

# Chapter 2: A review of some basic concepts

## 2.1 Light as an electromagnetic wave.

A review of complex numbers:

To add

# 2.1 Light as an electromagnetic wave.

Solution to the maxwell equations

To add

## 2.2 The monochromatic, time-harmonic plane wave

Construction of the plane-wave solution

$$\text{Speed of the EM wave} = \frac{c}{\sqrt{\epsilon\mu}} = \frac{c}{n}$$

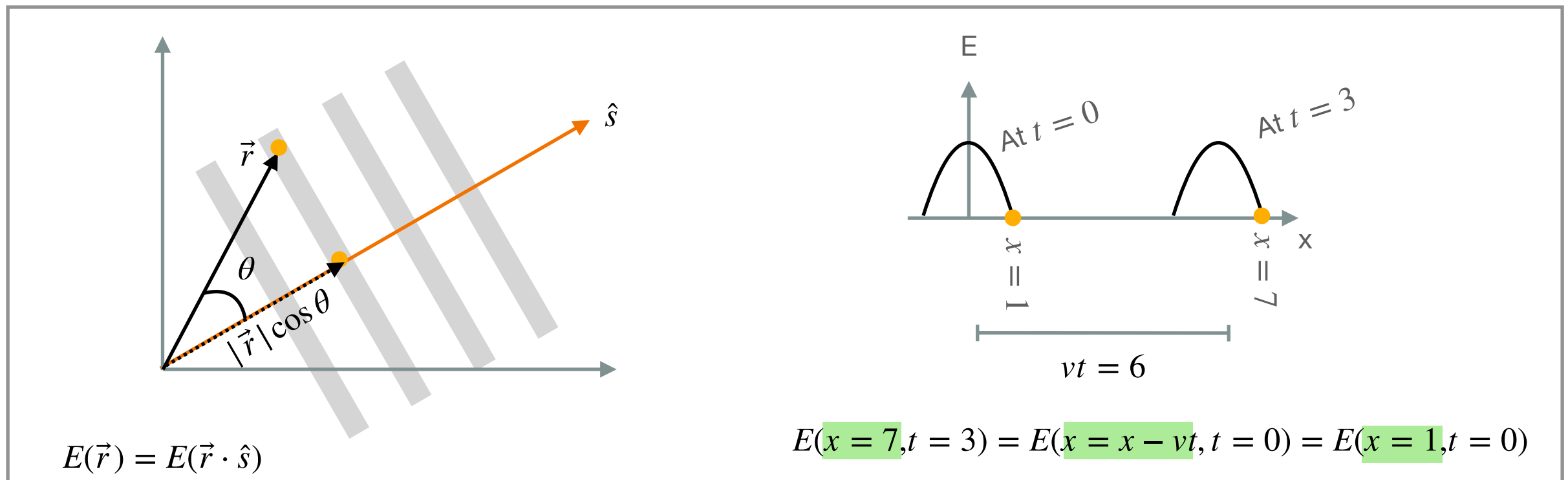
Wave vector:  $\vec{k} = \frac{2\pi}{\lambda} \hat{s}$

Direction of propagation

Wavelength

Angular frequency:  $\omega = 2\pi\nu = 2\pi \frac{v}{\lambda}$

Every point in space and time with the same  $\vec{u} = (\vec{r} \cdot \hat{s} - vt)\hat{s} = \frac{1}{k}(\vec{k} \cdot \vec{r} - \omega t)$  has the same electric field



Simple harmonic solution:

$$(2.7) \quad E_j(\vec{r}, t) = a_j e^{i(ku + \delta_j)}$$

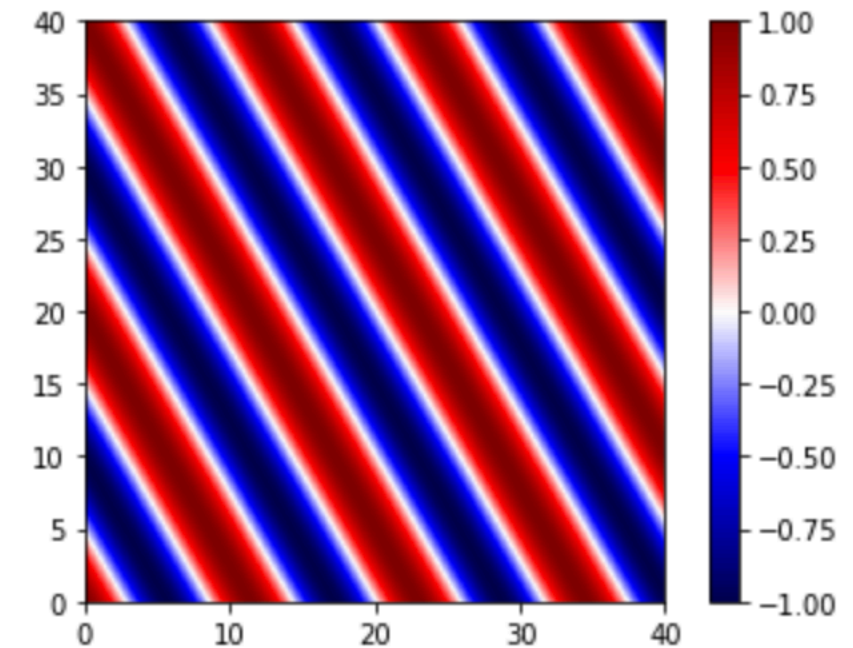
Replacing  $u$  in favor of  $\vec{k}$  and  $\omega$ :

$$(2.9) \quad E_j(\vec{r}, t) = a_j e^{i(\vec{k} \cdot \vec{r})} e^{-i(\omega t + \delta_j)}$$

## 2.2 The monochromatic, time-harmonic plane wave

$$(2.9) \quad E_j(\vec{r}, t) = a_j e^{i(\vec{k} \cdot \vec{r})} e^{-i(\omega t + \delta_j)}$$

Exercise in Python:



### ▼ 1. Illustration of the plane wave in space

As we have seen in class, a plane wave can be expressed as

$$(2.7) \quad E_j(\vec{r}, t) = a_j e^{i(\vec{k} \cdot \vec{r})} e^{-i(\omega t + \delta_j)}$$

Let's start with the spatial dependence of the plane wave. If we set  $t = 0$ , and we consider a plane wave for which the phase term  $\delta_j = 0$ , we simplify our expression below because the second exponential term is 1.0.

In the graph below, we will make a color map that illustrated the magnitude of the electric field as a function of position.

I have already defined two variables `xx` and `yy` which are 2D array that contains the x and y coordinates, respectively, of a cartesian grid of points.

1. Let's make the wavelength  $\lambda = 10$  distance units and let's make the light wave propagate in a direction  $\hat{s}$  that is oriented at  $30^\circ$  with

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i \vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i \vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

The volume energy density:

$$(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$$



## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$

The volume energy density:

$$(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$$

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$$

The volume energy density:

$$(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$

The volume energy density:

$$(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$

The volume energy density:  $(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$

The volume energy density:  $(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

## 2.3 The polarization tensor

Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

( $z = x + iy$ ,  $zz^* = x^2 + y^2 = |z|^2$ )

The volume energy density:  $(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.3 The polarization tensor

### Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

$$(z = x + iy, zz^* = x^2 + y^2 = |z|^2)$$

The volume energy density:  $(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.3 The polarization tensor

### Construction of the polarization tensor

Now let's make the wave propagate in the  $z$  direction

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

We cannot measure the electric field directly. Our detectors are sensitive to electromagnetic energy (averaged over the period of the wave)

Square of the magnitude of the electric field vector

( $z = x + iy$ ,  $zz^* = x^2 + y^2 = |z|^2$ )

The volume energy density:  $(2.12) \quad w = \frac{\epsilon}{8\pi} EE^*$

Our equations have separated the  $x$  and  $y$  components of the electric field. The magnitude of the electric vector (with time) will depend on how these two components are synchronized with each other.

e.g.: are they both zero at the same time? Is one zero while the other at its maximum?

