion estimation, in particular, have become ubiquitous in mobile robotics due to the availability of high-quality, compact, and inexpensive sensors that capture rich representations of the world. To remain computationally tractable, ‘classical’ visual-inertial pipelines (like visual odometry and visual SLAM) make simplifying assumptions that, while permitting reliable operation in ideal conditions, often lead to systematic error. In this thesis, we present several data-driven learned *pseudo-sensors* that serve to augment conventional pipelines by inferring latent information from the same sensor data. Our approach retains much of the benefits of traditional pipelines, while leveraging high-capacity hyper-parametric models to extract complementary information that can be used to improve uncertainty quantification, correct for systematic bias, and improve robustness to difficult-to-model deleterious effects. We validate our pseudo-sensors on several kilometres of sensor data collected in sundry settings such as urban roads, indoor labs, and planetary analogue sites in the Canadian high arctic.

Epigraph

A little learning is a dangerous thing;
drink deep, or taste not the Pierian
spring: there shallow draughts
intoxicate the brain, and drinking
largely sobers us again.

ALEXANDER POPE

The universe is no narrow thing and the
order within it is not constrained by any
latitude in its conception to repeat what
exists in one part in any other part.
Even in this world more things exist
without our knowledge than with it and
the order in creation which you see is
that which you have put there, like a
string in a maze, so that you shall not
lose your way. For existence has its own
order and that no man's mind can
compass, that mind itself being but a
fact among others.

CORMAC MCCARTHY

Elephants don't play chess.

RODNEY BROOKS

To all those who encouraged (or, at least, *never discouraged*) my intellectual wanderlust.

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Notation

- a : Symbols in this font are real scalars.
- \mathbf{a} : Symbols in this font are real column vectors.
- \mathbf{A} : Symbols in this font are real matrices.
- $\mathcal{N}(\boldsymbol{\mu}, \mathbf{R})$: Normally distributed with mean $\boldsymbol{\mu}$ and covariance \mathbf{R} .
- $E[\cdot]$: The expectation operator.
- $\underline{\mathcal{F}}_{\rightarrow a}$: A reference frame in three dimensions.
- $(\cdot)^\wedge$: An operator associated with the Lie algebra for rotations and poses. It produces a matrix from a column vector.
- $(\cdot)^\vee$: The inverse operation of $(\cdot)^\wedge$
- $\mathbf{1}$: The identity matrix.
- $\mathbf{0}$: The zero matrix.
- $\mathbf{p}_a^{c,b}$: A vector from point b to point c (denoted by the superscript) and expressed in $\underline{\mathcal{F}}_{\rightarrow a}$ (denoted by the subscript).
- $\mathbf{C}_{a,b}$: The 3×3 rotation matrix that transforms vectors from $\underline{\mathcal{F}}_{\rightarrow b}$ to $\underline{\mathcal{F}}_{\rightarrow a}$: $\mathbf{p}_a^{c,b} = \mathbf{C}_{a,b} \mathbf{p}_b^{c,b}$.
- $\mathbf{T}_{a,b}$: The 4×4 transformation matrix that transforms homogeneous points from $\underline{\mathcal{F}}_{\rightarrow b}$ to $\underline{\mathcal{F}}_{\rightarrow a}$: $\mathbf{p}_a^{c,a} = \mathbf{T}_{a,b} \mathbf{p}_b^{c,b}$.

Chapter 2

Mathematical Foundations

By relieving the brain of all unnecessary work, a good notation sets it free to concentrate on more advanced problems, and, in effect, increases the mental power of the race.

ALFRED NORTH WHITEHEAD

2.1 Coordinate Frames

Before we can present the main contributions of this thesis, it will be useful to first outline the notation and mathematical foundations that underly the work. Throughout this thesis, we largely follow the notation of [Barfoot \(2017\)](#) when dealing with three-dimensional rigid-body kinematics.

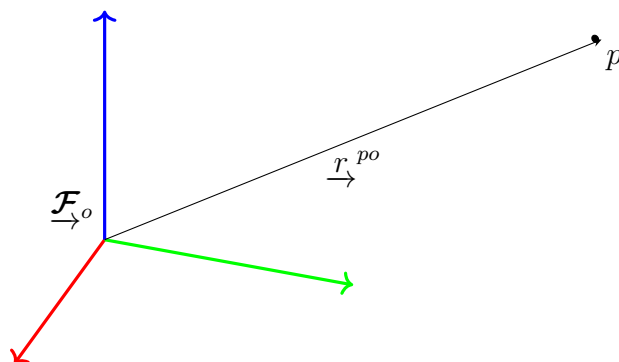


Figure 2.1: A position vector expressed in a coordinate frame.

