## Math Phys II HW5

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## 1 Problem 1

We investigate the solution of the 3D Helmholtz equation with the boundary that the solution represents a traveling wave as  $r \to \infty$ . The general solution is:

$$U(r,\theta,\phi) = \sum_{\ell mn} (a_m j_m(kr) + b_m n_m(kr)) Y_\ell^m(\theta,\phi)$$
 (1.1)

We look at the solution near  $\vec{x}' = \vec{x}$ , where we require  $u(r) = -\frac{1}{4\pi r}$ . We see that  $-\frac{1}{4\pi}j_0(x) = (-\frac{1}{4\pi})\frac{\cos x}{x}$  will satisfy this. Looking at the requirement that the exterior solution becomes a traveling wave as  $r \to \infty$ , we see that we can add  $n_0(x) = \frac{\sin x}{x}$  to the solution to get an exponential form (spherical Henkel function  $h_0^{(1)}$ ). With spherical symmetry our final Green's function is:

$$G(\vec{x}, \vec{x}') = -\frac{e^{ikr}}{4\pi r} \tag{1.2}$$

## 2 Problem 2

Our trial Green's function places an image point of  $\mathbf{x}$  at  $\mathbf{x}_1 = \alpha \mathbf{x}$ . Without loss of generality we examine the solution for  $\mathbf{x}$  lying on the x-axis in Cartesian coordinates. The intersection points with the sphere, where  $G(\mathbf{x}, \mathbf{x}') = 0$ ,

are at  $\pm a$ . We then have:

$$\beta = \frac{|\mathbf{x}' - \alpha \mathbf{x}|}{|\mathbf{x}' - \mathbf{x}|} \tag{2.1}$$

$$|\mathbf{x}'| = \pm a \tag{2.2}$$

$$\frac{|a - \alpha x|}{|a - x|} = \frac{|a + \alpha x|}{|a + x|} \tag{2.3}$$

$$\alpha = (\frac{a}{x})^2 \tag{2.4}$$

$$\beta = \frac{a + \frac{a^2}{x}}{a + x} \tag{2.5}$$

$$\beta = \frac{1 + \frac{a}{x}}{1 + \frac{x}{a}} = \frac{a}{x} \tag{2.6}$$

For a Dirichlet boundary condition  $f(\theta', \phi')$  defined at r' = a we need to compute only the surface integral contribution to the solution. We find the normal derivative of the Green's function using the geometric identity  $|\mathbf{x}' - \frac{a^2}{m^2}\mathbf{x}| = \frac{a}{n}|\mathbf{x}' - \mathbf{x}|$ :

$$G(\mathbf{x}', \mathbf{x}) = \frac{1}{4\pi} \left( -\frac{a}{r|\mathbf{x}' - \frac{a^2}{r^2}\mathbf{x}|} + \frac{1}{|\mathbf{x}' - \mathbf{x}|} \right)$$
(2.7)

$$\nabla G = \frac{1}{4\pi} \frac{\left(1 - \frac{a^2}{r^2}\right)\mathbf{x}'}{|\mathbf{x}' - \mathbf{x}|^3} \tag{2.8}$$

$$\frac{\mathbf{x}'}{|\mathbf{x}'|} \cdot \nabla G = \frac{1}{4\pi a} \frac{a^2 - r^2}{c^3} \tag{2.9}$$

We can now set up the solution as the integral on the surface of sphere r' = a.

$$u(r,\theta,\phi) = \frac{1}{4\pi a} \int_{\Omega'} \frac{a^2 - r^2}{|\mathbf{x}' - \mathbf{x}|^3} f(\theta',\phi') \ d\Omega'$$
 (2.10)

We compare this to the series solution of Laplace's equation subject to an inhomogenous boundary condition on the surface of the sphere. The regular series solution for Laplace's equation inside a sphere is:

$$u(r,\theta,\phi) = \sum_{\ell,m} A_{\ell m} r^{\ell} Y_{\ell}^{m}(\theta,\phi)$$
 (2.11)

Applying the surface boundary condition  $u(a, \theta, \phi) = f(\theta, \phi)$  we get an expansion of  $f(\theta, \phi)$  in spherical harmonics:

$$\sum_{\ell,m} A_{\ell m} a^{\ell} Y_{\ell}^{m}(\theta,\phi) = f(\theta,\phi)$$
(2.12)

$$\alpha_{\ell m} = A_{\ell m} a^{\ell} \tag{2.13}$$

$$\alpha_{\ell m} = \int_{\Omega} Y_{\ell}^{m}(\theta, \phi) f(\theta, \phi) d\Omega$$
 (2.14)

We can see a similarity between the two solutions if we expand  $|\mathbf{x}' - \mathbf{x}|$  in spherical harmonics.

## 3 Problem 3

We examine the inhomogenous wave equation with a point source moving on a trajectory. The retarded potential solution is:

$$\phi(\mathbf{x},t) = -\frac{c}{4\pi} \int d^3 \mathbf{x}' \ dt' \ f(\mathbf{x}',t') \frac{\delta \left[ |\mathbf{x} - \mathbf{x}'| - c(t-t') \right]}{|\mathbf{x} - \mathbf{x}'|}$$
(3.1)

With  $f(\mathbf{x}',t') = \delta\left[\mathbf{x}' - \xi(t')\right]$ , then  $\mathbf{x}' = \xi(t')$ . We then have:

$$\phi(\mathbf{x},t) = -\frac{c}{4\pi} \int dt' \, \frac{\delta\left[|\mathbf{x} - \xi(t')| - c(t-t')\right]}{|\mathbf{x} - \xi(t')|} \frac{d\xi}{dt}$$
(3.2)

$$|\mathbf{x} - \xi(t')| = c(t - t') \tag{3.3}$$

$$\phi(\mathbf{x},t) = -\frac{c}{4\pi} \int dt' \, \frac{\delta\left[|\mathbf{x} - \xi(t')| - c(t-t')\right]}{t-t'} \frac{d\xi}{dt}$$
(3.4)