



Institute of
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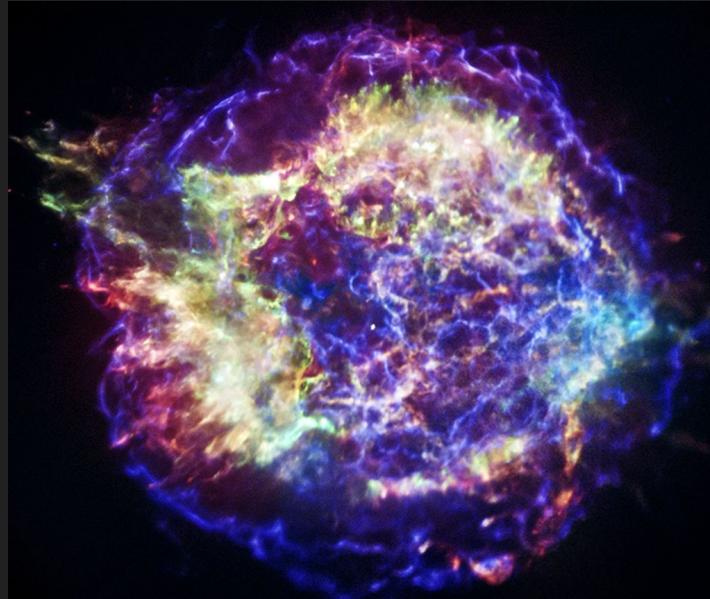


University of East Anglia,
May 9th 2023

Simulation-based inference (sbi) for pulsar population synthesis

Dr. Vanessa Gruber (gruber@ice.csic.es)

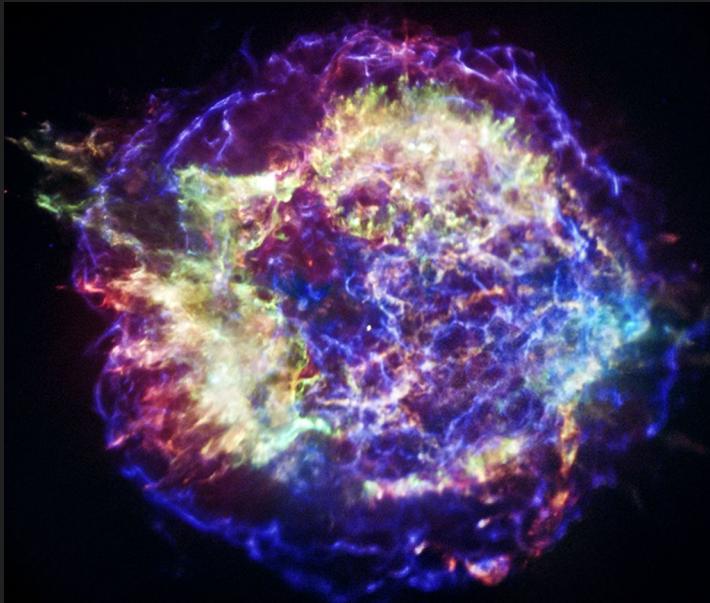
in collaboration with Michele Ronchi,
Celsa Pardo Araujo, and Nanda Rea



Cassiopeia A supernova remnant
(credit: NASA/CXC/SAO)

Outline

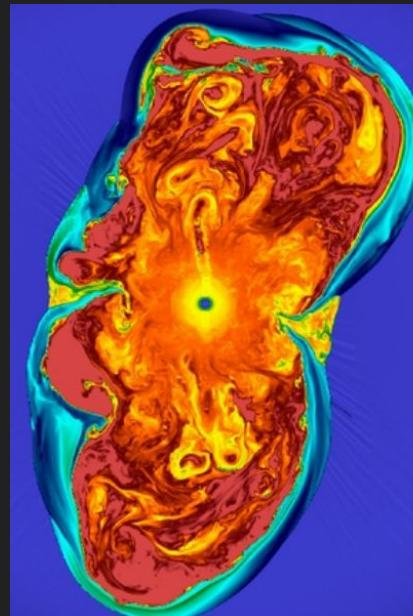
- **Neutron stars**
- **Pulsar population synthesis**
- **Machine learning and sbi**
- **Inference results**
- **Outlook**



Cassiopeia A supernova remnant
(credit: NASA/CXC/SAO)

Neutron-star formation

- Neutron stars are one of three types of **compact remnants**, created during the **final stages of stellar evolution**.
- When a **massive star of 8 - 25 solar masses** runs out of fuel, it collapses under its own gravitational attraction and **explodes in a supernova**.
- During the collapse, **electron capture** processes ($p + e^- \rightarrow n + \nu_e$) produce (a lot of) neutrons.



mass: $1.2 - 2.1 M_{\odot}$

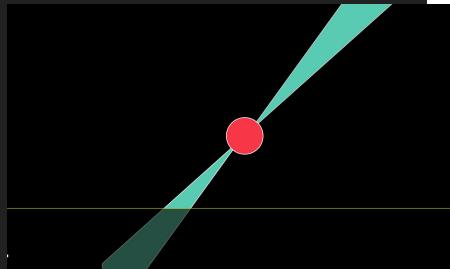
radius: 9 - 15 km

density: $10^{15} g/cm^3$

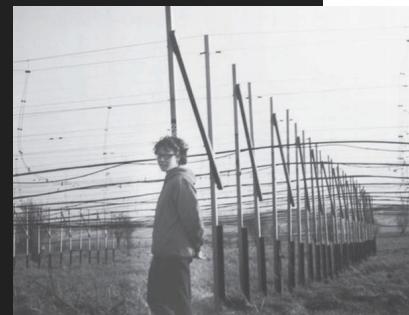
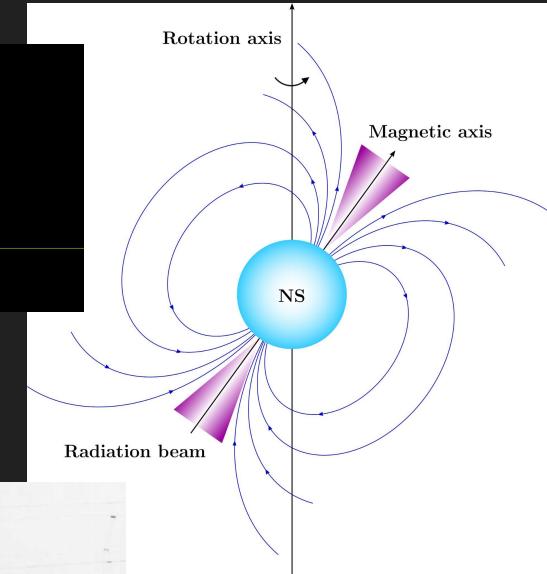
Snapshot of a 3D core-collapse supernova simulation (Mösta et al., 2014)

Lighthouse radiation

- Neutron stars have **extreme magnetic fields** between 10^8 - 10^{15} G. For comparison, the Earth's magnetic field is 0.5 G.
- Because rotation and magnetic axes are misaligned, neutron stars emit radio beams **like a lighthouse**.
- These pulses can be observed with radio telescopes. This is how neutron stars were first detected and why we call them **pulsars**.



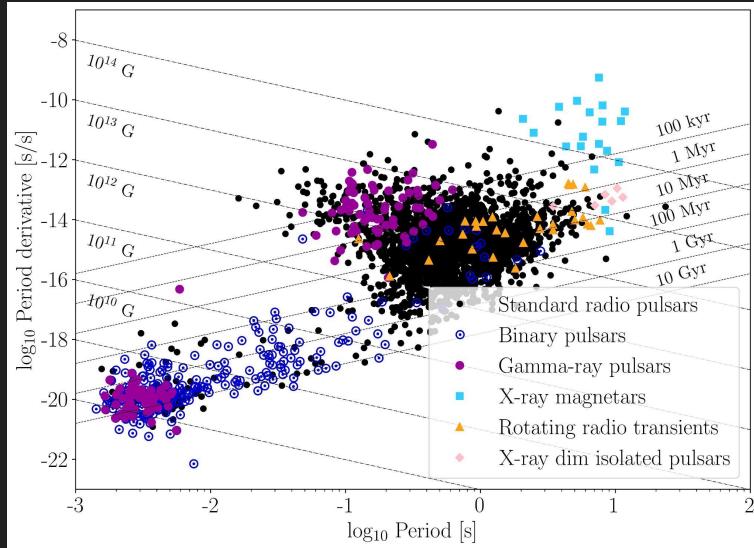
Credit: J. Christiansen



Sketch of the neutron-star exterior.

Dame Jocelyn Bell Burnell in front of her radio telescope in Cambridge, UK.

The neutron-star zoo



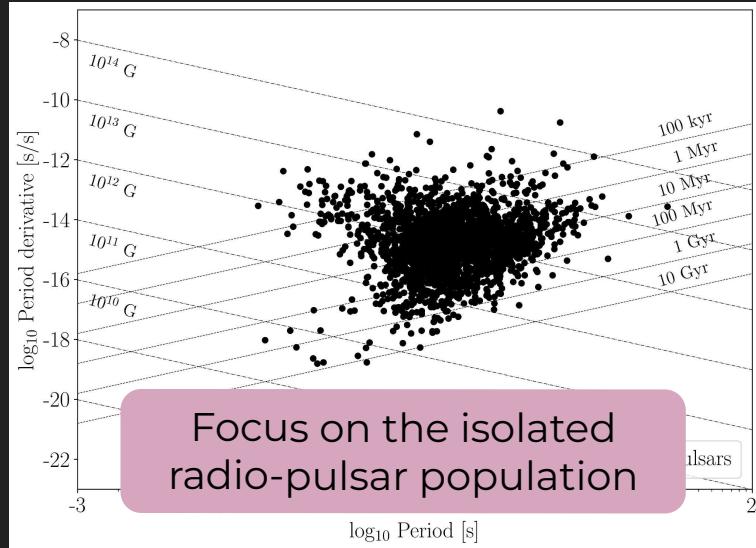
Period period-derivative plane for the pulsar population. Data taken from the ATNF Pulsar Catalogue (Manchester et al., 2005)

- Pulsars are **very precise clocks** and we time their pulses to **measure rotation periods P and derivatives \dot{P}** .
- We now observe neutron stars as pulsars **across the electromagnetic spectrum**.

~ **3,000 pulsars** are known to date

- Grouping neutron stars in the **$P\dot{P}$ -plane** according to their observed properties serves as a diagnostic tool to **identify different neutron-star classes**.

