

1 Earth's gravitational field and other perturbations

So far we have only considered the gravitational force acting between point masses. In reality, the Earth is not a point mass, neither a spherically symmetric mass distribution. In this section we will delve into the details of a more realistic model of the Earth's gravitational field.

1.1 Geopotential model

1.1.1 Continuous distribution of mass

In ?? we saw that the motion of a body orbiting another one can be described by a conic section. However, we have not been concerned about the mass distribution of the large body, in our case the Earth. In this section we will see that the motion of the smaller body, the satellite, is slightly perturbed by the mass distribution of the Earth as well as the presence of other forces such as atmospheric drag, solar radiation pressure, or the gravitational pull of the Moon and Sun, which we will discuss later on. Nevertheless, the perturbations are relatively small and the orbits of the satellites are still approximating ellipses, but as we will corroborate experimentally in ?? these perturbations are essential to obtain accurate results.

Consider a body confined in a compact region $\Omega \subseteq \mathbb{R}^3$ with a continuous density of mass $\rho : \Omega \rightarrow \mathbb{R}$. We would like to know the gravitational pull on a point mass m located at position $\mathbf{r} \in \Omega^c$ from the center of mass of the body. To do this, consider a covering of Ω in a set of disjoint cubes Q_i , $i = 1, \dots, N$, small enough to be considered as point masses and let $R_i := Q_i \cap \Omega$. Then, $\bigsqcup_{i=1}^N R_i = \Omega$. If each R_i has mass m_i , volume V_i , density ρ_i , and its center is located at $\mathbf{s}_i \in \mathbb{R}^3$, then the total gravitational acceleration \mathbf{g} exerted on m is the sum of the contributions of all the forces exerted by the N point masses, and it is given by:

$$\mathbf{g} = - \sum_{i=1}^N \frac{Gm_i}{\|\mathbf{r} - \mathbf{s}_i\|^3} (\mathbf{r} - \mathbf{s}_i) = - \sum_{i=1}^N \frac{G\rho_i}{\|\mathbf{r} - \mathbf{s}_i\|^3} (\mathbf{r} - \mathbf{s}_i) V_i \quad (1)$$

Note that Eq. (1) is a Riemann sum and so letting $N \rightarrow \infty$ we get:

$$\mathbf{g} = - \int_{\Omega} \frac{G\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|^3} (\mathbf{r} - \mathbf{s}) d^3\mathbf{s} \quad (2)$$

where $d^3\mathbf{s} := dx' dy' dz'$, if $\mathbf{s} = (x', y', z')$.

Theorem 1. Let Ω be a compact region in \mathbb{R}^3 with a continuous density of mass $\rho : \Omega \rightarrow \mathbb{R}$. Then, the gravitational acceleration field \mathbf{g} is conservative. That is, there exists a function $V : \mathbb{R}^3 \rightarrow \mathbb{R}$ such that $\mathbf{g} = -\nabla V$. We call V the *gravitational potential*.

Demostració. An easy computation shows that fixed $\mathbf{s} \in \mathbb{R}^3$ we have:

$$\nabla \left(\frac{1}{\|\mathbf{r} - \mathbf{s}\|} \right) = - \frac{1}{\|\mathbf{r} - \mathbf{s}\|^3} (\mathbf{r} - \mathbf{s}) \quad (3)$$

So we need to justify whether the following exchange between the gradient and the integral is correct:

$$\mathbf{g} = - \int_{\Omega} \frac{G\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|^3} (\mathbf{r} - \mathbf{s}) d^3\mathbf{s} = \int_{\Omega} G\rho(\mathbf{s}) \nabla \left(\frac{1}{\|\mathbf{r} - \mathbf{s}\|} \right) d^3\mathbf{s} = \nabla \int_{\Omega} \frac{G\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|} d^3\mathbf{s} \quad (4)$$

