

# Notes on the QED evolution to NLL accuracy

## Abstract

In these notes, we discuss the details of the implementation of the QED collinear evolution up to next-to-leading logarithmic (NLL) accuracy in the variable-flavour-number scheme (VFNS) in the presence of both charged leptons and quarks.

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## 1 The structure of the DGLAP equation

Suppose one wants to study the coupled collinear evolution in pure QED of the photon distribution  $\gamma$ , the  $i$ -th down-type quark distribution with a certain colour ( $rgb$ )  $d_i \in \{d^r, d^g, d^b, s^r, s^g, s^b, b^r, b^g, b^b\}$ , the  $k$ -th down-type anti-quark distribution  $\bar{d}_k \in \{\bar{d}^r, \bar{d}^g, \bar{d}^b, \bar{s}^r, \bar{s}^g, \bar{s}^b, \bar{b}^r, \bar{b}^g, \bar{b}^b\}$ , the  $j$ -th up-type quark distribution  $u_j \in \{u^r, u^g, u^b, c^r, c^g, c^b, t^r, t^g, t^b\}$ , the  $h$ -th up-type anti-quark distribution  $\bar{u}_h \in \{\bar{u}^r, \bar{u}^g, \bar{u}^b, \bar{c}^r, \bar{c}^g, \bar{c}^b, \bar{t}^r, \bar{t}^g, \bar{t}^b\}$ , the  $\alpha$ -th lepton distribution  $\ell_\alpha \in \{e^-, \mu^-, \tau^-\}$ , and the  $\beta$ -th anti-lepton distribution  $\bar{\ell}_\beta \in \{e^+, \mu^+, \tau^+\}$ . The most general form of the evolution equations reads:

$$\mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} \ell_\alpha \\ u_j \\ d_i \\ \gamma \\ \bar{d}_k \\ \bar{u}_h \\ \bar{\ell}_\beta \end{pmatrix} = \sum_{e,l,m,n,\gamma,\delta} \begin{pmatrix} \mathcal{P}_{\ell_\alpha \ell_\gamma} & \mathcal{P}_{\ell_\alpha u_e} & \mathcal{P}_{\ell_\alpha d_l} & \mathcal{P}_{\ell_\alpha \gamma} & \mathcal{P}_{\ell_\alpha \bar{d}_m} & \mathcal{P}_{\ell_\alpha \bar{u}_n} & \mathcal{P}_{\ell_\alpha \bar{\ell}_\delta} \\ \mathcal{P}_{u_j \ell_\gamma} & \mathcal{P}_{u_j u_e} & \mathcal{P}_{u_j d_l} & \mathcal{P}_{u_j \gamma} & \mathcal{P}_{u_j \bar{d}_m} & \mathcal{P}_{u_j \bar{u}_n} & \mathcal{P}_{u_j \bar{\ell}_\delta} \\ \mathcal{P}_{d_i \ell_\gamma} & \mathcal{P}_{d_i u_e} & \mathcal{P}_{d_i d_l} & \mathcal{P}_{d_i \gamma} & \mathcal{P}_{d_i \bar{d}_m} & \mathcal{P}_{d_i \bar{u}_n} & \mathcal{P}_{d_i \bar{\ell}_\delta} \\ \mathcal{P}_{\gamma \ell_\gamma} & \mathcal{P}_{\gamma u_e} & \mathcal{P}_{\gamma d_l} & \mathcal{P}_{\gamma \gamma} & \mathcal{P}_{\gamma \bar{d}_m} & \mathcal{P}_{\gamma \bar{u}_n} & \mathcal{P}_{\gamma \bar{\ell}_\delta} \\ \mathcal{P}_{\bar{d}_k \ell_\gamma} & \mathcal{P}_{\bar{d}_k u_e} & \mathcal{P}_{\bar{d}_k d_l} & \mathcal{P}_{\bar{d}_k \gamma} & \mathcal{P}_{\bar{d}_k \bar{d}_m} & \mathcal{P}_{\bar{d}_k \bar{u}_n} & \mathcal{P}_{\bar{d}_k \bar{\ell}_\delta} \\ \mathcal{P}_{\bar{u}_h \ell_\gamma} & \mathcal{P}_{\bar{u}_h u_e} & \mathcal{P}_{\bar{u}_h d_l} & \mathcal{P}_{\bar{u}_h \gamma} & \mathcal{P}_{\bar{u}_h \bar{d}_m} & \mathcal{P}_{\bar{u}_h \bar{u}_n} & \mathcal{P}_{\bar{u}_h \bar{\ell}_\delta} \\ \mathcal{P}_{\bar{\ell}_\beta \ell_\gamma} & \mathcal{P}_{\bar{\ell}_\beta u_e} & \mathcal{P}_{\bar{\ell}_\beta d_l} & \mathcal{P}_{\bar{\ell}_\beta \gamma} & \mathcal{P}_{\bar{\ell}_\beta \bar{d}_m} & \mathcal{P}_{\bar{\ell}_\beta \bar{u}_n} & \mathcal{P}_{\bar{\ell}_\beta \bar{\ell}_\delta} \end{pmatrix} \begin{pmatrix} \ell_\gamma \\ u_e \\ d_l \\ \gamma \\ \bar{d}_m \\ \bar{u}_n \\ \bar{\ell}_\delta \end{pmatrix}, \quad (1.1)$$