The neutron-star zoo



Period period-derivative plane for the pulsar population. Data taken from the ATNF Pulsar Catalogue (Manchester et al., 2005)

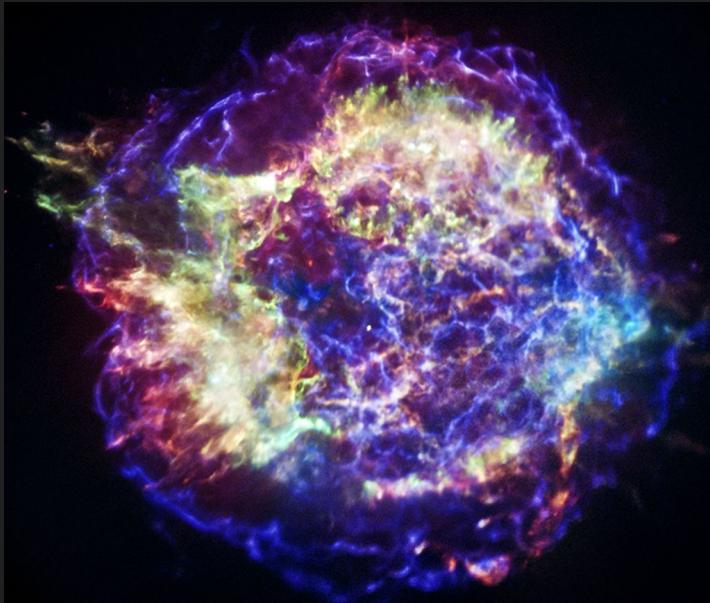
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General idea

- We can estimate the **total number of neutron stars in our Galaxy**

$$\text{CC supernova rate: } \sim 2 \text{ per century} \times \text{Galaxy age: } \sim 13.6 \text{ billion years} = \text{NS number: } \sim 2.8 \times 10^8$$

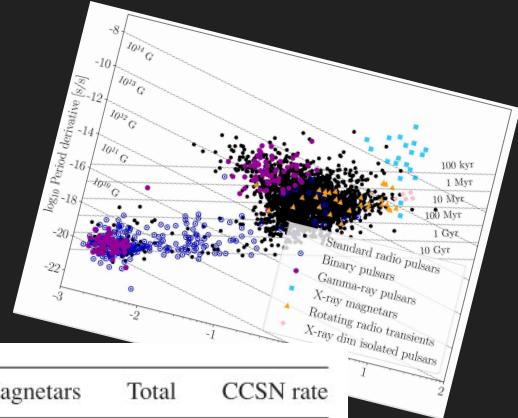
- We only **detect** a very **small fraction** of all neutron stars. Population synthesis bridges this gap focusing on the full population of neutron stars (e.g. Faucher-Giguère & Kaspi 2006, Lorimer et al. 2006, Gullón et al. 2014, Cieślar et al. 2020):



Goals

- Population synthesis allows us to **constrain the natal properties** of neutron stars and their **birth rates**.
- This is for example **relevant for**:
 - Massive star evolution
 - Gamma-ray bursts
 - Fast-radio bursts
 - Peculiar supernovae
- We can also learn about **evolutionary links between different neutron-star classes** (e.g., Viganó et al., 2013). This is important because estimates for the **Galactic core-collapse supernova rate** are **insufficient** for to explain the independent formation of different classes of pulsars (Keane & Kramer, 2008).

Estimated Galactic core-collapse supernova rate and birth rates for different pulsar classes (Keane & Kramer, 2008).



PSRs	RRATs	XDINSS	Magnetars	Total	CCSN rate
2.8 ± 0.5	$5.6^{+4.3}_{-3.3}$	2.1 ± 1.0	$0.3^{+1.2}_{-0.2}$	$10.8^{+7.0}_{-5.0}$	1.9 ± 1.1
1.4 ± 0.2	$2.8^{+1.6}_{-1.6}$	2.1 ± 1.0	$0.3^{+1.2}_{-0.2}$	$6.6^{+4.0}_{-3.0}$	1.9 ± 1.1
1.1 ± 0.2	$2.2^{+1.7}_{-1.3}$	2.1 ± 1.0	$0.3^{+1.2}_{-0.2}$	$5.7^{+4.1}_{-2.7}$	1.9 ± 1.1
1.6 ± 0.3	$3.2^{+2.5}_{-1.9}$	2.1 ± 1.0	$0.3^{+1.2}_{-0.2}$	$7.2^{+5.0}_{-3.4}$	1.9 ± 1.1
1.1 ± 0.2	$2.2^{+1.7}_{-1.3}$	2.1 ± 1.0	$0.3^{+1.2}_{-0.2}$	$5.7^{+4.1}_{-2.7}$	1.9 ± 1.1

Dynamical evolution I

- **Neutron stars are born in star-forming regions**, i.e., in the Galactic disk along the Milky Way's spiral arms, **and receive kicks** during the supernova explosions.
- We make the following assumptions:
 - Spiral-arm model (Yao et al., 2017) plus rigid rotation with $T = 250$ Myr
 - **Exponential disk model** with scale height h_c (Wainscoat et al., 1992)
 - Single-component **Maxwell kick-velocity distribution** with dispersion σ_k (Hobbs et al., 2005)
 - Galactic potential (Marchetti et al., 2019)

Artistic illustration of the Milky Way (credit: NASA JPL)



$$\mathcal{P}(z) = \frac{1}{h_c} e^{-\frac{|z|}{h_c}}$$

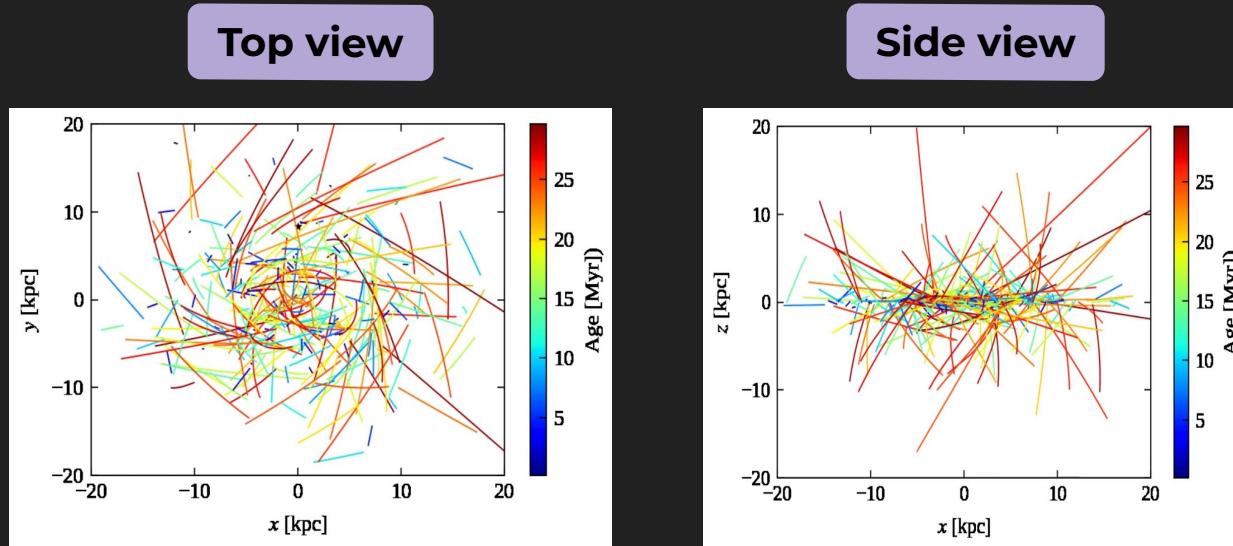
$$\mathcal{P}(v_k) = \sqrt{\frac{2}{\pi}} \frac{v_k^2}{\sigma_k^3} e^{-\frac{v_k^2}{2\sigma_k^2}}$$

For Monte-Carlo approach,
we **vary two uncertain parameters h_c and σ_k** .

Dynamical evolution II

- For our Galactic model Φ_{MW} , we evolve the stars' position & velocity by **solving Newtonian equations of motion** in cylindrical galactocentric coordinates:

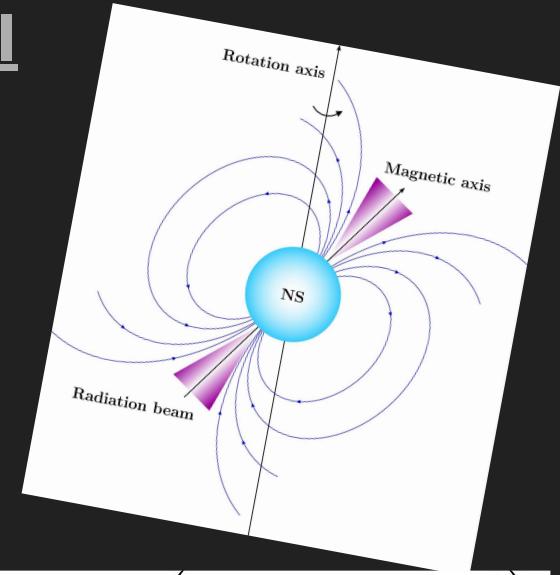
$$\ddot{\vec{r}} = -\vec{\nabla}\Phi_{\text{MW}}$$



Galactic evolution tracks for $h_c = 0.18$ kpc, $\sigma = 265$ km/s.

Magneto-rotational evolution I

- The neutron-star magnetosphere exerts a **torque onto the star**. This causes **spin-down** and **alignment of the magnetic and rotation axes**.
- Neutron star **magnetic fields decay** due to the Hall effect and Ohmic dissipation in the outer stellar layer (crust) (e.g., Viganó et al., 2013 & 2021).
- We make the following assumptions:
 - **Initial periods** follow a log-normal with μ_P and σ_P (Igoshev et al., 2022)
 - **Initial fields** follow a log-normal with μ_B and σ_B (Gullón et al., 2014)



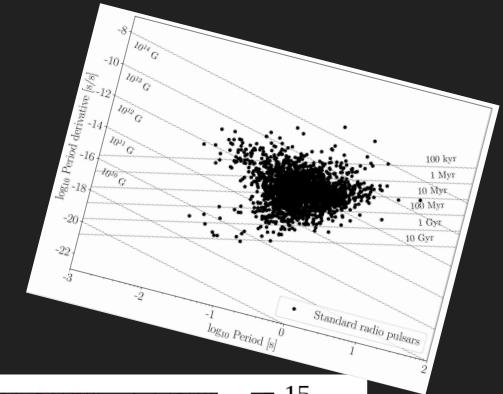
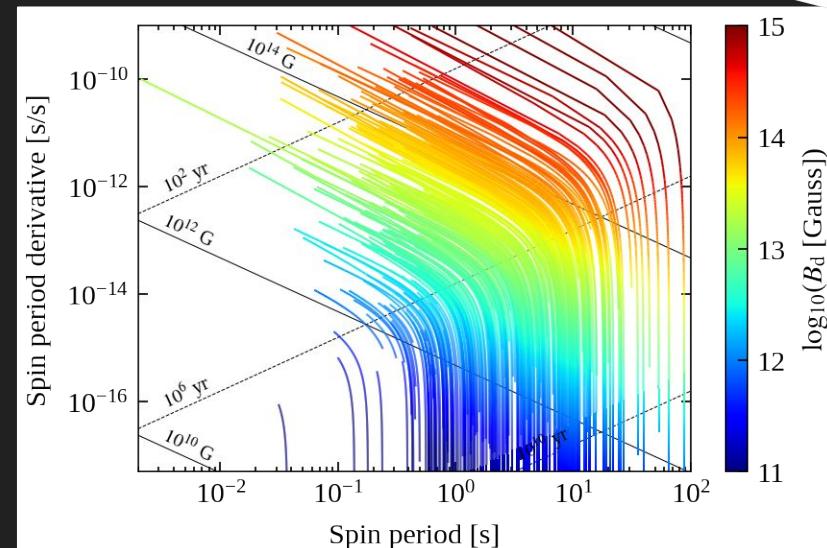
$$\mathcal{P}(P_0) = \frac{\log_{10}(e)}{\sqrt{2\pi}P_0\sigma_P} \exp\left(-\frac{[\log_{10}(P_0) - \mu_P]^2}{2\sigma_P^2}\right)$$

Here, we **vary** the four uncertain parameters μ_P , μ_B , σ_P and σ_B .