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

Complex: not measurable



## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} \boxed{E_x E_x^*} & E_x E_y^* \\ E_y E_x^* & \boxed{E_y E_y^*} \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} \boxed{a_x^2} & a_x a_y \boxed{e^{i\delta}} \\ a_x a_y \boxed{e^{-i\delta}} & \boxed{a_y^2} \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

As we can think of the electric vector as the addition of two vectors corresponding to the  $x$  and  $y$  components, it means that we can also think of the EM wave as being a sum of a wave in the  $x$ -direction and a wave in the  $y$ -direction. In this case,  $a_x$  and  $a_y$  corresponds to the amplitude of these two waves, and therefore  $a_x^2$  and  $a_y^2$  are proportional to the intensity of these two waves. As the intensity cannot be negative, it means that:

## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

As we can think of the electric vector as the addition of two vectors corresponding to the  $x$  and  $y$  components, it means that we can also think of the EM wave as being a sum of a wave in the  $x$ -direction and a wave in the  $y$ -direction. In this case,  $a_x$  and  $a_y$  corresponds to the amplitude of these two waves, and therefore  $a_x^2$  and  $a_y^2$  are proportional to the intensity of these two waves. As the intensity cannot be negative, it means that:

$$(2.17) \quad \text{Tr}(C) = a_x^2 + a_y^2 \geq 0$$

## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

As we can think of the electric vector as the addition of two vectors corresponding to the  $x$  and  $y$  components, it means that we can also think of the EM wave as being a sum of a wave in the  $x$ -direction and a wave in the  $y$ -direction. In this case,  $a_x$  and  $a_y$  corresponds to the amplitude of these two waves, and therefore  $a_x^2$  and  $a_y^2$  are proportional to the intensity of these two waves. As the intensity cannot be negative, it means that:

$$(2.17) \quad \text{Tr}(C) = a_x^2 + a_y^2 \geq 0$$

We can also see that the determinant of the coherence matrix (2.15) is zero, such that:

## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} \boxed{E_x E_x^*} & E_x E_y^* \\ E_y E_x^* & \boxed{E_y E_y^*} \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} \boxed{a_x^2} & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & \boxed{a_y^2} \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

As we can think of the electric vector as the addition of two vectors corresponding to the  $x$  and  $y$  components, it means that we can also think of the EM wave as being a sum of a wave in the  $x$ -direction and a wave in the  $y$ -direction. In this case,  $a_x$  and  $a_y$  corresponds to the amplitude of these two waves, and therefore  $a_x^2$  and  $a_y^2$  are proportional to the intensity of these two waves. As the intensity cannot be negative, it means that:

$$(2.17) \quad \text{Tr}(C) = \boxed{a_x^2} + \boxed{a_y^2} \geq 0$$

We can also see that the determinant of the coherence matrix (2.15) is zero, such that:

$$(2.18) \quad \det(C) = 0$$

## 2.3 The polarization tensor

Properties of the polarization tensor

Magnitude of the  $x$  component

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix}$$

Magnitude of the  $y$  component

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

Complex: not measurable

As we can think of the electric vector as the addition of two vectors corresponding to the  $x$  and  $y$  components, it means that we can also think of the EM wave as being a sum of a wave in the  $x$ -direction and a wave in the  $y$ -direction. In this case,  $a_x$  and  $a_y$  corresponds to the amplitude of these two waves, and therefore  $a_x^2$  and  $a_y^2$  are proportional to the intensity of these two waves. As the intensity cannot be negative, it means that:

$$(2.17) \quad \text{Tr}(C) = a_x^2 + a_y^2 \geq 0$$

We can also see that the determinant of the coherence matrix (2.15) is zero, such that:

$$(2.18) \quad \det(C) = 0$$

As we will see later, this is always true of mono-chromatic plane waves. But, this will not be always the case for polychromatic light. Light with  $\det(C) = 0$  is called totally (or completely) polarized light.

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

$$(2.10) \quad E_x = a_x \mathbf{e}^{i\vec{k} \cdot \vec{r}} \mathbf{e}^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y \mathbf{e}^{i\vec{k} \cdot \vec{r}} \mathbf{e}^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this



## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$



$$(2.10) \quad E_z = 0$$

Won't need this

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$



$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this



Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$



## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

This is how we get the cartesian circle equation from these:

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

This is how we get the cartesian circle equation from these:

$$\cos(\theta) = \frac{E_{x,R}}{a}$$

$$\sin(\theta) = \frac{E_{y,R}}{a}$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

This is how we get the cartesian circle equation from these:

$$\cos(\theta) = \frac{E_{x,R}}{a}$$

$$\sin(\theta) = \frac{E_{y,R}}{a}$$

We know that:

$$\cos^2(\theta) + \sin^2 \theta = 1$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

This is how we get the cartesian circle equation from these:

$$\cos(\theta) = \frac{E_{x,R}}{a}$$

$$\sin(\theta) = \frac{E_{y,R}}{a}$$

We know that:

$$\cos^2(\theta) + \sin^2 \theta = 1$$

## 2.3 The polarization tensor

Path of the electric vector in the  $E_x/E_y$  plane

We can describe the “path” of the electric vector of totally polarized light in the  $E_x$  vs  $E_y$  plane. This will be at a single point in space, as a function of time.

Set  $\vec{r} = 0$

$$(2.10) \quad E_x = a_x e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_x)}$$

$$(2.10) \quad E_y = a_y e^{i\vec{k} \cdot \vec{r}} e^{-i(\omega t - \delta_y)}$$

$$(2.10) \quad E_z = 0$$

Won't need this

Because I am looking for the path over a whole period of oscillation, I can recast the phase shift this way, without losing the information I am looking for (i.e. I am not interested in where the path starts at  $t = 0$ ).

$$E_x(t) = a_x e^{-i(\omega t)}$$

$$E_y(t) = a_y e^{-i(\omega t - \delta)}$$

$$(2.16) \quad \delta \equiv \delta_x - \delta_y$$

$$E_{x,R}(t) = a_x \cos(\omega t)$$

$$E_{y,R}(t) = a_y \cos(\omega t - \delta)$$

$$E_{x,R}(\theta) = a_x \cos(\theta)$$

$$E_{y,R}(\theta) = a_y \cos(\theta - \delta)$$

During one oscillation period,  $\omega t$  will go from 0 to  $2\pi$ , so I'll just substitute to an angle  $\theta$

A simple example: let's have  $a_x = a_y = a$  and  $\delta = \pi$ .

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \cos(\theta - \pi)$$

$$E_{x,R}(\theta) = a \cos(\theta)$$

$$E_{y,R}(\theta) = a \sin(\theta)$$

That is the x and y components of a circle of radius  $a$

This is how we get the cartesian circle equation from these:

$$\cos(\theta) = \frac{E_{x,R}}{a}$$

$$\sin(\theta) = \frac{E_{y,R}}{a}$$

We know that:

$$\cos^2(\theta) + \sin^2 \theta = 1$$

$$\frac{E_{x,R}^2}{a^2} + \frac{E_{y,R}^2}{a^2} = 1$$



## 2.3 The polarization tensor

A simple case with  $a_x = a_y = a$  and  $\delta = \pi/2$

### 2. Illustration of the electric field vector with time in a single plane.

In section 2.3, we have seen that we can express the electric field vector at a single point in space for a wave propagating in the z-direction as:

$$E_{x,R}(t) = a_x \cos(\omega t)$$

and

$$E_{y,R}(t) = a_y \cos(\omega t - \delta),$$

where  $\delta \equiv \delta_x - \delta_y$

Remember that as here we will be interested in the path taken by the electric vector in the  $E_{x,R} - E_{y,R}$  plane when  $\omega t$  goes from 0 to  $2\pi$ , so we do not care about the position of the electric field at  $t = 0$ .