Without loss of generality it suffices to justify that

$$\frac{\partial}{\partial x} \int_{\Omega} \frac{\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|} d^3\mathbf{s} = \int_{\Omega} \frac{\partial}{\partial x} \left(\frac{\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|} \right) d^3\mathbf{s} \quad (5)$$

assuming $\mathbf{r} = (x, y, z)$ and $\mathbf{s} = (x', y', z')$. In order to apply the theorem of derivation under the integral sign we need to control $\frac{\partial}{\partial x} \left(\frac{\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|} \right) = -\rho(\mathbf{s}) \frac{x - x'}{\|\mathbf{r} - \mathbf{s}\|^3}$ by an integrable function $h(\mathbf{s})$. Using spherical coordinates centered at \mathbf{r} and writing $(\mathbf{r} - \mathbf{s})_{\text{sph}} = (\rho_{\mathbf{r}}, \theta, \phi)$, the integrand to bound becomes (in spherical coordinates):

$$\left| -\rho(\mathbf{s}) \frac{x - x'}{\|\mathbf{r} - \mathbf{s}\|^3} \rho_{\mathbf{r}}^2 \sin \phi \right| = |\rho(\mathbf{s})| \left| \frac{\rho_{\mathbf{r}} \cos \theta \sin \phi}{\rho_{\mathbf{r}}^3} \rho_{\mathbf{r}}^2 \sin \phi \right| \leq |\rho(\mathbf{s})| \leq K \quad (6)$$

where the last inequality follows for certain $K \in \mathbb{R}$ by Weierstrass theorem (ρ is continuous and Ω is compact). Thus, since $h(\mathbf{s}) = K$ is integrable, because Ω is bounded, the equality of Eq. (5) is correct. \square

Physically speaking, the gravitational force \mathbf{F} being conservative means that the work W done by the force along a path C

$$W = \int_C \mathbf{F} \cdot d\mathbf{s} \quad (7)$$

depends only on the initial and final positions of it.

1.1.2 Laplace's equation for V

Theorem 2. Consider a distribution of matter of density ρ in a compact region Ω . Then, the gravitational potential V satisfies the Laplace equation

$$\Delta V = 0 \quad (8)$$

for all points outside Ω ¹.

Demostració. Recall that $\Delta V = \mathbf{div}(\nabla V)$. So since $\mathbf{g} = \nabla V$ it suffices to prove that $\mathbf{div}(\mathbf{g}) = 0$. Note that if $\mathbf{r} \in \Omega^c$, then $\exists \delta > 0$ such that $\|\mathbf{r} - \mathbf{s}\| \geq \delta > 0 \forall \mathbf{s} \in \Omega$ because Ω is closed. As a result, $\frac{\mathbf{r} - \mathbf{s}}{\|\mathbf{r} - \mathbf{s}\|^3}$ is differentiable and:

$$\begin{aligned} \mathbf{div} \left(\frac{\mathbf{r} - \mathbf{s}}{\|\mathbf{r} - \mathbf{s}\|^3} \right) &= \frac{\partial}{\partial x} \left(\frac{x - x'}{\|\mathbf{r} - \mathbf{s}\|^3} \right) + \frac{\partial}{\partial y} \left(\frac{y - y'}{\|\mathbf{r} - \mathbf{s}\|^3} \right) + \frac{\partial}{\partial z} \left(\frac{z - z'}{\|\mathbf{r} - \mathbf{s}\|^3} \right) = \\ &= \frac{\|\mathbf{r} - \mathbf{s}\|^2 - 3(x - x')^2}{\|\mathbf{r} - \mathbf{s}\|^5} + \frac{\|\mathbf{r} - \mathbf{s}\|^2 - 3(y - y')^2}{\|\mathbf{r} - \mathbf{s}\|^5} + \frac{\|\mathbf{r} - \mathbf{s}\|^2 - 3(z - z')^2}{\|\mathbf{r} - \mathbf{s}\|^5} = 0 \end{aligned}$$