Magneto-rotational evolution II

- To model the neutron stars' magneto-rotational evolution, we numerically **solve three coupled ordinary differential equations** for the period, the misalignment angle and the magnetic field strength (Aguilera et al., 2008; Philippov et al. 2014).
- This allows us to follow the stars' P and \dot{P} evolution in the \dot{P} -plane.

\dot{P} evolution tracks for $\mu_P = -0.6$, $\sigma_P = 0.3$, $\mu_B = 13.25$ and $\sigma_B = 0.75$.



Radio emission and detection

- The stars' **rotational energy E_{rot}** is converted into coherent radio emission. We assume that the corresponding **radio luminosity L_{radio}** is proportional to the loss of E_{rot} (Faucher- Giguère & Kaspi, 2006; Gullón et al., 2014). L_0 is taken from observations.
- As **emission is beamed**, ~ 90% of pulsars do not point towards us. For those intercepting our line of sight, compute **radio flux S_{radio}** & **pulse width W** .

$$L_{\text{radio}} = L_0 \left(\frac{\dot{P}}{P^3} \right)^{1/2} \propto \dot{E}_{\text{dot}}^{1/2}$$

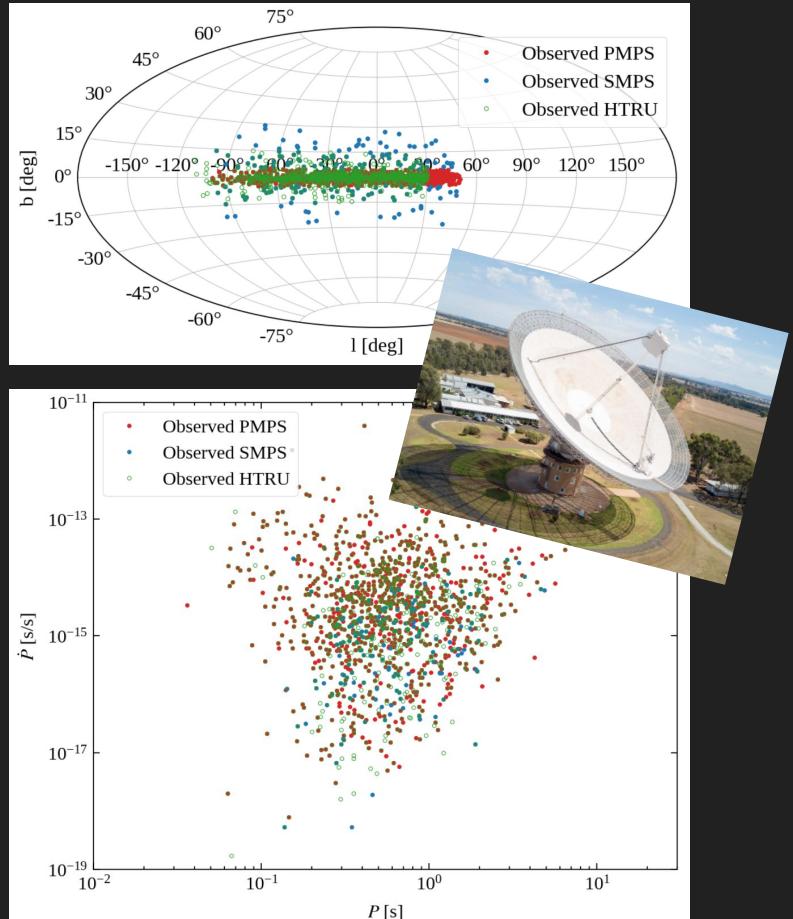
$$S_{\text{radio}} = \frac{L_{\text{radio}}}{\Omega_{\text{beam}} d^2}$$

A pulsar counts as detected, if it **exceeds the sensitivity threshold** for a survey recorded with a specific radio telescope.

Three pulsar surveys

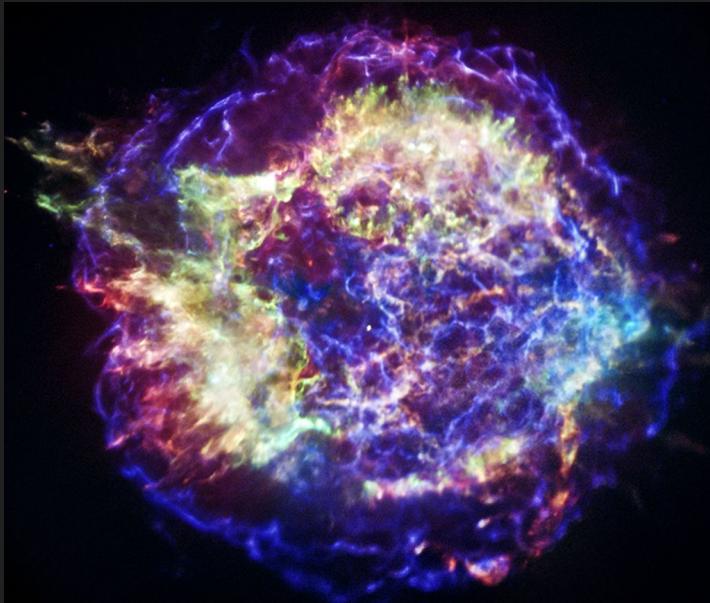
- We compare our simulated populations with three surveys from Murriyang (the Parkes Radio Telescope):
 - **Parkes Multibeam Pulsar Survey** (PMPS): 1,009 *isolated pulsars*
 - **Swinburne Parkes Multibeam Pulsar Survey** (SMPS): 218 *isol. p.*
 - **High Time Resolution Universe Survey** (HTRU): 1,023 *isol. pulsars*

Can we constrain birth properties by looking at a current snapshot of the pulsar population?



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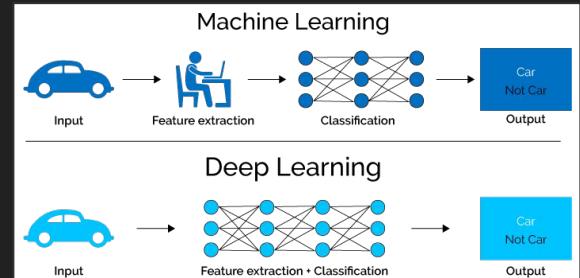
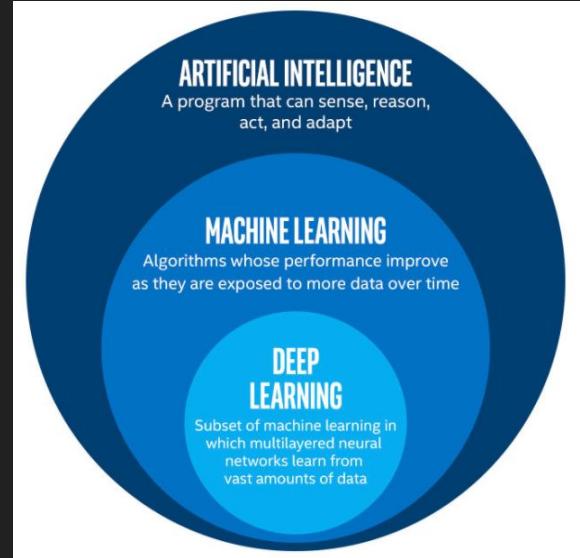
Comparing models and data

- Comparing observations to models and **constraining regions of the parameter space** that are **most probable given the data** is fundamental to many fields of science.
- Pulsar population synthesis is complex and has **many free parameters**. To compare synthetic simulations with observations, people have
 - Randomly sampled and then optimised ‘by eye’ (e.g., Gonthier et al., 2007)
 - Compared distributions of individual parameters using χ^2 - and KS-tests (e.g., Narayan & Ostriker, 1990; Faucher-Giguère & Kaspi, 2006)
 - Used annealing methods for optimisation (Gullón et al., 2014)
 - Performed Bayesian inference for simplified models (Cieślar et al., 2020)

These methods do not scale well and are **difficult to use** with the **multi-dimensional data** produced in population synthesis.

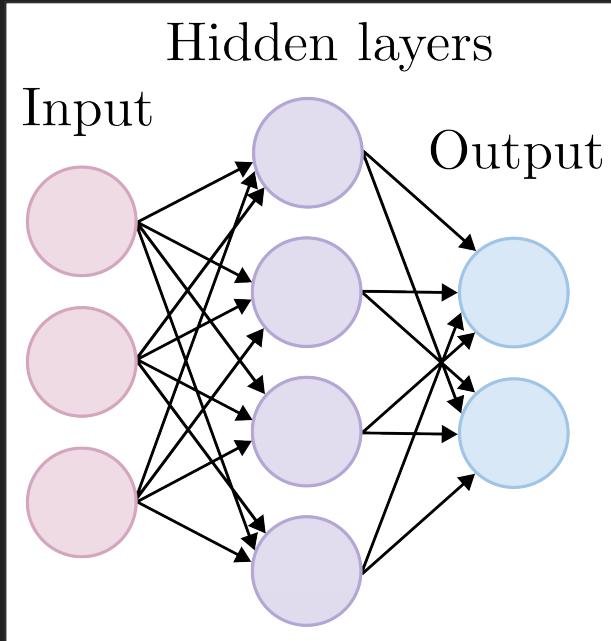
Deep learning

- Deep Learning is a subfield of Artificial Intelligence and Machine Learning. It focuses on using **multi-layered neural networks** to learn from large datasets. Different to classical ML approaches, deep learning does **not require external feature engineering**.
- **Recognising features in a hierarchical way** allows deep neural networks to model **complex non-linear relationships** for large input data. This makes deep learning powerful when **working with unstructured data such as images**, where the number of features / pixel can easily exceed millions.



Credit: www.bigdata-insider.de (top), Acheron Analytics (bottom)

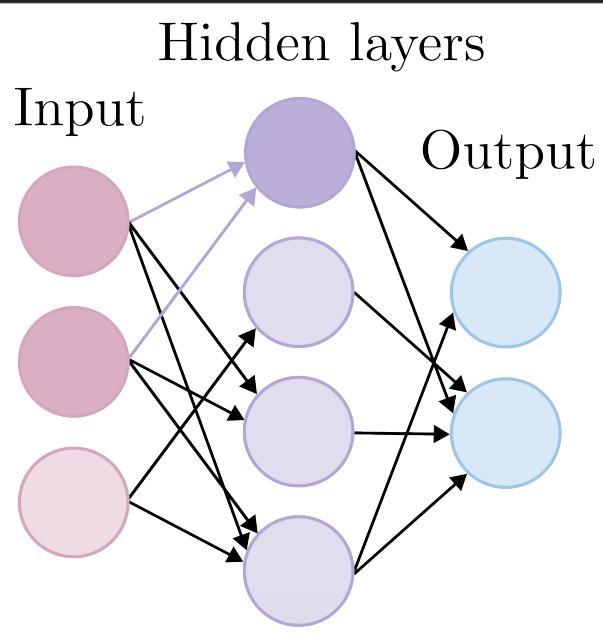
Convolutional neural networks (CNNs)



Sketch of a very simple
fully connected neural network.

