



EE6302E Dynamics of Electrical Machines (DEM)

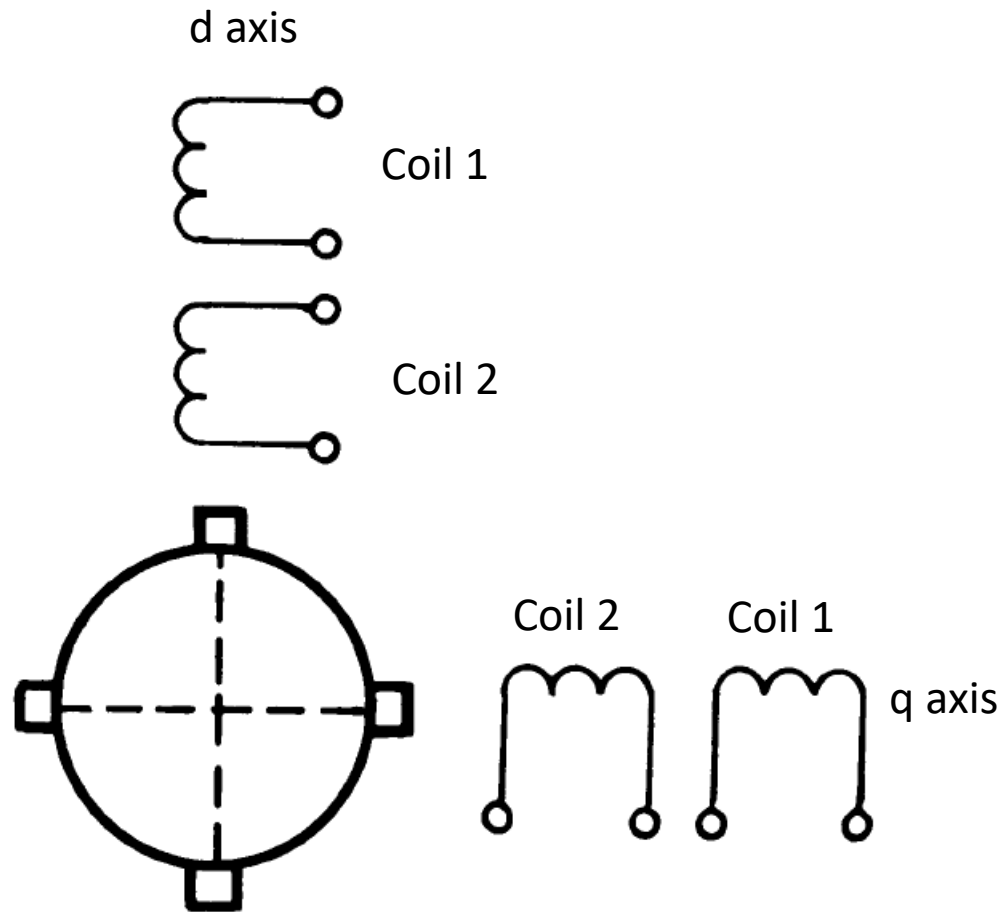
Module 2

References

1. D P Sengupta & J.B. Lynn, *Electrical Machine Dynamics*, The Macmillan Press Ltd., 1980.
2. R Krishnan, *Electric Motor Drives, Modeling, Analysis and Control*, Pearson Education, 2001.
3. P.C. Kraus, *Analysis of Electrical Machines*, McGraw Hill Book Company, 1987

PRIMITIVE MACHINE EQUATIONS ALONG STATIONARY AXES

Rotor Voltages



$$V_{dr} = R_{dr}i_{dr} + \frac{d\psi_{dr}}{dt} + B_{qr}p\theta$$

$$V_{qr} = R_{qr}i_{qr} + \frac{d\psi_{qr}}{dt} - B_{dr}p\theta$$

flux density terms

$$B_{qr} = L'_{qr}i_{qr} + M'_{q1}i_{qs1} + M'_{q2}i_{qs2}$$

$$B_{dr} = L'_{dr}i_{dr} + M'_{d1}i_{ds1} + M'_{d2}i_{ds2}$$

$$V_{dr} = R_{dr}i_{dr} + L_{dr}pi_{dr} + M_{d1}pi_{ds1} + M_{d2}pi_{ds2} \\ + L'_{qr}i_{qr}p\theta + M'_{q1}i_{qs1}p\theta + M'_{q2}i_{qs2}p\theta$$

$$V_{qr} = R_{qr}i_{qr} + L_{qr}pi_{q2} + M_{q1}pi_{qs1} + M_{q2}pi_{qs2} \\ - L'_{dr}i_{dr}p\theta - M'_{d1}i_{ds1}p\theta - M'_{d2}i_{ds2}p\theta$$

PRIMITIVE MACHINE EQUATIONS ALONG STATIONARY AXES

$$[\mathbf{V}] = [\mathbf{Z}][\mathbf{i}]$$

$$[\mathbf{V}] = [\mathbf{R}][\mathbf{i}] + [\mathbf{L}]p[\mathbf{i}] + [\mathbf{G}][\mathbf{i}]p\theta$$

$$\mathbf{V} = \mathbf{R}\mathbf{i} + \mathbf{L}p\mathbf{i} + \mathbf{G}ip\theta$$

$$\mathbf{R} =$$

	ds1	ds2	dr	qr	qs1	qs2
ds1	R_{ds1}					
ds2		R_{ds2}				
dr			R_r			
qr				R_r		
qs1					R_{qs1}	
qs2						R_{qs2}

$$\mathbf{L} =$$

	ds1	ds2	dr	qr	qs1	qs2
ds1	L_{ds1}	M_{d12}	M_{d1}			
ds2	M_{d12}	L_{ds2}	M_{d2}			
dr	M_{d1}	M_{d2}	L_{dr}			
qr				L_{qr}	M_{q1}	M_{q2}
qs1				M_{q1}	L_{qs1}	M_{q12}
qs2				M_{q2}	M_{q12}	L_{qs2}

$$\mathbf{G} =$$

	ds1	ds2	dr	qr	qs1	qs2
ds1						
ds2						
dr				L'_{qr}	M'_{q1}	M'_{q2}
qr	$-M'_{d1}$	$-M'_{d2}$	$-L'_{dr}$			
qs1						
qs2						

		ds1	ds2	dr	qr	qs1	qs2	
V_{ds1}	ds1	$R_{ds1} + L_{ds1}p$	$M_{d12}p$	$M_{d1}p$				i_{ds1}
V_{ds2}	ds2	$M_{d12}p$	$R_{ds2} + L_{ds2}p$	$M_{d2}p$				i_{ds2}
V_{dr}	dr	$M_{d1}p$	$M_{d2}p$	$R_r + L_{dr}p$	$L'_{qr}p\theta$	$M'_{q1}p\theta$	$M'_{q2}p\theta$	i_{dr}
V_{qr}	qr	$-M'_{d1}p\theta$	$-M'_{d2}p\theta$	$-L'_{dr}p\theta$	$R_r + L_{qr}p$	$M_{q1}p$	$M_{q2}p$	i_{qr}
V_{qs1}	qs1				$M_{q1}p$	$R_{qs} + L_{qs1}p$	$M_{q12}p$	i_{qs1}
V_{qs2}	qs2				$M_{q2}p$	$M_{q12}p$	$R_{qs2} + L_{qs2}p$	i_{qs2}

SQUIRREL CAGE INDUCTION MOTOR MODEL FROM PRIMITIVE MACHINE EQUATIONS

Step 1

Remove coil 2 and corresponding rows and column

Replace the operator p by $j\omega_1$

$$p\theta = \omega_r$$

$$\omega_r = \omega_1 (1 - s)$$

		ds1	dr	qr	qs1	
V_{ds1}	ds1	$R_{ds1} + j\omega_1 L_{ds1}$	$j\omega_1 M_{d1}$			i_{ds1}
V_{dr}	dr	$j\omega_1 M_{d1}$	$R_r + j\omega_1 L_{dr}$	$\omega_1 (1 - s) M_{q1}$	$\omega_1 (1 - s) M_{q1}$	i_{dr}
V_{qr}	qr	$-\omega_1 (1 - s) M_{d1}$	$-\omega_1 (1 - s) L_{dr}$	$R_r + j\omega_1 L_{qr}$	$j\omega_1 M_{q1}$	i_{qr}
V_{qs1}	qs1			$j\omega_1 M_{q1}$	$R_{qs1} + j\omega_1 L_{qs1}$	i_{qs1}

