

University of Minnesota
School of Physics and Astronomy

2025 Fall Physics 8011
Quantum Field Theory I
Assignment Solution

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Question 1

Problem 48.5

The charged pion π^- is represented by a complex scalar field φ , the muon μ^- by a Dirac field \mathcal{M} , and the muon neutrino ν_μ by a spin-projected Dirac field $P_L \mathcal{N}$, where $P_L = \frac{1}{2}(1 - \gamma_5)$. The charged pion can decay to a muon and a muon antineutrino via the interaction

$$\mathcal{L}_1 = c_1 G_F f_\pi \partial_\mu \varphi \bar{\mathcal{M}} \gamma^\mu P_L \mathcal{N} + h.c., \quad (1)$$

where c_1 is the cosine of the *Cabibbo angle*, G_F is the *Fermi constant*, and f_π is the *pion decay constant*.

- (a) Compute the charged pion decay rate Γ .
- (b) The charged pion mass is $m_\pi = 139.6$ MeV, the muon mass is $m_\mu = 105.7$ MeV, and the muon neutrino mass is massless. The Fermi constant is $G_F = 1.166 \times 10^{-5}$ GeV $^{-2}$, and the cosine of the Cabibbo angle is measured in nuclear beta decays to be $c_1 = 0.974$. The measured value of the charged pion life time is $\tau = 2.6033 \times 10^{-8}$ s. Determine the value of f_π in MeV. Your result is too large by 0.8%, due to neglect of electromagnetic loop corrections.
- (c) The previous parts assume π^- always decay into $\mu^- \bar{\nu}_\mu$, but actually π^- can also decay into $e^- \bar{\nu}_e$. The charged pion, electron by a Dirac field \mathcal{M}_e , and the electron neutrino by a spin-projected Dirac field $P_L \mathcal{N}_e$ have the form of interaction

$$\mathcal{L}_2 = c_2 G_F f_\pi \partial_\mu \varphi \bar{\mathcal{M}}_e \gamma^\mu P_L \mathcal{N}_e + h.c. \quad (2)$$

Given the decay branching ratio of $\pi^- \rightarrow e^- \bar{\nu}_e$ is 1.230×10^{-4} , the decay branching ratio of $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$ is 99.9877%. Find the value of c_2 . For example, the electronic decay branching ratio is

$$\text{Br}(\pi^- \rightarrow e^- \bar{\nu}_e) = \frac{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu) + \Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)}. \quad (3)$$

The coupling of pion-electron is similar with the coupling of pion-muon, why pion favoring decay into muon instead of electron? ($m_e = 0.511$ MeV.)

Answer

Question 2

Consider QED with both electron and muon:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \sum_{l=e,\mu} (i\bar{\Psi}_l \not{d} \Psi_l - m_l \bar{\Psi}_l \Psi_l + \frac{g}{2} \bar{\Psi}_l \gamma^\mu \Psi_l A_\mu), \quad (4)$$

where both Ψ_e and Ψ_μ are Dirac fields. Compute the $\langle |\mathcal{T}^2| \rangle$ for $e^+e^- \rightarrow \mu^+\mu^-$. Then, compute its cross section σ . Eq. (11.22) and Eq. (11.30) should be useful.

Answer

Question 3

Consider classical field theory with two real scalar fields in (3+1)-dimension spacetime:

$$\mathcal{L}(x) = \sum_{a=1}^2 \left(-\frac{1}{2} \partial^\mu \phi_a \partial_\mu \phi_a \right) - V(x), \quad (5)$$

$$V(x) = - \sum_{a=1}^2 \left(\frac{1}{2} \mu^2 \phi_a \phi_a \right) + \frac{\lambda}{4} (\phi_1^2 + \phi_2^2)^2, \quad (6)$$

where μ and λ are positive real constants.

- (a) Show that the Lagrangian has an $SO(2)$ transformation symmetry:

$$\phi_1(x) \rightarrow \phi'_1(x) = \phi_1(x) \cos \alpha_0 - \phi_2(x) \sin \alpha_0, \quad (7)$$

$$\phi_2(x) \rightarrow \phi'_2(x) = \phi_1(x) \sin \alpha_0 + \phi_2(x) \cos \alpha_0, \quad (8)$$

- (b) Find the conjugate momentum $\Pi_1(x)$, $\Pi_2(x)$ of $\phi_1(x)$, $\phi_2(x)$. Find the Hamiltonian density $\mathcal{H}(x)$ in the terms of $\phi_a(x)$, $\Pi_a(x)$, and $\partial_i \phi_a(x)$.

- (c) Find the ground state in the basis of $\{\phi_r(x), \phi_\theta(x)\}$ where

$$\phi_1(x) = \phi_r(x) \cos(\phi_\theta(x)), \quad (9)$$

$$\phi_2(x) = \phi_r(x) \sin(\phi_\theta(x)), \quad (10)$$

with $\phi_r(x) \geq 0$ and $\phi_\theta(x) \in [0, 2\pi]$. Is the Lagrangian \mathcal{L} invariant under a continuous shift symmetry of $\phi_\theta(x) \rightarrow \phi_\theta(x) + \alpha_0$?

Hint: In general, finding the ground state is to find $\phi(x)$ s.t. minimize $H = \int \mathcal{H}(x) d^3x$; but for this problem, finding $\phi(x)$ to minimize $\mathcal{H}(x)$ is the same. If you have trouble with the above procedure, given the Lagrangian of this problem, one can simply find $\phi(x)$ s.t. minimize $V(x)$, which is the same as minimizing \mathcal{H} for this problem.

- (d) Now let's study the system's dynamics around the ground state.

