

Statmech II HW6

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1 Problem 8.9

We calculate the second-order approximation to the specific heat of a Fermi gas using Sommerfeld's expansion of the Fermi integrals in powers of $\ln z$. To carry the whole solution to second-order we first extend the approximation of $\ln z$ derived from N/V , substituting ξ for $\ln z$:

$$\frac{N}{V} = \frac{h}{\lambda^3} f_{3/2}(z) \quad (1.1)$$

$$\frac{N}{V} \simeq \frac{4g}{3\sqrt{\pi}\lambda^3} (\xi)^{3/2} \left(1 + \frac{\pi^2}{8} \xi^{-2} - \frac{7\pi^4}{640} \xi^{-4} \right) \quad (1.2)$$

$$e_f = \left(\frac{3}{4\pi g} \right)^{2/3} \left(\frac{N}{V} \right)^{2/3} \frac{\hbar^2}{2m} \quad (1.3)$$

$$e_f \simeq \frac{\lambda^2}{\pi} \frac{\hbar^2}{2m} \xi \left(1 + \frac{\pi^2}{8} \xi^{-2} - \frac{7\pi^4}{640} \xi^{-4} \right)^{2/3} \quad (1.4)$$

Using a math package to help calculate the 2/3 power of the series we find:

$$e_f \simeq kT \ln z \left(1 + \frac{\pi^2}{8} (\ln z)^{-2} - \frac{\pi^4}{180} (\ln z)^{-4} \right) \quad (1.5)$$

We now approximate U/N :

$$\frac{U}{N} = \frac{3}{2} T \frac{f_{5/2}(z)}{f_{3/2}(z)} \quad (1.6)$$

$$f_{5/2}(z) \simeq \frac{8}{15\sqrt{\pi}} (\ln z)^{5/2} \left(1 + \frac{5\pi^2}{8} (\ln z)^{-2} - \frac{7\pi^4}{384} (\ln z)^{-4} \right) \quad (1.7)$$

$$\frac{U}{N} = \frac{3}{5} kT \ln z \left(1 + \frac{\pi^2}{2} (\ln z)^{-2} - \frac{11\pi^4}{120} (\ln z)^{-4} \right) \quad (1.8)$$

Using the provided result for $\mu = kT \ln z$:

$$\frac{U}{N} = \frac{3}{5} e_f \left(1 + \frac{\pi^2}{2} \left(\frac{kT}{e_f} \right)^2 - \frac{11\pi^4}{120} \left(\frac{kT}{e_f} \right)^4 \right) \times \left(1 - \frac{\pi^2}{12} \left(\frac{kT}{e_f} \right)^2 - \frac{\pi^4}{80} \left(\frac{kT}{e_f} \right)^4 \right) \quad (1.9)$$

$$\frac{U}{N} = \frac{3}{5} e_f \left(1 + \frac{5\pi^2}{12} \left(\frac{kT}{e_f} \right)^2 - \frac{7\pi^4}{48} \left(\frac{kT}{e_f} \right)^4 \right) \quad (1.10)$$

From the temperature-dependent part we find:

$$\frac{C_V}{Nk} = \frac{\pi^2}{2} \frac{kT}{e_f} - \frac{7\pi^4}{20} \left(\frac{kT}{e_f} \right)^3 \quad (1.11)$$

2 Problem 8.10

We first find the density of energy states for a system of dimension n with $e = p^s$:

$$a(e) = \frac{V}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)} p^{n-1} dp \quad (2.1)$$

$$p = e^{1/s} \quad (2.2)$$

$$p^{n-1} = e^{(n-1)/s} \quad (2.3)$$

$$dp = \frac{1}{s} e^{\frac{1}{s}-1} de \quad (2.4)$$

$$a(e) = \frac{V}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)} e^{(n-1)/s} \frac{1}{s} e^{\frac{1}{s}-1} de \quad (2.5)$$

$$a(e) = \frac{V}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} e^{(n/s)-1} de \quad (2.6)$$

We can now do the usual integrals for P/kT and U/N .

$$\frac{P}{kT} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} \int_0^\infty \ln(1 + ze^{-\beta e}) e^{(n/s)-1} de \quad (2.7)$$

$$\frac{P}{kT} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} (kT)^{n/s} \int_0^\infty \ln(1 + ze^{-x}) x^{(n/s)-1} dx \quad (2.8)$$

$$\frac{P}{kT} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} (kT)^{n/s} \Gamma(n/s) f_{n/s+1}(z) \quad (2.9)$$

$$\frac{N}{V} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} \int_0^\infty \frac{e^{(n/s)-1}}{z^{-1}e^{\beta e} + 1} de \quad (2.10)$$

$$\frac{N}{V} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} (kT)^{n/s} \int_0^\infty \frac{x^{(n/s)-1}}{z^{-1}e^x + 1} dx \quad (2.11)$$

$$\frac{N}{V} = \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} (kT)^{n/s} \Gamma(n/s) f_{n/s}(z) \quad (2.12)$$

Taking $U = kT^2 \left\{ \frac{\partial}{\partial T} (PV/kT) \right\}_{z,V}$ we have:

$$U = \frac{n}{s} V \frac{1}{h^n} \frac{2\pi^{n/2}}{\Gamma(n/2)s} \Gamma(n/s) (kT)^{n/s} f_{(n/s)+1}(z) \quad (2.13)$$

Comparing (13) and (9) we have shown that $U = \frac{n}{s} PV$ as required.
From (13) we have:

$$U = \frac{n}{s} N k T \frac{f_{(n/s)+1}(z)}{f_{n/s}(z)} \quad (2.14)$$

We will the derivative of z wrt T:

$$\frac{1}{z} \frac{\partial z}{\partial T} = -\frac{n}{s} \frac{1}{T} \frac{f_{n/s}(z)}{f_{(n/s)-1}(z)} \quad (2.15)$$

3 Problem 8.13