SQUIRREL CAGE INDUCTION MOTOR MODEL FROM PRIMITIVE MACHINE EQUATIONS

Step 2

$$\begin{aligned} R_{ds1} &= R_{qs1} = R_1, & R_r &= R_2 \\ L_{ds1} &= L_{qs1} = L_1, & L_{dr} &= L_{qr} = L_2 \\ M_{d1} &= M_{q1} = M \\ M'_{d1} &= M_{d1} = M \\ M'_{q1} &= M_{q1} = M \end{aligned}$$

stator coils which are symmetrically distributed

air gap is uniform

flux wave is sinusoidally distributed in space and hence the coefficients of mutual inductance for transformer and generated voltages are the same

		ds1	dr	qr	qs1	
V_{ds1}	ds1	$R_1 + jX_1$	jX_m			i_{ds1}
V_{dr}	dr	jX_m	$R_2 + jX_2$	$(1-s)X_2$	$(1-s)X_m$	i_{dr}
V_{qr}	qr	$-(1-s)X_m$	$-(1-s)X_2$	$R_2 + jX_2$	jX_m	i_{qr}
V_{qs1}	qs1			jX_m	$R_1 + jX_1$	i_{qs1}

SQUIRREL CAGE INDUCTION MOTOR MODEL FROM PRIMITIVE MACHINE EQUATIONS

Step 3

		ds1	dr	qr	qs1	
V_1	ds1	$R_1 + jX_1$	jX_m			I_1
0	dr	jX_m	$R_2 + jX_2$	$(1-s)X_2$	$(1-s)X_m$	I_2
0	qr	$-(1-s)X_m$	$-(1-s)X_2$	$R_2 + jX_2$	jX_m	$-jI_2$
$-jV_1$	qs1			jX_m	$R_1 + jX_1$	$-jI_1$

net m.m.f. they produce rotates at synchronous speed

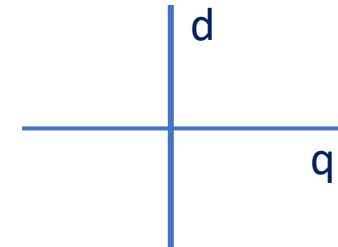
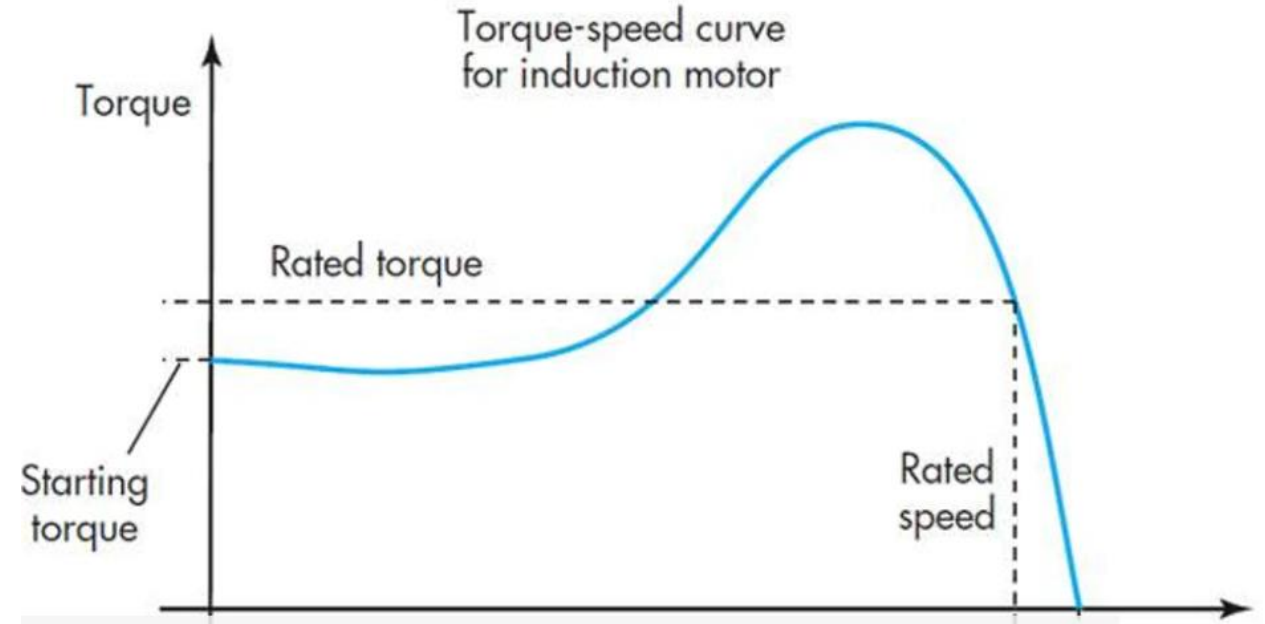
During balanced operation, net MMF $N_1(I_1^2 \sin^2 \omega_1 t + I_1^2 \cos^2 \omega_1 t)^{\frac{1}{2}}$

rotor voltages are zero since the rotor coils in an induction motor are short-circuited

SQUIRREL CAGE INDUCTION MOTOR MODEL FROM PRIMITIVE MACHINE EQUATIONS

$$T = -\mathbf{i}^* \mathbf{G} \mathbf{i}$$

$$= -i_{qr}^* M_d i_{ds} + i_{dr}^* M_q i_{qs}$$



Squirrel Cage Induction Motor: Stability Analysis

Assignment Q1

Marks: 10

Last date: before 12.12.2023

Develop the state space model for the induction machine with following parameters. Estimate eigen values and eigen vector using any simulation platform. Infer on the stability

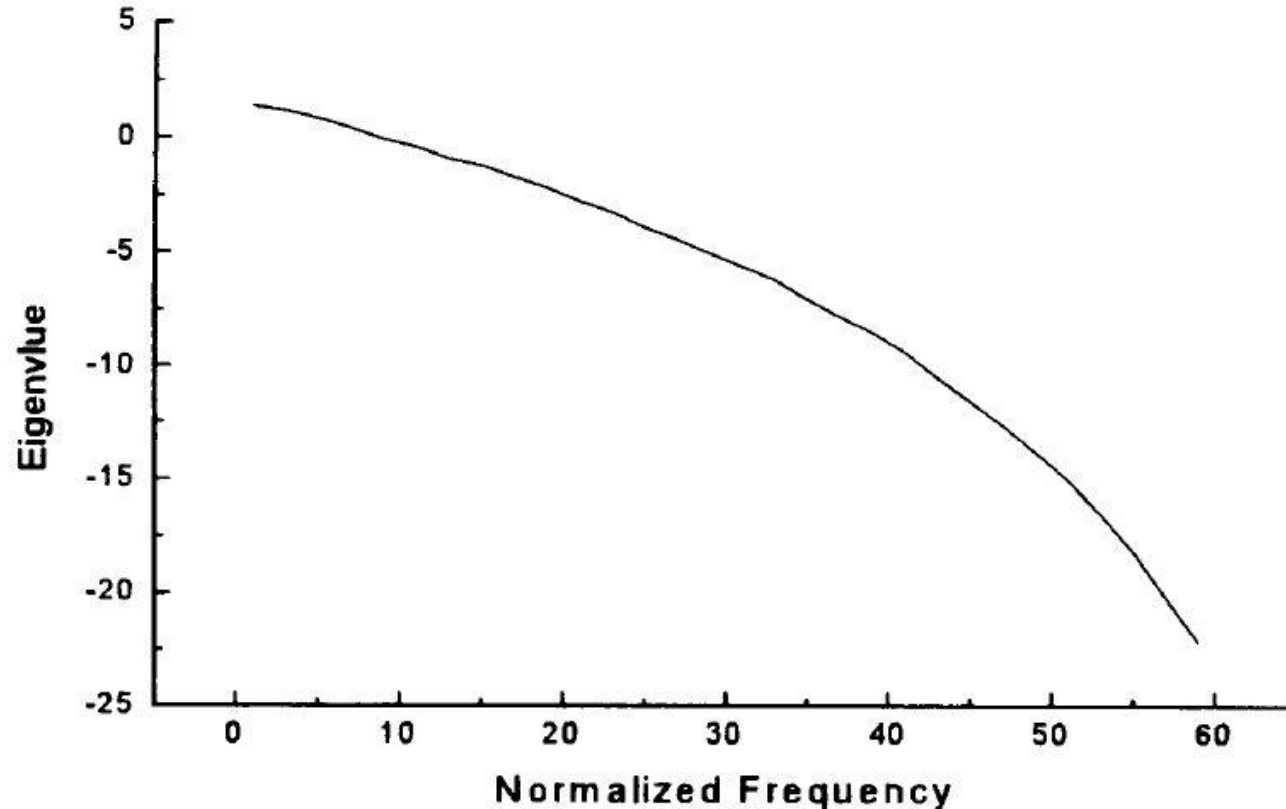
Machines Parameters	Value	Per Unit Value
Horse Power (Hp)	50 hp	-
Voltage (V_L)	460 V	-
Frequency (Hz)	60 Hz	-
Stator Resistance (r_s)	0.087 Ω	0.015336
Stator Reactance (X_{ls})	0.302 Ω	0.053235
Mutual Reactance (X_M)	13.08 Ω	2.30569
Equivalent Rotor Resistance (r'_r)	0.302 Ω	0.040191
Equivalent Rotor Reactance (X'_{lr})	0.228 Ω	0.053235
Moment of Inertia (J)	1.662 Ω	-

Ref. Section 5.10 of "Electrical Machine Dynamics" by Sengupta for State space model example.