We refer to a three-dimensional position vector, \underline{r}^{po} , as one that originates at the origin of a coordinate reference frame, \mathcal{F}_o , and terminates at the point p . This geometric quantity has

the numerical coordinates \mathbf{r}_o^{po} when expressed in \mathcal{F}_o . Often, we will refer to two reference frames such as a world or *inertial* frame, \mathcal{F}_i , and a vehicle frame, \mathcal{F}_v . Rotation matrices or rigid-body transformations that convert coordinates from \mathcal{F}_i to \mathcal{F}_v will be represented as \mathbf{T}_{vi} , and \mathbf{C}_{vi} ¹, respectively.

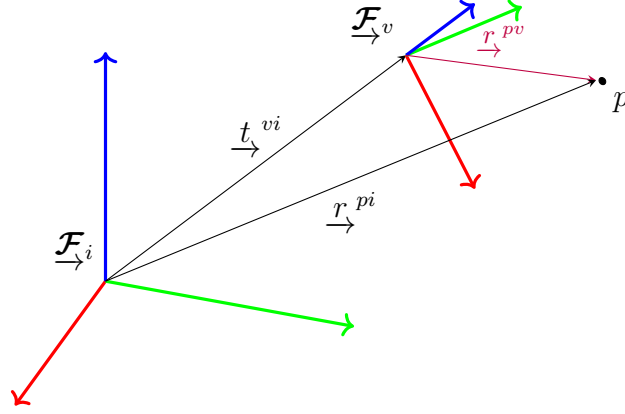


Figure 2.2: Two common reference frames used throughout this thesis.

2.2 Rotations

The rotation matrix \mathbf{C} is a member of the matrix Lie group $\text{SO}(3)$ (the Special Orthogonal group) and can be defined as a matrix as follows:

$$\text{SO}(3) = \{\mathbf{C} \in \mathbb{R}^{3 \times 3} \mid \mathbf{C}^T \mathbf{C} = \mathbf{1}, \det \mathbf{C} = 1\}. \quad (2.1)$$

Active vs. Passive

An active (or *alibi*) rotation changes the coordinates of a position directly while implicitly assuming that the reference frame is fixed. A passive (or *alias*) rotation rotates the reference frame. Following Barfoot (2017), all rotation matrices in this dissertation are passive unless otherwise noted.

Exponential and Logarithmic Maps

Since rotations form a matrix Lie group (we refer the reader to Solà et al. (2018) and Barfoot (2017) for a thorough treatment of Lie groups for state estimation), we can define a surjective

¹We use \mathbf{C} and not \mathbf{R} for rotation matrices to avoid confusion with common notation for measurement model covariance.

exponential map² from three axis-angle parameters, $\phi = \phi \mathbf{a}$, $\phi \in \mathbb{R}$, $\mathbf{a} \in S^2$, to a rotation matrix, \mathbf{C} :

$$\mathbf{C} = \text{Exp}(\phi) = \exp(\phi^\wedge) = \sum_{n=0}^{\infty} \frac{1}{n!} (\phi^\wedge)^n \quad (2.2)$$

$$= \cos \phi \mathbf{1} + (1 - \cos \phi) \mathbf{a} \mathbf{a}^T + \sin \phi \mathbf{a}^\wedge, \quad (2.3)$$

where the wedge operator $(\cdot)^\wedge$ ³ is defined as

$$\mathbf{a}^\wedge = \begin{bmatrix} a_0 \\ a_1 \\ a_2 \end{bmatrix}^\wedge = \begin{bmatrix} 0 & -a_2 & a_1 \\ a_2 & 0 & -a_0 \\ -a_1 & a_0 & 0 \end{bmatrix}. \quad (2.4)$$