- A neural network is composed of layers, which represent **stacks of neurons** (objects holding a single numerical value). Each layer encodes a simplified representation of the input data.
- A deep-learning **algorithm learns more and more about the input** as the data is passed through successive network layers.
- The **Multilayer Perceptron** is the simplest set-up where input and output are **fully connected**. In a CNN, not all nodes are connected, which **reduces the number of trainable parameters** and allows more flexibility for training.

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Convolutional and max pooling layers

- Besides fully connected layers, CNNs are composed of two types of filters:

Convolutional filters



$$\begin{bmatrix} -1 & -1 & -1 \\ -1 & 8 & -1 \\ -1 & -1 & -1 \end{bmatrix}$$



Max-pooling layers

3	1	0	9
8	4	7	3
6	5	0	4
1	2	9	0

2×2

8	9
6	9

These filters recognise features, such as detecting edges of an object in an image.

These filters extract the most relevant features, helping to speed up the training process.

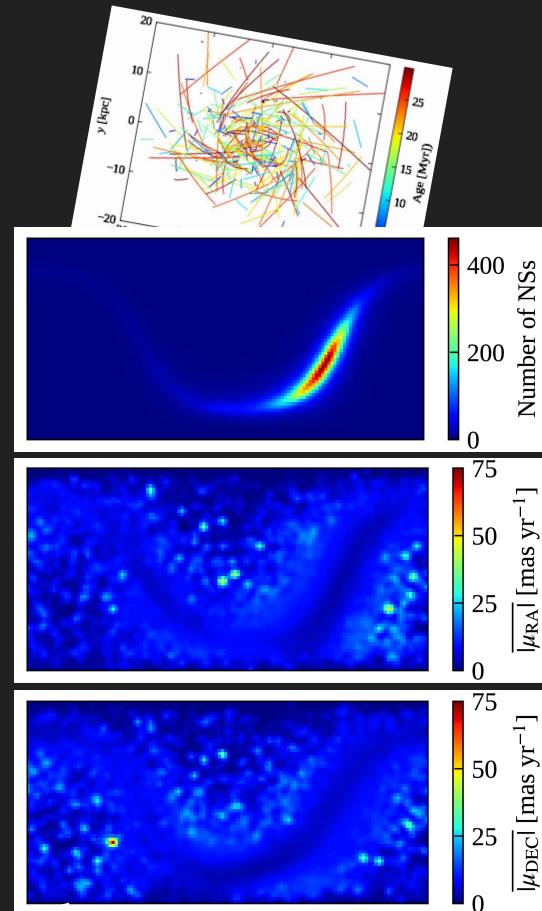
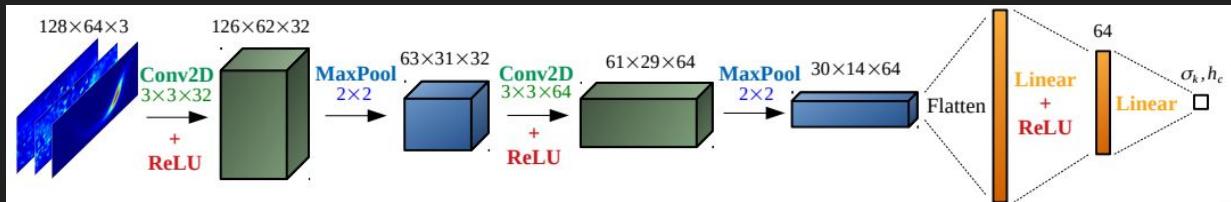
Proof of concept study I

- In Ronchi et al. (2021), we focused on the dynamical evolution and **simulated a database of 128 x 128 (=16,384) synthetic neutron-star populations.**

Vary **scale height**
 h_c in range
[0.02-2] kpc

Vary **dispersion** σ_k of
kick distribution
between [1-700] km/s

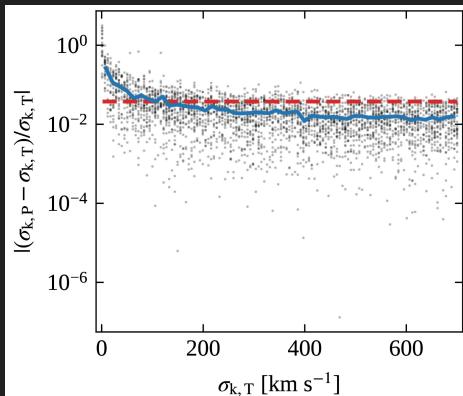
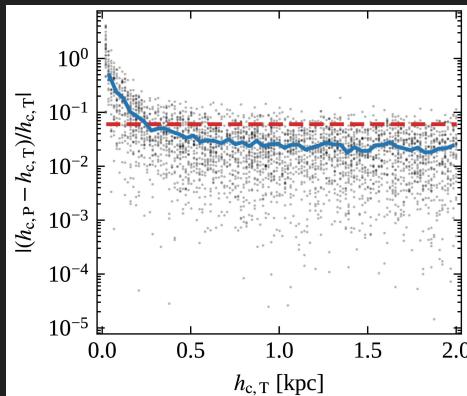
- We **perform supervised ML** and train a CNN to extract labels h_c and σ from position / velocity maps:



Stellar density and velocity maps in ICRS coordinates.

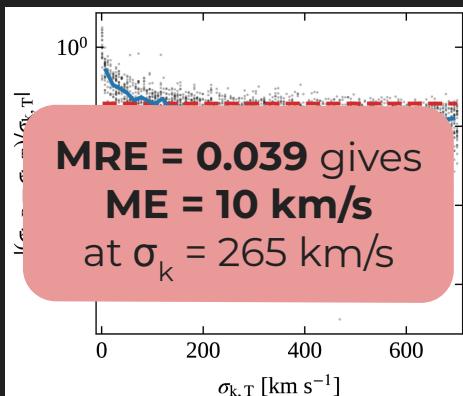
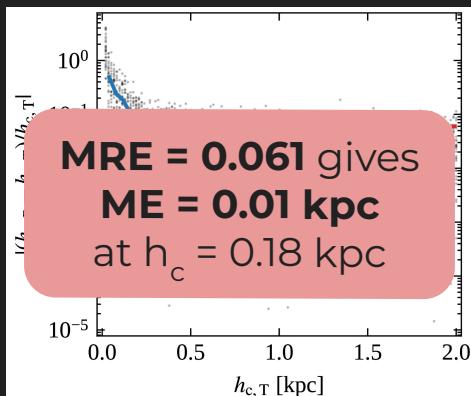
Proof of concept study II

- **Training info:** We use the root mean square error as the loss function and validation metric, Kaiming initialisation, Adam for gradient-descent optimisation, and apply a 80 / 20% split of the full dataset for training and validation.
- The **CNN recovers the input values** very well. To visualise this, we can look at the **relative error between target and predicted labels**



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We did **not include observational biases** and assumed all simulated stars are detectable! **216 pulsars have measured proper motions**, insufficient for this precision.

Statistical inference

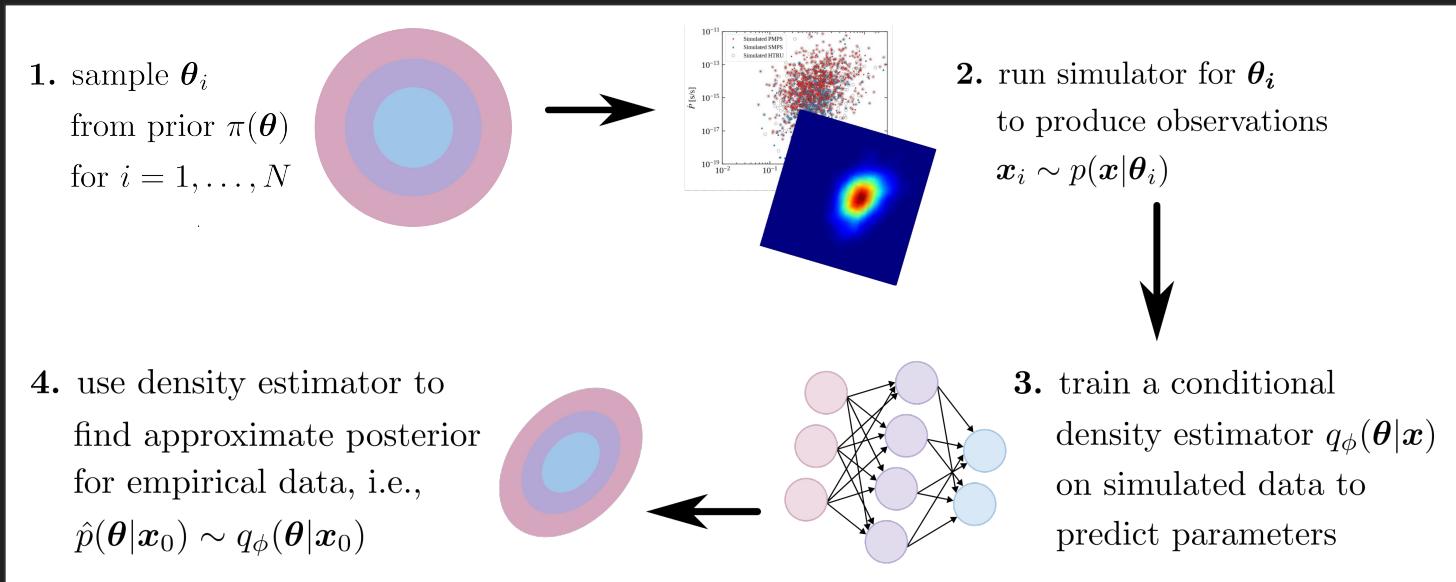
- Our initial study focused on **deducing point estimates**. However, we often do not require exact estimates but **knowledge of probable regions** is sufficient.
- This is where **Bayesian inference** comes in: based on some prior knowledge $\pi(\theta)$, a stochastic model and some observation x , we want to infer the most likely distribution $P(\theta|x)$ for our model parameters θ given the data x . This is **encoded in Bayes' Theorem**:

$$\underbrace{P(\theta|x)}_{\text{posterior}} = \frac{\overbrace{P(\theta)}^{\text{prior } \pi} \overbrace{P(x|\theta)}^{\text{likelihood } \mathcal{L}}}{\underbrace{P(x)}_{\text{evidence}}} = \frac{P(\theta) \int P(x,z|\theta) dz}{\int P(x|\theta') P(\theta') d\theta'}$$

For complex simulators, the **likelihood is defined implicitly and often intractable**. This is overcome with **simulation-based** (likelihood-free) **inference** (see e.g. Cranmer et al., 2020).

