

Statmech II HW6

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1 Problem 8.15

a. We prove the relation:

$$X = \frac{2n\mu^{*2}}{(\frac{\partial\mu_0}{\partial x})|_{x=1/2}} = \frac{n\mu^{*2}}{kT} \frac{f_{1/2}(z)}{f_{5/2}(z)} \quad (1.1)$$

We start with the relation:

$$\mu_0(xN) = kT \ln\left(\frac{xN\lambda^3}{V}\right) = kT \ln z \quad (1.2)$$

$$\frac{\partial\mu_0}{\partial x} = kT \frac{\partial \ln z}{\partial x} \quad (1.3)$$

With $f_{3/2}(z) \simeq z$ we identify $f_{3/2}(z) = \frac{xN\lambda^3}{V}$ we find an expression for $\frac{\partial \ln z}{\partial x}$:

$$\frac{\partial f_{3/2}(z)}{\partial \ln z} \frac{\partial \ln z}{\partial x} = \frac{N\lambda^3}{V} \quad (1.4)$$

$$\frac{\partial \ln z}{\partial x} = \frac{f_{3/2}(z)}{x f_{1/2}(z)} \quad (1.5)$$

Using 1.3 and 1.6 we can now write 1.1 as:

$$X = \frac{2n\mu^{*2}}{kT/x}|_{x=1/2} \frac{f_{1/2}(z)}{f_{3/2}(z)} \quad (1.6)$$

$$X = \frac{n\mu^{*2}}{kT} \frac{f_{1/2}(z)}{f_{3/2}(z)} \quad (1.7)$$

At high temperatures $z \ll 1$ and (keeping terms to first order in z):

$$\frac{f_{1/2}(z)}{f_{3/2}(z)} = \frac{z - z^2 2^{-1/2} + \dots}{z - z^2 2^{-1/2} + \dots} \quad (1.8)$$

$$\frac{f_{1/2}(z)}{f_{3/2}(z)} \simeq \frac{1 - z 2^{-1/2}}{1 - z 2^{-3/2}} \quad (1.9)$$

$$\frac{f_{1/2}(z)}{f_{3/2}(z)} \simeq 1 - z 2^{-3/2} \quad (1.10)$$

Where we have used the fact that $2^{-1/2} - 2^{-3/2} = 2^{-3/2}$. Using the high temperature expression $z = \frac{n\lambda^3}{2}$ we can write the susceptibility:

$$X = \frac{n\mu^{*2}}{kT} \left(1 - \frac{n\lambda^3}{2} 2^{-3/2} \right) \quad (1.11)$$

$$X = \frac{n\mu^{*2}}{kT} \left(1 - \frac{n\lambda^3}{2^{5/2}} \right) \quad (1.12)$$

With $X_0 \equiv \frac{n\mu^{*2}}{kT}$ we have proved the provided relation.

At low temperatures we use the Sommerfeld expansions of the Fermi integrals:

$$\frac{f_{1/2}(z)}{f_{3/2}(z)} = \frac{3}{2} \frac{1}{\ln z} \left(1 - \frac{\pi^2}{6} (\ln z)^{-2} + \dots \right) \quad (1.13)$$

Using the low-temperature approximation $\ln z = \frac{e_f}{kT} \left(1 - \frac{\pi^2}{12} \left(\frac{kT}{e_f} \right)^2 \right)$:

$$X = \frac{n\mu^{*2}}{kT} \frac{f_{1/2}(z)}{f_{3/2}(z)} \quad (1.14)$$

$$X = \frac{n\mu^{*2}}{kT} \frac{3}{2} \frac{kT}{e_f} \left(1 - \frac{\pi^2}{6} \left(\frac{kT}{e_f} \right)^{-2} + \dots \right) \left(1 + \frac{\pi^2}{12} \left(\frac{kT}{e_f} \right)^2 + \dots \right) \quad (1.15)$$

$$X = \frac{3n\mu^{*2}}{2e_f} \left(1 - \frac{\pi^2}{12} \left(\frac{kT}{e_f} \right)^{-2} + \dots \right) \quad (1.16)$$

2 Problem 8.19

$$e = mc^2 \left(\sqrt{1 + \left(\frac{p}{mc} \right)^2} - 1 \right) \quad (2.1)$$

$$p = mc \sqrt{\left(\frac{e}{mc^2} + 1 \right)^2 - 1} \quad (2.2)$$

$$p = \sqrt{\frac{e^2}{c^2} + 2me} \quad (2.3)$$

$$\frac{dp}{de} = \frac{1}{c} \left(\frac{e}{mc^2} + 1 \right) \left(\left(\frac{e}{mc^2} + 1 \right)^2 - 1 \right)^{-1/2} \quad (2.4)$$

$$p^2 \frac{dp}{de} = m^2 c \left(\frac{e}{mc^2} + 1 \right) \left(\left(\frac{e}{mc^2} + 1 \right)^2 - 1 \right)^{1/2} \quad (2.5)$$

3 Problem 7.3

Using the relation from note 6 in chapter 7 of Pathria:

$$\frac{g_{3/2}(z)}{g_{3/2}(1)} = \left(\frac{T_c}{T} \right)^{3/2} \quad (3.1)$$