Equation (2.3) is known as the Euler-Rodriguez formula and it can also be derived geometrically, starting from Euler's theorem that any rotation can be expressed as an axis of rotation and an angle of rotation about that axis. Although the map in Equation (2.2) is surjective, we can define an inverse map if we restrict its domain to $0 \leq \phi < \pi$:

$$\phi = \text{Log}(\mathbf{C}) = \log(\mathbf{C})^\vee = \frac{\phi(\mathbf{C} - \mathbf{C}^T)^\vee}{2 \sin \phi}, \quad (2.5)$$

where $\phi = \arccos \frac{\text{tr}(\mathbf{C})-1}{2}$ and the *vee* operator, $(\cdot)^\vee : \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}^3$, is defined as the unique inverse of the wedge operator $(\cdot)^\wedge$. Note Equation (2.5) is undefined at both $\phi = 0$ and at $\phi = \pi$. In the former case, we can use a small-angle approximation and define

$$\text{Log}(\mathbf{C}) \approx (\mathbf{C} - \mathbf{1})^\vee \text{ when } \phi \approx 0. \quad (2.6)$$

The latter case (when $\phi = \pi$) defines the *cut locus* of the space where $\text{Exp}(\cdot)$ is not a covering map and both $+\phi$ and $-\phi$ map to the same \mathbf{C} . This *cut locus* is related to the idea that any three parameterization of $\text{SO}(3)$ will have singularities associated with it.

2.2.1 Unit Quaternions

Another way (and historically, the original way) to represent a general rotation is to use a unit quaternion, \mathbf{q} . A unit quaternion has four parameters, a scalar value q_ω , and a three-dimensional vector component, \mathbf{q}_v :

²We follow Solà et al. (2018) and also define *capitalized* map for notational clarity.

³This operator is sometimes also expressed as $(\cdot)^\times$ or $[\cdot]_\times$ and is known as the skew-symmetric operator.

$$\mathbf{q} = \begin{bmatrix} q_\omega \\ \mathbf{q}_v \end{bmatrix} \in S^3, \quad (\|\mathbf{q}\| = 1). \quad (2.7)$$

Unit quaternions also form a Lie group (Solà et al., 2018) and lie on a three-dimensional unit sphere within \mathbb{R}^4 . This manifold represents a double cover of $\text{SO}(3)$ (since both \mathbf{q} and $-\mathbf{q}$ represent the same rotation). As with rotation matrices, we can define a surjective map from three parameters to the group itself,

$$\mathbf{q} = \text{Exp}(\phi) = \begin{bmatrix} \cos \phi/2 \\ \mathbf{a} \sin \phi/2 \end{bmatrix}. \quad (2.8)$$

Similarly, we can also define a logarithmic map as

$$\phi = \text{Log}(\mathbf{q}) = 2\mathbf{q}_v \frac{\arctan(\|\mathbf{q}_v\|, q_\omega)}{\|\mathbf{q}_v\|}. \quad (2.9)$$

To avoid issues with the double cover, we replace \mathbf{q} with $-\mathbf{q}$ if q_ω is negative before evaluating Equation (2.9). Also note again that Equation (2.9) is undefined when $\phi = 0$, but, importantly, we do not face any issues when $\phi = \pi$ due to the half angle. As with rotation matrices, we can use small angle approximations to define:

$$\text{Log}(\mathbf{q}) \approx \frac{\mathbf{q}_v}{q_\omega} \left(1 - \frac{\|\mathbf{q}_v\|^2}{3q_\omega^2} \right) \quad \text{when } \phi \approx 0. \quad (2.10)$$

A fantastic summary of the history of rotation parameterizations, unit quaternions and the story of Hamilton and Rodriguez can be found in Altmann (1989).

2.3 Spatial Transforms

The rigid body transform \mathbf{T} is also a member of the matrix Lie group, the Special Euclidian group $\text{SE}(3)$ and can be defined as a 4×4 matrix as follows:

$$\text{SE}(3) = \left\{ \mathbf{T} = \begin{bmatrix} \mathbf{C} & \mathbf{t} \\ \mathbf{0}^T & 1 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \mid \mathbf{C} \in \text{SO}(3), \mathbf{t} \in \mathbb{R}^3 \right\}. \quad (2.11)$$

As a member of a matrix Lie group, it also admits a surjective exponential map,

$$\mathbf{T} = \text{Exp}(\xi) = \exp(\xi^\wedge) = \sum_{n=0}^{\infty} \frac{1}{n!} (\xi^\wedge)^n \quad (2.12)$$

where $\xi = \begin{bmatrix} \rho \\ \phi \end{bmatrix} \in \mathbb{R}^6$ and the wedge operator is overloaded (following Barfoot (2017)) as follows:

$$\xi^\wedge \triangleq \begin{bmatrix} \rho \\ \phi \end{bmatrix}^\wedge = \begin{bmatrix} \phi^\wedge & \rho \\ \mathbf{0}^T & 0 \end{bmatrix}. \quad (2.13)$$