where the index  $l$  runs over the down-type quarks, the index  $m$  over the down-type anti-quarks, the index  $e$  runs over the up-type quarks, the index  $n$  over the up-type anti-quarks, the index  $\gamma$  over the leptons, and the index  $\delta$  over the anti-leptons. In addition, the Mellin convolution integral between the splitting-function matrix and the vector of distributions in the r.h.s. is understood. Despite in general all entries of the splitting-function matrix are different from zero, in massless QED one can identify the following equalities for the splitting functions that that involve a photon:

$$\begin{aligned} \mathcal{P}_{\ell_\alpha \gamma} &= \mathcal{P}_{\bar{\ell}_\beta \gamma} \equiv \mathcal{P}_{\ell \gamma}, \\ \mathcal{P}_{u_j \gamma} &= \mathcal{P}_{\bar{u}_h \gamma} \equiv \mathcal{P}_{u \gamma}, \\ \mathcal{P}_{d_i \gamma} &= \mathcal{P}_{\bar{d}_k \gamma} \equiv \mathcal{P}_{d \gamma}, \end{aligned} \quad (1.2)$$

and:

$$\begin{aligned} \mathcal{P}_{\gamma \ell_\gamma} &= \mathcal{P}_{\gamma \bar{\ell}_\delta} \equiv \mathcal{P}_{\gamma \ell}, \\ \mathcal{P}_{\gamma u_e} &= \mathcal{P}_{\gamma \bar{u}_n} \equiv \mathcal{P}_{\gamma u}, \\ \mathcal{P}_{\gamma d_l} &= \mathcal{P}_{\gamma \bar{d}_m} \equiv \mathcal{P}_{\gamma d}. \end{aligned} \quad (1.3)$$

The splitting functions that only involve fermions (quarks and/or leptons) instead obey the following decompositions:

$$\begin{aligned}
\mathcal{P}_{\ell_\alpha \ell_\gamma} &= \mathcal{P}_{\ell_\alpha \bar{\ell}_\gamma} \equiv \delta_{\alpha\gamma} \mathcal{P}_{\ell\ell}^V + \mathcal{P}_{\ell\ell}^S, \\
\mathcal{P}_{\ell_\beta \ell_\gamma} &\equiv \mathcal{P}_{\ell_\beta \bar{\ell}_\gamma} = \delta_{\beta\gamma} \mathcal{P}_{\ell\ell}^V + \mathcal{P}_{\ell\ell}^S, \\
\mathcal{P}_{u_j u_e} &= \mathcal{P}_{\bar{u}_j \bar{u}_e} \equiv \delta_{je} P_{uu}^V + P_{uu}^S, \\
\mathcal{P}_{\bar{u}_h u_e} &= \mathcal{P}_{u_h \bar{u}_e} \equiv \delta_{he} P_{u\bar{u}}^V + P_{u\bar{u}}^S, \\
\mathcal{P}_{d_i d_l} &= \mathcal{P}_{\bar{d}_i \bar{d}_l} \equiv \delta_{il} P_{dd}^V + P_{dd}^S, \\
\mathcal{P}_{\bar{d}_k d_l} &= \mathcal{P}_{d_k \bar{d}_l} \equiv \delta_{kl} P_{d\bar{d}}^V + P_{d\bar{d}}^S, \\
\mathcal{P}_{\ell_\alpha u_e} &= \mathcal{P}_{\ell_\alpha \bar{u}_e} \equiv \mathcal{P}_{\ell u}^S, \\
\mathcal{P}_{\ell_\alpha d_l} &= \mathcal{P}_{\ell_\alpha \bar{d}_l} \equiv \mathcal{P}_{\ell d}^S, \\
\mathcal{P}_{\ell_\alpha \bar{u}_e} &= \mathcal{P}_{\ell_\alpha u_e} \equiv \mathcal{P}_{\ell \bar{u}}^S, \\
\mathcal{P}_{\ell_\alpha \bar{d}_l} &= \mathcal{P}_{\ell_\alpha d_l} \equiv \mathcal{P}_{\ell \bar{d}}^S, \\
\mathcal{P}_{u_e \ell_\alpha} &= \mathcal{P}_{\bar{u}_e \bar{\ell}_\alpha} \equiv \mathcal{P}_{u\ell}^S, \\
\mathcal{P}_{d_l \ell_\alpha} &= \mathcal{P}_{\bar{d}_l \bar{\ell}_\alpha} \equiv \mathcal{P}_{d\ell}^S, \\
\mathcal{P}_{\bar{u}_e \ell_\alpha} &= \mathcal{P}_{u_e \bar{\ell}_\alpha} \equiv \mathcal{P}_{\bar{u}\ell}^S, \\
\mathcal{P}_{\bar{d}_l \ell_\alpha} &= \mathcal{P}_{d_l \bar{\ell}_\alpha} \equiv \mathcal{P}_{\bar{d}\ell}^S, \\
\mathcal{P}_{u_l d_k} &= \mathcal{P}_{d_l u_k} = \mathcal{P}_{\bar{u}_l \bar{d}_k} = \mathcal{P}_{\bar{d}_l \bar{u}_k} \equiv \mathcal{P}_{ud}^S, \\
\mathcal{P}_{u_l \bar{d}_k} &= \mathcal{P}_{\bar{d}_l u_k} = \mathcal{P}_{\bar{u}_l d_k} = \mathcal{P}_{d_l \bar{u}_k} \equiv \mathcal{P}_{u\bar{d}}^S.
\end{aligned} \tag{1.4}$$

Each one of splitting functions on the r.h.s. of the  $\equiv$  symbol above admits a perturbative expansion that truncated to next-to-leading order (as appropriate to achieve NLL accuracy at the level of the solution of the evolution equation), reads:

$$\mathcal{P} = \left(\frac{\alpha}{4\pi}\right) \mathcal{P}^{(0)} + \left(\frac{\alpha}{4\pi}\right)^2 \mathcal{P}^{(1)}. \tag{1.5}$$

