	CONTENTS		10	Total delay in dependence of the discretization parameter $\Delta_d$ for 9 different problem
List o	of Tables	1		instances with up to in 12 flights and 25 conflicts
List o	of Figures	1	11	Optimal total delay found by using the Isoenergetic Cluster Method (ICM) at
IIntro	duction	2		fixed time step $\Delta_d$ as a function of numbers of flight within each connected component.
I <b>P</b> robl	em specification	2		ICM was unable to find solutions for connected component with more than 12 flights.
I <b>I</b> hstai	nces	4	12 13	Penalty weight phase diagram  Number of physical qubits versus the num-
∏ <b>D</b> iscr	etizing the configuration space	6		ber of logical qubits after embedding of QUBO instances for the departure delay
Wapr	oing to QUBO	7		model
ABinary encoding		7	14	Average success probability for QUBO in-
Boftening the constraints		7		stances in dependence of the number of flights $N_f$ and the number of conflicts
Muan	tum Annealing	8		$N_c$ . The error bars indicate the stan-
Wuantum Annealing		8		dard deviation. We used 10000 annealing
AEmbedding				runs for each instance and penalty weights
Buccess Probability 8 Which clusions 9				$\lambda = \lambda_{\text{conflict}} = \lambda_{\text{unique}} \in \{0.5, 1, 2\}. \dots$
Which is the second of the seco		9		
V <b>AL</b> kno	owledgements	9		
<b>A</b> Gener	ral QUBO mapping	9		
1Global trajectory modifications		9		
2Local trajectory modifications		10		
aExc	clusive avoidance	10		
bFle	xible avoidance	10		
cInterstitial delays		10		
References		10		
	LIST OF TABLES			
I	Parameters of the largest embeddable instances for the D-Wave 2X	8		
	LIST OF FIGURES			
1	Conflict example	3		
$\frac{2}{3}$	Conflict preprocessing  Number of non-trivial connected compo-	4		
	nents vs. $d_{\text{max}}$	4		
4	Histogram of connected component sizes	5		
5	Histogram of degrees	5		
6	Power-law exponent vs. $d_{\text{max}}$	5		
7	Correlation between connected component			
	size and treewidth	5		
8	Treewidth-size correlation coefficient vs.			
	$d_{\max}\ldots\ldots\ldots$	6		

Effect of discretization on solution quality . 6

# Quantum Annealing for Air Traffic Management

Tobias Stollenwerk German Aerospace Center, Linder Höhe, 51147 Cologne, Germany

## Bryan O'Gorman

University of California, Berkeley

NASA Ames Research Center Quantum Artificial Intelligence Laboratory (QuAIL), Mail Stop 269-1, 94035 Moffett Field CA and Stinger Ghaffarian Technologies Inc., 7701 Greenbelt Rd., Suite 400, Greenbelt, MD 20770

### Davide Venturelli

NASA Ames Research Center Quantum Artificial Intelligence Laboratory (QuAIL), Mail Stop 269-1, 94035 Moffett Field CA and USRA

#### Eleanor G. Rieffel

NASA Ames Research Center Quantum Artificial Intelligence Laboratory (QuAIL), Mail Stop 269-1, 94035 Moffett Field CA

## Salvatore Mandrà

NASA Ames Research Center Quantum Artificial Intelligence Laboratory (QuAIL), Mail Stop 269-1, 94035 Moffett Field CA and Stinger Ghaffarian Technologies Inc., 7701 Greenbelt Rd., Suite 400, Greenbelt, MD 20770

Olga Rodionova, Hok K. Ng and Banavar Sridhar tbd (Dated: May 19, 2017)

In this paper we present the mapping of air traffic management (ATM) problem on quadratic unconstrained boolean optimization (QUBO) problem. After the representation of the ATM problem in terms of a conflict graph, where nodes of the graph represent flights and edges represent a potential conflict between flights, we proceed by discretize the ATM problem and then mapping it in binary variables. As part of our study, we tested the QUBO formulation of the ATM problem using both classical solvers and the D-Wave 2X quantum chip.

# I. INTRODUCTION

There is an overall increase in air traffic over the last decades and this trend is believed to continue. As a result, the air traffic control system is increasingly under pressure. The current system works with pre-planned routes for the flights and more or less manual conflict avoidance with the help of air traffic controllers. With the limited airspace available, novel approaches are necessary to meet the demand of even more air traffic in the coming decades. A promising approach to solve this problem is using wind-optimal trajectories beyond the predefined air traffic network [1]. However, the conflicts in the wind-optimal trajectories need to be avoided, or "deconflicted" [2].

Quantum annealing is a promising computational method which became increasingly important in recent years. This development is driven also by first commercially available quantum annealing device by the company D-Wave Systems. In addition to studying the fundamental properties of quantum annealing, it is imperative to find possible real world application for this technology. Hard operational planning problems are a promising candidate for the latter [3–5].

In this work, we investigate the feasibility of applying quantum annealing to the deconfliction of wind-optimal trajectories. In the course of this, we use wind-optimal trajectory data to extract realistic problem instances. To be amenable to a D-Wave Quantum annealer, the problem has to be formulated as a quadratic unconstrained binary optimization (QUBO) problem. We demonstrate the mapping of deconfliction problem to a QUBO formulation. Since restrictions to the configurations space are necessary for the reformulation of the problem as a QUBO, we investigate the influence of this restriction on the solution quality. Moreover, we study the embeddability and solution quality using both classical solvers and quantum annealing runs.

# II. PROBLEM SPECIFICATION

The basic input of the deconflicting problem is a set of ideal flight trajectories (space-time paths). These ideal trajectories are specified by the individual flight operators. Each ideal trajectory represents some independent optimization from the operator's perspective, especially minimizing fuel costs given expected wind conditions between the desired origin and destination at the desired times; for this reason, they are called the "wind-optimal" trajectories. Because of the number of such trajectories and the correlation between them, these trajectories are likely to conflict; that is, two or more aircraft are likely to get dangerously close to each other if their ideal trajec-

tories are followed without modification. The goal thus is to modify the trajectories to avoid such conflicts.

In theory, the configuration space consists of all physically realistic trajectories; in practice, computational bounds constrain us to consider perturbations of the ideal trajectories. The simplest such perturbation is a departure delay, which is the main focus of the present work. Previous work [2] additionally considered a global perturbation by which a trajectory is sinusoidally shifted parallel to the Earth's surface. We focus instead on local perturbations to the trajectories, in which a modification to the trajectory is parameterized by some choice of active maneuvers near a potential conflict; such a modification does not affect the preceding part of the trajectory and only affects the subsequent part by the additional delay it introduces.

A full accounting of the cost of such modifications would take into account the cost of departure delays, the change in fuel cost due to perturbing the trajectories, the relative importance of each flight, and many other factors. As in previous work, we consider only the total, unweighted arrival delay, aggregated equally over all of the flights.

Formally, each ideal trajectory  $\mathbf{x}_i = (x_{i,t})_{t=t_{i,0}}^{t_{i,1}}$  is specified as a time-discretized path from the departure point  $x_{i,t_{i,0}}$  at time  $t_{i,0}$  to the arrival point  $x_{i,t_{i,1}}$  at time  $t_{i,1}$ . For each flight i, the geographical coordinates  $x_{i,t}$  (as latitude, longitude, and altitude) are specified at every unit of time (i.e. one minute) between  $t_{i,0}$  and  $t_{i,1}$ ; we call this interval  $T_i = (t_{i,0}, t_{i,0} + 1, \dots, t_{i,1})$ .

For notational simplicity, suppose momentarily that each trajectory  $\mathbf{x}_i$  is modified only by introducing delays between time steps. Let  $\delta_{i,t}$  be the accumulated delay of flight i at the time that it reaches the point  $x_{i,t}$ , and let  $\delta_{i,t}^*$  be the maximum such delay.

A pair of flights (i,j) are in conflict with each other if any pair of points from their respective trajectories is in conflict. (The trajectories are reasonably assumed to be sufficiently time-resolved so that if the continuously interpolated trajectories conflict than there is a pair of discrete trajectory points that conflict.) A pair of trajectory points  $(x_{i,s}, x_{j,t})$  conflict if their spatial and temporal separations are both within the respective mandatory separation standards  $\Delta_x$  and  $\Delta_t$  (i.e. 3 nautical miles and 3 minutes):

$$||x_{i,s} - x_{i,t}|| < \Delta_x \text{ and } |(s + \delta_{i,s}) - (t + \delta_{i,t})|$$
 (1)

. The latter condition can be met for some  $(\delta_{i,s},\delta_{j,t})\in[0,\delta_{i,s}^*]\times[0,\delta_{j,t}^*]$  if and only if

$$\max\left\{\delta_{i,s}^*, \delta_{j,t}^*\right\} + \Delta_t > |s - t|,\tag{2}$$

in which case we call the pair of trajectory points potentially conflicting. The set C of such pairs of potentially conflicting trajectory points contains strongly correlated clusters. To simplify the constraints, we enumerate all such clusters and refer to them simply as the conflicts. That is, we partition the potentially conflicting pairs of

trajectory points into disjoint sets,

$$C = \bigcup_{k} C_k, \tag{3}$$

such that if  $\{(i,s),(j,t)\},\{(i',s'),(j',t')\}\in C_k$  for some k then i=i'< j=j' and for all  $s''\in [\min\{s,s'\},\max\{s,s'\}]$  there exists some  $t''\in [\min\{t,t'\},\max\{t,t'\}]$  such that  $\{(i,s''),(j,t'')\}\in C_k$ . Thus every conflict k is associated with a pair of flights  $I_k=\{i,j\}$ . Let  $K_i=\{k|i\in I_k\}$  be the set of conflicts to which flight i is associated,  $N_c$  the number of conflicts, and  $N_f$  the number of flights.

Having identified disjoint sets of conflicts, we relax the supposition that the trajectory modifications only introduce delays between time steps. Instead, we consider modifications to the trajectories that introduce delays local to particular conflicts. Specifically, the configuration space consists of the departure delays  $\mathbf{d} = (d_i)_{i=1}^{N_{\mathrm{f}}}$  and the set of local maneuvers  $\mathbf{a_k} = (\mathbf{a_k})_k$ , where  $\mathbf{a_k}$  represents some parameterization of the local maneuvers used to avoid conflict k. Let  $d_{i,k}(\mathbf{d}, \mathbf{a_k})$  be the delay introduced to flight i at conflict k, as a function of the departure delays and local maneuvers. With this notation, we can write the total delay as

$$D = \sum_{i=1}^{N_{\rm f}} \left( d_i + \sum_{k \in K_i} d_{i,k} \right).$$
 (4)

This is the quantity we wish to minimize subject to avoiding all potential conflicts.

