

# PHYS 502: Mathematical Physics II

## Winter 2015, Homework #4

(Due February 25, 2015)

1. (a) Show explicitly from the series solutions that

$$\begin{aligned}J_{1/2}(x) &= Ax^{-1/2} \sin x \\J_{-1/2}(x) &= Bx^{-1/2} \cos x.\end{aligned}$$

Hence, taking  $A = B = 1$  and using the recurrence relations, write down expressions for  $J_{3/2}(x)$  and  $J_{5/2}(x)$ .

- (b) A function  $f(x)$  is expressed as a Bessel series

$$f(x) = \sum_{n=1}^{\infty} a_n J_m(\alpha_{mn} x),$$

where  $\alpha_{mn}$  is the  $n$ -th root of  $J_m$ . Prove the Parseval relation

$$\int_0^1 [f(x)]^2 x dx = \frac{1}{2} \sum_{n=1}^{\infty} a_n^2 [J_{m+1}(\alpha_{mn})]^2.$$

2. The curved surface of a long cylinder of radius  $b$  is kept at a constant temperature  $T = 0$ . Initially the cylinder is at a uniform temperature  $T_0 > 0$ . Derive an expression for the temperature at the center of the cylinder at any time  $t > 0$ , and write down a simplified solution (not  $T = 0$ !) valid in the limit  $t \gg b^2/\kappa$ , where  $\kappa$  is the heat diffusion coefficient of the cylinder.
3. (a) Two hemispherical shells each of radius  $a$  are fitted together, insulated around their circle of contact, and kept at potentials  $\pm V_0$ , respectively. Find the potential  $\Phi(r, \theta)$  inside the resulting sphere, where  $\nabla^2 \Phi = 0$  inside the sphere,  $\Phi = \pm V_0$  on the two hemispheres, and the polar axis  $\theta = 0$  is the axis of symmetry.
- (b) Now suppose that the potentials of the hemispheres in part (a) oscillate in time, with  $V(t) = \pm V_0 e^{-i\omega t}$ . Find the *exterior* solution  $\Phi(r, \theta, t)$  to the wave equation

$$\nabla^2 \Phi = \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2}$$

for  $r > a$ , with  $\Phi = \pm V(t)$  on the hemispheres, that describes an *outgoing* wave as  $r \rightarrow \infty$ .

[Recall that the asymptotic forms of the spherical Bessel functions  $j_l(x)$  and  $n_l(x)$  as  $x \rightarrow \infty$  are  $j_l(x) \sim \frac{1}{x} \cos\{x - \frac{\pi}{2}(l+1)\}$  and  $n_l(x) \sim \frac{1}{x} \sin\{x - \frac{\pi}{2}(l+1)\}$ .]

4. Each of the two 1S electrons in a helium atom may be described by a hydrogenic wave function

$$\psi(\mathbf{r}) = \left( \frac{8}{\pi a_0^3} \right)^{1/2} e^{-2r/a_0}$$

in the absence of the other electron. Here,  $a_0 = \hbar^2/me^2$  is the Bohr radius. Use the expansion

$$\frac{1}{r_{12}} = \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{4\pi}{2n+1} [Y_n^m(\theta_1, \phi_1)]^* Y_n^m(\theta_2, \phi_2) \begin{cases} \frac{r_1^n}{r_2^{n+1}}, & |\mathbf{r}_1| < |\mathbf{r}_2| \\ \frac{r_2^n}{r_1^{n+1}}, & |\mathbf{r}_1| > |\mathbf{r}_2| \end{cases}$$

to find the mutual electrostatic potential energy of the two electrons

$$U = \int \psi^*(\mathbf{r}_1) \psi^*(\mathbf{r}_2) \frac{e^2}{r_{12}} \psi(\mathbf{r}_1) \psi(\mathbf{r}_2) d^3r_1 d^3r_2,$$

where  $r_i = |\mathbf{r}_i|$ ,  $d^3r_i = r_i^2 dr_i \sin \theta_i d\theta_i d\phi_i$  and  $r_{12} = |\mathbf{r}_1 - \mathbf{r}_2|$ .