Statmech II HW1

Vince Baker

October 4, 2015

1 Problem 1

We write the wavefunctions corresponding to the two polarizations as (1,0), (0,1). In this basis we can write the density matrix as:

$$\rho = \begin{bmatrix} f & 0 \\ 0 & 1 - f \end{bmatrix}
\tag{1.1}$$

2 Problem 2

Since the $|n\rangle$ are a complete orthonormal basis, each diagonal element of $\rho\hat{A}$ is:

$$(\rho \hat{A})_n = e^{-\langle n \rangle} \frac{\langle n \rangle^n}{n!} \langle n | \hat{A} | n \rangle$$
 (2.1)

a) For $\hat{a} | n \rangle = \sqrt{n} | n - 1 \rangle$, each term contains a dot product $\langle n | n - 1 \rangle = 0$ because the eigenvectors are orthogonal, so $\langle \hat{a} \rangle = 0$. For $\langle a^{\dagger} a \rangle$ the diagonal terms are:

$$(\rho a^{\dagger} a)_n = e^{-\langle n \rangle} \frac{\langle n \rangle^n}{n!} \langle n | n | n \rangle$$
 (2.2)

$$(\rho a^{\dagger} a)_n = e^{-\langle n \rangle} \frac{\langle n \rangle^n}{(n-1)!}$$
 (2.3)

We and pull out one $\langle n \rangle$ from the numerator and use the series definition of e:

$$Tr(\rho a^{\dagger} a) = \langle n \rangle e^{-\langle n \rangle} \sum_{j=0}^{\infty} \frac{\langle n \rangle^{j-1}}{(j-1)!}$$
 (2.4)

$$Tr(\rho a^{\dagger} a) = \langle n \rangle e^{-\langle n \rangle} e^{\langle n \rangle}$$
 (2.5)

$$Tr(\rho a^{\dagger}a) = \langle n \rangle$$
 (2.6)

b) To find $\langle (a^{\dagger}a)^2 \rangle$ we expand the exponential in its series:

$$(\rho a^{\dagger} a)_n = e^{-\langle n \rangle} \frac{\langle n \rangle^n}{n!} \langle n | n^2 | n \rangle$$
 (2.7)

$$(\rho(a^{\dagger}a)^2)_n = e^{-\langle n \rangle} \frac{n \langle n \rangle^n}{(n-1)!}$$
(2.8)

$$Tr((a^{\dagger}a)^2) = \langle n \rangle e^{-\langle n \rangle} \sum_{j=0}^{\infty} \frac{j \langle n \rangle^{j-1}}{(j-1)!}$$
 (2.9)

$$Tr((a^{\dagger}a)^{2}) = \langle n \rangle \left(\sum_{i=0}^{\infty} \frac{\langle n \rangle^{i}}{i!} \right)^{-1} \sum_{j=0}^{\infty} \frac{j \langle n \rangle^{j-1}}{(j-1)!}$$
 (2.10)

$$Tr((a^{\dagger}a)^{2}) = \frac{\langle n \rangle}{\langle n \rangle} \left(\sum_{i=0}^{\infty} \frac{i \langle n \rangle^{i-1}}{(i-1)!} \right)^{-1} \sum_{i=0}^{\infty} \frac{j \langle n \rangle^{j-1}}{(j-1)!}$$
(2.11)

$$Tr((a^{\dagger}a)^2) = 1$$
 (2.12)

So we find that the variance is $1 - \langle n \rangle^2$.

3 Problem 3

a) The Hamiltonian for a electron with magnetic moment μ_B in a magnetic field B along the z axis is:

$$H = -\mu_B \mathbf{S} \cdot \mathbf{B} \tag{3.1}$$

$$H = -\mu_B B \sigma_z \tag{3.2}$$

The eigenvalues of σ_z are $\pm \frac{\hbar}{2}$ in the basis where σ_z is diagonalized. Defining $\kappa \equiv \mu_b B^{\frac{\hbar}{2}}$ we find the probabilities of the two states:

$$\rho_{\uparrow} = \frac{e^{\beta \kappa}}{e^{-\beta \kappa} + e^{\beta \kappa}} = \frac{1}{1 + e^{-2\beta \kappa}} \tag{3.3}$$

$$\rho_{\downarrow} = \frac{e^{-\beta\kappa}}{e^{-\beta\kappa} + e^{\beta\kappa}} = \frac{1}{1 + e^{2\beta\kappa}} \tag{3.4}$$

Knowing that the eigenvectors of σ_z are (1,0) and (0,1) we can directly write the density matrix:

$$\rho = \begin{bmatrix} \frac{1}{1 + e^{-2\beta\kappa}} & 0\\ 0 & \frac{1}{1 + e^{2\beta\kappa}} \end{bmatrix}$$
(3.5)

b) We compute $\langle \sigma_z \rangle$ as $Tr(\rho \sigma_z)$:

$$\langle \sigma_z \rangle = Tr(\rho \sigma_z) \tag{3.6}$$

$$\langle \sigma_z \rangle = Tr(\rho \sigma_z) \tag{3.6}$$

$$\langle \sigma_z \rangle = Tr \left(\begin{bmatrix} \frac{1}{1 + e^{-2\beta \kappa}} & 0\\ 0 & \frac{1}{1 + e^{2\beta \kappa}} \end{bmatrix} \begin{bmatrix} 1 & 0\\ 0 & -1 \end{bmatrix} \right) \tag{3.7}$$

$$\langle \sigma_z \rangle = \frac{e^{\beta \kappa} - e^{-\beta \kappa}}{e^{\beta \kappa} + e^{-\beta \kappa}} \tag{3.8}$$

$$\langle \sigma_z \rangle = \frac{e^{\beta \kappa} - e^{-\beta \kappa}}{e^{\beta \kappa} + e^{-\beta \kappa}} \tag{3.8}$$