Simulation-based inference I

- To perform **Bayesian inference for any kind of (stochastic) forward model** (e.g. those specified by simulators), we use the following approach:



Simulation-based inference II

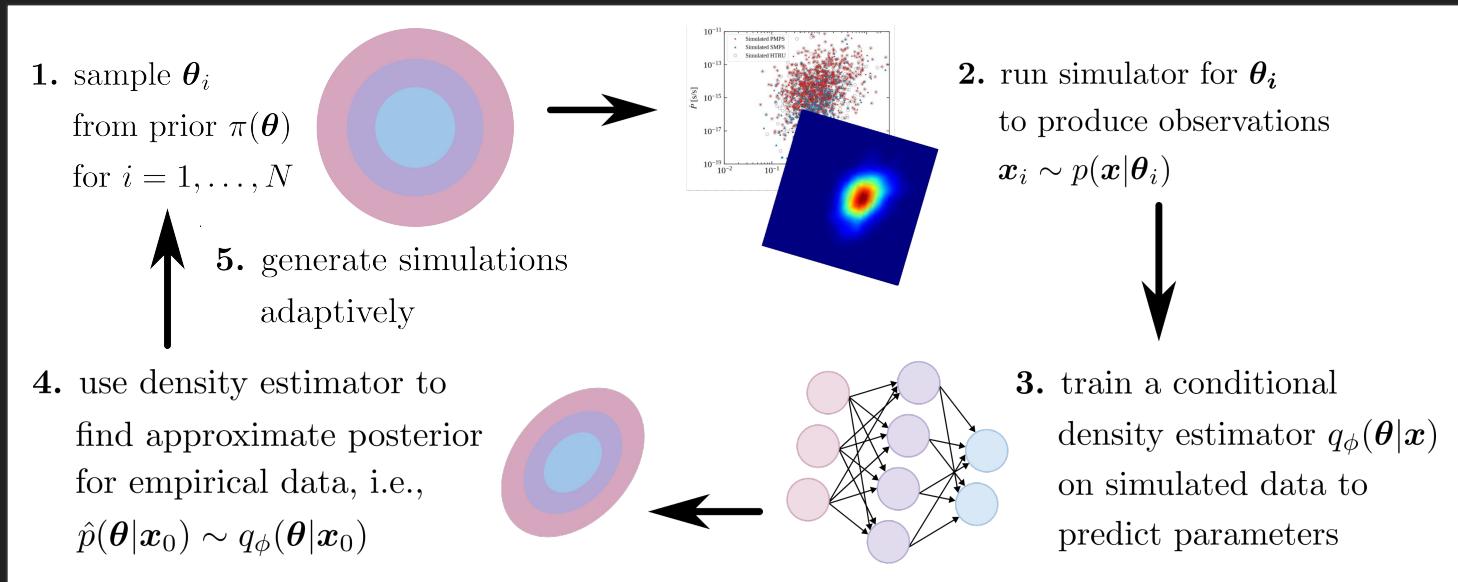
- Different approaches (all relying on deep learning) exist to **learn a probabilistic association** between the simulated data and the underlying parameters. These algorithms essentially focus on different pieces of Bayes' theorem:
 - Neural Posterior Estimation (NPE) (e.g., Papamakarios & Murray, 2016)
 - Neural Likelihood Estimation (NLE) (e.g., Papamakarios et al., 2019)
 - Neural Ratio Estimation (NRE) (e.g., Hermans et al., 2020; Delaunoy et al., 2022)

We focus on NPE. This allows us to **directly learn the posterior distribution**. In contrast, NLE and NRE need an extra (potentially time consuming) MCMC sampling step to construct a posterior.

- All methods exist in **sequential form** (SNPE, SNLE, SNRE), **which adds a fifth step to workflow**. Instead of sampling from the prior, we adaptively generate simulations from the posterior. This **typically requires fewer simulations**.

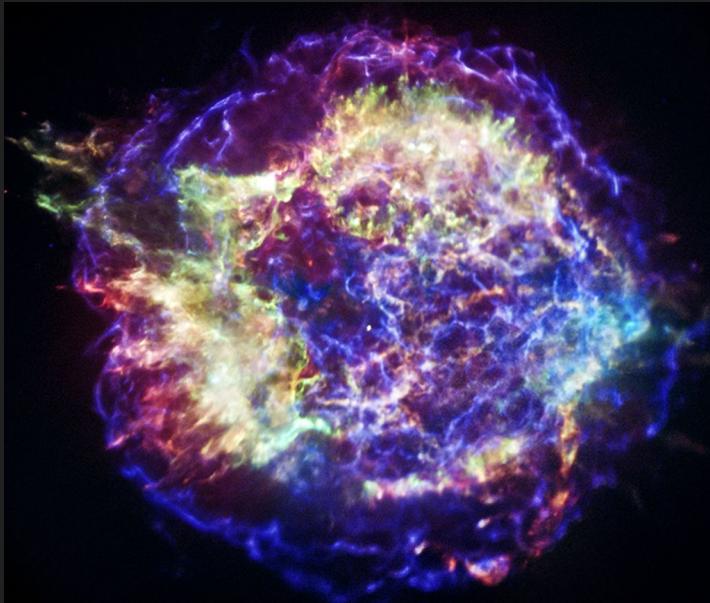
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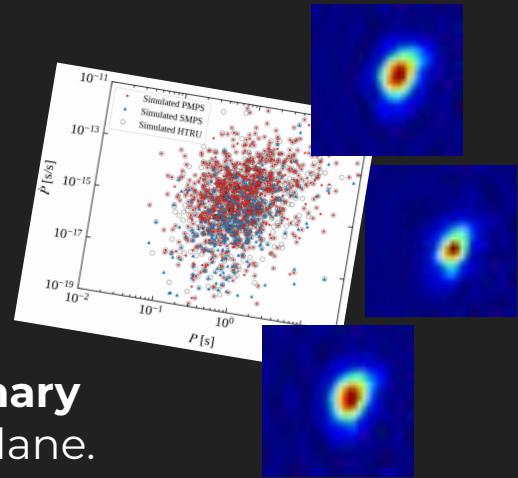
Workflow

- With our complex population synthesis simulator, we fix the dynamics to a fiducial model and **focus on the magneto-rotational evolution**.
- From our simulated populations, we **generate summary statistics**: density maps for three surveys in the PP-plane.
- To perform the inference, we use the **PyTorch package sbi** (Tejero-Cantero et al., 2020; <https://www.mackelab.org/sbi/>). Our trainable neural network has two parts:
 - CNN** (see Ronchi et al., 2021): compresses the data into a latent vector.
 - Mixture density network**: our posterior is approximated by a mixture of 10 Gaussians components; we learn the means, stds and coefficients.
- We initialise the CNN with Kaiming, use 89% of data for training, 10% for validation and 1% for testing, set the batch size to 8, and learning rate to 5×10^{-4} .

Varying the four parameters μ_p , μ_B , σ_p and σ_B , we simulate 360,000 synthetic populations over 6 weeks.

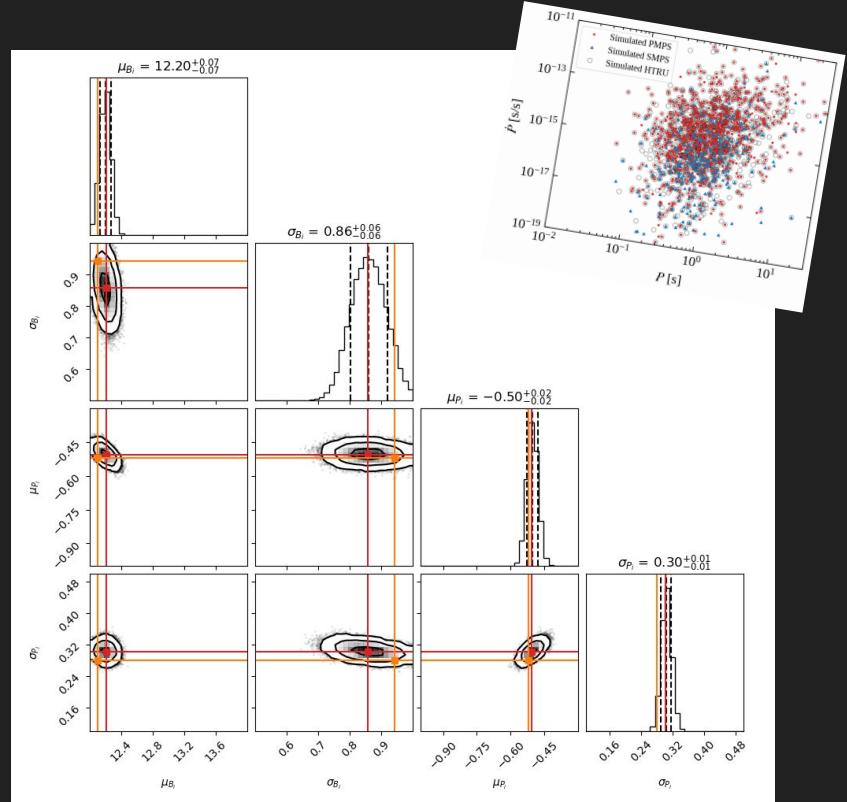
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Posterior distributions for test sample

- As our conditional density estimator is represented by a neural network, we can **directly evaluate the posterior distributions for a given observation**.
- We recover **narrow, well-defined posteriors** for all four parameters that contain the ground truth (parameters used for the forward simulation) at the 2σ levels.



Posterior distributions for observed population

- With our optimised neural network, we can also **infer the posteriors** for the **pulsar population recorded in our three surveys** and recover the following constraints:

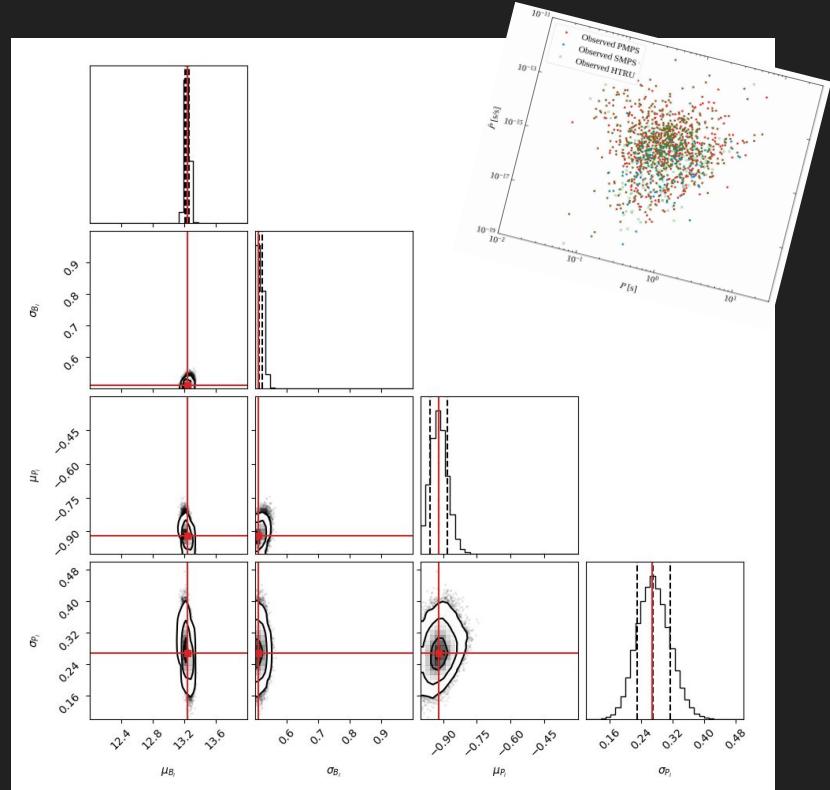
$$\mu_P = -0.92^{+0.04}_{-0.04}$$

$$\mu_B = 13.23^{+0.03}_{-0.03}$$

$$\sigma_P = 0.27^{+0.04}_{-0.04}$$

$$\sigma_B = 0.51^{+0.01}_{-0.01}$$

$$\mathcal{P}(P_0) = \frac{\log_{10}(e)}{\sqrt{2\pi}P_0\sigma_P} \exp\left(-\frac{[\log_{10}(P_0) - \mu_P]^2}{2\sigma_P^2}\right)$$



