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February 9, 2015

Abstract

Quantum 1 HW 3

1 Problem 1

To calculate the Compton effect we look at relativistic scattering of a photon from an electron following notes from Dr. McCray's Classical Dynamics course. We use both conservation of energy and conservation of momentum. We define the momentum of the incident photon as p_i and the momentum of the scattered photon as p_s . The electron is initially at rest, and after scattering has momentum p_e . Prior to scattering, the photon has energy $p_i c$ and the electron has energy $m_e c^2$. After scattering the scattered photon and electron obey conservation of energy.

$$p_i c + m_e c^2 = p_s c + \sqrt{(p_e^2 c^2 + m_e^2 c^4)} \quad (1.1)$$

$$p_e^2 c^2 = (p_i - p_s)^2 c^2 + 2c(p_i - p_e)m_e c^2 \quad (1.2)$$

$$p_e^2 = (p_i - p_s)^2 + 2c(p_i - p_s)m_e \quad (1.3)$$

We also use conservation of momentum. The electron initially has no momentum. The vector sum of the scattered photon momentum and the electron momentum must equal the original photon momentum. Using the law of cosines with θ the scattering angle of the photon:

$$p_e^2 = p_i^2 + p_s^2 - 2p_i p_s \cos(\theta) \quad (1.4)$$

We can now use 1.3 and 1.4 together:

$$p_i^2 + p_s^2 - 2p_i p_s \cos(\theta) = (p_i - p_s)^2 + 2c(p_i - p_s)m_e = p_i^2 - 2p_i p_s + p_s^2 + 2cm_e(p_i - p_s) \quad (1.5)$$

$$-2p_i p_s \cos(\theta) = -2p_i p_s + 2cm_e(p_i - p_s) \quad (1.6)$$

$$p_i p_s (1 - \cos(\theta)) = cm_e(p_i - p_s) \quad (1.7)$$

The photon momentum is equal to $\frac{\hbar}{\lambda}$, so we can write:

$$\frac{\hbar^2}{\lambda_i \lambda_s} (1 - \cos(\theta)) = m_e c \hbar \left(\frac{1}{\lambda_i} - \frac{1}{\lambda_s} \right) \quad (1.8)$$

$$\lambda_s - \lambda_i = \frac{\hbar}{m_e c} (1 - \cos(\theta)) \quad (1.9)$$

A photon that scatters from an electron will see an average change in wavelength of $\frac{\hbar}{m_e c}$, the Compton wavelength of the electron.

2 Problem 2

The energy to assemble a charge distribution is:

$$W = \frac{1}{2} \int \sigma V da \quad (2.1)$$

The potential at the surface of the uniformly charged sphere is $\frac{1}{4\pi\epsilon_0} \frac{-e}{R}$.

$$W = \frac{1}{8\pi\epsilon_0} \frac{-e}{R} \int \sigma da = \frac{1}{8\pi\epsilon_0} \frac{e^2}{R} \quad (2.2)$$

We now equate this energy to the rest mass of the electron and calculate the classical electron radius.

$$\frac{1}{8\pi\epsilon_0} \frac{e^2}{R} = m_e c^2 \quad (2.3)$$

$$R_e = \frac{1}{8\pi\epsilon_0} \frac{e^2}{m_e c^2} \quad (2.4)$$

3 Problem 3

Trying simple combinations of two of the three physical constants, we have:

$$QM \bigcup EM : \frac{\hbar^2}{e^2} = ML \quad (3.1)$$

$$QM \bigcup SR : \frac{\hbar}{c} = ML \quad (3.2)$$

$$EM \bigcup SR : \frac{e^2}{c^2} = ML \quad (3.3)$$

These are the Bohr radius of hydrogen, the Compton wavelength of the electron, and the classical radius of the electron. Seeking a dimensionless combination of the fundamental constants we find the fine structure constant:

$$\alpha = \frac{e^2}{\hbar c} \text{ (dimensionless)} \quad (3.4)$$

Multiplying each length scale by the fine structure constant moves to the next lower length scale.

4 Problem 4

The Schrodinger equation of a fixed rotor with angular momentum I_0 is:

$$\frac{-\hbar^2}{2I_0} \frac{d^2 \Psi}{d\phi^2} = E \Psi \quad (4.1)$$

$$\frac{d^2 \Psi}{d\phi^2} + \frac{2I_0 E}{\hbar^2} \Psi = 0 \quad (4.2)$$

The solutions are $\Psi_m = e^{im\phi}$, $m = 0, \pm 1, \pm 2, \dots$, where the restriction to integer values of m comes from the requirement that $\Psi(\phi)$ be single-valued. The separation constant $m^2 = \frac{2I_0 E}{\hbar^2}$ allows us to calculate the energy levels:

$$E_m = \frac{m^2 \hbar^2}{2I_0}, \quad m = 0, \pm 1, \pm 2, \dots \quad (4.3)$$

$E_m \leq 100 \frac{\hbar^2}{2I_0}$ when $|m| \leq 10$. There are therefore 21 states with $E_m \leq 100 \frac{\hbar^2}{2I_0}$.

With a variable axis we are now working with a full spherical coordinate system. The Schrodinger equation is:

$$\frac{-\hbar^2}{2I_0} \nabla^2 \Psi = E \Psi \quad (4.4)$$

$$\nabla^2 \Psi + \frac{2I_0 E}{\hbar^2} \Psi = 0 \quad (4.5)$$

With a fixed rotor length we can ignore the radial part of this solution. The angular part of the solution is a spherical harmonic $Y_\ell^m(\theta, \phi)$. This will consist of an exponential part in ϕ and an associated Legendre function in θ . The exponential part gives rise to the constraint that m is an integer.

The angular part will be equal to a separation constant, which we set to $-\ell(\ell+1)$ by convention. We can then use the separation constant to calculate the energy levels.

$$\nabla^2 Y_\ell^m(\theta, \phi) = -\ell(\ell+1) Y_\ell^m(\theta, \phi) \quad (4.6)$$

$$E = \frac{\ell(\ell+1) \hbar^2}{2I_0} \quad (4.7)$$

We apparently recognize that this is $(2\ell+1)$ -fold degenerate. The maximum value of ℓ that satisfies $E_\ell \leq 100 \frac{\hbar^2}{2I_0}$ is 9. With 9 energy levels and $(2\ell+1)$ -fold degeneracy there are 99 states with energy $E_\ell \leq 100 \frac{\hbar^2}{2I_0}$.

5 Problem 5

We put the energy levels of the fixed rotor (4.3) into the partition function.

$$\zeta(\beta) = \sum_{m=-\infty}^{+\infty} e^{-m^2 \beta \frac{\hbar^2}{2I_0}}, \quad \beta = \frac{1}{kT} \quad (5.1)$$

$$= 1 + 2 \sum_{m=1}^{+\infty} e^{-m^2 \beta \frac{\hbar^2}{2I_0}} \quad (5.2)$$

$$= 1 + 2 \sum_{m=1}^{+\infty} (e^{-\beta \frac{\hbar^2}{2I_0}})^{m^2} \quad (5.3)$$

With $\frac{\hbar^2}{2I_0} \ll kT$, $\frac{\hbar^2}{2I_0} \gg \beta$. This guarantees the term $e^{-\beta \frac{\hbar^2}{2I_0}}$ will be small enough to quickly go to 0 for $m > 1$. Therefore, $\zeta(\beta) \approx 1 + 2e^{-\beta \frac{\hbar^2}{2I_0}}$.

We put the energy of the rotor with variable axis(4.7) into the partition function.

$$\zeta(\beta) = \sum_{\ell=0}^{+\infty} e^{-\ell(\ell+1)\beta \frac{\hbar^2}{2I_0}}, \quad \beta = \frac{1}{kT} \quad (5.4)$$

$$= 1 + \sum_{\ell=1}^{+\infty} e^{-\ell(\ell+1)\beta \frac{\hbar^2}{2I_0}} \quad (5.5)$$

$$= 1 + \sum_{\ell=1}^{+\infty} (e^{-\beta \frac{\hbar^2}{2I_0}})^{\ell(\ell+1)} \quad (5.6)$$

Using the same approximation as above, we find $\zeta(\beta) \approx 1 + e^{-3\beta \frac{\hbar^2}{2I_0}}$.

6 Problem 6

We find the Dunham coefficients Y_{ij} of NaCl^1 :

Table 1: Dunham coefficients of NaCl

Constant	Na^{35}Cl	Na^{37}Cl
Y_{10}	364.68	360.75
Y_{20}	-1.7761	-1.7366
$Y_{30} \times 10^3$	5.9369	5.452(1)
$Y_{40} \times 10^5$	-1.231(2)	-
Y_{01}	0.21806	0.21339
$Y_{11} \times 10^3$	-1.6248	-1.5716
$Y_{21} \times 10^6$	5.1543	4.9913
$Y_{02} \times 10^7$	-3.1190	-2.9520
$Y_{12} \times 10^{10}$	6.397(2)	4.019(5)
$Y_{22} \times 10^{12}$	6.397(21)	-
$Y_{03} \times 10^{14}$	8.96(29)	-

These are coefficients of the Dunham expansion of the total energy of a rotating vibrator. The Dunham expansion provides an approximation that incorporates both the translational and rotational energy.

$$E(v, J) = \sum_{k, \ell} Y_{k, \ell} \left(v + \frac{1}{2}\right)^k [J(J+1)]^\ell \quad (6.1)$$

Where v and J are the quantum rotational and translational numbers and the $Y_{k, \ell}$ are the Dunham coefficients.

¹R. S. Ram, M. Dulick, B. Guo, K.-Q. Zhang, and P. F. Bernath, Journal of Molecular Spectroscopy, **183**, 360-373