Truncating Pathria D.9 two two terms and instering into 1:

$$g_{3/2}(e^{-a}) = \frac{\Gamma(-1/2)}{a^{-1/2}} + \xi(\frac{3}{2}) + \dots \quad (3.2)$$

$$g_{3/2}(e^{-a}) = \xi(\frac{3}{2}) - 2\sqrt{\pi}a^{1/2} \quad (3.3)$$

$$1 - \frac{2\sqrt{\pi}a^{1/2}}{\xi(\frac{3}{2})} = \left(\frac{T}{T_c}\right)^{3/2} \quad (3.4)$$

$$a^{1/2} = \frac{\xi(\frac{3}{2})}{2\sqrt{\pi}} \left(1 - \left(\frac{T}{T_c}\right)^{3/2}\right) \quad (3.5)$$

Now make a Taylor expansion of $1 - \left(\frac{T}{T_c}\right)^{3/2}$.

$$1 - \left(\frac{T}{T_c}\right)^{3/2} \simeq (1 - 1) - \frac{3}{2} \frac{1}{T_c^{3/2}} T_c^{1/2} (T - T_c) \quad (3.6)$$

$$1 - \left(\frac{T}{T_c}\right)^{3/2} \simeq -\frac{3}{2} \frac{T - T_c}{T_c} \quad (3.7)$$

Inserting the approximation into 5 and squaring both sides:

$$a \simeq \frac{1}{\pi} \left(\frac{3\xi(3/2)}{4}\right)^2 \left(\frac{T - T_c}{T_c}\right)^2 \quad (3.8)$$

4 Problem 7.5

a) We prove the following relations for the isothermal compressibility and adiabatic compressibility of an ideal Bose gas:

$$\kappa_T = \frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_T \quad (4.1)$$

$$\kappa_S = \frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_S \quad (4.2)$$

Using the expressions for P and n and taking the derivatives with T held constant and with z held constant:

$$P = kT \frac{1}{\lambda^3} g_{5/2}(z) \quad (4.3)$$

$$\frac{\partial P}{\partial z} = \frac{kT}{\lambda^3} \frac{1}{z} g_{3/2}(z) \quad (4.4)$$

$$\frac{\partial P}{\partial T} = \frac{5}{2} \frac{(2\pi m)^{3/2}}{h^3} k^{5/2} T^{3/2} g_{5/2}(z) \quad (4.5)$$

$$n = \frac{1}{\lambda^3} g_{3/2}(z) \quad (4.6)$$

$$\frac{\partial n}{\partial z} = \frac{1}{\lambda^3} \frac{1}{z} g_{1/2}(z) \quad (4.7)$$

$$\frac{\partial n}{\partial T} = \frac{3}{2} \frac{(2\pi m k)^{3/2}}{h^3} T^{1/2} g_{3/2}(z) \quad (4.8)$$

Writing the compressibility expressions as functions of n and using the appropriate derivatives, we show:

$$\kappa_T = \frac{V}{N} \left(\frac{\partial(\frac{N}{V})}{\partial P} \right)_T = \frac{1}{n} \left(\frac{\partial n}{\partial P} \right)_T \quad (4.9)$$

$$\kappa_T = \frac{1}{nkT} \frac{g_{1/2}(z)}{g_{3/2}(z)} \quad (4.10)$$

$$\kappa_S = \frac{V}{N} \left(\frac{\partial(\frac{N}{V})}{\partial P} \right)_S = \frac{1}{n} \left(\frac{\partial n}{\partial P} \right)_S \quad (4.11)$$

$$\kappa_S = \frac{3}{5nkT} \frac{g_{3/2}(z)}{g_{5/2}(z)} \quad (4.12)$$

b) We now derive the relations:

$$\gamma = \frac{C_p}{C_v} = 1 + \frac{4}{9} \frac{C_v}{Nk} \frac{g_{1/2}(z)}{g_{3/2}(z)} \quad (4.13)$$

$$= \frac{5}{3} \frac{g_{5/2}(z) g_{1/2}(z)}{(g_{3/2}(z))^2} \quad (4.14)$$

We first note that $\frac{C_p - C_v}{C_v} = \gamma - 1$, so $\gamma = 1 + \frac{C_p - C_v}{C_v}$. We calculate $\frac{\partial P}{\partial T}|_V$:

$$P = \frac{2U}{3V} \quad (4.15)$$

$$\left(\frac{\partial P}{\partial T} \right)|_V = \frac{2}{3V} \left(\frac{\partial U}{\partial T} \right)_V \quad (4.16)$$

$$\left(\frac{\partial P}{\partial T} \right)|_V = \frac{2}{3V} C_V \quad (4.17)$$

Using the provided relation for $C_P - C_V$ and plugging in the expression for $\left(\frac{\partial P}{\partial T}\right)|_V$:

$$\frac{C_P - C_V}{C_V} = \frac{4T}{9V} C_V \frac{1}{nkT} \frac{g_{1/2}(z)}{g_{3/2}(z)} \quad (4.18)$$

$$\frac{C_P - C_V}{C_V} = \frac{4}{9} \frac{C_V}{Nk} \frac{g_{1/2}(z)}{g_{3/2}(z)} \quad (4.19)$$

$$\gamma = 1 + \frac{4}{9} \frac{C_V}{Nk} \frac{g_{1/2}(z)}{g_{3/2}(z)} \quad (4.20)$$

Using the relation $\frac{C_P}{C_V} = \frac{\kappa_T}{\kappa_S}$ and substituting our previous values for the compressibilities:

$$\gamma = \frac{C_P}{C_V} = \frac{5}{3} \frac{g_{5/2}(z)g_{1/2}(z)}{(g_{3/2}(z))^2} \quad (4.21)$$