Hence, as in Theorem 1, we have that:

$$\left| \rho(\mathbf{s}) \frac{\|\mathbf{r} - \mathbf{s}\|^2 - 3(x - x')^2}{\|\mathbf{r} - \mathbf{s}\|^5} \right| \leq \frac{4|\rho(\mathbf{s})|}{\|\mathbf{r} - \mathbf{s}\|^3} \leq \frac{4|\rho(\mathbf{s})|}{\delta^3} \quad (9)$$

which is integrable by Weierstrass theorem. Thus, by the theorem of derivation under the integral sign:

$$\mathbf{div}(\mathbf{g}) = -\mathbf{div} \int_{\Omega} \frac{G\rho(\mathbf{s})}{\|\mathbf{r} - \mathbf{s}\|^3} (\mathbf{r} - \mathbf{s}) d^3\mathbf{s} = -\int_{\Omega} G\rho(\mathbf{s}) \mathbf{div} \left(\frac{\mathbf{r} - \mathbf{s}}{\|\mathbf{r} - \mathbf{s}\|^3} \right) d^3\mathbf{s} = 0 \quad (10)$$

\square

So far we have seen that the gravitational potential V satisfies the Laplace equation. If, moreover, we choose the origin of potential to be at the infinity, that is, if we impose $\lim_{\|\mathbf{r}\| \rightarrow \infty} V = 0$, then the gravitational potential created by a distribution of mass in a compact region Ω is a solution of the following exterior Dirichlet problem:

$$\begin{cases} \Delta V = 0 & \text{in } \Omega^c \\ V = f & \text{on } \partial\Omega \\ \lim_{\|\mathbf{r}\| \rightarrow \infty} V = 0 \end{cases} \quad (11)$$

If Ω represents the Earth, then $f = f(\theta, \phi)$ is the boundary condition concerning the gravitational potential at the surface of the Earth as a function of the longitude θ and colatitude ϕ .

We will see now that Eq. (11) has a unique solution. To do that we invoke the maximum principle, which we will not prove here (see [Eva10] for more details).

Theorem 3 (Maximum principle). Let $U \subset \mathbb{R}^n$ be open and bounded and $u \in \mathcal{C}^2(U) \cap \mathcal{C}(\overline{U})$. Suppose that u is harmonic within U , that is, $\Delta u = 0$ in U . Then, $\max_{\overline{U}} u = \max_{\partial U} u$.

¹It can be seen that V satisfies in fact the Poisson equation $\Delta V = 4\pi G\rho$ for any point $\mathbf{r} \in \mathbb{R}^3$, which reduces to Laplace equation when $\mathbf{r} \in \Omega^c$, because there we have $\rho(\mathbf{r}) = 0$.

Theorem 4. The Dirichlet problem of Eq. (11) has a unique solution.

Demostració. Suppose we have two solutions V_1, V_2 of Eq. (11). Then, $W := V_1 - V_2$ is harmonic in Ω^c , $W = 0$ on $\partial\Omega$ and $\lim_{\|\mathbf{r}\| \rightarrow \infty} W = 0$. So $\forall \varepsilon > 0$, $\exists n \in \mathbb{N}$ large enough such that $\Omega \subseteq B(0, n)$ and $|W| \leq \varepsilon$ on $\mathbb{R}^3 \setminus \overline{B(0, n)}$. Thus, by the maximum principle, $|W| \leq \varepsilon$ on $\overline{B(0, n)} \cap \Omega^c$. Since the ε is arbitrary, we must have $W = 0$ on Ω^c , that is, $V_1 = V_2$. \square

Now that we know that the Dirichlet problem has a unique solution, we can proceed to find it. In the next section we will construct an explicit solution for the gravitational potential created by the Earth.

1.2 Spherical harmonics

1.2.1 Legendre polynomials, regularity and orthonormality

In this section we aim to introduce a class of functions that will appear later on in the general solution of the Laplace equation (see Section 1.2.2). To accomplish this, we need first to introduce the Legendre polynomials. There are several ways to define them, but the most convenient one for our purposes is from the following differential equation. Consider the following second-order differential equation called *Legendre differential equation*:

$$(1 - x^2)y'' - 2xy' + \lambda y = 0 \quad (12)$$

for $\lambda \in \mathbb{R}$. This equation can be rewritten as:

$$((1 - x^2)y')' + \lambda y = 0 \quad (13)$$

Seeking for analytic solutions of this equation using the power series method [Mez], i.e. looking for solutions of the form $y(x) = \sum_{j=0}^{\infty} a_j x^j$, we see that:

$$\begin{aligned} 0 &= (1 - x^2) \sum_{j=2}^{\infty} a_j(j-1)jx^{j-2} - 2x \sum_{j=1}^{\infty} a_j j x^{j-1} + \lambda \sum_{j=0}^{\infty} a_j x^j = \sum_{j=0}^{\infty} a_{j+2}(j+1)(j+2)x^j - \\ &\quad - \sum_{j=0}^{\infty} a_j(j-1)jx^j - \sum_{j=0}^{\infty} 2a_j j x^j + \sum_{j=0}^{\infty} \lambda a_j x^j = \sum_{j=0}^{\infty} [a_{j+2}(j+1)(j+2) - a_j(j(j+1) - \lambda)]x^j \end{aligned} \quad (14)$$

Equating the general term of the series to 0 we obtain this recursion:

$$a_{j+2} = \frac{j(j+1) - \lambda}{(j+1)(j+2)} a_j, \quad j = 0, 1, 2, \dots \quad (15)$$

From here we can obtain two independent solutions by setting the initial conditions a_0 and a_1 of the iteration. For example, setting $a_1 = 0$ we obtain a series that has only even powers of x . On the other hand, setting $a_0 = 0$ we obtain a series that has only odd powers of x . These two series converge on the interval $(-1, 1)$ by the ratio test (by looking at Eq. (15)) and can be expressed compactly as [Mez]:

$$y_e(x) = a_0 \sum_{j=0}^{\infty} \left[\prod_{k=0}^{j-1} (2k(2k+1) - \lambda) \right] \frac{x^{2j}}{(2j)!} \quad y_o(x) = a_1 \sum_{j=0}^{\infty} \left[\prod_{k=0}^{j-1} ((2k+1)(2k+2) - \lambda) \right] \frac{x^{2j+1}}{(2j+1)!} \quad (16)$$

Here, the empty product (that is, when k ranges from 0 to -1) is defined to be 1. However, for each $\lambda \in \mathbb{R}$ either one of these series or both diverge at $x = \pm 1$, as they behave as the harmonic series in a neighborhood of $x = \pm 1$. We are interested, though, in the solutions that remain bounded on the whole interval $[-1, 1]$. Looking at the expressions of Eq. (16) one can check that the only possibility to make the series converge in $[-1, 1]$ is when $\lambda = n(n+1)$, $n \in \mathbb{N} \cup \{0\}$. In this case, for each $n \in \mathbb{N} \cup \{0\}$ exactly one of the series is in fact a polynomial of degree n . If, furthermore, we choose a_0 or a_1 be such that the polynomial evaluates to 1 at $x = 1$, these polynomials are called *Legendre polynomials*, and they are denoted by $P_n(x)$. The other (divergent) series is usually denoted in the literature by $Q_n(x)$ (check [RHB99; Mez]) and it is independent of $P_n(x)$. Thus, the general solution of Eq. (13) for $\lambda = n(n+1)$ can be expressed as a linear combination of P_n and Q_n , because the space of solutions form a vector space of dimension 2.

n	$P_n(x)$
0	1
1	x
2	$\frac{1}{2}(3x^2 - 1)$
3	$\frac{1}{2}(5x^3 - 3x)$
4	$\frac{1}{8}(35x^4 - 30x^2 + 3)$
5	$\frac{1}{8}(63x^5 - 70x^3 + 15x)$
6	$\frac{1}{16}(231x^6 - 315x^4 + 105x^2 - 5)$
7	$\frac{1}{16}(429x^7 - 693x^5 + 315x^3 - 35x)$