#### a. The simple case where $a_x = a_y = a$ and $\delta = \pi/2$

For this case, we have seen that the electric field simplifies to:

$$E_{x,R}(\theta) = a \cos(\theta)$$

and

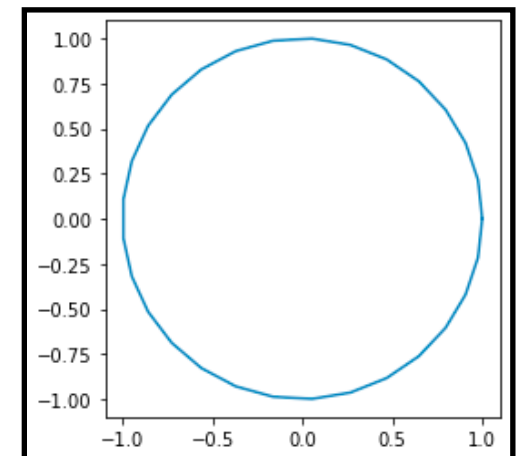
$$E_{y,R}(\theta) = a \sin(\theta)$$

In the code cell below, I have already define an array of  $\theta$  values.

Make a graph of the path of the electric field vector in the  $E_{x,R} - E_{y,R}$  plane

```
[4] fig, ax = plt.subplots(1,1)
    ax.set_aspect('equal')

    a = 1.0
    theta = np.linspace(0,2*np.pi, 30)
```

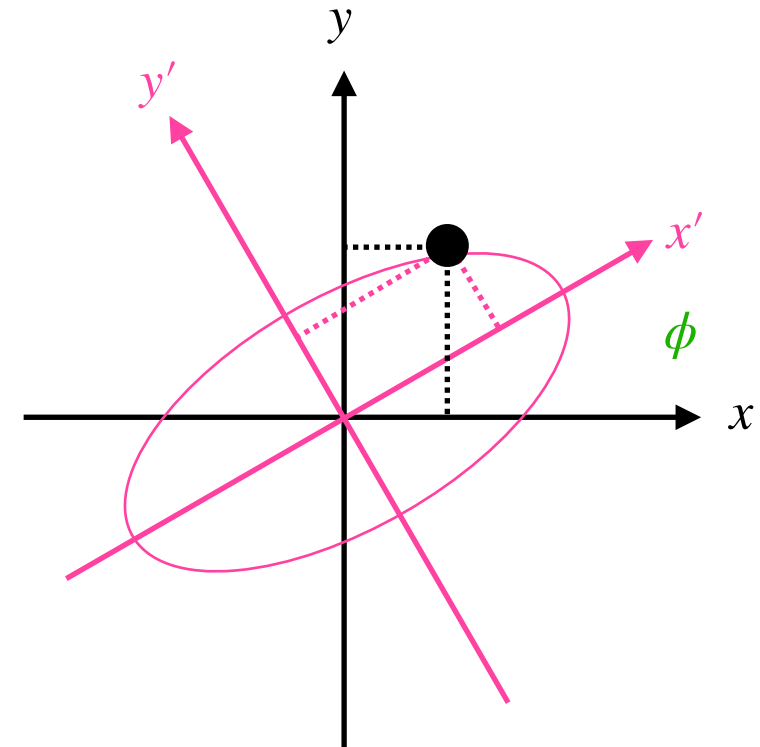


## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$



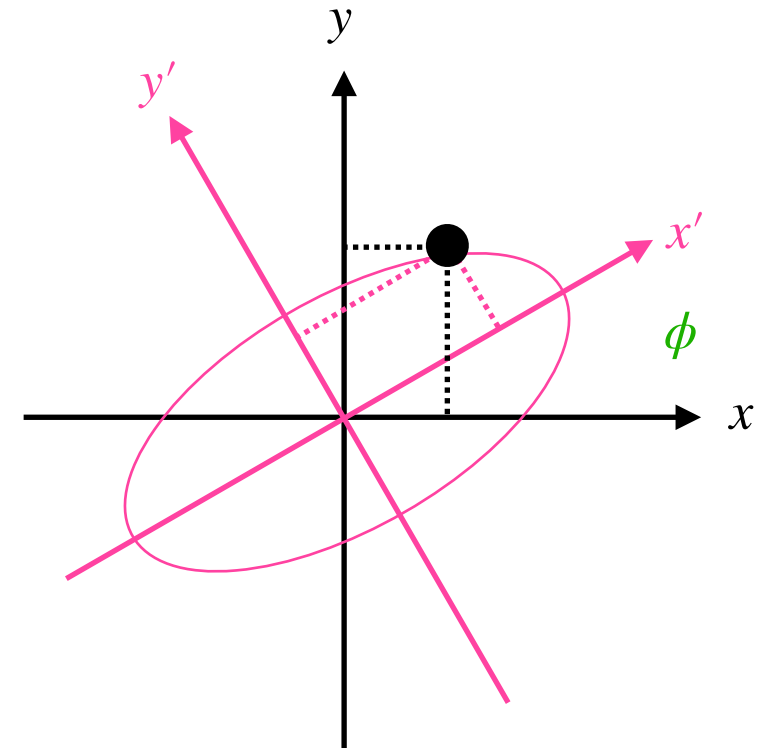
## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$



## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

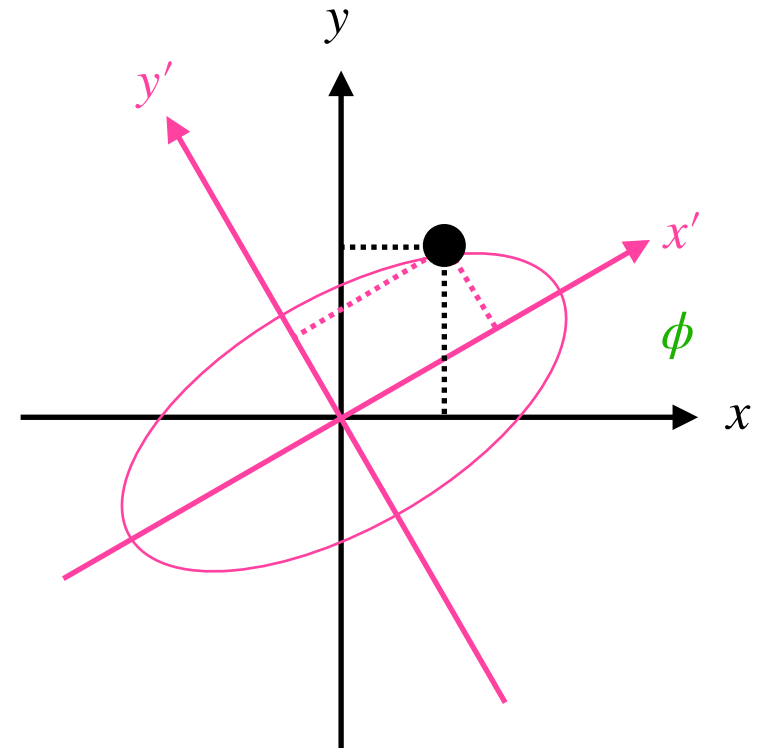
Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$



## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

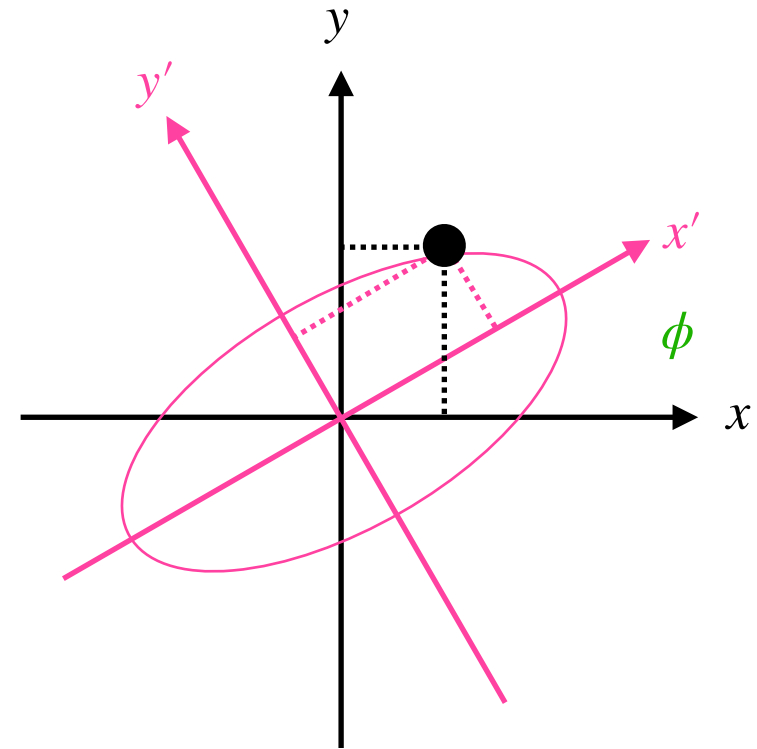
$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$

$$\frac{(x \cos \phi + y \sin \phi)^2}{a^2} + \frac{(y \cos \phi - x \sin \phi)^2}{b^2} = 1$$



## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

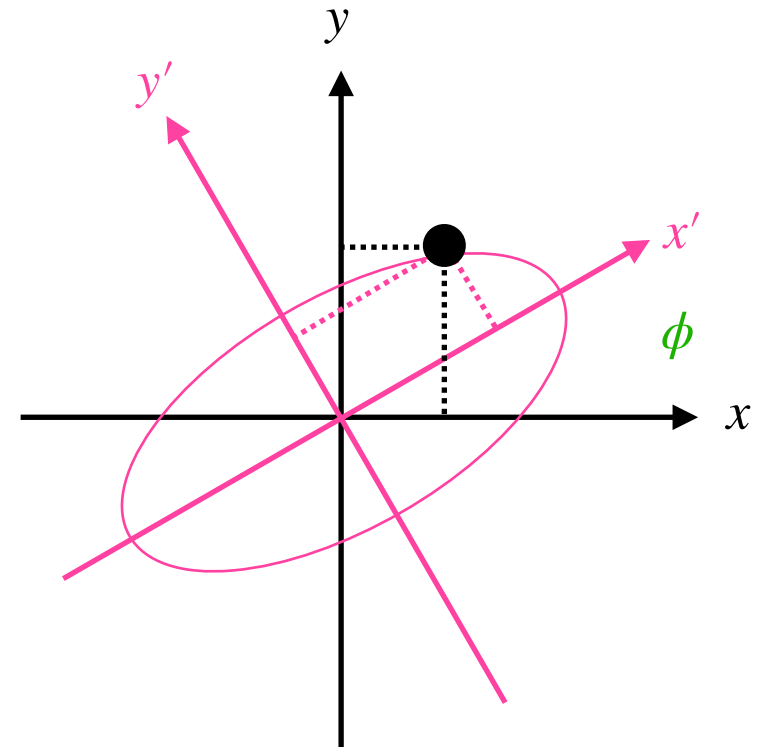
$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$

$$\frac{(x \cos \phi + y \sin \phi)^2}{a^2} + \frac{(y \cos \phi - x \sin \phi)^2}{b^2} = 1$$

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$



## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

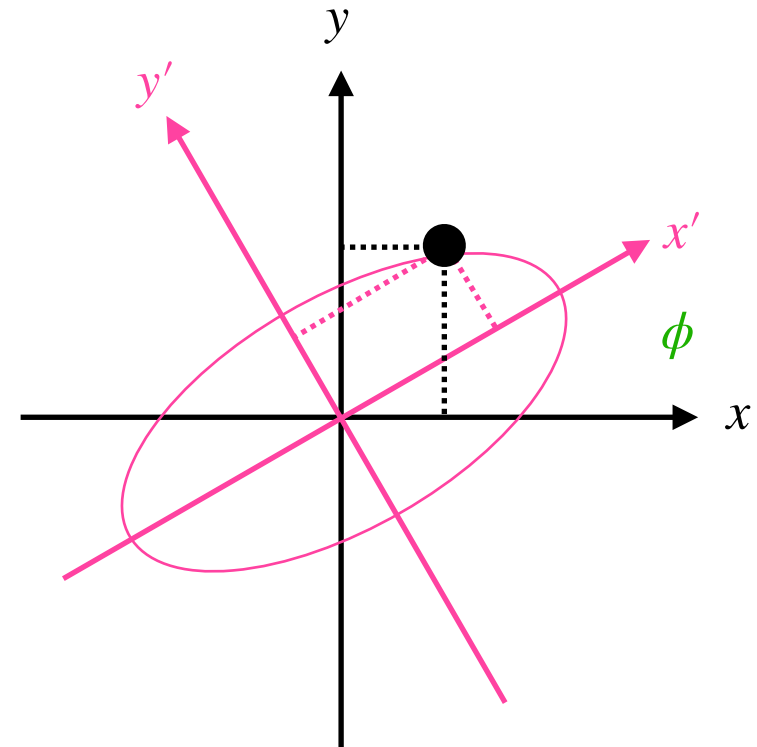
$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$

$$\frac{(x \cos \phi + y \sin \phi)^2}{a^2} + \frac{(y \cos \phi - x \sin \phi)^2}{b^2} = 1$$

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$



## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

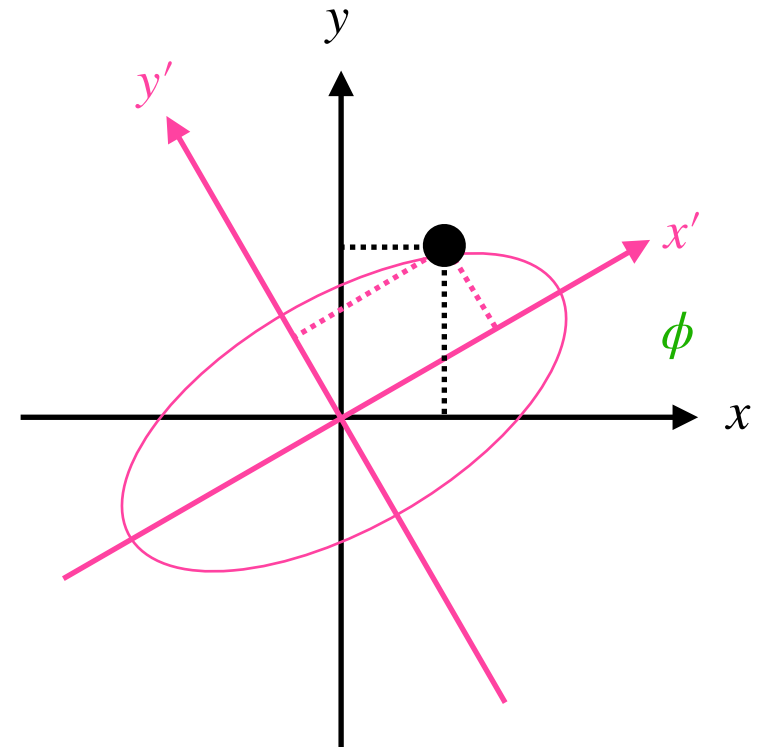
$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$

$$\frac{(x \cos \phi + y \sin \phi)^2}{a^2} + \frac{(y \cos \phi - x \sin \phi)^2}{b^2} = 1$$

$$\boxed{x^2} \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + \boxed{y^2} \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2\boxed{x}\boxed{y} \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$