In practice, we can evaluate the exponential map through the Euler-Rodriguez formula (Equation (2.3)) and by computing the left-Jacobian of $\text{SO}(3)$, \mathbf{J} ,

$$\mathbf{T} = \text{Exp} \left(\begin{bmatrix} \rho \\ \phi \end{bmatrix} \right) = \begin{bmatrix} \mathbf{C}(\phi) & \mathbf{J}(\phi)\rho \\ \mathbf{0}^T & 1 \end{bmatrix}, \quad (2.14)$$

where

$$\mathbf{J}(\phi) = \frac{\sin \phi}{\phi} \mathbf{1} + \left(1 - \frac{\sin \phi}{\phi}\right) \mathbf{a}\mathbf{a}^T + \frac{1 - \cos \phi}{\phi} \mathbf{a}^\wedge. \quad (2.15)$$

2.3.1 Applying Transforms

Applying our notation for coordinate frames (and referring back to Section 2.1), a transform, \mathbf{T}_{vi} can be expressed as

$$\mathbf{T}_{vi} = \begin{bmatrix} \mathbf{C}_{vi} & \mathbf{t}_v^{iv} \\ \mathbf{0}^T & 1 \end{bmatrix}. \quad (2.16)$$

This allows us to use the homogenous point representation for \mathbf{r}_i^{pi} and express the following relation:

$$\begin{bmatrix} \mathbf{r}_v^{pi} \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} \mathbf{C}_{vi} & \mathbf{t}_v^{iv} \\ \mathbf{0}^T & 1 \end{bmatrix}}_{\mathbf{T}_{vi}} \begin{bmatrix} \mathbf{r}_i^{pi} \\ 1 \end{bmatrix} \quad (2.17)$$

which is numerically equivalent to

$$\mathbf{r}_v^{pi} = \mathbf{C}_{vi}\mathbf{r}_i^{pi} + \mathbf{t}_v^{iv}. \quad (2.18)$$

2.4 Perturbations

When solving optimization problems that involve rotations or rigid-body transforms, it is often useful to consider a small *perturbation* about an operating point. By leveraging a core property of Lie groups (that they are locally ‘Euclidian’), we can convert difficult non-linear problems into ones that have local linear approximations.

Using rotations as an example, we can perturb an operating point, $\bar{\mathbf{C}} \triangleq \text{Exp}(\bar{\boldsymbol{\phi}})$, in three different ways:

$$\mathbf{C} = \begin{cases} \text{Exp}(\delta\boldsymbol{\phi}^\ell) \bar{\mathbf{C}} & \text{left perturbation,} \\ \text{Exp}(\bar{\boldsymbol{\phi}} + \delta\boldsymbol{\phi}^m) & \text{middle perturbation,} \\ \bar{\mathbf{C}} \text{Exp}(\delta\boldsymbol{\phi}^r) & \text{right perturbation.} \end{cases} \quad (2.19)$$

We can relate all the left and middle perturbations through the left Jacobian of $\text{SO}(3)$ with the following useful identity,

$$\text{Exp}((\boldsymbol{\phi} + \delta\boldsymbol{\phi}^m)) \approx \text{Exp}(\mathbf{J}(\boldsymbol{\phi})\delta\boldsymbol{\phi}^m) \text{Exp}(\boldsymbol{\phi}). \quad (2.20)$$

From this it follows that $\delta\boldsymbol{\phi}^\ell \approx \mathbf{J}(\boldsymbol{\phi})\delta\boldsymbol{\phi}^m$ and elucidates why \mathbf{J} is called the *left* Jacobian.

In this thesis, we will use the left and middle perturbations when appropriate. Using small angle approximations, the Euler-Rodriguez formula (Equation (2.3)) yields $\text{Exp}(\delta\boldsymbol{\phi}) \approx \mathbf{1} + \delta\boldsymbol{\phi}^\wedge$, which allows us to write the useful formula for the left perturbation:

$$\mathbf{C} = (\mathbf{1} + (\delta\boldsymbol{\phi}^\ell)^\wedge) \bar{\mathbf{C}}. \quad (2.21)$$