Posterior distributions for observed population

- With our optimised neural network, we can also **infer the posteriors** for the **pulsar population recorded in our three surveys** and the following constraints:

$$\mu_P = -0.92^{+0.08}_{-0.07}$$

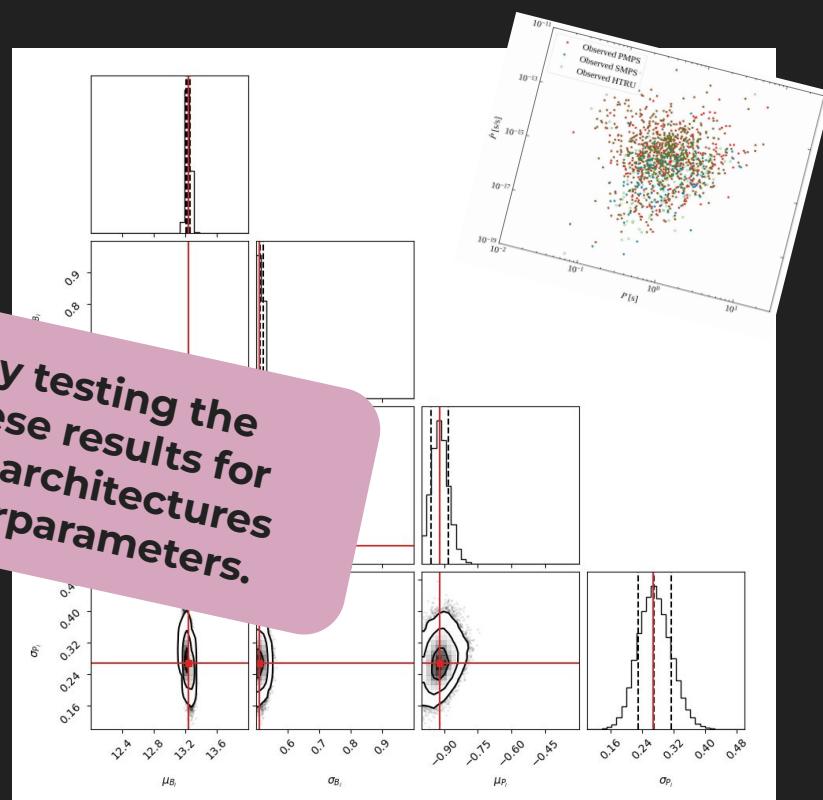
$$\sigma_P = 0.27^{+0.09}_{-0.07}$$

$$\mu_B$$

$$\sigma_B = 0.51^{+0.02}_{-0.02}$$

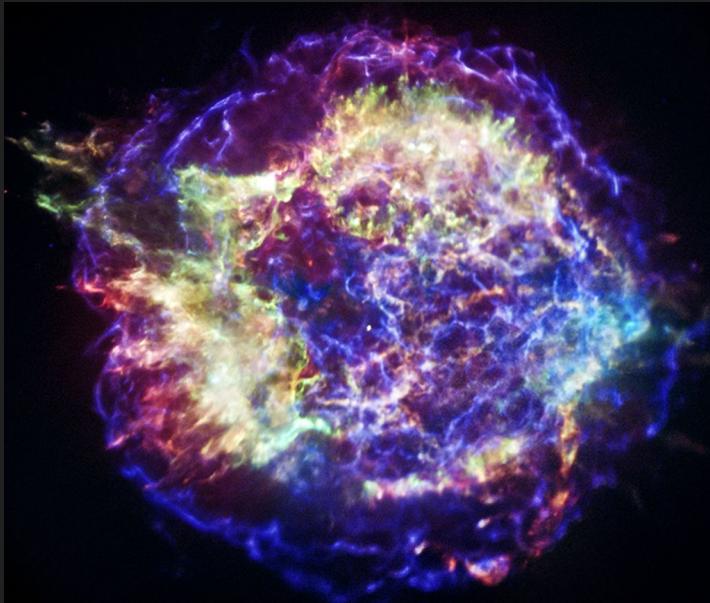
We are currently testing the robustness of these results for different network architectures and training hyperparameters.

$$\mathcal{P}(P_0) = \frac{\log_{10}(e)}{\sqrt{2\pi}P_0\sigma_P} \exp\left(-\frac{[\log_{10}(P_0) - \mu_P]^2}{2\sigma_P^2}\right)$$



Outline

- **Neutron stars**
- **Pulsar population synthesis**
- **Machine learning and sbi**
- **Inference results**
- **Summary and outlook**



Cassiopeia A supernova remnant
(credit: NASA/CXC/SAO)

Take-home points

- Neutron stars are **compact remnants** that **emit pulsed radiation** across the electromagnetic spectrum.
- **Standard radio pulsars** constitutes the largest class of observed neutron star.

- Pulsar **population synthesis** bridges gap between known pulsars and the invisible population.
- It **allows us to constrain birth rates** of different neutron star classes **and birth properties**.

- **Deep learning** with neural networks is ideal to **analyse high-dimensional astrophysical data**.
- **Simulation-based inference** has opened up the possibility for statistical inference **for complex simulators**.

Outlook

- There are **several directions** that we have started to look into:

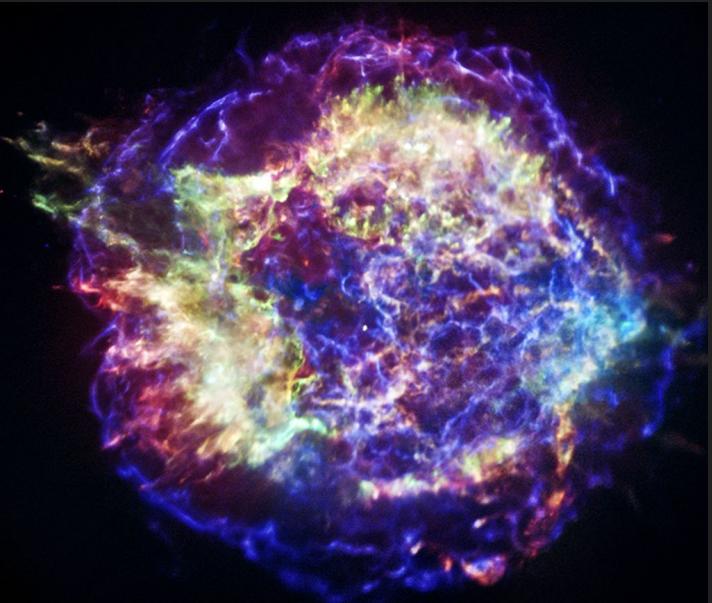
IMPROVING THE SIMULATOR

- Explore **different assumptions on initial period** and magnetic-field **distributions**
- Extend framework to model also gamma-ray and X-ray emission and **predict the multi-wavelength emission**

IMPROVING SBI

- **Test other approaches**
- Expand the approach to **active learning** and derive posteriors sequentially **using SNPE**, etc.

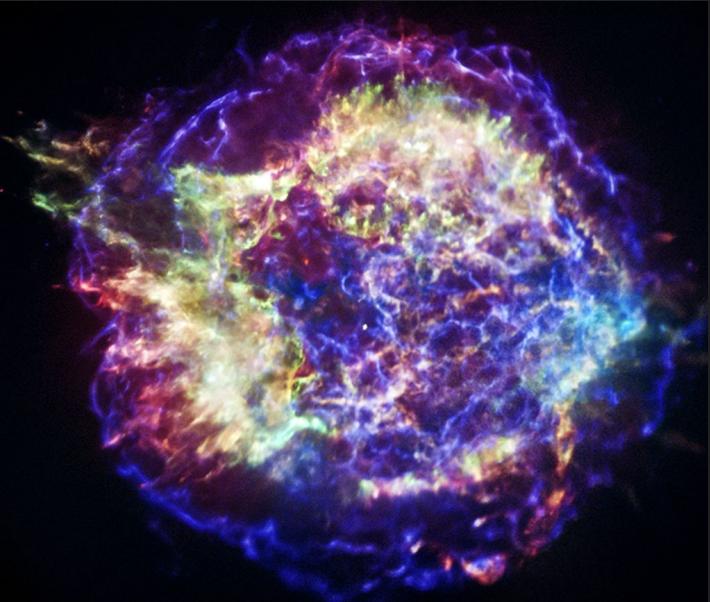
THANK YOU



Cassiopeia A supernova remnant
(credit: NASA/CXC/SAO)

Back-up

Slides



Cassiopeia A supernova remnant
(credit: NASA/CXC/SAO)

Magneto-rotational evolution - analytics

- For a given initial rotation period, magnetic field strength and inclination angle (uniformly distributed along the sphere), the evolution of a neutron star is determined by **three ordinary differential equations**. The first two are

$$\dot{P} = \frac{\pi^2}{c^3} \frac{B^2 R^6}{IP} (\kappa_0 + \kappa_1 \sin^2 \chi) \quad \dot{\chi} = -\frac{\pi^2}{c^3} \frac{B^2 R^6}{IP^2} (\kappa_2 \sin \chi \cos \chi)$$

The κ are determined from simulations. For a **realistic pulsar magnetosphere** filled with plasma we have $\kappa_0 \approx \kappa_1 \approx \kappa_2 \approx 1$ (Spitkovsky 2006; Philippov et al., 2014).