Various currents and rotor speed are to be taken as State variables.

M. B. Uddin, M. N. Pramanik and S. A. Reza, "Low frequency stability study of a three-phase induction motor," 2007 7th International Conference on Power Electronics, 2007, pp. 1115-1120

Squirrel Cage Induction Motor: Stability Analysis

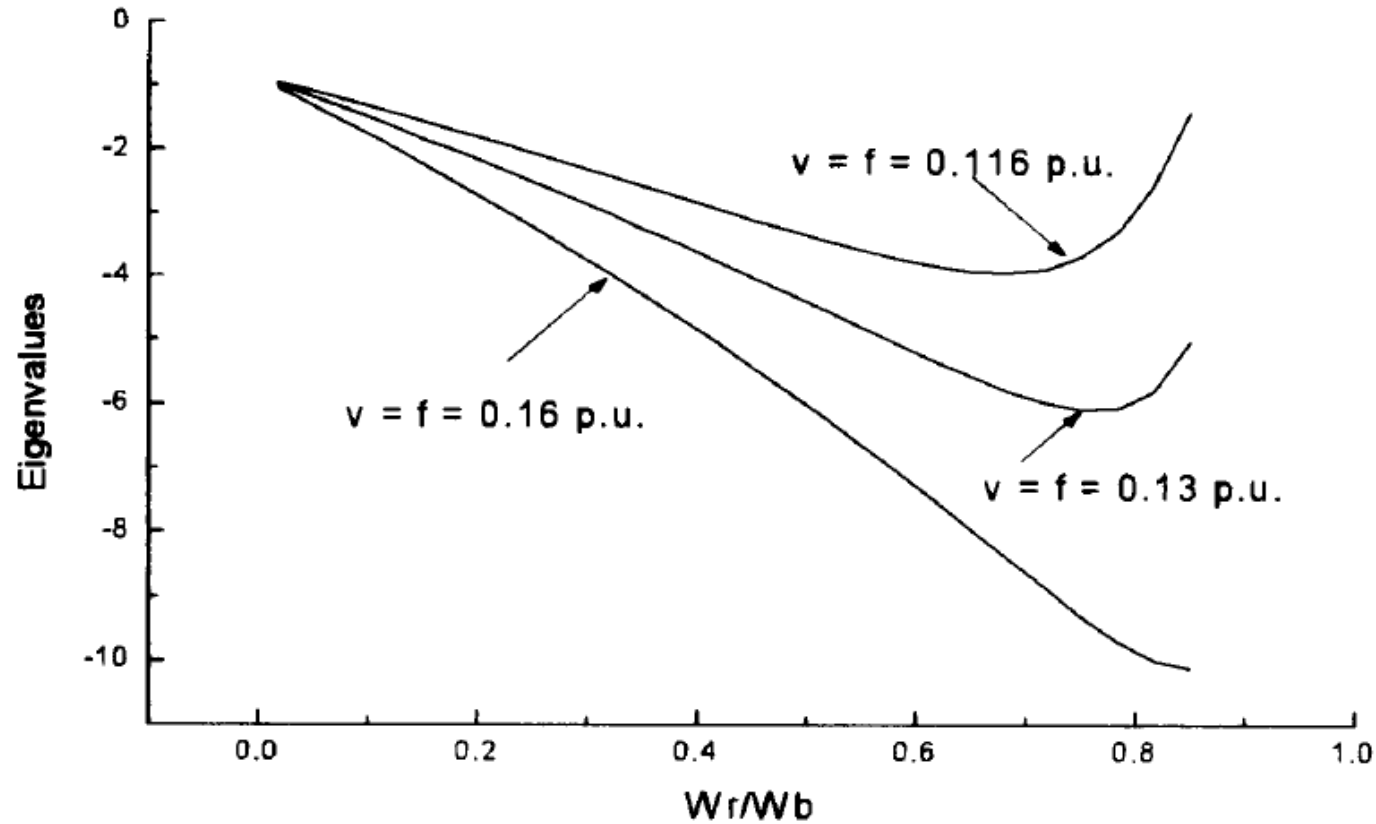


The excursion of eigen values for the rated operating condition of induction motor are depicted in Fig. Which indicate that operation becomes unstable at lower frequency. In each step of computation the applied voltage is decreased linearly with frequency. The eigen values are found to cross the boundary at low frequency of 0.116 p. u. (7 Hz) while the corresponding stator voltage is also 0.116 p. u. indicating unstable operation below this frequency.

The straight line parallel to the x-axis passing through zero value of the ordinate forms the boundary between the stable and unstable region of operation. The lower part of the boundary represents the stable region and the upper part signifies the unstable region of operation.

M. B. Uddin, M. N. Pramanik and S. A. Reza, "Low frequency stability study of a three-phase induction motor," *2007 7th International Conference on Power Electronics*, 2007, pp. 1115-1120

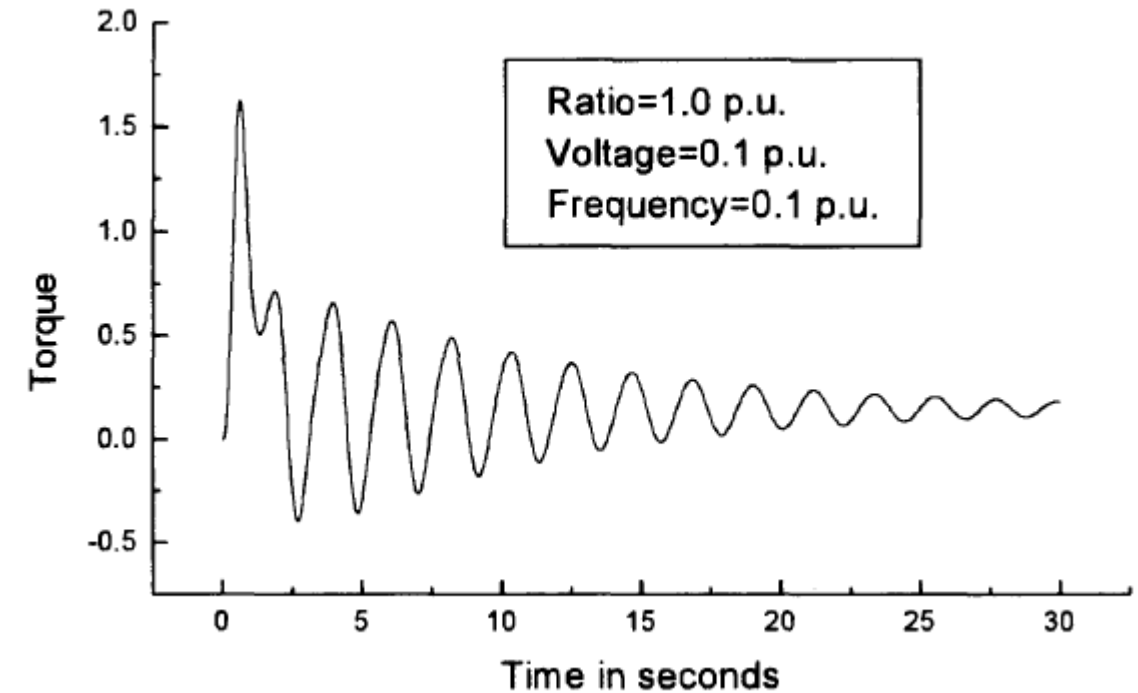
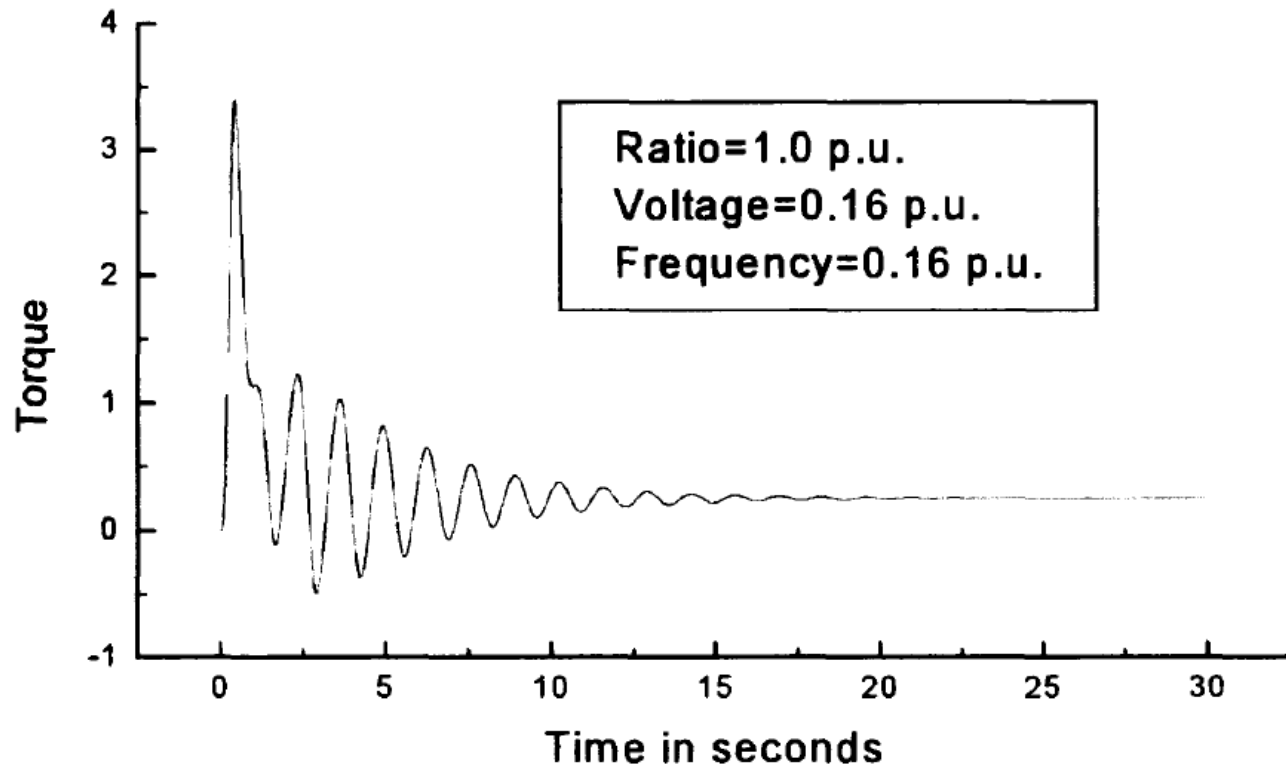
Squirrel Cage Induction Motor: Stability Analysis