## 2.3 The polarization tensor

The cartesian formula for a tilted ellipse

Useful for the next slide: The cartesian formula for a tilted ellipse

$$\frac{x'^2}{a^2} + \frac{y'^2}{b^2} = 1$$

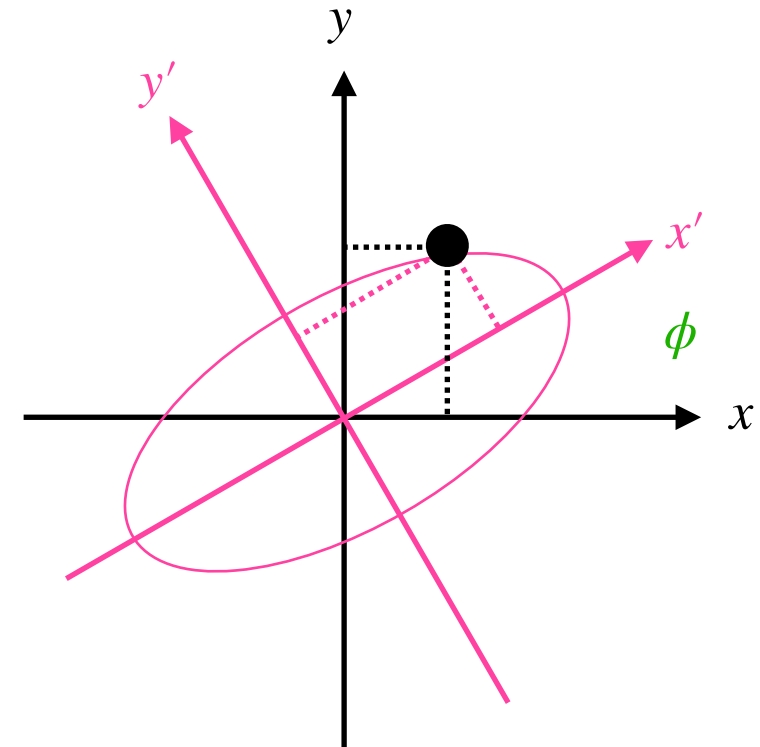
$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$x' = x \cos \phi + y \sin \phi$$

$$y' = y \cos \phi - x \sin \phi$$

$$\frac{(x \cos \phi + y \sin \phi)^2}{a^2} + \frac{(y \cos \phi - x \sin \phi)^2}{b^2} = 1$$

$$\boxed{x^2} \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + \boxed{y^2} \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2\boxed{x}\boxed{y} \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$



If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

## 2.3 The polarization tensor

The general case for the path of the electric vector

$$E_{x,R} = a_x \cos(\theta)$$

$$E_{y,R} = a_y \cos(\theta - \delta)$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$

$$E_{y,R} = a_y \cos(\theta - \delta)$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

$$E_{y,R} = a_y \cos(\theta - \delta)$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta] \quad \rightarrow \quad \frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute



## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta] \Rightarrow \frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

This equation has the form of a tilted ellipse

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

This equation has the form of a tilted ellipse

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

This equation has the form of a tilted ellipse

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

## 2.3 The polarization tensor

The general case for the path of the electric vector

1

$$E_{x,R} = a_x \cos(\theta)$$



$$\cos(\theta) = \frac{E_{x,R}}{a_x} \text{ and } \sin(\theta) = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2}$$

2

$$E_{y,R} = a_y \cos(\theta - \delta)$$

$$E_{y,R} = a_y [\cos \theta \cos \delta + \sin \theta \sin \delta]$$



$$\frac{E_{y,R}}{a_y} = [\cos \theta \cos \delta + \sin \theta \sin \delta]$$

Substitute

$$\frac{E_{y,R}}{a_y} = \frac{E_{x,R}}{a_x} \cos \delta + \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right]^{1/2} \sin \delta$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \cos \delta \right]^2 = \left[ 1 - \frac{E_{x,R}^2}{a_x^2} \right] \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} \cos^2 \delta - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta - \frac{E_{x,R}^2}{a_x^2} \sin^2 \delta$$

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} (\cos^2 \delta + \sin^2 \delta) - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

This equation has the form of a tilted ellipse

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

## 2.3 The polarization tensor

### ▼ b. The general case

Now make a graph for the general case, where  $a_x$ ,  $a_y$ , and  $\delta$  can vary.

Then, let's play a detective game, using your code. By playing with the parameters, figure out what is needed for the path of the electric field to be

1. A perfect circle

Answer:

2. A line that is tilted:

Answer:

3. A line that is purely vertical or horizontal

Answer:

4. An ellipse that is not tilted (so horizontal or vertical)

Answer:

✓  
0s



```
fig, ax = plt.subplots(1,1)
ax.set_aspect('equal')

a_x = 1.0
a_y = 1.5
delta = np.pi/3
theta = np.linspace(0,2*np.pi, 30)
```

```
##### Solution #####
```

```
##### Solution for the part below about the Stokes parameters.
```

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$



## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

If  $\delta = 0$  (or similarly with  $\delta = \pi$ ):

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

If  $\delta = 0$  (or similarly with  $\delta = \pi$ ):

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} = 0$$

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

If  $\delta = 0$  (or similarly with  $\delta = \pi$ ):

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} = 0$$
$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \right]^2 = 0$$

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

If  $\delta = 0$  (or similarly with  $\delta = \pi$ ):

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} = 0$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \right]^2 = 0$$

$$E_{y,R} = \frac{a_y}{a_x} E_{x,R}$$

## 2.3 The polarization tensor

Explaining analytically the empirical results from the notebook exercise

$$(2.19) \quad \frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} \cos \delta = \sin^2 \delta$$

If  $\delta = \pi/2$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

$$x^2 \left[ \frac{\cos^2 \phi}{a^2} - \frac{\sin^2 \phi}{b^2} \right] + y^2 \left[ \frac{\sin^2 \phi}{a^2} - \frac{\cos^2 \phi}{b^2} \right] + 2xy \cos \phi \sin \phi \left[ \frac{1}{a} - \frac{1}{b} \right] = 1$$

If  $\phi = 0$ , the  $xy$  cross term vanishes,  
and we recover the normal ellipse equation

If  $\delta = 0$  (or similarly with  $\delta = \pi$ ):

$$\frac{E_{y,R}^2}{a_y^2} + \frac{E_{x,R}^2}{a_x^2} - 2 \frac{E_{x,R} E_{y,R}}{a_x a_y} = 0$$

$$\left[ \frac{E_{y,R}}{a_y} - \frac{E_{x,R}}{a_x} \right]^2 = 0$$

$$E_{y,R} = \frac{a_y}{a_x} E_{x,R}$$

Equation of a line with slope  $a_y/a_x$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

(2.21)  $C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$\begin{aligned} I &\equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) & \begin{pmatrix} \boxed{C_{11}} + C_{21} \\ C_{12} + \boxed{C_{22}} \end{pmatrix} & \begin{pmatrix} \boxed{a_x^2} & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + \boxed{a_y^2} \end{pmatrix} \\ Q &\equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) & \begin{pmatrix} \boxed{C_{11}} - C_{21} \\ C_{12} - \boxed{C_{22}} \end{pmatrix} & \begin{pmatrix} \boxed{a_x^2} & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - \boxed{a_y^2} \end{pmatrix} \\ U &\equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta & \begin{pmatrix} C_{11} + \boxed{C_{21}} \\ \boxed{C_{12}} + C_{22} \end{pmatrix} & \begin{pmatrix} a_x^2 & \boxed{a_x a_y e^{i\delta}} \\ \boxed{a_x a_y e^{-i\delta}} + a_y^2 \end{pmatrix} \end{aligned} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \cos \delta = \frac{e^{i\delta} + e^{-i\delta}}{2}$$