The coefficients  $\mathcal{P}^{(0)}$  and  $\mathcal{P}^{(1)}$  undergo further simplifications. Let us start with the leading order one  $\mathcal{P}^{(0)}$ . At this order all of the pure-singlet contributions vanish, that is  $\mathcal{P}_{xy}^{(0),S} = 0$  for all pairs  $xy$ , and so do the valence fermion/anti-fermion splitting functions vanish, *i.e.*  $\mathcal{P}_{\ell\ell}^{(0),V} = \mathcal{P}_{u\bar{u}}^{(0),V} = \mathcal{P}_{d\bar{d}}^{(0),V} = 0$ . In addition, the splitting functions that survive obey the following equalities:

$$\begin{aligned}
\mathcal{P}_{\ell\gamma}^{(0)} &= e_\ell^2 P_{f\gamma}^{(0)}, \\
\mathcal{P}_{u\gamma}^{(0)} &= e_u^2 P_{f\gamma}^{(0)}, \\
\mathcal{P}_{d\gamma}^{(0)} &= e_d^2 P_{f\gamma}^{(0)}, \\
\mathcal{P}_{\gamma\ell}^{(0)} &= e_\ell^2 P_{\gamma f}^{(0)}, \\
\mathcal{P}_{\gamma u}^{(0)} &= e_u^2 P_{\gamma f}^{(0)}, \\
\mathcal{P}_{\gamma d}^{(0)} &= e_d^2 P_{\gamma f}^{(0)}, \\
\mathcal{P}_{\ell\ell}^{V,(0)} &= e_\ell^2 P_{ff}^{V,(0)}, \\
\mathcal{P}_{u\bar{u}}^{V,(0)} &= e_u^2 P_{ff}^{V,(0)}, \\
\mathcal{P}_{d\bar{d}}^{V,(0)} &= e_d^2 P_{ff}^{V,(0)},
\end{aligned} \tag{1.6}$$

where  $e_\ell^2 = 1$ ,  $e_u^2 = 4/9$ , and  $e_d^2 = 1/9$  are the electric charges of up- and down-type quarks, respectively.

Let us now move to the next-to-leading order coefficient  $\mathcal{P}^{(1)}$ . At this order, the pure-single splitting functions are different from zero but all the same up to an overall factor:

$$\mathcal{P}_{xy}^{S,(1)} = e_x^2 e_y^2 P_{ff}^{S,(1)} \tag{1.7}$$

where  $x$  and  $y$  are either a lepton, or a down-type quark, or an up-type quark. For the other splitting functions

one finds:

$$\begin{aligned}
\mathcal{P}_{\ell\gamma}^{(1)} &= e_\ell^4 P_{f\gamma}^{(1)}, \\
\mathcal{P}_{u\gamma}^{(1)} &= e_u^4 P_{f\gamma}^{(1)}, \\
\mathcal{P}_{d\gamma}^{(1)} &= e_d^4 P_{f\gamma}^{(1)}, \\
\mathcal{P}_{\ell\ell}^{(1)} &= e_\ell^4 P_{\gamma f}^{(1)} + e_\ell^2 P_{\gamma f,n}^{(1)}, \\
\mathcal{P}_{\gamma u}^{(1)} &= e_u^4 P_{\gamma f}^{(1)} + e_u^2 P_{\gamma f,n}^{(1)}, \\
\mathcal{P}_{\gamma d}^{(1)} &= e_d^4 P_{\gamma f}^{(1)} + e_d^2 P_{\gamma f,n}^{(1)}, \\
\mathcal{P}_{\ell\ell}^{V,(1)} &= e_\ell^4 P_{ff}^{V,(1)} + e_\ell^2 P_{ff,n}^{V,(1)}, \\
\mathcal{P}_{uu}^{V,(1)} &= e_u^4 P_{ff}^{V,(1)} + e_u^2 P_{ff,n}^{V,(1)}, \\
\mathcal{P}_{dd}^{V,(1)} &= e_d^4 P_{ff}^{V,(1)} + e_d^2 P_{ff,n}^{V,(1)}, \\
\mathcal{P}_{\ell\bar{\ell}}^{V,(1)} &= e_\ell^4 P_{f\bar{f}}^{V,(1)}, \\
\mathcal{P}_{u\bar{u}}^{V,(1)} &= e_u^4 P_{f\bar{f}}^{V,(1)}, \\
\mathcal{P}_{d\bar{d}}^{V,(1)} &= e_d^4 P_{f\bar{f}}^{V,(1)}.
\end{aligned} \tag{1.8}$$