The effect of amplitude of the stator voltage on stability is illustrated in Fig. These curves are obtained by decreasing the stator voltage and frequency keeping volt/Hz ratio constant. The reduction of voltage and frequency simultaneously is indicating the induction motor trends to be unstable at lower frequency

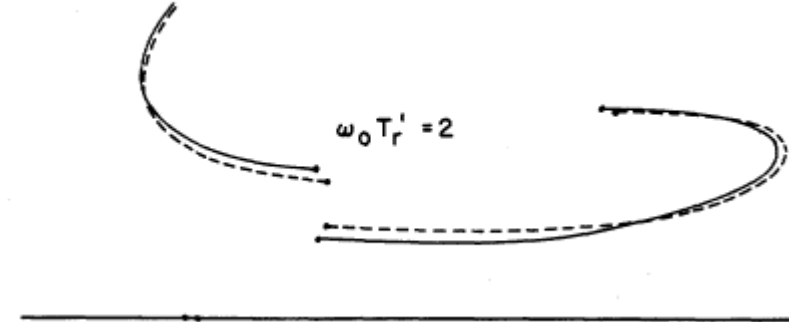
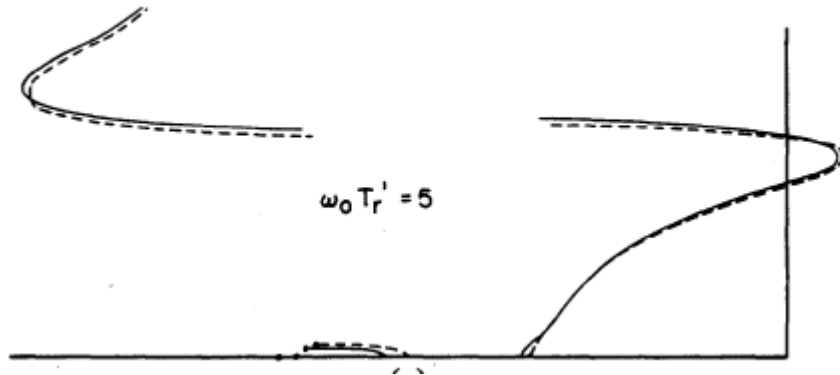
M. B. Uddin, M. N. Pramanik and S. A. Reza, "Low frequency stability study of a three-phase induction motor," *2007 7th International Conference on Power Electronics*, 2007, pp. 1115-1120

Squirrel Cage Induction Motor: Stability Analysis

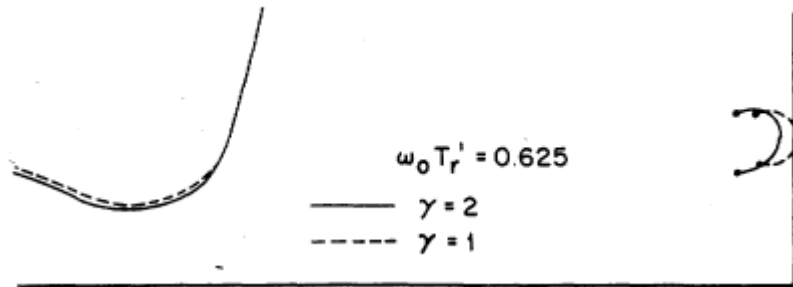


M. B. Uddin, M. N. Pramanik and S. A. Reza, "Low frequency stability study of a three-phase induction motor," *2007 7th International Conference on Power Electronics*, 2007, pp. 1115-1120

Squirrel Cage Induction Motor: Stability Analysis



The leftward shift of the overall root locus in the region of low damping is perhaps the most significant effect of transient saturation.

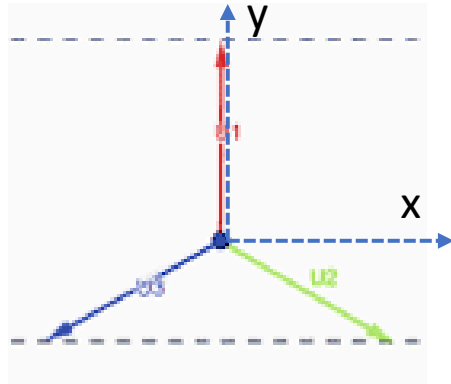
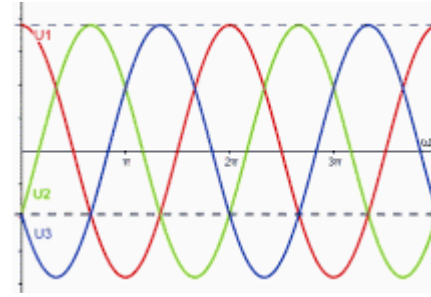


Modern high efficiency machines which have increased pu magnetizing reactance and are designed to operate at reduced flux levels are likely to have larger and stronger regions of instability and to exhibit poorer damped transient response than conventional machines.

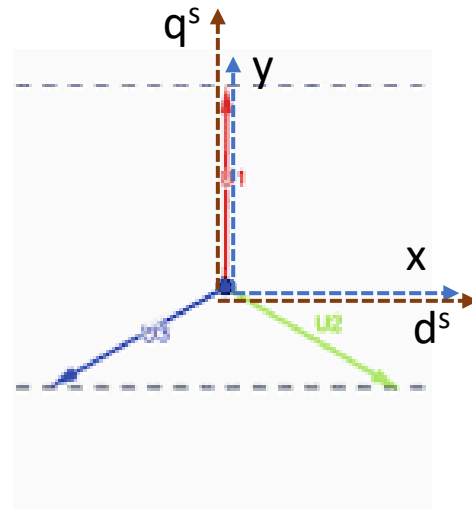


Reference Frame

A three phase voltage (or current or any balanced sinusoidal oscillation) can be represented as rotating vectors in multiple ways (different perspectives)



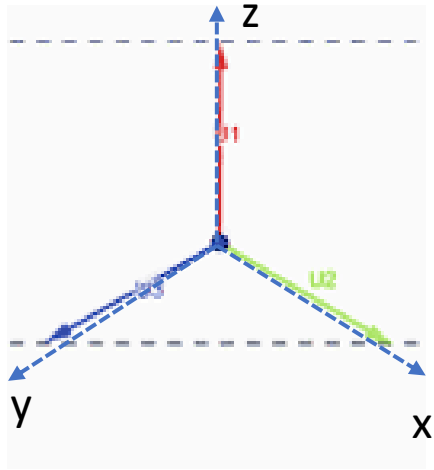
Two stationary axes (x and y) and three rotating vectors



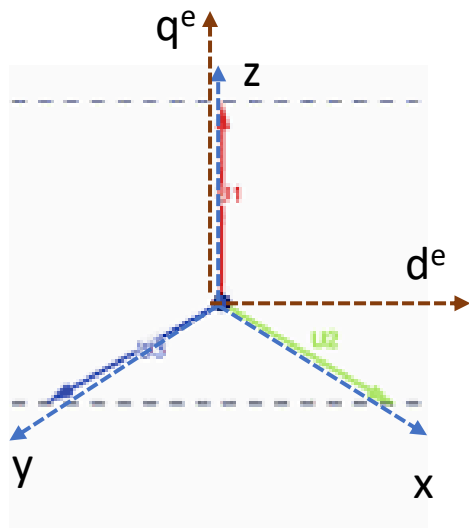
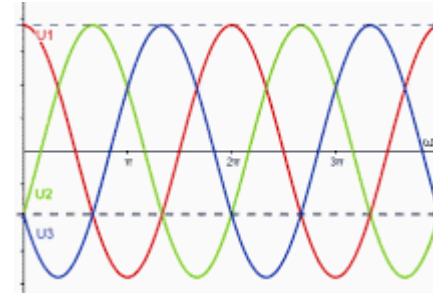
- The three rotating vectors can be mapped to two stationary axes (d^s and q^s). Since axes are stationary, the mapped quantities will be sinusoidal.
- That is, three-phase voltage (V_R, V_Y, V_B) can be represented in their equivalent two phase components (V_d^s, V_q^s phase shifted by 90 deg). This mapping is called Clarke's transformation.
- By changing the magnitudes of direct and quadrature axis components (the mapped components), we could realize three-phase components of any magnitude and phase angle at any particular instant (i.e. by increasing quadrature component magnitude alone in this case, we could generate phase lead to three phase set in comparison with original set. Also change in magnitude).

- In developing control for three phase system, this approach facilitates easier implementation than using three-phase quantities directly (since number of variables reduces)

Reference Frame



Three rotating axes
and three stationary
vectors in those axes



- The **three stationary vectors (constant)** in **three rotating axes** can be mapped to **two rotating axes** (d^e and q^e). Since axes are rotating, the **mapped quantities also will be constant (dc)** (assuming d^e - q^e axes rotating in synchronism with that of three-phase).
- That is, three-phase voltage (V_R , V_Y , V_B) can be represented in their equivalent two phase components (V_d^e , V_q^e , phase shifted by 90 deg). This mapping is called Park's transformation.
- By changing the magnitudes of direct and quadrature axis components (the mapped components), we could realize three-phase components of any magnitude and phase angle at any particular instant (i.e. by increasing quadrature component magnitude alone in this case, we could generate phase lead to three phase set in comparison with original set. Also change in magnitude).

- In developing control for three phase system, this approach facilitates easier implementation than Clarke's transformation (since number of variables reduces and becomes DC (lower bandwidth requirement))

Fortescue's Transformation

- This transformation is known as the method of symmetrical components and developed by Fortescue.
- This transformation states that N unbalanced phasors can be represented by N systems of N balanced phasors.
- It uses a complex transformation to decouple the abc phase variables.
- The method of symmetrical components is used to simplify analysis of unbalanced three phase power systems under both normal and abnormal conditions.
- It is used to decouple an unbalanced three-phase network into three simpler sequence (zero, positive and negative) networks.

Fortescue's Transformation

- The method of symmetrical components is expressed as follows

$$[\mathbf{f}_{012}] = [\mathbf{T}_{012}][\mathbf{f}_{abc}]$$

$$[\mathbf{f}_{abc}] = [\mathbf{T}_{012}]^{-1} [\mathbf{f}_{012}]$$

$$[\mathbf{f}_{012}] = \begin{bmatrix} f_0 \\ f_1 \\ f_2 \end{bmatrix}$$

$$[\mathbf{f}_{abc}] = \begin{bmatrix} f_a \\ f_b \\ f_c \end{bmatrix}$$

- Variable f may be the currents, voltages or fluxes and the transformation and its inverse are given by

$$[\mathbf{T}_{012}] = \frac{1}{3} \begin{bmatrix} 1 & 1 & 1 \\ 1 & a & a^2 \\ 1 & a^2 & a \end{bmatrix}$$

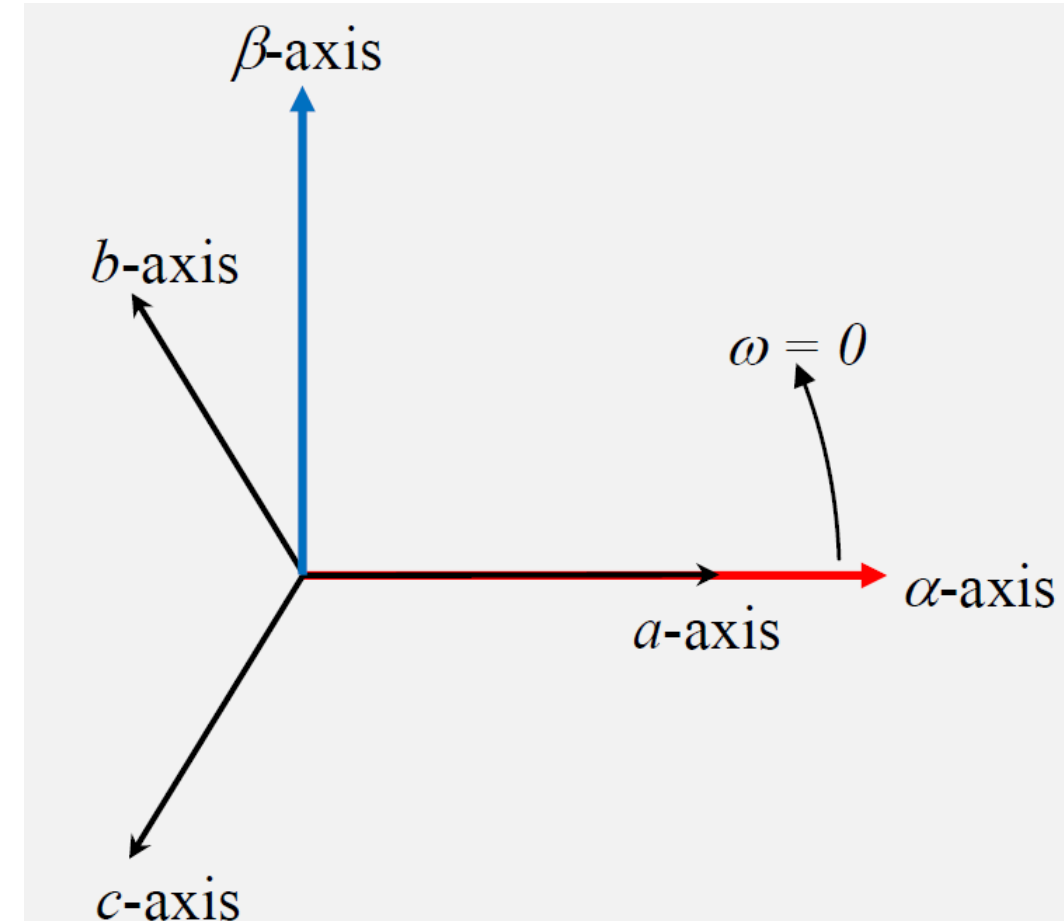
$$[\mathbf{T}_{012}]^{-1} = \begin{bmatrix} 1 & 1 & 1 \\ 1 & a^2 & a \\ 1 & a & a^2 \end{bmatrix}$$

where

$$a = e^{j\frac{2\pi}{3}}$$

Clarke's Transformation

- The stationary two-phase variables of Clarke's transformation are denoted as α and β .
- As shown below, the α -axis coincides with the phase a-axis and the α -axis leads the β -axis by $\pi/2$.
- A third variable known as the zero-sequence component is also included.
- Clarke's transformation is not power-invariant (i.e. the values of power before and after the transformation are not the same.)



Clarke's Transformation

Clarke's transformation is expressed as follows

$$\begin{bmatrix} \mathbf{f}_{\alpha\beta 0} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{\alpha\beta 0} \end{bmatrix} \begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} = \begin{bmatrix} \mathbf{T}_{\alpha\beta 0} \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{f}_{\alpha\beta 0} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{f}_{\alpha\beta 0} \end{bmatrix} = \begin{bmatrix} f_{\alpha} \\ f_{\beta} \\ f_0 \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{f}_{abc} \end{bmatrix} = \begin{bmatrix} f_a \\ f_b \\ f_c \end{bmatrix}$$

Similarly variable f may be the currents, voltages or fluxes and the transformation and its inverse are given by

$$\begin{bmatrix} \mathbf{T}_{\alpha\beta 0} \end{bmatrix} = \frac{2}{3} \begin{bmatrix} 1 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & \frac{\sqrt{3}}{2} & -\frac{\sqrt{3}}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{T}_{\alpha\beta 0} \end{bmatrix}^{-1} = \begin{bmatrix} 1 & 0 & 1 \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} & 1 \\ -\frac{1}{2} & -\frac{\sqrt{3}}{2} & 1 \end{bmatrix}$$

Clarke's Transformation

Power-Invariant Property

$$\begin{cases} v_a = V_m \cos(\omega t) \\ v_b = V_m \cos(\omega t - 2\pi / 3) \\ v_c = V_m \cos(\omega t - 4\pi / 3) \end{cases}$$

$$\begin{cases} i_a = I_m \cos(\omega t) \\ i_b = I_m \cos(\omega t - 2\pi / 3) \\ i_c = I_m \cos(\omega t - 4\pi / 3) \end{cases}$$

at $\omega t = 0$

$$\begin{cases} v_a = V_m \\ v_b = \frac{-1}{2} V_m \\ v_c = \frac{-1}{2} V_m \end{cases}$$

$$\begin{cases} i_a = I_m \\ i_b = \frac{-1}{2} I_m \\ i_c = \frac{-1}{2} I_m \end{cases}$$

Three Phase Power

$$P = v_a i_a + v_b i_b + v_c i_c = \frac{3}{2} V_m I_m$$

Clarke's Transformed quantities

$$\begin{cases} v_\alpha = V_m \\ v_\beta = 0 \\ v_0 = 0 \end{cases}$$

$$\begin{cases} i_\alpha = I_m \\ i_\beta = 0 \\ i_0 = 0 \end{cases}$$

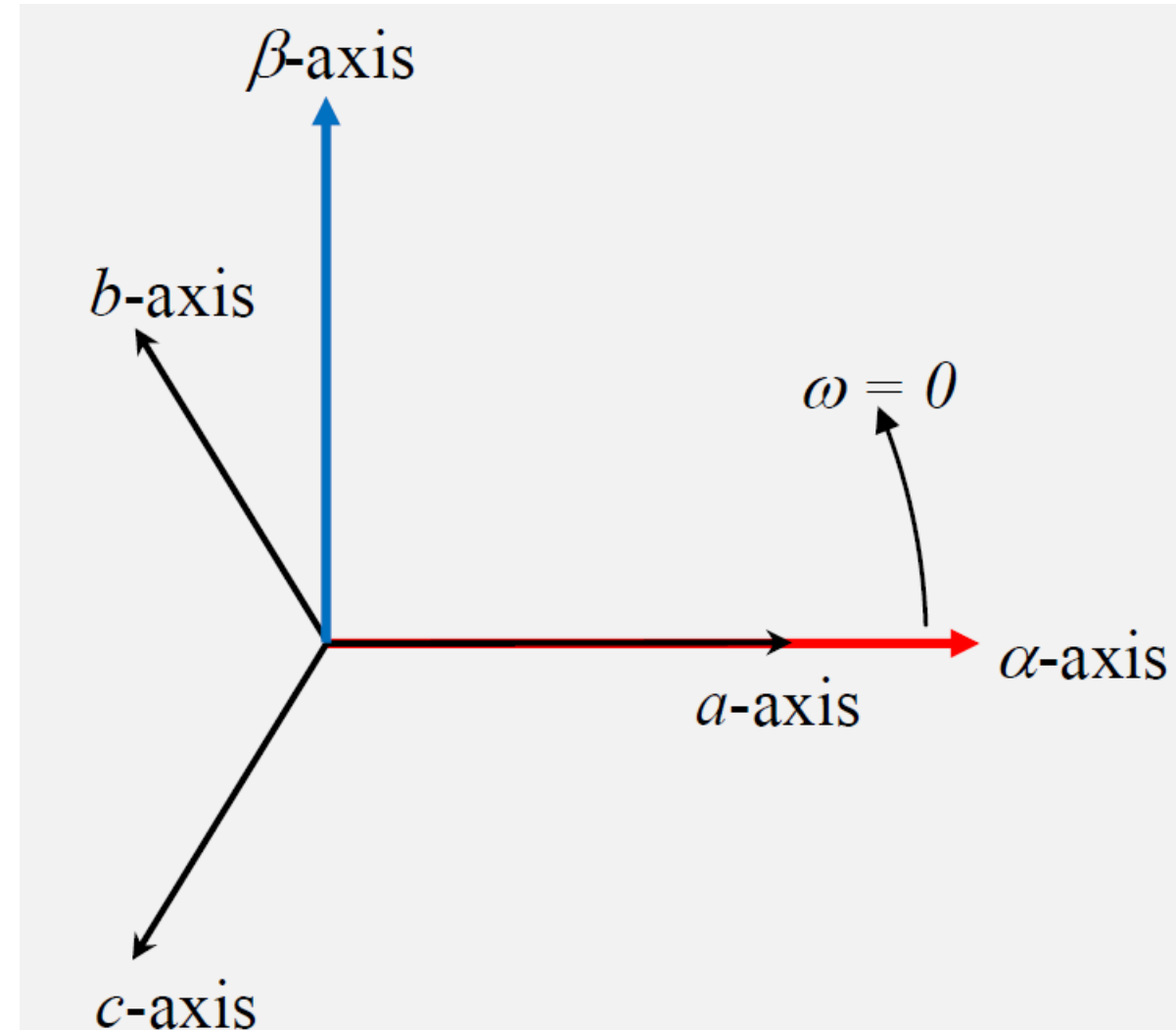
Clarke's transformation is not power-invariant

Power in $\alpha\beta$ frame

$$P = v_\alpha i_\alpha + v_\beta i_\beta + v_0 i_0 = V_m I_m$$

Concordia's Transformation

- Concordia's transformation is similar to Clarke's transformation.
- The only difference is that Concordia's transformation is power-invariant (i.e. the values of power before and after the transformation are identical).
- To have the power-invariant property, the transformation matrix must be orthogonal.
- A matrix is orthogonal if its inverse and its transpose are the same



Concordia's Transformation

$$[\mathbf{f}_{\alpha\beta 0}] = [\mathbf{T}_{\alpha\beta 0}] [\mathbf{f}_{abc}]$$

$$[\mathbf{f}_{abc}] = [\mathbf{T}_{\alpha\beta 0}]^{-1} [\mathbf{f}_{\alpha\beta 0}]$$

$$[\mathbf{f}_{\alpha\beta 0}] = \begin{bmatrix} f_{\alpha} \\ f_{\beta} \\ f_0 \end{bmatrix}$$

$$[\mathbf{f}_{abc}] = \begin{bmatrix} f_a \\ f_b \\ f_c \end{bmatrix}$$

$$[\mathbf{T}_{\alpha\beta 0}] = \sqrt{\frac{2}{3}} \begin{bmatrix} 1 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & \frac{\sqrt{3}}{2} & -\frac{\sqrt{3}}{2} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

$$[\mathbf{T}_{\alpha\beta 0}]^{-1} = \sqrt{\frac{2}{3}} \begin{bmatrix} 1 & 0 & \frac{1}{\sqrt{2}} \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} & \frac{1}{\sqrt{2}} \\ -\frac{1}{2} & -\frac{\sqrt{3}}{2} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

Concordia's Transformation

Power-Invariant Property

$$\begin{cases} v_a = V_m \cos(\omega t) \\ v_b = V_m \cos(\omega t - 2\pi / 3) \\ v_c = V_m \cos(\omega t - 4\pi / 3) \end{cases}$$

$$\begin{cases} i_a = I_m \cos(\omega t) \\ i_b = I_m \cos(\omega t - 2\pi / 3) \\ i_c = I_m \cos(\omega t - 4\pi / 3) \end{cases}$$

at $\omega t = 0$

$$\begin{cases} v_a = V_m \\ v_b = \frac{-1}{2} V_m \\ v_c = \frac{-1}{2} V_m \end{cases}$$

$$\begin{cases} i_a = I_m \\ i_b = \frac{-1}{2} I_m \\ i_c = \frac{-1}{2} I_m \end{cases}$$

Three Phase Power

$$P = v_a i_a + v_b i_b + v_c i_c = \frac{3}{2} V_m I_m$$

Concordia's Transformed quantities

$$\begin{cases} v_\alpha = \sqrt{\frac{3}{2}} V_m \\ v_\beta = 0 \\ v_0 = 0 \end{cases}$$

$$\begin{cases} i_\alpha = \sqrt{\frac{3}{2}} I_m \\ i_\beta = 0 \\ i_0 = 0 \end{cases}$$

Concordia's transformation is power-invariant

Power in $\alpha\beta$ frame

$$P = v_\alpha i_\alpha + v_\beta i_\beta + v_0 i_0 = \frac{3}{2} V_m I_m$$

Induction Motor Model in α - β frame

Voltage Equations in the stationary reference frame

$$v_{s\alpha} = R_s i_{s\alpha} + s\varphi_{s\alpha}$$

$$v_{s\beta} = R_s i_{s\beta} + s\varphi_{s\beta}$$

$$v_{r\alpha} = 0 = R_r i_{r\alpha} + s\varphi_{r\alpha} + \omega_r \varphi_{r\beta}$$

$$v_{r\beta} = 0 = R_r i_{r\beta} + s\varphi_{r\beta} - \omega_r \varphi_{r\alpha}$$

where 's' indicates the differential operator (d/dt)
[NOTE: its instead of usual p operator. not $j\omega$].

stator and rotor fluxes equations

$$\varphi_{s\alpha} = L_s i_{s\alpha} + L_m i_{r\alpha}$$

$$\varphi_{s\beta} = L_s i_{s\beta} + L_m i_{r\beta}$$

$$\varphi_{r\alpha} = L_r i_{r\alpha} + L_m i_{s\alpha}$$

$$\varphi_{r\beta} = L_r i_{r\beta} + L_m i_{s\beta}$$

In these equations, R_s , R_r , L_s , and L_r are, respectively, the resistors and the inductances of the stator windings and the rotor windings, L_m is the mutual inductance and $\omega_r = p \cdot \Omega_r$ is the rotor speed (with p is the pairs poles number). Additionally, ω_s is the synchronous frequency.

Induction Motor Model in α - β frame

Electromagnetic torque $\Gamma_{em} = \frac{3}{2}p (\varphi_{s\alpha} i_{s\beta} - \varphi_{s\beta} i_{s\alpha})$

For the complete model of the induction machine, the flux expressions are replaced in the voltage equations. We obtain a mechanical equation and four electrical equations in terms of the stator currents, rotor fluxes components, and the electric speed of induction machine as well

$$\frac{di_{s\alpha}}{dt} = -\frac{1}{\sigma L_s} \left(R_s + \frac{1}{T_r} \frac{L_m^2}{L_r} \right) i_{s\alpha} + \frac{1}{\sigma L_s} \left(\frac{L_m}{L_r} \frac{1}{T_r} \right) \varphi_{r\alpha} + \frac{1}{\sigma L_s} \left(\frac{L_m}{L_r} \right) \omega_r \varphi_{r\beta}$$

$$\frac{di_{s\beta}}{dt} = -\frac{1}{\sigma L_s} \left(R_s + \frac{1}{T_r} \frac{L_m^2}{L_r} \right) i_{s\beta} - \frac{1}{\sigma L_s} \left(\frac{L_m}{L_r} \right) \omega_r \varphi_{r\alpha} + \frac{1}{\sigma L_s} \left(\frac{L_m}{L_r} \frac{1}{T_r} \right) \varphi_{r\beta}$$

$$\frac{d\varphi_{r\alpha}}{dt} = \frac{L_m}{T_r} i_{s\alpha} - \frac{1}{T_r} \varphi_{r\alpha} - \omega_r \varphi_{r\beta}$$

$$\frac{d\varphi_{r\beta}}{dt} = \frac{L_m}{T_r} i_{s\beta} + \omega_r \varphi_{r\alpha} - \frac{1}{T_r} \varphi_{r\beta}$$

$$\omega_m = p \Omega_m; \omega_r = [\omega_s - \omega_m]; \sigma = 1 - \frac{L_m^2}{L_s L_r}; T_r = \frac{L_r}{R_r}; T_s = \frac{L_s}{R_s}$$

Induction Motor Model in α - β frame

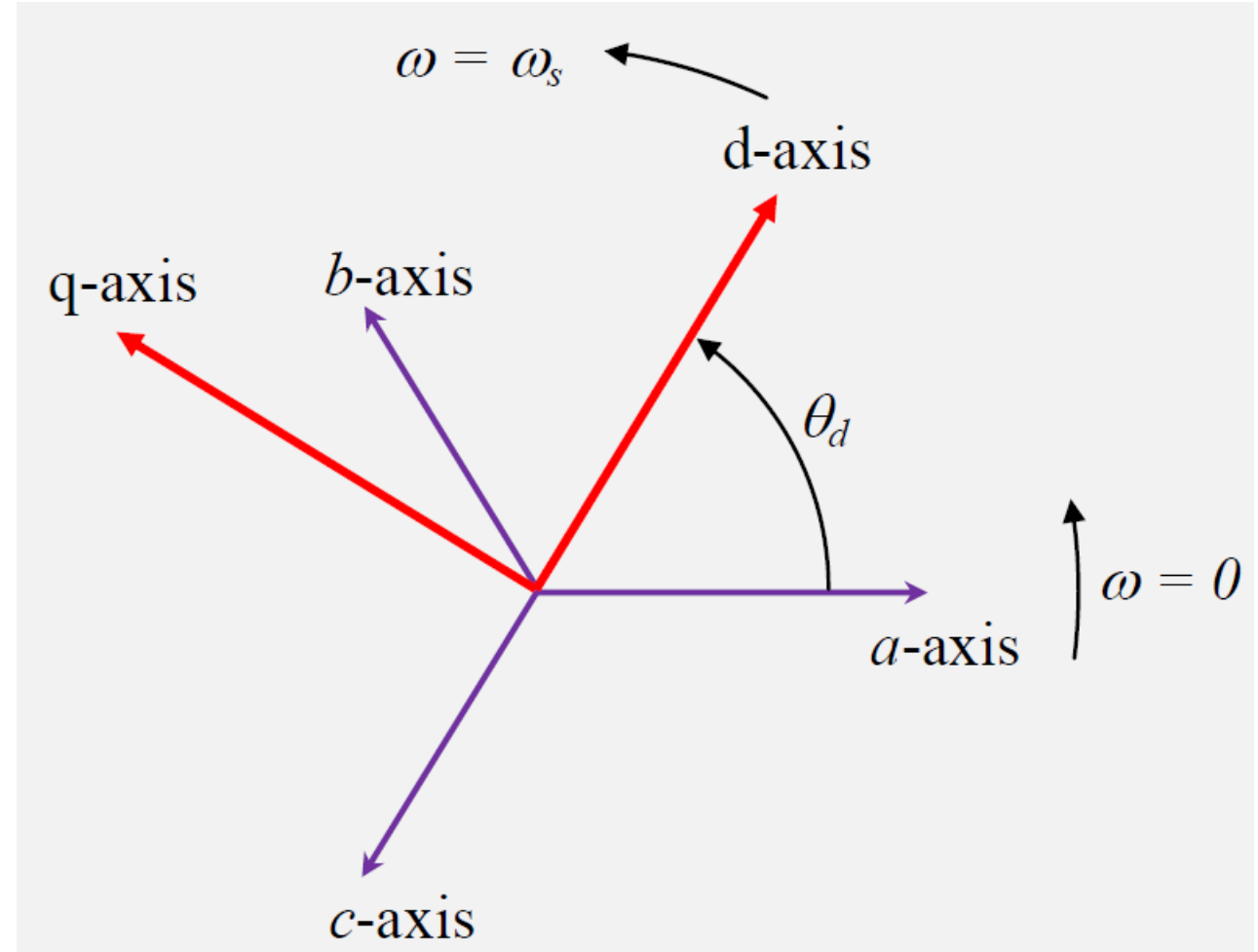
$$\Gamma_{em} - \Gamma_r = f\Omega_m + J \frac{d\Omega_m}{dt}$$

where Γ_{em} is the electromagnetic torque [N.m] and Γ_r is the resistive torque imposed by the machine shaft [N. m]. Ω : Mechanical speed, f : damping coefficient

Inference: Modeling the machine in this way (α - β frame) that reduces the number of quantities that we need to know in order to simulate machine operation. In fact, only the instantaneous values of the stator voltages and the resistive torque must be determined in order to impose them on the machine. Therefore, we do not need to know the stator frequency value, or the slip as in the case of the model whose equations are written in the reference frame rotating in synchronism

Park's Transformation

The d-axis is leading the q-axis by 90 electrical degrees; and the angle between the d-axis w.r.t. the a-axis is used



Park's Transformation

$$[\mathbf{f}_{dq0}] = [\mathbf{T}_{dq0}(\theta_d)] [\mathbf{f}_{abc}]$$

$$[\mathbf{f}_{dq0}] = \begin{bmatrix} f_d \\ f_q \\ f_0 \end{bmatrix}$$

$$[\mathbf{f}_{abc}] = \begin{bmatrix} f_a \\ f_b \\ f_c \end{bmatrix}$$

$$[\mathbf{T}_{dq0}(\theta_d)] = \frac{2}{3} \begin{bmatrix} \cos \theta_d & \cos(\theta_d - 2\pi/3) & \cos(\theta_d + 2\pi/3) \\ -\sin \theta_d & -\sin(\theta_d - 2\pi/3) & -\sin(\theta_d + 2\pi/3) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\theta_d = \omega t + \theta_0$$

Park's transformation is not power-invariant. But it can be made power invariant by replacing $2/3$ with $\sqrt{2/3}$

Induction Motor Model in d-q frame (Park's Frame) Synchronized to Rotating Field

Voltage Equations

$$v_{sd} = R_s i_{sd} + \frac{d\varphi_{sd}}{dt} - \omega_s \varphi_{sq}$$

$$v_{sq} = R_s i_{sq} + \frac{d\varphi_{sq}}{dt} + \omega_s \varphi_{sd}$$

$$v_{rd} = 0 = R_r i_{rd} + \frac{d\varphi_{rd}}{dt} - \omega_r \varphi_{rq}$$

$$v_{rq} = 0 = R_r i_{rq} + \frac{d\varphi_{rq}}{dt} + \omega_r \varphi_{rd}$$

ω_s = Synchronous speed in electrical rad/sec
 ω_r = slip speed in electrical rad/sec

Induction Motor Model in d-q frame (Park's Frame) Synchronized to Machine Rotor

Voltage Equations

$$v_{sd} = R_s i_{sd} + \frac{d\varphi_{sd}}{dt} - \omega_s \varphi_{sq}$$

$$v_{sq} = R_s i_{sq} + \frac{d\varphi_{sq}}{dt} + \omega_s \varphi_{sd}$$

$$v_{rd} = 0 = R_r i_{rd} + \frac{d\varphi_{rd}}{dt}$$

$$v_{rq} = 0 = R_r i_{rq} + \frac{d\varphi_{rq}}{dt}$$

$\omega_s = \omega_m$ = rotor speed in electrical rad/sec

Voltage Equations in the stationary reference frame

$$v_{s\alpha} = R_s i_{s\alpha} + s\varphi_{s\alpha}$$

$$v_{s\beta} = R_s i_{s\beta} + s\varphi_{s\beta}$$

$$v_{r\alpha} = 0 = R_r i_{r\alpha} + s\varphi_{r\alpha} + \omega_r \varphi_{r\beta}$$

$$v_{r\beta} = 0 = R_r i_{r\beta} + s\varphi_{r\beta} - \omega_r \varphi_{r\alpha}$$

Voltage Equations in the dq frame synchronized to mag field

$$v_{sd} = R_s i_{sd} + \frac{d\varphi_{sd}}{dt} - \omega_s \varphi_{sq}$$

$$v_{sq} = R_s i_{sq} + \frac{d\varphi_{sq}}{dt} + \omega_s \varphi_{sd}$$

$$v_{rd} = 0 = R_r i_{rd} + \frac{d\varphi_{rd}}{dt} - \omega_r \varphi_{rq}$$

$$v_{rq} = 0 = R_r i_{rq} + \frac{d\varphi_{rq}}{dt} + \omega_r \varphi_{rd}$$

$$v_{sd} = R_s i_{sd} + \frac{d\varphi_{sd}}{dt} - \omega_s \varphi_{sq}$$

$$v_{sq} = R_s i_{sq} + \frac{d\varphi_{sq}}{dt} + \omega_s \varphi_{sd}$$

$$v_{rd} = 0 = R_r i_{rd} + \frac{d\varphi_{rd}}{dt}$$

$$v_{rq} = 0 = R_r i_{rq} + \frac{d\varphi_{rq}}{dt}$$

Voltage Equations in the dq frame synchronized to rotor

Tutorial #1

An induction motor specifications are as follows; 5 hp, 200V, 3-phase star connected 4 pole 60 Hz. $R_s=0.277\Omega$, $R_r=0.183\Omega$, $L_m=0.0538\text{ H}$, $L_s=0.0553\text{ H}$, $L_r=0.056\text{ H}$. Calculate stator and rotor currents using stator-reference frame model. $\omega_r=140\text{ rad/sec}$

$$[\mathbf{T}_{\alpha\beta 0}] = \sqrt{\frac{2}{3}} \begin{bmatrix} 1 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & \frac{\sqrt{3}}{2} & -\frac{\sqrt{3}}{2} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

$$[\mathbf{T}_{\alpha\beta 0}]^{-1} = \sqrt{\frac{2}{3}} \begin{bmatrix} 1 & 0 & \frac{1}{\sqrt{2}} \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} & \frac{1}{\sqrt{2}} \\ -\frac{1}{2} & -\frac{\sqrt{3}}{2} & \frac{1}{\sqrt{2}} \end{bmatrix}$$

$$v_{s\alpha} = R_s i_{s\alpha} + s\varphi_{s\alpha}$$

$$v_{s\beta} = R_s i_{s\beta} + s\varphi_{s\beta}$$

$$v_{r\alpha} = 0 = R_r i_{r\alpha} + s\varphi_{r\alpha} + \omega_r \varphi_{r\beta}$$

$$v_{r\beta} = 0 = R_r i_{r\beta} + s\varphi_{r\beta} - \omega_r \varphi_{r\alpha}$$

Calculate electromagnetic torque generated in an induction motor with specifications as follows; 2000 hp, 1400 RPM 2300V, 3-phase star connected 4 pole 60 Hz, full load slip of 0.03746. $R_s=0.02\Omega$, $R_r=0.12\Omega$, $X_m=50\Omega$, $X_{ls}=X_{lr}=0.32\Omega$. Use synchronous reference frame

$$[\mathbf{T}_{dq0}(\theta_d)] = \frac{2}{3} \begin{bmatrix} \cos \theta_d & \cos(\theta_d - 2\pi/3) & \cos(\theta_d + 2\pi/3) \\ -\sin \theta_d & -\sin(\theta_d - 2\pi/3) & -\sin(\theta_d + 2\pi/3) \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\theta_d = \omega t + \theta_0$$

$$T_e = \frac{3}{2} p (\lambda_{ds} i_{qs} - \lambda_{qs} i_{ds})$$

$$p = P/2 = \text{Pole pairs}$$

$$v_{sd} = R_s i_{sd} + \frac{d\varphi_{sd}}{dt} - \omega_s \varphi_{sq}$$

$$v_{sq} = R_s i_{sq} + \frac{d\varphi_{sq}}{dt} + \omega_s \varphi_{sd}$$

$$v_{rd} = 0 = R_r i_{rd} + \frac{d\varphi_{rd}}{dt} - \omega_r \varphi_{rq}$$

$$v_{rq} = 0 = R_r i_{rq} + \frac{d\varphi_{rq}}{dt} + \omega_r \varphi_{rd}$$

Nonlinearities In Machine Equations

Saturation

Saturation of machine causes lower inductance and hence lower flux linkage. It can also generate more higher order harmonics due to current waveform distortion. The inductance and torque matrices have been assumed to have constant coefficients in Park's axes, but they are no longer constant when we consider saturation, their values depending on various currents. During disturbances the currents may vary over wide ranges. In computation we have to start with a certain set of inductances and then alter their values in short time steps as the calculation proceeds, according to the saturation characteristic of the magnetic path.

Space Harmonics

In most studies relating to synchronous machines, it is assumed that the flux wave is sinusoidally distributed in space. This assumption is reasonable for induction motors and cylindrical-rotor synchronous machines of good design. It is not always a valid assumption, however, and the effects of space harmonics, which produce voltage time-harmonics, electrical noise and parasitic torques, must in certain cases be taken into account.