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \cos \delta = \frac{e^{i\delta} + e^{-i\delta}}{2}$$

$$V \equiv i\kappa(C_{21} - C_{12}) = 2\kappa a_x a_y \sin \delta \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \cos \delta = \frac{e^{i\delta} + e^{-i\delta}}{2}$$

$$V \equiv i\kappa(C_{21} - C_{12}) = 2\kappa a_x a_y \sin \delta \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow \quad (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \cos \delta = \frac{e^{i\delta} + e^{-i\delta}}{2}$$

$$V \equiv i\kappa(C_{21} - C_{12}) = 2\kappa a_x a_y \sin \delta \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix} \quad \sin \delta = \frac{e^{i\delta} - e^{-i\delta}}{2i}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.14) \quad C \equiv \begin{pmatrix} E_x E_x^* & E_x E_y^* \\ E_y E_x^* & E_y E_y^* \end{pmatrix} \quad \begin{pmatrix} C_{11} & C_{21} \\ C_{12} & C_{22} \end{pmatrix}$$

$$(2.15) \quad C \equiv \begin{pmatrix} a_x^2 & a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} & a_y^2 \end{pmatrix}$$

$$I \equiv \kappa(C_{11} + C_{22}) = \kappa(a_x^2 + a_y^2) \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \Rightarrow (2.21) \quad C \equiv \begin{pmatrix} I + Q & \\ & I - Q \end{pmatrix}$$

$$Q \equiv \kappa(C_{11} - C_{22}) = \kappa(a_x^2 - a_y^2) \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix}$$

$$U \equiv \kappa(C_{12} + C_{21}) = 2\kappa a_x a_y \cos \delta \quad \begin{pmatrix} C_{11} + C_{21} \\ C_{12} + C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 + a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} + a_y^2 \end{pmatrix} \quad \cos \delta = \frac{e^{i\delta} + e^{-i\delta}}{2}$$

$$V \equiv i\kappa(C_{21} - C_{12}) = 2\kappa a_x a_y \sin \delta \quad \begin{pmatrix} C_{11} - C_{21} \\ C_{12} - C_{22} \end{pmatrix} \quad \begin{pmatrix} a_x^2 - a_x a_y e^{i\delta} \\ a_x a_y e^{-i\delta} - a_y^2 \end{pmatrix} \quad \sin \delta = \frac{e^{i\delta} - e^{-i\delta}}{2i} \quad \Rightarrow (2.21) \quad C \equiv \begin{pmatrix} & U + iV \\ U - iV & \end{pmatrix}$$

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

In the following sets of slides, we have seen that the Stokes parameters are defined as:

$$\begin{aligned}I &= \kappa(a_x^2 + a_y^2) \\Q &= \kappa(a_x^2 - a_y^2) \\U &= 2\kappa a_x a_y \cos \delta \\V &= 2\kappa a_x a_y \sin \delta\end{aligned}$$

In your code above, calculate the Stokes parameters, and display the values of  $Q/I$ ,  $U/I$ , and  $V/I$  in the graph.

Then, let's play the detective again, using your code. For the cases we discussed above, what are the values of the Stokes parameters?

1. A perfect circle

Answer:

2. A line that is tilted:

Answer:

3. A line that is purely vertical or horizontal

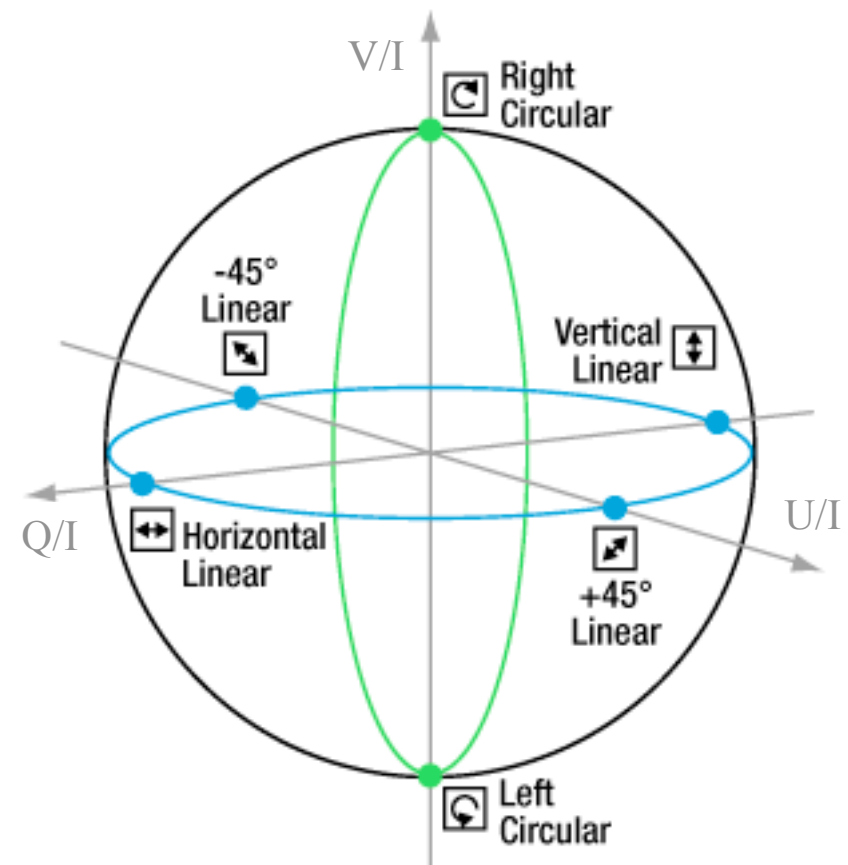
Answer:

4. An ellipse that is not tilted (so horizontal or vertical)

Answer:

## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

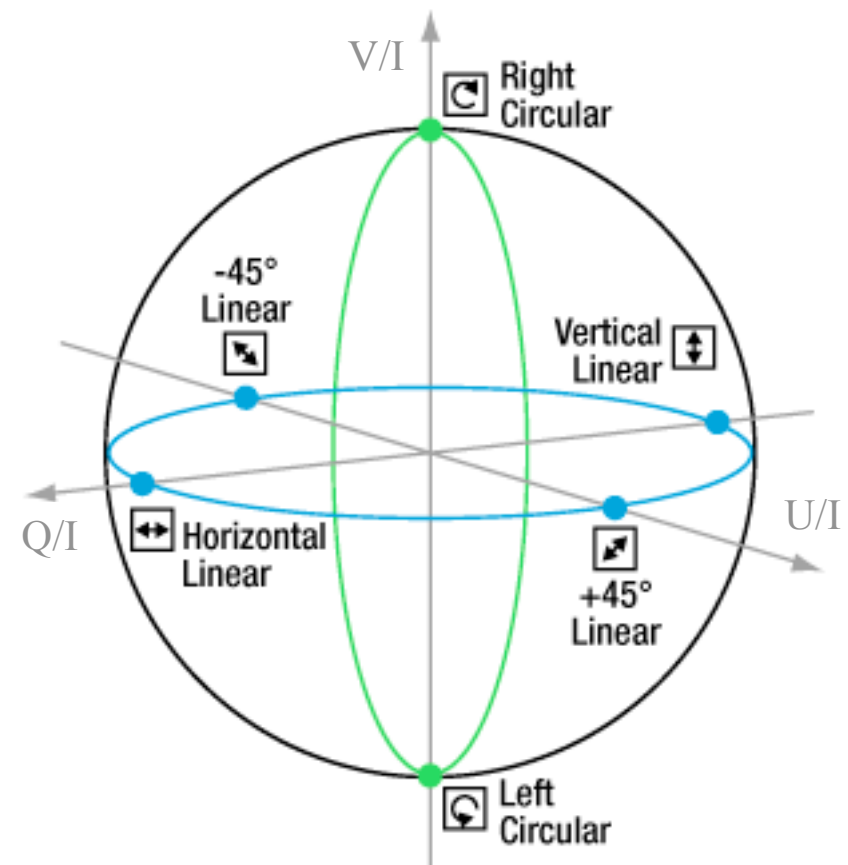
$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