We are now in a position to rewrite the evolution system in Eq. (1.1) exploiting the equalities identified above:

$$\begin{aligned}
\mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} \ell_\alpha^+ \\ u_j^+ \\ d_i^+ \\ \gamma \end{pmatrix} &= \left[ \left( \frac{\alpha}{4\pi} \right) \begin{pmatrix} e_\ell^2 P_{ff}^{V,(0)} & 0 & 0 & 2e_\ell^2 P_{f\gamma}^{(0)} \\ 0 & e_u^2 P_{ff}^{V,(0)} & 0 & 2e_u^2 P_{f\gamma}^{(0)} \\ 0 & 0 & e_d^2 P_{ff}^{V,(0)} & 2e_d^2 P_{f\gamma}^{(0)} \\ e_\ell^2 P_{\gamma f}^{(0)} & e_u^2 P_{\gamma f}^{(0)} & e_d^2 P_{\gamma f}^{(0)} & P_{\gamma\gamma}^{(0)} \end{pmatrix} \begin{pmatrix} \ell_\alpha^+ \\ u_j^+ \\ d_i^+ \\ \gamma \end{pmatrix} \right. \\
&+ \left( \frac{\alpha}{4\pi} \right)^2 \begin{pmatrix} e_\ell^4 P_{ff}^{V,(1)} + e_\ell^2 P_{ff,n}^{V,(1)} + e_\ell^4 P_{f\bar{f}}^{V,(1)} & 0 & 0 & 2e_\ell^4 P_{f\gamma}^{(1)} \\ 0 & e_u^4 P_{ff}^{V,(1)} + e_u^2 P_{ff,n}^{V,(1)} + e_u^4 P_{f\bar{f}}^{V,(1)} & 0 & 2e_u^4 P_{f\gamma}^{(1)} \\ 0 & 0 & e_d^4 P_{ff}^{V,(1)} + e_d^2 P_{ff,n}^{V,(1)} + e_d^4 P_{f\bar{f}}^{V,(1)} & 2e_d^4 P_{f\gamma}^{(1)} \\ e_\ell^4 P_{\gamma f}^{(1)} + e_\ell^2 P_{\gamma f,n}^{(1)} & e_u^4 P_{\gamma f}^{(1)} + e_u^2 P_{\gamma f,n}^{(1)} & e_d^4 P_{\gamma f}^{(1)} + e_d^2 P_{\gamma f,n}^{(1)} & P_{\gamma\gamma}^{(1)} \end{pmatrix} \begin{pmatrix} \ell_\alpha^+ \\ u_j^+ \\ d_i^+ \\ \gamma \end{pmatrix} \\
&\left. + \left( \frac{\alpha}{4\pi} \right)^2 2P_{ff}^{S,(1)} \begin{pmatrix} e_\ell^4 & e_\ell^2 e_u^2 & e_\ell^2 e_d^2 & 0 \\ e_u^2 e_\ell^2 & e_u^4 & e_u^2 e_d^2 & 0 \\ e_d^2 e_\ell^2 & e_d^2 e_u^2 & e_d^4 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \Sigma_\ell \\ \Sigma_u \\ \Sigma_d \\ \gamma \end{pmatrix} \right], \tag{1.9}
\end{aligned}$$

$$\begin{aligned}
\mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} d_i^- \\ u_j^- \\ \ell_\alpha^- \end{pmatrix} &= \left[ \left( \frac{\alpha}{4\pi} \right) \begin{pmatrix} e_d^2 P_{ff}^{V,(0)} & 0 & 0 \\ 0 & e_u^2 P_{ff}^{V,(0)} & 0 \\ 0 & 0 & e_\ell^2 P_{ff}^{V,(0)} \end{pmatrix} \begin{pmatrix} d_i^- \\ u_j^- \\ \ell_\alpha^- \end{pmatrix} \right. \\
&+ \left( \frac{\alpha}{4\pi} \right)^2 \begin{pmatrix} e_d^4 P_{ff}^{V,(1)} + e_d^2 P_{ff,n}^{V,(1)} - e_d^4 P_{f\bar{f}}^{V,(1)} & 0 & 0 \\ 0 & e_u^4 P_{ff}^{V,(1)} + e_u^2 P_{ff,n}^{V,(1)} - e_u^4 P_{f\bar{f}}^{V,(1)} & 0 \\ 0 & 0 & e_\ell^4 P_{ff}^{V,(1)} + e_\ell^2 P_{ff,n}^{V,(1)} - e_\ell^4 P_{f\bar{f}}^{V,(1)} \end{pmatrix} \begin{pmatrix} d_i^- \\ u_j^- \\ \ell_\alpha^- \end{pmatrix} \left. \right], \tag{1.10}
\end{aligned}$$

where we have used the following definitions:

$$\begin{aligned}
\ell_\alpha^\pm &= \ell_\alpha \pm \bar{\ell}_\alpha \\
u_j^\pm &= u_j \pm \bar{u}_j \\
d_i^\pm &= d_i \pm \bar{d}_i
\end{aligned} \tag{1.11}$$

and also introduced:

$$\Sigma_\ell = \sum_{\alpha=e,\mu,\tau} \ell_\alpha^+ \quad \Sigma_u = \sum_{k=i}^{N_c n_u} u_k^+ \quad \Sigma_d = \sum_{k=i}^{N_c n_d} d_k^+. \tag{1.12}$$

For completeness we also introduce:

$$V_\ell = \sum_{\alpha=e,\mu,\tau} \ell_\alpha^- \quad V_u = \sum_{k=i}^{N_c n_u} u_k^- \quad V_d = \sum_{k=i}^{N_c n_d} d_k^-. \tag{1.13}$$

Notice that quarks are summed also over their colour degree of freedom.

In order to diagonalise as much as possible the evolution matrix in the presence of QED corrections avoiding unnecessary couplings between parton distributions, we propose the following evolution basis:

$$\begin{array}{ll}
1) \ \gamma & 11) \ V_\ell \\
2) \ \Sigma_\ell & 12) \ V_u \\
3) \ \Sigma_u & 13) \ V_d \\
4) \ \Sigma_d & 14) \ T_2^\ell = \ell_e^+ + \ell_\mu^+ - 2\ell_\tau^+ \\
5) \ T_1^\ell = \ell_e^+ - \ell_\mu^+ & 15) \ V_2^\ell = \ell_e^- + \ell_\mu^- - 2\ell_\tau^- \\
6) \ V_1^\ell = \ell_e^- - \ell_\mu^- & 16) \ V_1^u = u^- - c^- \\
7) \ T_1^u = u^+ - c^+ & 17) \ V_2^u = u^- + c^- - 2t^- \\
8) \ T_2^u = u^+ + c^+ - 2t^+ & 18) \ V_1^d = d^- - s^- \\
9) \ T_1^d = d^+ - s^+ & 19) \ V_2^d = d^- + s^- - 2b^- \\
10) \ T_2^d = d^+ + s^+ - 2b^+ &
\end{array} \tag{1.14}$$

we then introduce the following splitting functions:

$$\begin{aligned}
P_p^\pm &= \left(\frac{\alpha}{4\pi}\right) e_p^2 P_{ff}^{V,(0)} + \left(\frac{\alpha}{4\pi}\right)^2 \left[ e_p^4 P_{ff}^{V,(1)} + e_p^2 P_{ff,n}^{V,(1)} \pm e_p^4 P_{f\bar{f}}^{V,(1)} \right] \\
P_{p\gamma} &= \left(\frac{\alpha}{4\pi}\right) e_p^2 P_{f\gamma}^{V,(0)} + \left(\frac{\alpha}{4\pi}\right)^2 e_p^4 P_{f\gamma}^{V,(1)} \\
P_{\gamma p} &= \left(\frac{\alpha}{4\pi}\right) e_p^2 P_{\gamma f}^{V,(0)} + \left(\frac{\alpha}{4\pi}\right)^2 \left[ e_p^4 P_{\gamma f}^{V,(1)} + e_p^2 P_{\gamma f,n}^{V,(1)} \right] \\
P_{pp'}^{\text{PS}} &= \left(\frac{\alpha}{4\pi}\right)^2 2N_c^{(p)} n_p e_p^2 e_{p'}^2 P_{ff}^{S,(1)} \\
P_{pp} &= P_p^+ + P_{pp}^{\text{PS}} \\
P_{\gamma\gamma} &= \left(\frac{\alpha}{4\pi}\right) P_{\gamma\gamma}^{V,(0)} + \left(\frac{\alpha}{4\pi}\right)^2 P_{\gamma\gamma}^{V,(1)}
\end{aligned} \tag{1.15}$$

with  $p, p' = \ell, u, d$ , and  $N_c^{(\ell)} = 1$  and  $N_c^{(u)} = N_c^{(d)} = N_c = 3$ . With these definitions at hand we can finally write the full set of systems of evolution equation using the basis given in Eq. (1.14):

$$\mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} \gamma \\ \Sigma_\ell \\ \Sigma_u \\ \Sigma_d \end{pmatrix} = \begin{pmatrix} P_{\gamma\gamma} & P_{\gamma\ell} & P_{\gamma u} & P_{\gamma d} \\ 2n_\ell P_{\ell\gamma} & P_{\ell\ell} & P_{\ell u}^{\text{PS}} & P_{\ell d}^{\text{PS}} \\ 2N_c n_u P_{u\gamma} & P_{u\ell}^{\text{PS}} & P_{uu} & P_{ud}^{\text{PS}} \\ 2N_c n_d P_{d\gamma} & P_{d\ell}^{\text{PS}} & P_{du}^{\text{PS}} & P_{dd} \end{pmatrix} \begin{pmatrix} \gamma \\ \Sigma_\ell \\ \Sigma_u \\ \Sigma_d \end{pmatrix} \tag{1.16}$$

$$\mu^2 \frac{\partial V_p}{\partial \mu^2} = P_p^- V_p \tag{1.17}$$

$$\mu^2 \frac{\partial T_{1,2}^p}{\partial \mu^2} = P_p^+ T_{1,2}^p \tag{1.18}$$

$$\mu^2 \frac{\partial V_{1,2}^p}{\partial \mu^2} = P_p^- T_{1,2}^p \tag{1.19}$$

## 2 Threshold crossing

In order to implement the evolution in the variable-flavour-number scheme (VFNS), it is necessary to introduce thresholds. Thresholds need to be ordered in value. To do so we use the PDG mass values:

$$\begin{aligned}
 m_e &= 0.5109989461 \cdot 10^{-3} \text{ GeV} \\
 m_u &= 2.16 \cdot 10^{-3} \text{ GeV} \\
 m_d &= 4.67 \cdot 10^{-3} \text{ GeV} \\
 m_s &= 0.093 \text{ GeV} \\
 m_\mu &= 0.1056583745 \text{ GeV} \\
 m_c &= 1.27 \text{ GeV} \\
 m_\tau &= 1.77686 \text{ GeV} \\
 m_b &= 4.18 \text{ GeV} \\
 m_t &= 172.76 \text{ GeV}
 \end{aligned} \tag{2.1}$$