- From the induction equation, we deduce for the B-field (Aguilera et al., 2008):

$$\dot{B} = -\frac{c^2 B}{4\pi\sigma L^2} - \frac{cB^2}{4\pi e n_e L^2} \quad \text{with} \quad \tau_{\text{Ohm}} = \frac{4\pi\sigma L^2}{c^2} \sim 10^6 \text{ yr}, \quad \tau_{\text{Hall}} = \frac{4\pi e n_e L^2}{cB} \sim 10^4 B_{12} \text{ yr}$$

Improved B-field prescription

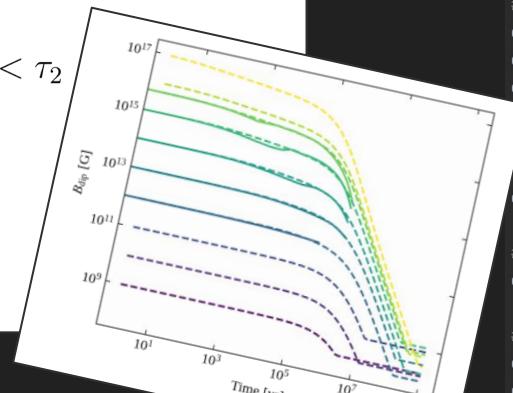
- To provide a more **realistic description of the magnetic field decay**, we **fit an analytical function** to B-field evolution curves obtained from 2D magneto-thermal simulations (Viganó et al., 2021) for various fields. As they are only valid up to 10^6 yr, we extend the **late-time behaviour with a power law** as follows:

$$B(t) = B_0 \left(1 + \frac{t}{\tau_1}\right)^{-a_1} \left(1 + \frac{t}{\tau_2}\right)^{a_1-a_2} \left(1 + \frac{t}{t_{\text{tr}}}\right)^{a_2-a_{\text{late}}} , \quad \text{if } \tau_1 < \tau_2 < t_{\text{tr}}$$

$$B(t) = B_0 \left(1 + \frac{t}{\tau_1}\right)^{-a_1} \left(1 + \frac{t}{t_{\text{tr}}}\right)^{a_2-a_{\text{late}}} , \quad \text{if } \tau_1 < t_{\text{tr}} < \tau_2$$

$$B(t) = B_0 \left(1 + \frac{t}{t_{\text{tr}}}\right)^{-a_{\text{late}}} , \quad \text{if } t_{\text{tr}} < \tau_1 < \tau_2$$

with $\tau_1 = A_1 B_0^{-b_1}$, $\tau_2 = A_2 B_0^{-b_2}$



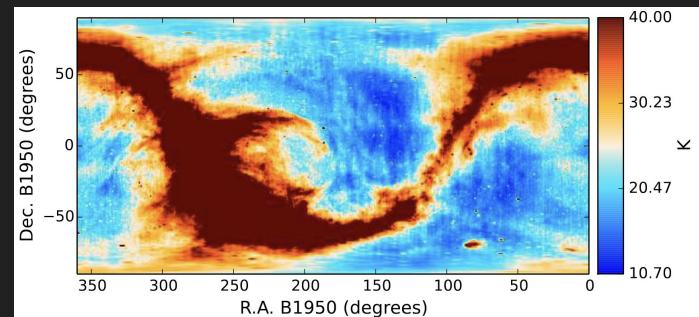
```
# Power-law indices.  
cfg["a1"] = -0.13  
cfg["a2"] = -3.0  
  
# Timescale parameters, normalization  
cfg["A1"] = 1.0e14  
cfg["b1"] = -0.8  
cfg["A2"] = 6.0e8  
cfg["b2"] = -0.2  
  
# Timescale in [yr] when transitioning  
cfg["tau_late"] = 2.0e6  
  
# Late time power-law index.  
cfg["a_late"] = -2.0  
  
# Parameters for a log-normal distri  
cfg["B_millisec_mean"] = 8.5  
cfg["B_millisec_sigma"] = 0.5
```

Radiometer equation

$$S_{\text{radio}} = \frac{L_{\text{radio}}}{\Omega_{\text{beam}} d^2}$$

- Following Lorimer & Kramer (2005), we convert the intrinsic (bolometric) radio flux S_{radio} into a flux density measured at a given frequency. We then account for the fact that the intrinsic pulse width is broadened to w_{eff} (due to interstellar scattering, dispersion and finite sampling resolution of a detector) and compute the period-averaged flux S_{mean} observed by a radio telescope.
- With these estimates and information specific to the PMPS, SMPS and HTRU surveys, we **determine the signal-to-noise-ratio** for a given simulated pulsar using the **radiometer equation**:

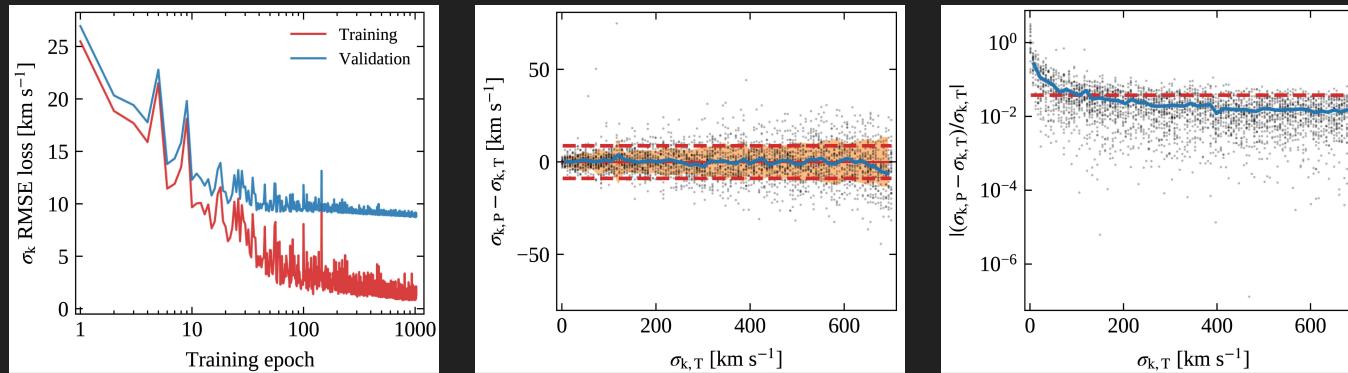
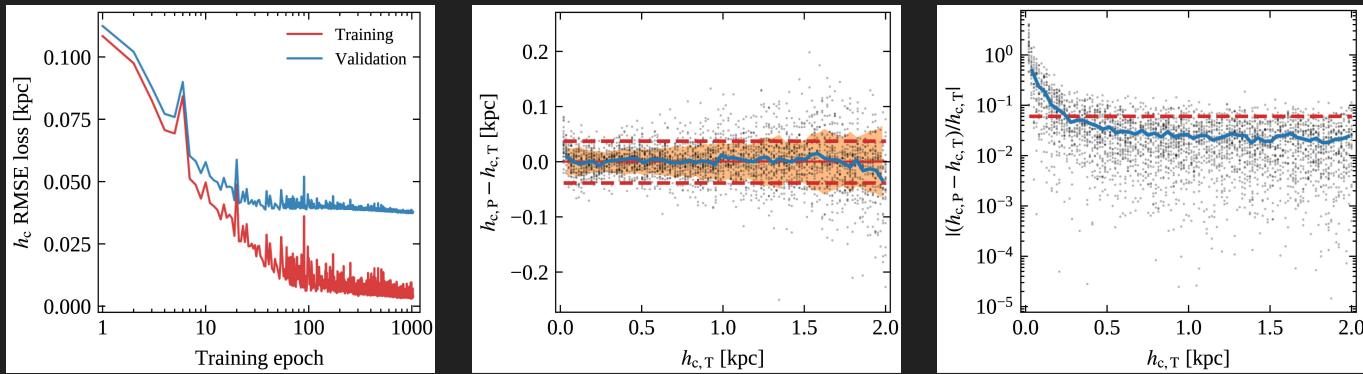
$$S/N = \frac{S_{\text{mean}} G \sqrt{N_{\text{pol}} \Delta\nu \Delta t_{\text{obs}}}}{\beta (T_{\text{rec}} + T_{\text{sky}}(l, b))} \sqrt{\frac{P - w_{\text{eff}}}{w_{\text{eff}}}}$$



Haslam sky temperature map taken from Remazeilles et al., (2015).

Ronchi et al. (2021) - Training Results

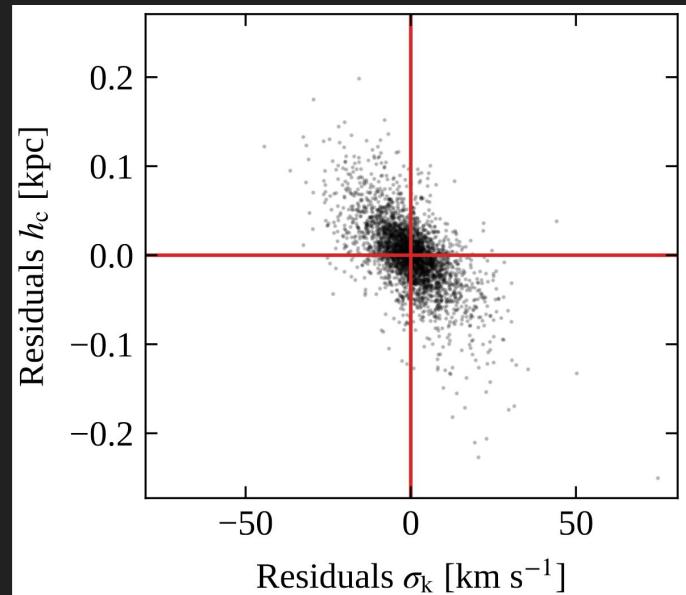
CNN validation
results for h_c :
 $\text{RMSE} = 0.038 \text{ kpc}$
& $\text{MRE} = 0.061$.



CNN validation
results for σ_k :
 $\text{RMSE} = 8.8 \text{ km/s}$
& $\text{MRE} = 0.039$.

Ronchi et al. (2021) - Degeneracy

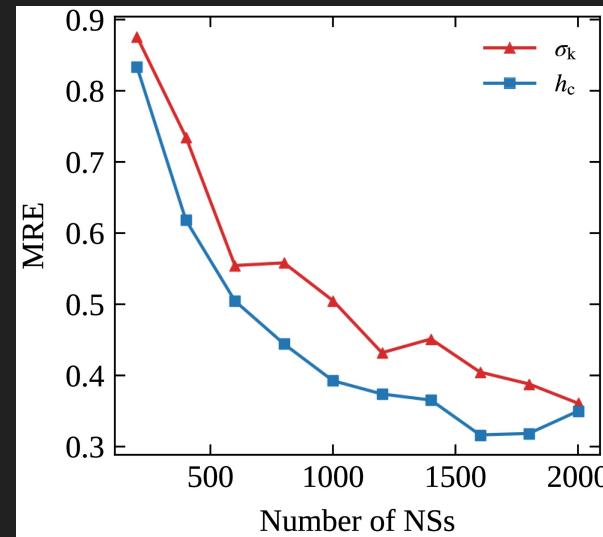
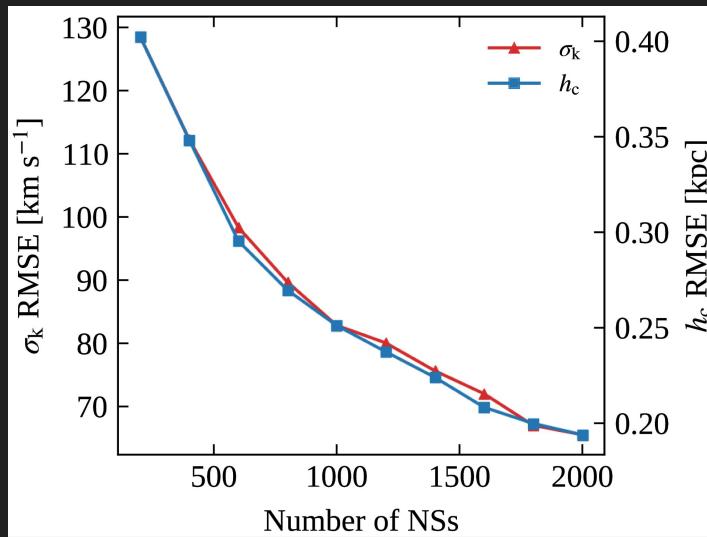
- Histogramming **distances from the Galactic plane** for current mock pulsar populations, we note a **degeneracy between h_c and σ_k** : large scale heights combined with small velocity dispersions lead to same outcomes as small scale heights with large velocity dispersions.
- **CNN recovers the degeneracy!**



We find an anticorrelation between the residuals in h_c and σ_k .