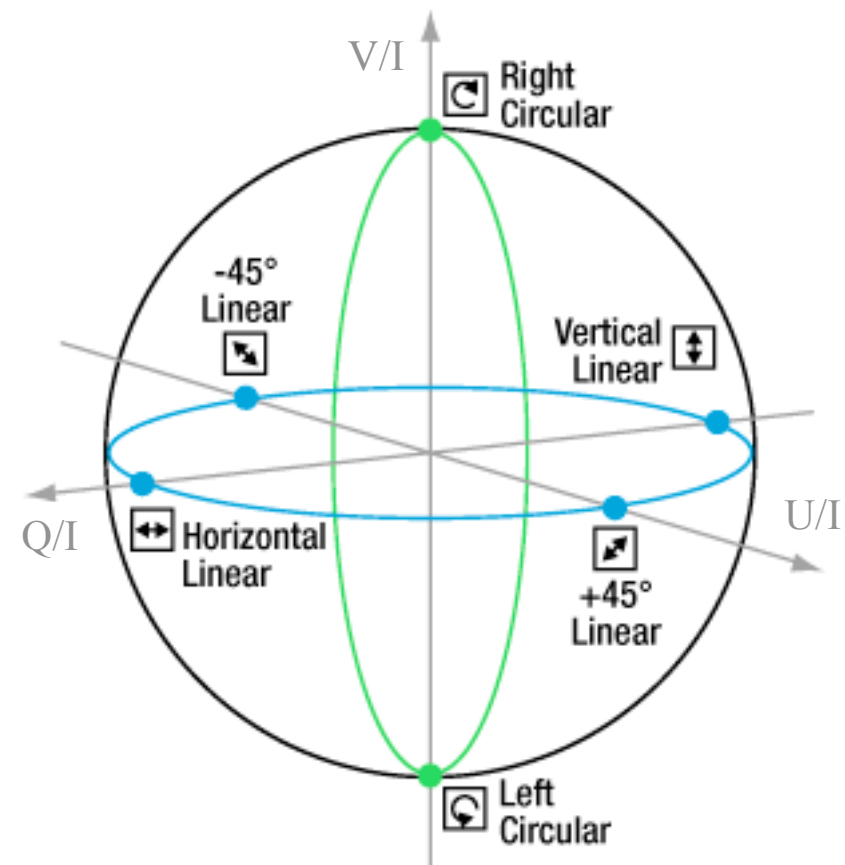


## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$





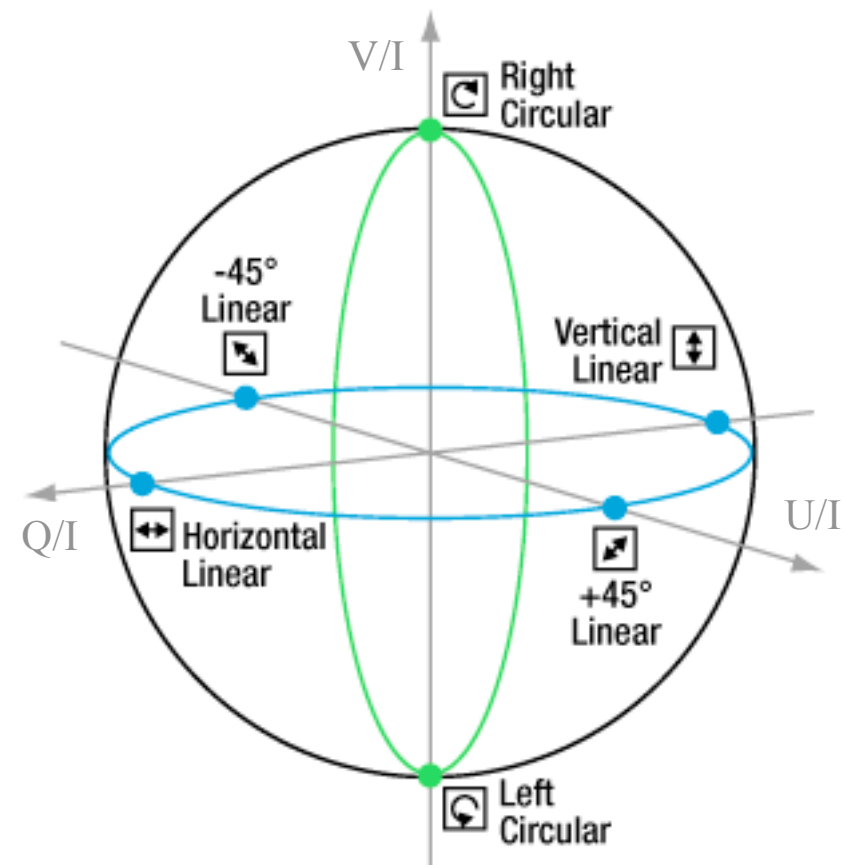
## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$

Therefore:



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

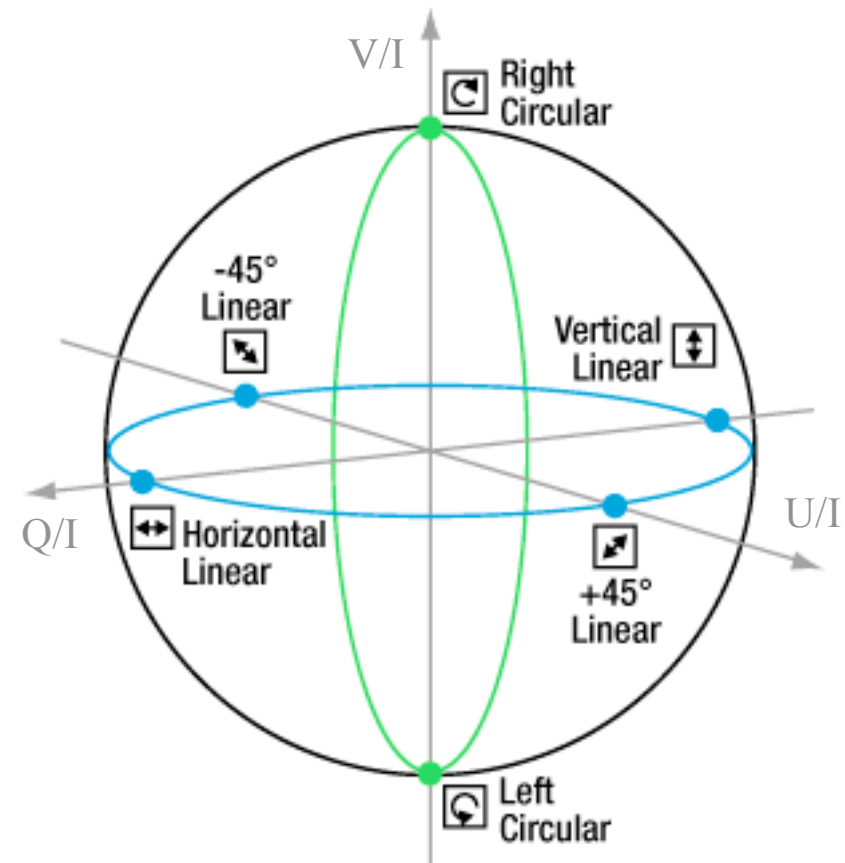
$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$