Taulla 1: First eight Legendre polynomials

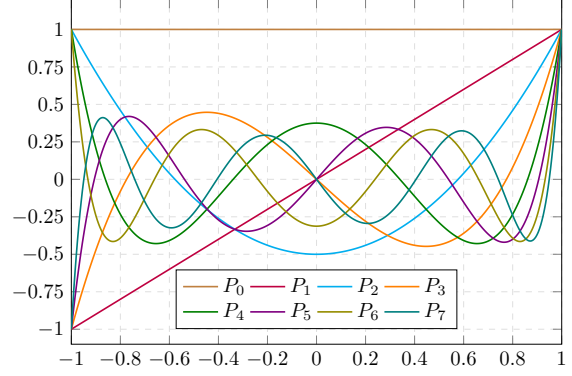


Figure 1: Graphic representation of the first eight Legendre polynomials.

The following proposition will be of our interest in the next section [RHB99].

Proposition 5. Let $y(x)$ be a solution to the Legendre differential equation. Then, $\forall m \in \mathbb{N} \cup \{0\}$ the function

$$w_m(x) = (1 - x^2)^{m/2} \frac{d^m y(x)}{dx^m} \quad (17)$$

solves the *general Legendre differential equation*:

$$(1 - x^2)y'' - 2xy' + \left(\lambda - \frac{m^2}{1 - x^2}\right)y = 0 \quad (18)$$

In particular if $\lambda = n(n + 1)$ for $n \in \mathbb{N} \cup \{0\}$, then $w_m(x)$ is denoted as

$$P_{n,m}(x) := (1 - x^2)^{m/2} \frac{d^m P_n}{dx^m} \quad (19)$$

and it is called the *associated Legendre polynomial* of degree n and order m .

Note that although we opted to call the functions $P_{n,m}$ as polynomials, they are only true polynomials when m is even. But we have opted to call them in that manner as it is the common practice in the literature (see [Wei; RHB99; Mez]).

Moreover, from the definition of $P_{n,m}$, we can see that $P_{n,0} = P_n$ and that $P_{n,m} = 0$ if $m > n$. So we can restrict the domain of m to the set $\{0, 1, \dots, n\}$.

n	$P_{n,1}(x)$	n	$P_{n,2}(x)$
1	$\sqrt{1 - x^2}$	2	$3(1 - x^2)$
2	$3x\sqrt{1 - x^2}$	3	$15x(1 - x^2)$
3	$\frac{3}{2}(5x^2 - 1)\sqrt{1 - x^2}$	4	$\frac{15}{2}(7x^2 - 1)(1 - x^2)$
4	$\frac{5}{2}x(7x^2 - 3)\sqrt{1 - x^2}$	5	$\frac{105}{2}x(3x^2 - 1)(1 - x^2)$
5	$\frac{15}{8}(21x^4 - 14x^2 + 1)\sqrt{1 - x^2}$	6	$\frac{105}{8}(33x^4 - 18x^2 + 1)(1 - x^2)$

Taulla 2: First associated Legendre polynomials for $m = 1$ and $m = 2$.

Definition 6. Let $n \in \mathbb{N} \cup \{0\}$ and $m \in \{0, 1, \dots, n\}$. We define the *real spherical harmonics* $Y_{n,m}^c$ and $Y_{n,m}^s$ as:

$$Y_{n,m}^c(\theta, \phi) = \sqrt{(2 - \delta_{0,m})(2n + 1) \frac{(n - m)!}{(n + m)!}} P_{n,m}(\cos \phi) \cos m\theta \quad (20)$$

$$Y_{n,m}^s(\theta, \phi) = \sqrt{(2 - \delta_{0,m})(2n + 1) \frac{(n - m)!}{(n + m)!}} P_{n,m}(\cos \phi) \sin m\theta \quad (21)$$

Here the coordinates $\theta \in [0, 2\pi)$ and $\phi \in [0, \pi]$ are the spherical coordinates such that any point on the sphere can be written uniquely as $(x, y, z) = (\sin \phi \cos \theta, \sin \phi \sin \theta, \cos \phi)$. The factor $N_{n,m} := \sqrt{(2 - \delta_{0,m})(2n + 1) \frac{(n - m)!}{(n + m)!}}$ is called the *normalization factor* of the spherical harmonics and $\delta_{0,m}$ is the Kronecker delta. Its appearance here will become clear in the next section.