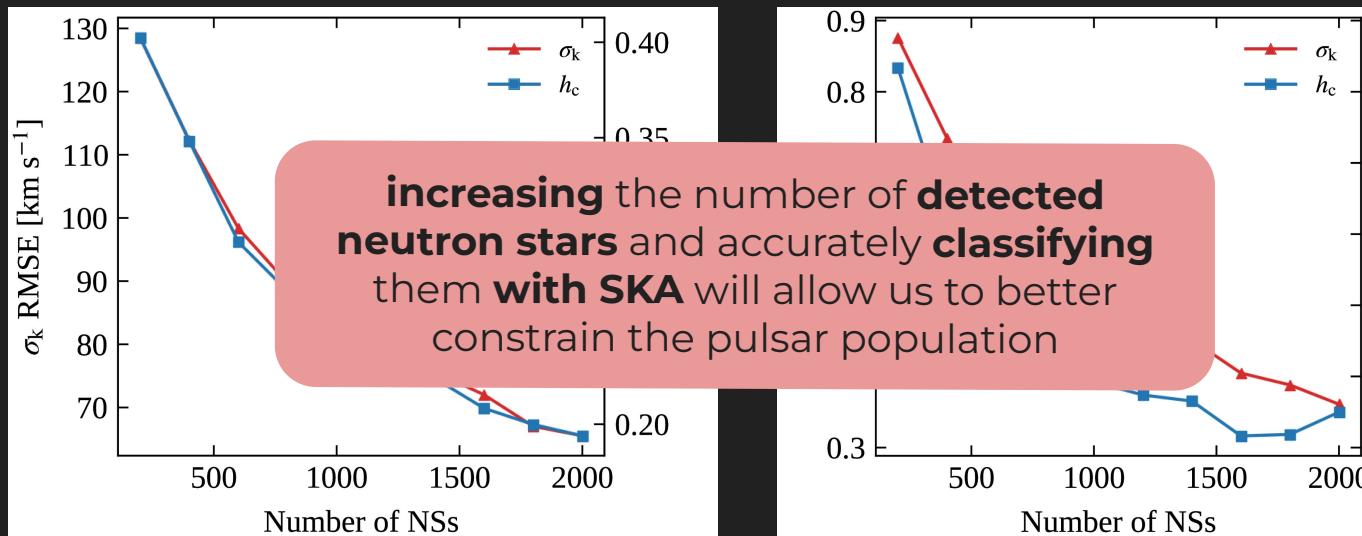
Ronchi et al. (2021) - Selection Biases

- Analyse the **CNN's predictive power** as a function of available data points (i.e. number of neutron stars) by resampling our fiducial simulation to **incorporate selection biases** from pulsars with **proper-motion measurements**.



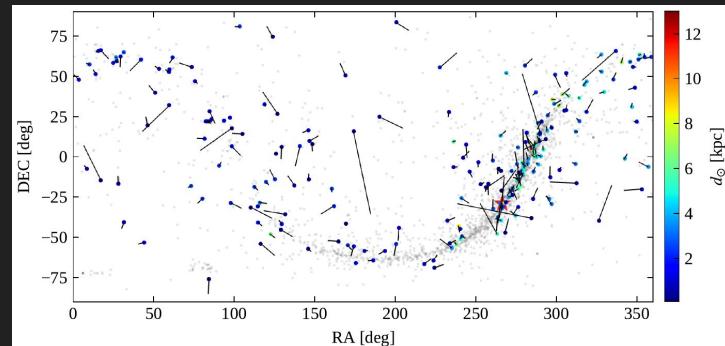
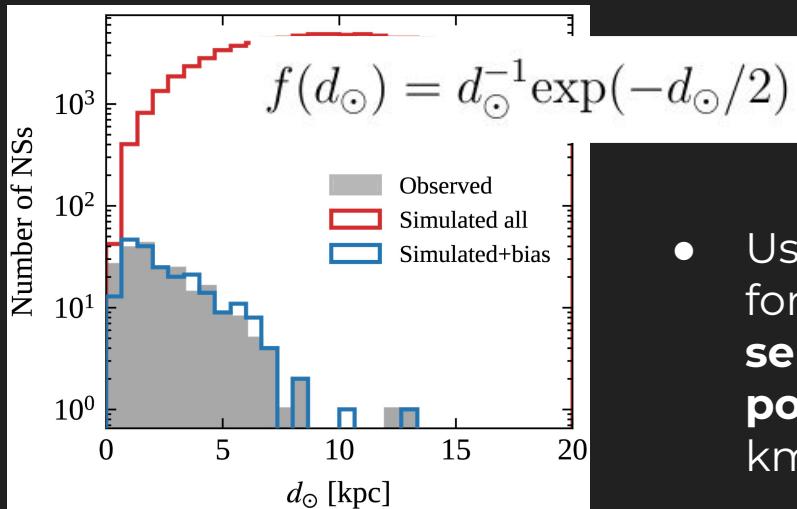
Ronchi et al. (2021) - Selection Biases

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Ronchi et al. (2021) - Selection Function

- To incorporate selection effects & observational biases, we use a **phenomenological approach** to reanalyse the CNN performance.

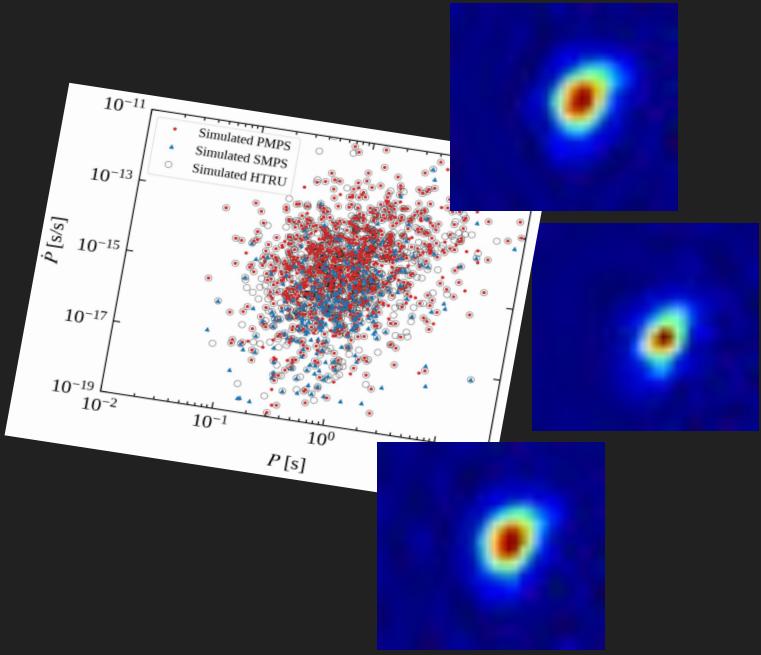


216 isolated NSs with proper motions.

- Use proper motion and distance estimates for 216 isolated pulsars to deduce **empirical selection function** $f(d_{\odot})$ and **resample population** with $h_c = 0.18$ kpc and $\sigma_k = 265$ km/s.

Simulated training dataset for sbi

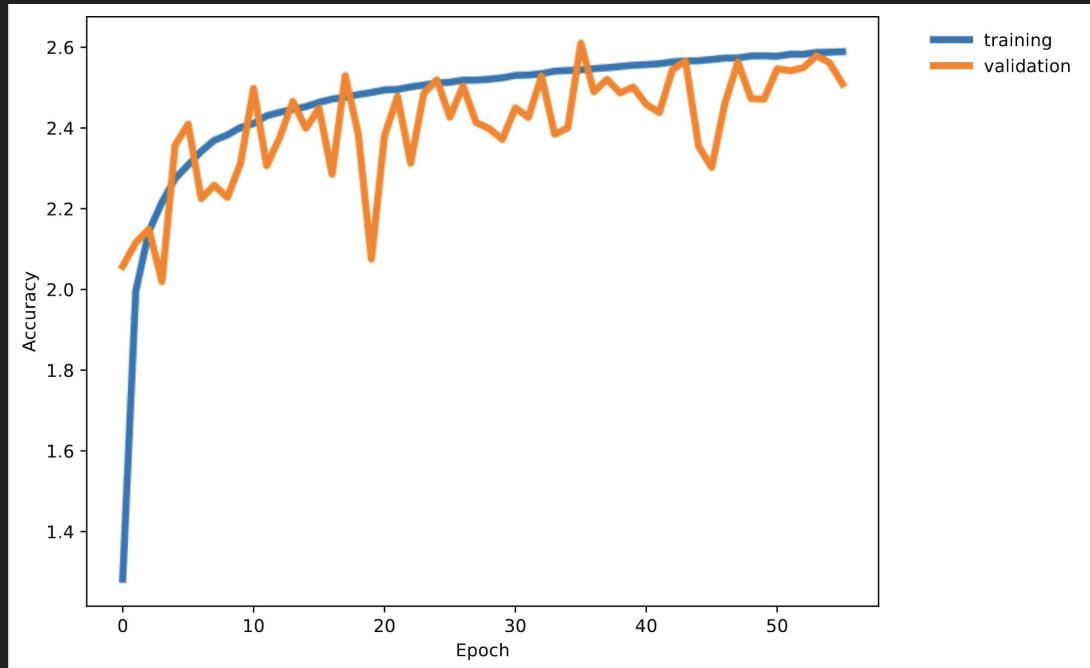
- We vary our **four parameters** for the initial period and magnetic field log-normal distributions as follows:
 - $\mu_P \in \text{Uniform}(-1, -0.3)$
 - $\sigma_P \in \text{Uniform}(0.1, 0.5)$
 - $\mu_B \in \text{Uniform}(12.0, 14.0)$
 - $\sigma_B \in \text{Uniform}(0.5, 1.0)$
- These also correspond to our **flat priors** for the simulation-based inference experiments.



$$\mathcal{P}(P_0) = \frac{\log_{10}(e)}{\sqrt{2\pi}P_0\sigma_P} \exp\left(-\frac{[\log_{10}(P_0) - \mu_P]^2}{2\sigma_P^2}\right)$$

Sbi training behaviour

- **Typical learning behaviour** for our simulation-based inference experiments.
- The **log-probability increases** as a function of training epochs for the training and validation loss.
- We see some variation in the validation loss but **little overfitting**.



Quality checks

- To assess the quality of our neural density estimator, we **perform simulation-based calibration** to assess how well our posterior approximates the true posterior by **plotting the 1D rank statistics** for our four parameters.
- For a well calibrated **unbiased posterior, ranks follow a uniform distribution**:

