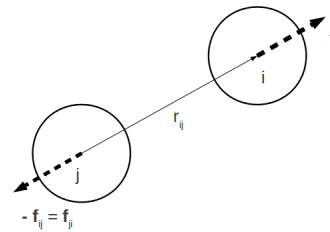


## “Tricks” of the Trade

- *Efficiency* is the word
- But the straightforward implementation of algorithms to solve eq. of motion may not be the most efficient way
- There are ways of skipping some calculations and still get the “right” answer
- Force calculations are expensive: can we avoid parts of it, square roots? (show bench)

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## Force calculation shortcuts



- Newton's 3rd: once calculate  $F_{ij}$ , have  $F_{ji}$
- If  $V=V(r_{ij})$  is an even power, then because

$$F_{ij} = -\frac{1}{r_{ij}} \left( \frac{dV(r_{ij})}{dr_{ij}} \right) r_{ij}$$

only need  $r_{ij}^2$  for  $F$ , thus no sqrt()

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## Potential Interpolation

One way to not calculate a complicated potential at every step is to use interpolations:

$$\text{e.g. } V(r_{ij}) = A \exp(-C \cdot r_{ij}) - \frac{B}{r_{ij}^6} \quad \text{Barker, et.al. 1971 For Argon}$$

- Use tables to evaluate the potential explicitly for only a small number of distance values
- For arbitrary distances, interpolate the potential from bracketing values in the table

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## Potential Interpolation

Newton-Gregory forward difference method:

- Use  $s=r_{ij}^2$  and calculate  $V_1 = V(s_1)$ ,  $V_2 = V(s_2)$ , ... at equal intervals  $\delta s$ .
- Define 1<sup>st</sup> and 2<sup>nd</sup> differences:
  - $\delta V_k = V_{k+1} - V_k$
  - $\delta^2 V_k = \delta V_{k+1} - \delta V_k$
- For a value  $s$  between  $s_k$  and  $s_{k+1}$ , *interpolate*

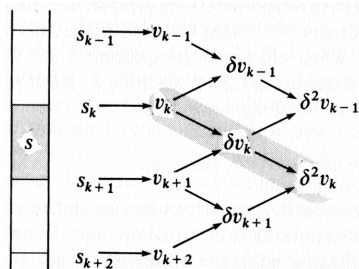
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## Potential Interpolation

$$V(s) \approx V_k + \xi \delta V_k + \frac{1}{2} \xi (\xi - 1) \delta^2 V_k$$

- where

$$\xi = (s - s_k) / \delta s$$



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## Potential Interpolation

For the force, using

$$f_{ij} = -\frac{1}{r_{ij}} \left( \frac{dV(r_{ij})}{dr_{ij}} \right) r_{ij} = -\frac{w(r_{ij})}{r_{ij}^2} r_{ij}$$

and noting that

$$\frac{w(r_{ij}^2)}{r_{ij}^2} = \frac{w(s)}{s} = 2 \frac{dV}{ds}$$

Thus, for  $f$  can just differentiate

$$V(s) \approx V_k + \xi \delta V_k + \frac{1}{2} \xi (\xi - 1) \delta^2 V_k$$

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## Potential Interpolation

- An alternate method. For each interval ( $s_k, s_{k+1}$ ) represent potential by 5<sup>th</sup> order polynomial:

$$V(s) \approx c_0 + c_1 \delta s + c_2 \delta s^2 + c_3 \delta s^3 + c_4 \delta s^4 + c_5 \delta s^5$$

- Where  $\delta s = s - s_k$ .
- $c_i$  coeffs are determined by the exact values of  $V(s)$ ,  $dV(s)/ds$ , and  $d^2V(s)/ds^2$  evaluated at the two ends of the interval.
- Advantage: the  $s_k$  need not be evenly spaced.

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## Not only by algorithms...

Compiler optimizations can help:

- O3 (illustrate)

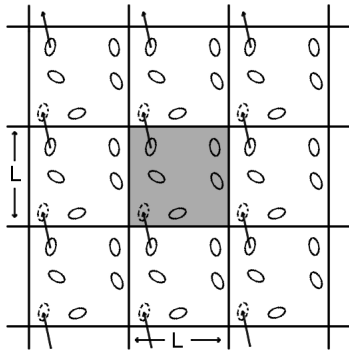
Hardware compiler type and version:

- commercial

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## Problems with PBC: time to pay...

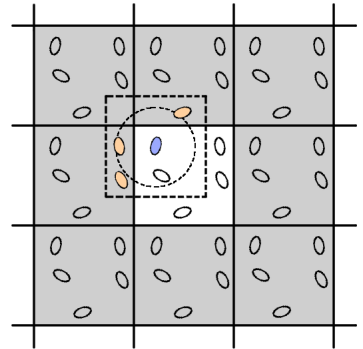
- PBC introduces the problem that all image pair interactions should be considered (infinite).



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## PBC: still hope...

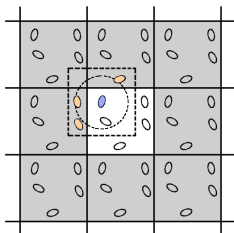
- For short-ranged interactions could possibly skip far-away interactions.
- Approximate to only consider closest real or image particles: minimum image convention.
- Pairwise interactions only require  $N(N-1)/2$  terms.



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## Avoid $N^2$ calculation with $r_c$

- Minimum Image convention still  $O(N^2)$
- What if for short-ranged interactions could possibly skip far-away interactions – **do less!**
- WARNING: This introduces an error and discontinuity in force and potential at cutoff  $r_c$
- Thus, energy will not be conserved for truncated interactions

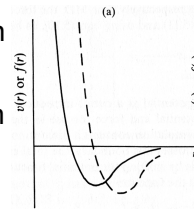


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## Avoid the $N^2$ calculation

Solution:

- Truncate interaction at  $r_c$
- Shift potential by an amount  $V_c = V(r_c)$



$$V^T(r_{ij}) = \begin{cases} V(r_{ij}) - V_c & r_{ij} \leq r_c \\ 0 & r_{ij} > r_c \end{cases}$$

Corresponds to (a)

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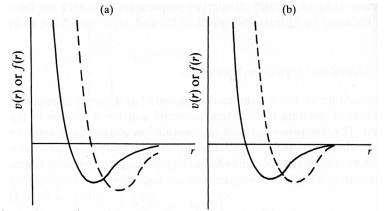
## Avoid the $N^2$ calculation

- Additional term  $V_c$  is constant and does not affect force calculation nor eq. of motion
- However, contribution to total energy varies from step to step. Have to account for it in the energy
- Force is still discontinuous at  $r=r_c$
- Solution: add yet another term

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## Avoid the $N^2$ calculation

- Additional term is linear such that derivative is zero at cutoff distance
- This is the 'shifted-force potential'



$$V^T(r_{ij}) = \begin{cases} V(r_{ij}) - V_c - \left( \frac{dV(r_{ij})}{dr_{ij}} \right)_{r_{ij}=r_c} (r_{ij} - r_c) & r_{ij} \leq r_c \\ 0 & r_{ij} > r_c \end{cases}$$

Corresponds to (b)

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## Avoid the $N^2$ calculation

### Caveats:

- Discontinuity now shifts to the gradient of the force.
- The 'shifted-force potential' does not correspond anymore to desired model potential.
- However, thermodynamics are still very similar to original problem.

### Alternate Route:

- Introduce a “switching” function to smoothly taper potential to zero at large  $r$ .

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## In sum: on reducing the distance “checking”

### Can do:

- Reduced use of expensive functions
- Could use potential interpolations
- Get good compiler and use optimization flags
- PBC with minimum image convention
- Add distance cutoff  $r_c$  and correct energy

### BUT:

- Still have to check (but not compute) all pairs -  $O(N^2)$

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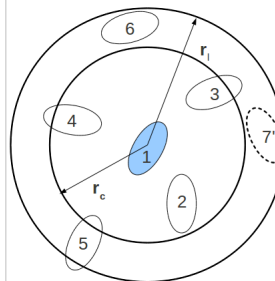
## Neighbor Lists

### Solution:

- Again rely on short-range interactions and keep a list of only neighbors
- Only update this list occasionally
- Between updates, calculate interactions with all neighbors in the list
- Can do:
  - Verlet neighbor list
  - Cell structures

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## Verlet Neighbor List

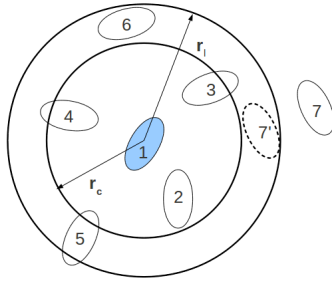


- Surround  $r_c$  by a “skin” of  $r_i$
- Loop over all particles inside  $r_i$  (Note, more interactions than with pure  $r_c$ )
- But number of pairs  $\sim O(N)$

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## Verlet Neighbor List

- Once in a while, update list (depends on size of  $r_l$ )
- $r_l$  should be big enough so as to prevent particle 7 into penetrating within  $r_c$  in between updates
- 10-20 updates typical, proportional to  $r_l$
- Can do automatic updates



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## Verlet Neighbor List

### Caveats:

- As the size of the system increases, the total size of the neighbor lists also increases thus affecting storage.
- Update (of the lists) is still  $N^2$   
[Show bench](#)
- Next: Use alternative method of cell structures.

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