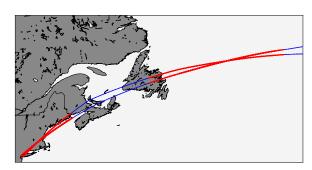


FIG. 1. Example of two parallel potential conflicts between two transatlantic flights starting from the east cost of the USA.

A conflict can be avoided locally by introducing earlier delays differentially, thereby increasing the temporal separation; by some active maneuver of one or both of the flights; or by some combination thereof. We focus on the former case. Let

$$D_{i,k} = d_i + \sum_{k' \in K_i | k' < k} d_{i,k'}$$
 (5)

be the accumulated delay of flight i by the time it reaches conflict k. In this case,  $\delta_{i,s} = D_{i,k}$  for all s associated

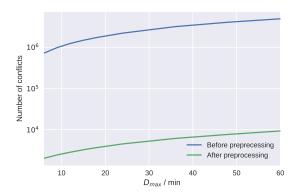


FIG. 2. Preprocessing: Reduction in the number of potential conflicts for various upper delay bounds  $D_{\text{max}}$ .

with flight i in conflict k,  $\{s|(s,t) \in T_k\}$ . We assume that the set of conflicts  $K_i$  associated with flight i is indexed in temporal order, i.e. if k' < k and  $k, k' \in K_i$ , then flight i reaches conflict k' before conflict k. The pairs of conflicting trajectory points associated with conflict k are given by

$$T_k = \{(s,t)|\{(i,s),(j,t)\} \in C_k, i < j\}.$$
 (6)

Thus the potential conflict is avoided only if

$$D_k = D_{i,k} - D_{i,k} \notin B_k \tag{7}$$

where

$$B_k = \bigcup_{(s,t)\in T_k} (-\Delta_t + t - s, \Delta_t + t - s) = [\Delta_k^{\min}, \Delta_k^{\max}],$$

$$\Delta_k^{\min} = 1 - \Delta_t + \min_{(s,t) \in T_k} \{t - s\},\tag{9}$$

$$\Delta_k^{\max} = \Delta_t - 1 + \max_{(s,t) \in T_k} \{t - s\}.$$
 (10)

In the remainder of this paper, we focus on the restricted problem in which only departure delays are allowed. In this simplified case, the configuration space is simply  $\mathbf{d} = (d_i)_{i=1}^{N_{\mathrm{f}}}$ , the cost function simply  $D = \sum_{i=1}^{N_{\mathrm{f}}} d_i$ , and the constraints simply  $d_i - d_j \notin B_k$  for all k.

## III. INSTANCES

To assess our methods on realistic instances of the problem, we use the actual wind-optimal trajectories for transatlantic flights on July 29, 2012, as was done in previous work [2]. In these trajectories, each flight has a constant (cruising) altitude and constant speed, to within (classical) machine precision, though our methods generalize to instances without these special properties.

One perspective into the nature of an instance of the deconflicting problem is the *conflict graph*, whose vertices

correspond to flights and which has an edge between a pair of vertices if there is at least one potential conflict between the corresponding flights. Note that the conflict graph for a given set of trajectories depends on the parameters of the problem. In the case of only departure delays, whether or not a potential conflict, and thus an edge in the conflict graph, exists between two flights is a function of the maximum allowable departure delay  $d_{\text{max}}$ . For a certain value of  $d_{\text{max}}$ , the conflict graph may contain several connected components, which can be considered as smaller, independent instances. Figure 3 shows this dependence of the number of connected components (both including and excluding trivial connected components, i.e. those containing a single vertex) on the maximum delay  $d_{\text{max}}$ , and Figure 4 shows the distribution of the sizes of the connected components for various values of  $d_{\text{max}}$ . Interestingly, most of the connected components are very small; for example, with  $d_{\text{max}} = 60$  minutes, approximately 75% of the connected components contain no more than 10 flights.

In the remainder of this paper, we consider sets of smaller instances corresponding to the connected components of the conflict graph from the larger single instance for various values of  $d_{\rm max}$ . Let  $\mathcal{I}_{d_{\rm max}}$  be the set of such instances for a particular value of  $d_{\rm max}$ , excluding trivial instances. We say that an instance is trivial if there are no conflicts when all flights therein depart without delay; in particular, this includes instances containing only a single flight.

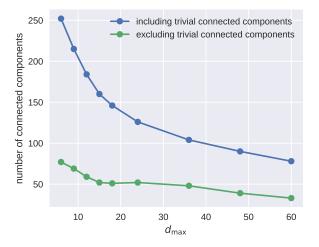


FIG. 3. Number of connected components versus  $d_{\text{max}}$ .

As part of our analysis, we also studied the probability distribution of the degree, namely the number of flights for which a given flight share a potential conflict with. Figure (5) shows the distribution of degrees in the conflict graph for  $d_{\text{max}} = 60$ , which seem to be approximately distributed according to a power law, i.e. the number of vertices with degree d is proportional to  $d^{\alpha}$ . This is consistent with a so-called "small-world" model believed

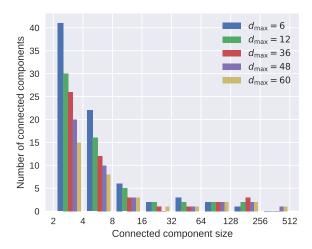


FIG. 4. Histogram of the connected component size for various values of the maximum delay time  $d_{\text{max}}$ .

to be typical of many real-world graphs [6], which are generated by preferential attachment and resultingly contain a few number of highly-connected hubs, as is the case with air traffic. Figure (6) shows the dependence of this empirical power-law exponent  $\alpha$  as a function of  $d_{\rm max}$ . As  $d_{\rm max}$  increases, the exponent decreases. The larger the delay, the less the structure of the trajectories matters and the flatter the distribution of degrees in the conflict graph.

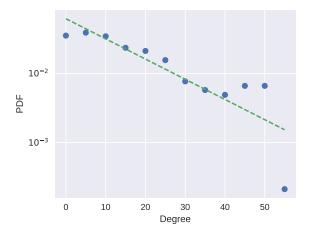


FIG. 5. Histogram of the degrees of vertices in the conflict graph for  $d_{\text{max}} = 60$ . The distribution of the degrees approximately follows a power law, with the exponent depending on  $d_{\text{max}}$ .

In many cases, generally hard problems are easy when restricted to tree-like instances [7, 8]. For example, if the conflict graph here is a tree, then the optimum could be easily found by propagating the delays along the tree; on

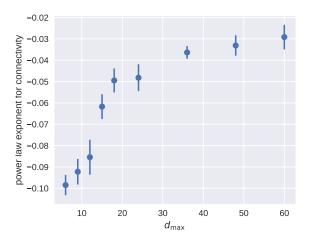


FIG. 6. Empirical power-law exponent versus  $d_{\text{max}}$ .

the other hand, if the conflict graph is a complete graph, finding the optimum is much harder. The tree-width of a graph formalizes this notion of tree-likeness, ranging from 1 for a tree to n-1 for fully connected graph. We examine the treewidth of the connected components as a proxy for the hardness of the instances they represent.

Figure 7 shows that the treewidth of a connected component scales approximately linearly with its size. This indicates that realistic instances of the deconflicting are indeed hard, and not restricted to easier (bounded treewidth) instances of the generally hard problem. Moreover, the correlation  $\gamma$  between the tree-width of a connected component and its size increases with  $d_{\rm max}$ , as shown in Figure 8. The larger  $d_{\rm max}$ , the more potential conflicts there are; restricting  $d_{\rm max}$  also restricts the number of conflicts.

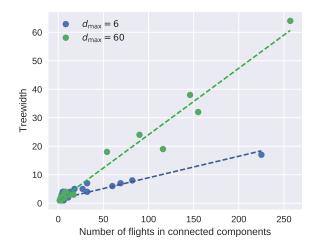


FIG. 7. The treewidths of connected components versus their sizes for various values of  $d_{\rm max}$ . The correlation is approximately linear, with a slope  $\gamma$  that depends on  $d_{\rm max}$ .

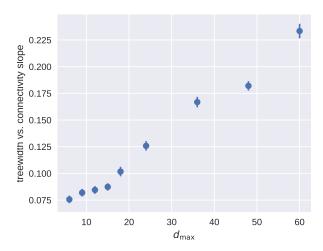


FIG. 8. Slope  $\gamma$  as a function of the maximum delay time.

# IV. DISCRETIZING THE CONFIGURATION SPACE

To apply quantum annealing to the deconflicting problem, we must encode the configuration space  $\mathbf{d}$  in binary-valued variables. To do so, we must first discretize and bound the allowed values. Let  $\Delta_d$  be the resolution of the allowed delays and  $d_{\max} = N_d \Delta_d$  the maximum allowed delay, so that  $d_i \in \{\Delta_d l | l \in [0,1,\ldots,N_d]\}$ . The larger the configuration space, the more qubits needed to encode it, and so determining the effect of this discretization on solution quality is integral to the effective use of quantum annealing. To do so, we solve the deconflicting problem with departure delays only for various delay resolutions and upper bounds and compare the various optima to the continuous problem without restrictions (other than non-negativity) on the delays.

We consider two sets of instances,  $\mathcal{I}_{18}$  and  $\mathcal{I}_{60}$ . For  $\mathcal{I}_{18}$ , the exact optima are found by modeling the problem as a constraint satisfaction problem [9]; the largest instance in  $\mathcal{I}_{18}$  has 50 flights and 104 potential conflicts.

The instances in  $\mathcal{I}_{60}$  are much larger and harder; we solved them by mapping to QUBO (as described in the next section) and then using the Isoenergetic Cluster Method (a rejection-free cluster algorithm for spin glasses that greatly improves thermalization) [10], which has been shown to be one of the fastest classical heuristic to optimize QUBO problems [11]. Because ICM is a classical method, the penalty weights can be set arbitrarily large, ensuring that the desired constraints are satisfied. While not guaranteed to return the global optimum in general, for the sizes of instances to which we applied ICM the results are sufficiently well converged to conclude that the solution found is indeed globally optimal with exceedingly high probability.

Figure 9 shows the minimum total delay of a problem

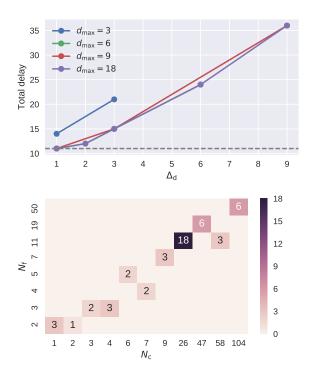


FIG. 9. Top: Minimum total delay of a problem instance with 19 flights and 47 conflicts for various values of  $\Delta_d$  and  $d_{\text{max}}$ . Bottom: Minimum  $d_{\text{max}}$  necessary to obtain same optimum as that without bounding the delay for instances in  $\mathcal{I}_{18}$ .

instance with 19 flights and 47 potential conflicts for various values of  $\Delta_d$  and  $d_{\text{max}}$ . With the exception of the small maximum delay  $d_{\text{max}} = 3$ , the total delay of the solutions is nearly independent of the maximum delay. The total delay is non-decreasing with respect to the coarseness  $\Delta_d$  of the discretization for a fixed maximum delay  $d_{\text{max}}$ , and non-increasing with respect to  $d_{\text{max}}$  for a fixed  $\Delta_d$ . Since the original data is discretized in time in units of 1 minute,  $\Delta_d = 1$  yield the same result as a continuous variable with the same upper bound. Above some threshold value  $d_{\text{max}}^0$ , further increasing the maximum delay does not decrease the minimum total delay. With one exception, we found that for all the investigated problem instances  $d_{\text{max}}^0 \leq 6$  minutes (see Figure 9). Therefore we conclude, that a moderate maximum delay is sufficient even for larger problem instances. On the other hand, the delay discretization should be as fine as possible to obtain a high quality solutions.

Figure (10) shows the dependence of the total delay time optimized by ICM on the delay discretization  $\Delta_d$  for various problem instances extracted from the connected components of the conflict graph. Results are for maximum delay of 60 minutes. As expected, the total delay decreases by decreasing  $\Delta_d$ . This is consistent with the idea that smaller  $\Delta_d$  allows a finer optimization of the delays of the flights.

In Figure (11) we show the optimal delay time found by ICM as a function of the number of the flights in

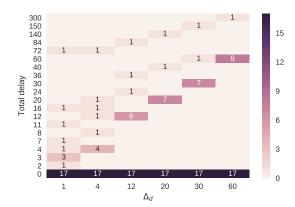


FIG. 10. Total delay in dependence of the discretization parameter  $\Delta_d$  for 9 different problem instances with up to in 12 flights and 25 conflicts.

the connected components. Results are for a maximum delay of 60 minutes. Unfortunately, ICM was unable to optimize connected components with more than 12 flights. This can be explained by recalling that ICM works the best for almost-planar problem while the its performance quickly decreases for fully-connected problems. Indeed, as shown in Section III, the underlying graph of connected components look more like a fully-connected graph rather than a tree graph by increasing the number of flights inside the connected component.

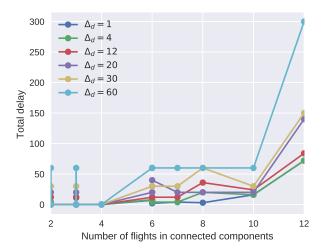


FIG. 11. . Optimal total delay found by using the Isoenergetic Cluster Method (ICM) at fixed time step  $\Delta_d$  as a function of numbers of flight within each connected component. ICM was unable to find solutions for connected component with more than 12 flights.

# V. MAPPING TO QUBO

In this section, we describe how to map to QUBO from the deconflicting problem restricted to only departure delays; a more general mapping is found in the appendix.

## A. Binary encoding

Having suitably discretized the configuration space, we must then encode it into binary-valued variables. The value of  $d_i$  is encoded in  $N_d+1$  variables  $d_{i,0},\ldots,d_{i,N_d+1}\in\{0,1\}$  using a one-hot encoding:

$$d_{i,\alpha} = \begin{cases} 1, & d_i = \alpha, \\ 0, & d_i \neq \alpha; \end{cases} \qquad d_i = \Delta_d \sum_{l=0}^{N_d} d_{i,l}.$$
 (11)

To enforce this encoding, we add the penalty function

$$f_{\text{encoding}} = \lambda_{\text{encoding}} \sum_{i=1}^{n} \left( \sum_{l=0}^{N_d} d_{i,l} - 1 \right)^2,$$
 (12)

where  $\lambda_{\text{encoding}}$  is a penalty weight sufficiently large to ensure that any cost minimizing state satisfies  $f_{\text{encoding}} = 0$ . In terms of these binary variables, the cost function is

$$f_{\text{delay}} = \Delta_d \sum_{i=1}^{n} \sum_{l=0}^{N_d} d_{i,l},$$
 (13)

Lastly, actualized conflicts are penalized by

$$f_{\text{conflict}} = \lambda_{\text{conflict}} \sum_{k} \sum_{\substack{l,l' \mid \Delta_d(l-l') \in D_k \\ i,j \in I_k \mid i < j}} d_{i,l} d_{j,l}, \qquad (14)$$

where again  $\lambda_{\text{conflict}}$  is a sufficiently large penalty weight. The overall cost function to be minimized is

$$f = f_{\text{encoding}} + f_{\text{delay}} + f_{\text{conflict}}.$$
 (15)

# B. Softening the constraints

In the QUBO formalism, there are no hard constraints; thus the use of penalty functions in the previous section. For sufficiently large penalty weights, the optimum will satisfy the desired constraints. However, precision is a limited resource in quantum annealing [?]; therefore, we would like to determine the smallest sufficient penalty weights.

For a given instance, we say that a pair of penalty weights ( $\lambda_{\text{conflict}}$ ,  $\lambda_{\text{encoding}}$ ) is valid if the minimum of the total cost function satisfies both the conflict and encoding constraints when using those weights. Figure 12 shows the phase space of these penalty weights for a single instance with 7 flights and 9 conflicts. The box-like boundary between valid and invalid penalty weights suggests that

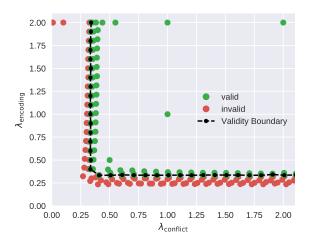


FIG. 12. Validity of exact solution to a QUBO extracted from a problem instance with  $N_f = 7$  flights and  $N_c = 9$  conflicts in dependence on the choice of the penalty weights,  $\lambda_{\text{encoding}}$  and  $\lambda_{\text{conflict}}$ . Here,  $\Delta_t = 6$  and  $d_{\text{max}} = 18$ .

the validity of the two penalty weights is independent; this box-like boundary is found for all of our instances with up to 7 flights and 9 conflicts.

One can give an upper bound for the sufficiently large penalty weights by the following considerations. A minimal violation of the hard constraints yield an additional contribution to the QUBO cost function of  $\lambda_{\text{encoding}}$  or  $\lambda_{\text{conflict}}$ , respectively. Such a violation would correspond to single bit flip in the binary delay variables  $d_{i\alpha}$ . Therefore the contribution from (??) would be reduced maximally by

$$\min_{\alpha} \frac{-\alpha}{d_{\max}} = -1$$

Hence, sufficiently large penalty weights must fulfill the following conditions

$$\lambda_{\text{encoding}} > 1$$
 $\lambda_{\text{conflict}} > 1$ .

This corresponds to a box like shape as it appears in figure 12.

# VI. QUANTUM ANNEALING

In this section we report on our efforts to solve problem instances from the departure delay model from Section ?? with a D-Wave 2X quantum annealer. We restricted ourselves to instances with  $d_{\text{max}} = D_{\text{max}} = 18$  and  $\Delta_d \in \{3,6,9\}$ .

## A. Embedding

In order to make a QUBO amenable for a D-Wave 2X quantum annealer, it has to obey certain hardware

constraints. For instance the connections between the binary variables are restricted to the so called Chimera graph [3] However, it is possible to map every QUBO to another QUBO which obeys the Chimera architecture while increasing the number of binary variables used by a so called minor-embedding technique [12]

$\Delta_d$	3	6	9
Number of flights $N_f$	13	19	50
Number of conflicts $N_c$	27	47	104
Number of logical qubits	91	76	150
Average number of physical qubits	631	395	543

TABLE I. Parameters of the largest embeddable instances for the D-Wave  $2\mathrm{X}$ 

We found, that non of the instances were suitable for direct calculation on the D-Wave machine. Therefore we used D-Wave's heuristic embedding algorithm [13] to embed instances with up to  $N_f=50$  and  $N_c=104$  depending on discretization (cf. Table I). We used up to 5 different embeddings for each QUBO instance. In figure 13 on can see the dependence of the number of physical qubits on the number of logical qubits.

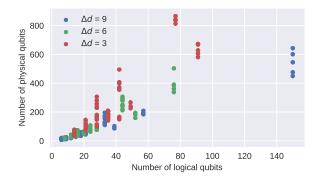
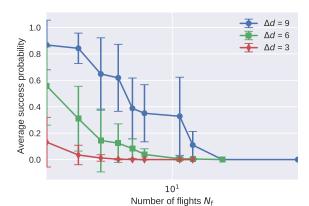


FIG. 13. Number of physical qubits versus the number of logical qubits after embedding of QUBO instances for the departure delay model.

## B. Success Probability

In order to investigate the performance of the D-Wave 2X machine, we compared its results to the ones of an exact solver. We used an exact Max-SAT solver [?] after we mapped the QUBOs to Max-SAT [?]. For each QUBO instance, we ran the annealing process  $N_r=10000$  times. The success probability is then given by the ratio of the number of annealing solutions which are equal to exact solution and the  $N_r$ .

In figure 14 the dependence of the average success probability on the number of flights and the number of conflicts is shown. The average was taken over all successfully embedded QUBO instances with the same



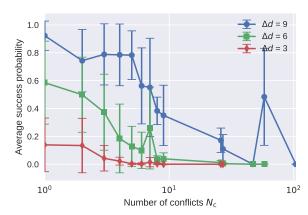


FIG. 14. Average success probability for QUBO instances in dependence of the number of flights  $N_f$  and the number of conflicts  $N_c$ . The error bars indicate the standard deviation. We used 10000 annealing runs for each instance and penalty weights  $\lambda = \lambda_{\text{conflict}} = \lambda_{\text{unique}} \in \{0.5, 1, 2\}$ .

number of flights and number of conflicts, respectively. On can see, that the success probability decreases for larger problem instances as well as for finer discretizations. This is a result of the limited precision in the specification of a QUBO on the D-Wave 2X machine [?].

## VII. CONCLUSIONS

TODO

#### VIII. ACKNOWLEDGEMENTS

# Appendix A: General QUBO mapping

In this section we describe a mapping to QUBO of a more general version of the deconflicting problem than that covered in the main text.

## 1. Global trajectory modifications

Consider the case in which each trajectory can be modified by a departure delay and some parameterized spatial transformation, i.e. for each flight i there is a variable  $d_i$  and some parameter  $\theta_i$ . For example, Rodionova et al. consider a single angle  $\theta_i$  that determines a sinusoidal transformation of the trajectory. To map to QUBO, we require that these variables be allowed to take on values from some finite set, so that are QUBO variables are  $\{d_{i,\alpha}\}$  and  $\{\boldsymbol{\theta}_{i,\phi}\}$ , where  $d_{i,\alpha}=1$   $(d_{i,\alpha})$  indicates that  $d_i = \alpha \ (d_i \neq \alpha)$  and similarly for  $\theta_{i,\phi}$ . For every pair of flights i < j, we can efficiently (in time and space polynomial in the size of the input) compute whether the corresponding trajectories conflict when modified according to  $d_i$ ,  $d_j$ ,  $\theta_i$  and  $\theta_j$ . Let  $B_{i,j}$  be the set of values of  $(d_i, \boldsymbol{\theta}_i, d_j, \boldsymbol{\theta}_j)$  such that the modified trajectories conflict. Lastly, let  $d_{(i,\alpha),(j,\beta)}$  indicate that  $d_i = \alpha$  and  $d_j = \beta$ , and similarly for  $\boldsymbol{\theta}_{(i,\phi),(j,\psi)}$ . The overall cost function is

$$f_{\text{global}}\left(\left(d_{i,\alpha}\right)_{i,\alpha}\left(d_{(i,\alpha),(j,\beta)}\right)_{i,j,\alpha,\beta}\left(\boldsymbol{\theta}_{(i,\phi),(j,\psi)}\right)_{i,j,\phi,\psi}\right) = f_{\text{encoding}} + f_{\text{co}}$$
(A1)

where

$$f_{\text{encoding}} = \lambda_{\text{encoding}} \left[ \sum_{i} \left( \sum_{\alpha} d_{i,\alpha} - 1 \right)^{2} + \sum_{i} \left( \sum_{\phi} \theta_{i,\phi} - 1 \right)^{2} \right]$$
(A2)

ensures that the values of  $d_i$  and  $\boldsymbol{\theta}_i$  are uniquely encoded;

$$f_{\text{consistency}} = \lambda_{\text{consistency}} \left[ \sum_{i < j, \alpha, \beta} s \left( d_{i,\alpha}, d_{j,\beta}, d_{(i,\alpha),(j,\beta)} \right) + \sum_{i < j, \phi, \psi} s \left( \boldsymbol{\theta}_{i,\phi}, \boldsymbol{\theta}_{j,\psi}, \boldsymbol{\theta}_{(i,\phi),(j,\psi)} \right) \right]$$
(A3)

ensures consistency between the values of  $d_{i,\alpha}$ ,  $d_{j,\beta}$ , and  $d_{(i,\alpha),(j,\beta)}$ ;

is a non-negative penalty function that is zero if and only if z = xy;

$$s(x, y, z) = 3z + xy - 2xz - 2yz$$
(A4)
$$f_{\text{delay}} = \sum_{i, \alpha} \alpha d_{i, \alpha}$$
(A5)

is the cost function to be minimized; and

$$f_{\text{conflict}} = \lambda_{\text{conflict}} \sum_{i < j} \sum_{(\alpha, \phi, \beta, \psi) \in B_{i,j}} d_{(i,\alpha),(j,\beta)} \boldsymbol{\theta}_{(i,\phi),(j,\psi)}$$
(A6)

penalize conflicts.

## 2. Local trajectory modifications

Alternatively, we can consider modifications to the trajectory only near conflicts. We describe a few special models and their mapping to QUBO, though many more such ways of doing so, and we leave a full accounting for future work.

## a. Exclusive avoidance

Suppose for every conflict k and associated pair of flights i < j, there is a way for either flight to go around the trajectory of the other, introducing some delay  $d_{i,k}$  to flight i or  $d_{j,k}$  to flight j depending on which trajectory is changed. Let  $a_k = a_{i,k} = 1$   $(a_{i,k} = 0)$  indicate that flight i's trajectory is changed, and for convenience let  $a_{j,k} = a_{i,k} - 1$ , though only one (qu)bit will be used per conflict. Adding in the departure delay, we have the total cost function

$$f_{\text{exclusive}} = f_{\text{delay}} + f_{\text{encoding}},$$
 (A7)

where

$$f_{\text{delay}} = \sum_{i} \left[ \sum_{\alpha} \alpha d_{i,\alpha} + \sum_{k \in K_i} d_{i,k} a_{i,k} \right]$$
 (A8)

and  $f_{\text{encoding}}$  is as in (12). This assumes that the trajectory modifications don't introduce potential conflicts with other flights; this assumption can be partially relaxed by adding penalty terms of the form  $a_{i,k}a_{j,k'}$  or  $d_{i,\alpha}a_{j,k}$  as appropriate.

## b. Flexible avoidance

Exclusive requires that one or the other flight is delayed at each conflict. We can relax this by accounting for the fact that if the flights arriving at a potential conflict are already relatively delayed, the conflict could be passively avoided. Let  $D_{k,\gamma} = 1$   $(D_{k,\gamma} = 0)$  indicate that  $D_k = \gamma$   $(D_k \neq \gamma)$ , where  $D_k$  is the difference in the accumulated delays at conflict k as defined in (7).

$$f_{\text{flexible}} = f_{\text{encoding}} + f_{\text{delay}} + f_{\text{consistency}} + f_{\text{conflict}}$$
 (A9)

$$f_{\text{encoding}} = \lambda_{\text{encoding}} \left[ \sum_{i} \left( \sum_{\alpha} d_{i,\alpha} - 1 \right)^{2} + \sum_{k} \left( \sum_{\gamma} D_{k,\gamma} - 1 \right)^{2} \right]$$
(A10)

$$f_{\text{consistency}} = \sum_{k} \left( \sum_{\alpha} \alpha d_{i,\alpha} + \sum_{k' \in K_i | k' < k} d_{i,k'} a_{i,k'} - \sum_{\alpha} \alpha d_{j,\alpha} - \sum_{k' \in K_i | k' < k} d_{i,k'} a_{i,k'} \right) \right)$$
(A11)

and  $f_{\text{delay}}$  as in (A8) but where  $a_{i,k}$  and  $a_{j,k}$  are separate bits.

$$f_{\text{conflict}} = \lambda_{\text{conflict}} \sum_{k} \sum_{\gamma \in B_k} \left[ D_{k,\gamma} \left( 1 - a_{i,k} - a_{j,k} \right) + 2 a_{i,k} a_{j,k} \right]$$
(A12)

If we want to allow both flights to be delayed at conflict  $a_{i,k} = a_{j,k} = 1$ , we must introduce an ancillary bit  $a_k$  that indicates whether at least one flight is delayed at conflict k, adding

$$\lambda_{\text{consistency}} \sum_{k} \left[ \left( a_{i,k} + a_{j,k} \right) \left( 1 - 2a_k \right) + a_{i,k} a_{j,k} \right]$$
 (A13)

to  $f_{\text{consistency}}$ , and replacing  $f_{\text{conflict}}$  with

$$\sum_{k} \sum_{\gamma \in B_k} D_{k,\gamma} (1 - a_k). \tag{A14}$$

c. Interstitial delays

- H. K. Ng, B. Sridhar, and S. Grabbe, Optimizing Aircraft Trajectories with Multiple Cruise Altitudes in the Presence of Winds, Journal of Aerospace Information Systems 11, 35 (2014), URL http://dx.doi.org/10.2514/ 1.1010084.
- [2] O. Rodionova, D. Delahaye, B. Sridhar, and H. Ng., Deconflicting wind-optimal aircraft trajectories in north atlantic oceanic airspace, Proceedings of Advanced Aircraft Efficiency in a Global Air Transport System (AEGATS'16) Conference (2016).
- [3] E. G. Rieffel, D. Venturelli, B. O'Gorman, M. B. Do,
- E. M. Prystay, and V. N. Smelyanskiy, A case study in programming a quantum annealer for hard operational planning problems, Quantum Information Processing 14, 1 (2015), ISSN 1573-1332, URL http://dx.doi.org/10.1007/s11128-014-0892-x.
- [4] E. Rieffel, D. Venturelli, M. Do, I. Hen, and J. Frank, Parametrized families of hard planning problems from phase transitions (2014), URL https://www.aaai.org/ ocs/index.php/AAAI/AAAI14/paper/view/8614/8697.
- [5] D. Venturelli, D. J. J. Marchand, and G. Rojo, Quantum Annealing Implementation of Job-Shop Scheduling,

- arXiv:1506.08479 [quant-ph] (2015), arXiv: 1506.08479, URL http://arxiv.org/abs/1506.08479.
- [6] R. Albert, H. Jeong, and A.-L. Barabási, Internet: Diameter of the world-wide web, nature 401, 130 (1999).
- [7] U. Bertele and F. Brioschi, Nonserial dynamic programming (Academic Press, 1972).
- [8] R. Halin, S-functions for graphs, Journal of geometry 8, 171 (1976).
- [9] E. Hebrard, E. O'Mahony, and B. O'Sullivan, Constraint Programming and Combinatorial Optimisation in Numberjack (Springer Berlin Heidelberg, Berlin, Heidelberg, 2010), pp. 181–185, ISBN 978-3-642-13520-0, URL http://dx.doi.org/10.1007/978-3-642-13520-0\_22.
- [10] Z. Zhu, A. J. Ochoa, and H. G. Katzgraber, Efficient

- cluster algorithm for spin glasses in any space dimension, Physical review letters 115, 077201 (2015).
- [11] S. Mandrà, Z. Zhu, W. Wang, A. Perdomo-Ortiz, and H. G. Katzgraber, Strengths and weaknesses of weakstrong cluster problems: A detailed overview of state-ofthe-art classical heuristics versus quantum approaches, Physical Review A 94, 022337 (2016).
- [12] V. Choi, [0804.4884] Minor-Embedding in Adiabatic Quantum Computation: I. The Parameter Setting Problem, URL http://arxiv.org/abs/0804.4884.
- [13] J. Cai, W. G. Macready, and A. Roy, A practical heuristic for finding graph minors, arXiv:1406.2741 [quant-ph] (2014), arXiv: 1406.2741, URL http://arxiv.org/abs/ 1406.2741.