Therefore:

$$(2.22) \quad I^2 - Q^2 - U^2 - V^2 = 0$$



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

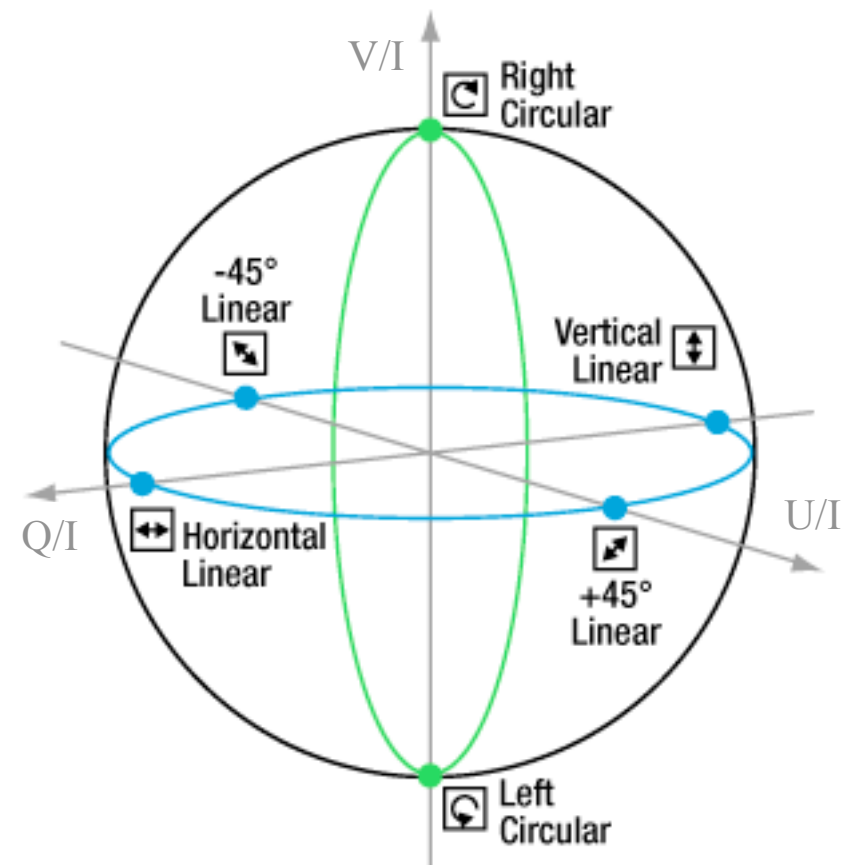
In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$

Therefore:

$$(2.22) \quad I^2 - Q^2 - U^2 - V^2 = 0$$

$$(2.24) \quad \frac{Q^2}{I^2} + \frac{U^2}{I^2} + \frac{V^2}{I^2} = 1$$



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

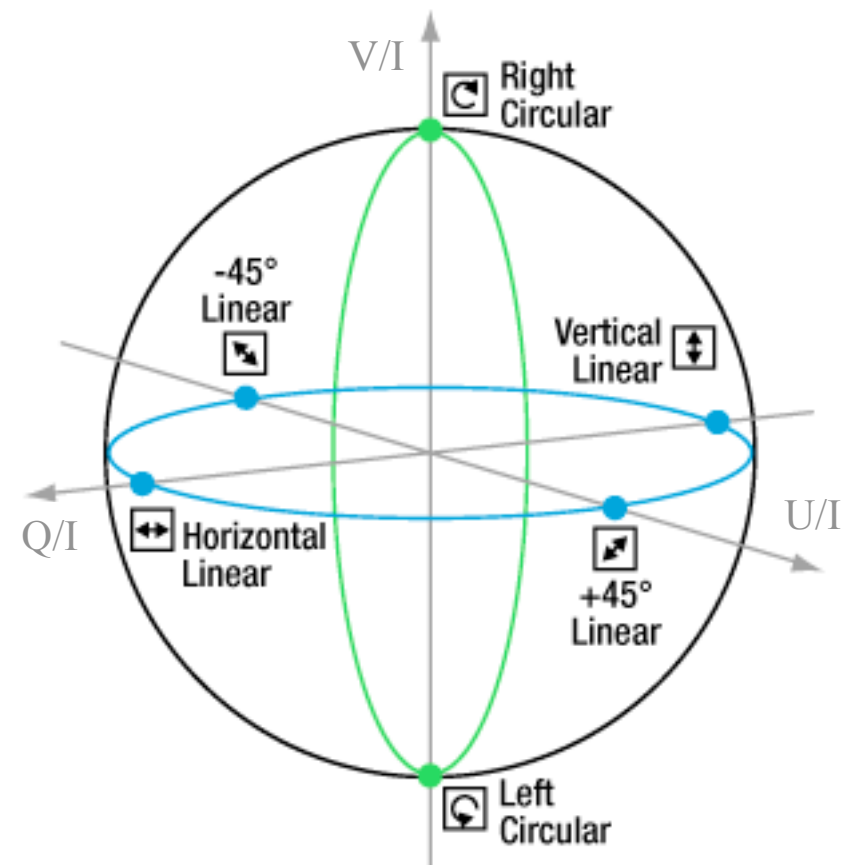
$$(2.18) \quad \det(C) = 0$$

Therefore:

$$(2.22) \quad I^2 - Q^2 - U^2 - V^2 = 0$$

$$(2.24) \quad \frac{Q^2}{I^2} + \frac{U^2}{I^2} + \frac{V^2}{I^2} = 1$$

Stokes V is the 'roundness' of the path



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$

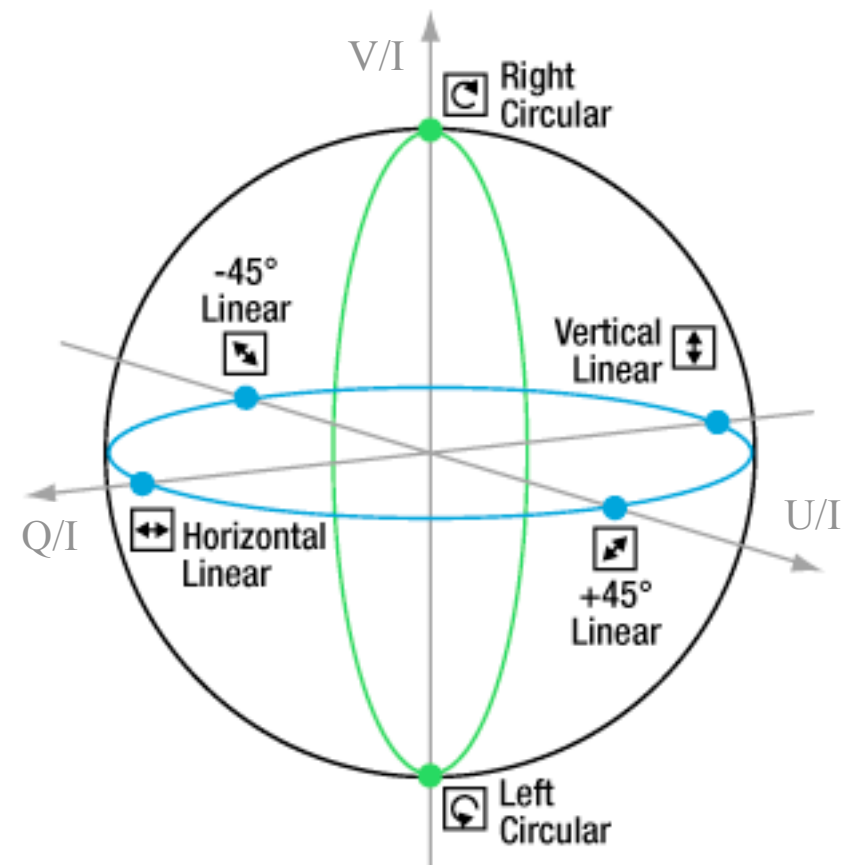
Therefore:

$$(2.22) \quad I^2 - Q^2 - U^2 - V^2 = 0$$

$$(2.24) \quad \frac{Q^2}{I^2} + \frac{U^2}{I^2} + \frac{V^2}{I^2} = 1$$

Stokes V is the 'roundness' of the path

Q is the 'horizontalness/verticalness' of the path



## 2.4 The Stokes parameters of a monochromatic, time-harmonic plane wave

$$(2.21) \quad C \equiv \begin{pmatrix} I+Q & U+iV \\ U-iV & I-Q \end{pmatrix}$$

In the previous section, we have seen that for a monochromatic plane wave:

$$(2.18) \quad \det(C) = 0$$

Therefore:

$$(2.22) \quad I^2 - Q^2 - U^2 - V^2 = 0$$

$$(2.24) \quad \frac{Q^2}{I^2} + \frac{U^2}{I^2} + \frac{V^2}{I^2} = 1$$

Stokes V is the 'roundness' of the path

Q is the 'horizontalness/verticalness' of the path

U is the 'tiltness' of the path (with max tilt = 45°)