$\phi_r(x)$ should fluctuate around $\sqrt{\frac{\mu^2}{\lambda}}$: $\phi_r(x) = \sqrt{\frac{\mu^2}{\lambda}} + f_r(x)$. $\phi_\theta(x)$ should fluctuate within $[0, 2\pi]$.

Show that $f_r(x)$ is a massive field and find its mass. Taking $f_\theta(x) \equiv \sqrt{\frac{\mu^2}{\lambda}} \phi_\theta(x)$ as the other scalar field, does $f_\theta(x)$ have a mass? Does \mathcal{L} have a continuous shift symmetry of $f_\theta(x) \rightarrow f_\theta(x) + \Lambda_0$?

Remark: This problem paves the road for your understanding of spontaneous symmetry breaking. We also see again that the symmetry groups of $SO(2)$ and $U(1)$ are isomorphic.

Remark: More to think about after solving the problems above: Note that we reparametrized the field into a non-linear realization, where you see the $U(1)$ symmetry explicitly. How do you interpret the kinetic term? How do you interpret the $f_r(x)$ field-dependent kinetic terms for $f_\theta(x)$? Is it canonically normalized? How does the field $f_\theta(x)$ relate to the original $SO(2)$ field $\phi_a(x)$? And again, is the ratio of

the field a linear redefinition of the field configuration? It is a non-linear realization because all powers of $f_\theta(x)/\sqrt{\frac{\mu^2}{\lambda}}$ need to enter. There is only a region of validity, that is $f_r(x) \ll \sqrt{\frac{\mu^2}{\lambda}}$

Answer

Question 4

Problem 66.3

Use the result of problem 66.2 to compute the anomalous dimension of m and the beta function for e in spinor electrodynamics in R_ξ gauge. You should find that the results are independent of ξ .

Remark:

$$\tilde{\Delta}^{\mu\nu}(k) = \frac{g^{\mu\nu} + (\xi - 1)k^\mu k^\nu/k^2}{k^2 - i\epsilon} \quad (11)$$

The book only choose the Feynman gauge ($\xi = 1$) to show the loop calculation and get $Z_{1,2,3,m}$. For arbitrary gauge choice ξ , we can repeat the calculation and get:

$$Z_3 = 1 - \frac{e^2}{6\pi^2} \left(\frac{1}{\epsilon} + \text{finite} \right) + \mathcal{O}(e^4), \quad \text{derived from photon propagator loop correction} \quad (12)$$

$$Z_2 = 1 - \xi \frac{e^2}{8\pi^2} \left(\frac{1}{\epsilon} + \text{finite} \right) + \mathcal{O}(e^4), \quad \text{derived from fermion propagator loop correction} \quad (13)$$

$$Z_m = 1 - (3 + \xi) \frac{e^2}{8\pi^2} \left(\frac{1}{\epsilon} + \text{finite} \right) + \mathcal{O}(e^4), \quad \text{derived from fermion mass loop correction} \quad (14)$$

$$Z_1 = 1 - \xi \frac{e^2}{8\pi^2} \left(\frac{1}{\epsilon} + \text{finite} \right) + \mathcal{O}(e^4), \quad \text{derived from vertex loop correction} \quad (15)$$

Use the above to finish this problem.

Answer

Question 5

Consider the following theory:

$$\mathcal{L} = \mathcal{L}_\phi^0 + \mathcal{L}_\Psi^0 + \mathcal{L}_A^0 + \mathcal{L}_I \quad (16)$$

$$= -\frac{1}{2}\partial^\mu\phi\partial_\mu\phi - \frac{1}{2}m_\phi^2\phi^2 + \bar{\Psi}(iD^\mu - m_\Psi) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + y\phi\bar{\Psi}\Psi. \quad (17)$$

The Dirac field Ψ is charged under a $U(1)$ gauge symmetry with a charge Q , and the gauge interaction strength is e . The $U(1)$ gauge field is A_μ , whose kinetic term is $\mathcal{L}_A^0 = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu}$. (This is part of the real-world calculation for the discovery mode for the Higgs boson, which gone through heroic phenomenological studies on predicting the Higgs properties.)

- (a) Draw the leading diagrams that enable $\phi \rightarrow \gamma\gamma$ decay. (The gauge field A_μ is identified as the photon field γ .)
 - (b) In the ϕ rest frame, write down the amplitude in the general d dimension. No need to carry out the loop integral at this point, but need to simplify the trace. (Notice that $k_\mu\epsilon^\mu(k) = 0$ in Lorenz gauge.)
 - (c) Does the integral have a UV divergence in $d = 4$ dimension (loop momentum goes to ∞)? Answer Yes or No with a few lines of argument.
 - (d) Does the integral have a singularity in $d = 4$ dimension when the Euclidean loop momentum squared \bar{q}^2 go to $-D$? Answer Yes or No with a few lines of argument. (For simplicity, assume that D is real and can be zero for some configuration of x_1, x_2, x_3 .)
 - (e) For $m_\Phi = 0$, calculate using dimensional regularization in $d = 4 - \epsilon$. Write down your final answer in the simplest form. (The final answer would be short.)
 - (f) Carry out the full calculation of the amplitude in Part b using dimensional regularization in $d = 4 - \epsilon$. Write down your final answer in the simplest form. (The full answer would be a long calculation.)
- Hint:** The following few equations, identities, and tricks, and the discussion around them might be helpful for you: Eq. (62.18), Eq. (47.18), Eq. (67.2).
- Remark:** No need to answer this, but one can think about it for fun. Recall that taking $\epsilon \rightarrow 0$ (from plus or minus direction?) get you back to $d = 4$. In such a limit, contrast your result in Part f and Part c and think about why.

Answer