Similarly, we can write analogous expressions for a rigid body transform, $\mathbf{T} \in \text{SE}(3)$, as composed of an operating point $\bar{\mathbf{T}} \triangleq \text{Exp}(\bar{\boldsymbol{\xi}})$, and a small perturbation about that operating point:

$$\mathbf{T} = \begin{cases} \text{Exp}(\delta\boldsymbol{\xi}^\ell) \bar{\mathbf{T}} & \text{left perturbation,} \\ \text{Exp}(\bar{\boldsymbol{\xi}} + \delta\boldsymbol{\xi}^m) & \text{middle perturbation,} \\ \bar{\mathbf{T}} \text{Exp}(\delta\boldsymbol{\xi}^r) & \text{right perturbation.} \end{cases} \quad (2.22)$$

Now, we can also note a similar identity for $\text{SE}(3)$,

$$\text{Exp}((\boldsymbol{\xi} + \delta\boldsymbol{\xi}^m)) \approx \text{Exp}((\mathcal{J}(\boldsymbol{\xi})\delta\boldsymbol{\xi}^m)) \text{Exp}(\boldsymbol{\xi}), \quad (2.23)$$

where \mathcal{J} is the left Jacobian of $\text{SE}(3)$ and defined as

$$\mathcal{J}(\boldsymbol{\xi}) \triangleq \begin{bmatrix} \mathbf{J}(\boldsymbol{\phi}) & \mathbf{Q}(\boldsymbol{\xi}) \\ \mathbf{0} & \mathbf{J}(\boldsymbol{\phi}) \end{bmatrix}, \quad (2.24)$$

where $\mathbf{Q}(\boldsymbol{\xi})$ can be evaluated analytically (see Barfoot (2017)). This again allows us to write $\delta\boldsymbol{\xi}^\ell \approx \mathcal{J}(\boldsymbol{\xi})\delta\boldsymbol{\xi}^m$ and form a similar expression,

$$\mathbf{T} = (\mathbf{1} + (\delta\boldsymbol{\xi}^\ell)^\wedge) \bar{\mathbf{T}}. \quad (2.25)$$

To derive locally linear systems from sets of rigid-body transforms, or ‘poses’, we can apply Equation (2.25). To update an operating point, we solve for $\delta\boldsymbol{\xi}^\ell$ and then use the constraint-sensitive update $\mathbf{T} \leftarrow \text{Exp}(\delta\boldsymbol{\xi}^\ell) \bar{\mathbf{T}}$.

2.5 Uncertainty

We can also use perturbation theory to implicitly define uncertainty on constrained manifolds (see [Barfoot and Furgale \(2014\)](#) for a thorough discussion).

Given a concentrated normal density, $\delta\boldsymbol{\xi} \sim \mathcal{N}(\mathbf{0}, \boldsymbol{\Sigma}_{6 \times 6})$, we can *inject* this unconstrained density onto the Lie group through left perturbations about some mean:

$$\mathbf{T} = \text{Exp}(\delta\boldsymbol{\xi}) \bar{\mathbf{T}} \quad (2.26)$$

This allows us to keep track of a random variable, \mathbf{T} , by keeping its mean in group form, $\bar{\mathbf{T}}$, while its second statistical moment is stored as a standard 6×6 covariance matrix, $\boldsymbol{\Sigma}$.

We can define an analogous density for rotation matrices given normal densities over rotation perturbations $\delta\boldsymbol{\phi} \sim \mathcal{N}(\mathbf{0}, \boldsymbol{\Sigma}_{3 \times 3})$,

$$\mathbf{C} = \text{Exp}(\delta\boldsymbol{\phi}) \bar{\mathbf{C}}, \quad (2.27)$$

and also, for unit quaternions,

$$\mathbf{q} = \text{Exp}(\delta\boldsymbol{\phi}) \otimes \bar{\mathbf{q}} \quad (2.28)$$

where \otimes refers to the standard quaternion product operator [Sola \(2017\)](#).

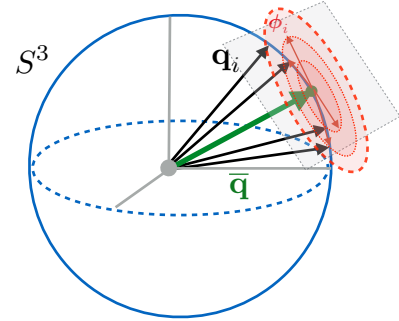


Figure 2.3: We can define uncertainty in the left tangent space of a mean element of a Lie group (here illustrated for unit quaternions).

Appendices

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