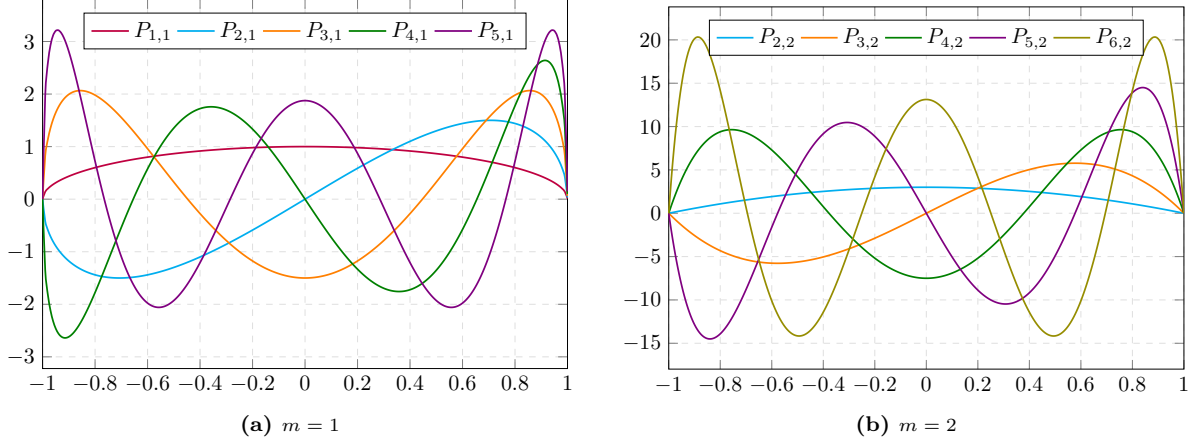


Figure 2: Graphic representation of the first associated Legendre polynomials for $m = 1$ and $m = 2$.

n	m	$Y_{n,m}^c(\theta, \phi)$	n	m	$Y_{n,m}^c(\theta, \phi)$
0	0	1	2	2	$\frac{\sqrt{15}}{2}(\sin \phi)^2 \cos 2\theta$
1	0	$\sqrt{3} \cos \phi$	3	0	$\frac{\sqrt{7}}{2} \cos \phi (5(\cos \phi)^2 - 3)$
1	1	$\sqrt{3} \sin \phi \cos \theta$	3	1	$\frac{\sqrt{42}}{4} (5(\cos \phi)^2 - 1) \sin \phi \cos \theta$
2	0	$\frac{\sqrt{5}}{2} (3(\cos \phi)^2 - 1)$	3	2	$\frac{\sqrt{105}}{2} (\sin \phi)^2 \cos \phi \cos 2\theta$
2	1	$\sqrt{15} \sin \phi \cos \phi \cos \theta$	3	3	$\frac{\sqrt{70}}{4} (\sin \phi)^3 \cos 3\theta$

Taula 3: First cosine spherical harmonics.

The associated Legendre polynomials satisfy an orthogonality relation:

Lemma 7. Let $n_1, n_2 \in \mathbb{N} \cup \{0\}$ and $m \leq \min\{n_1, n_2\}$. Then:

$$\int_0^1 P_{n_1,m}(x) P_{n_2,m}(x) dx = \frac{2}{2n_1 + 1} \frac{(n_1 + m)!}{(n_1 - m)!} \delta_{n_1, n_2} \quad (22)$$

where δ_{n_1, n_2} denotes the Kronecker delta.

Similarly, it can be shown that the spherical harmonics form an orthonormal family of functions:

Proposition 8. The family of spherical harmonics $\{Y_{n,m}^c(\theta, \phi), Y_{n,m}^s(\theta, \phi) : n \in \mathbb{N} \cup \{0\}, m \leq n\}$ is orthonormal in the following sense:

$$\frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi Y_{n_1, m_1}^i(\theta, \phi) Y_{n_2, m_2}^j(\theta, \phi) d\Omega = \delta_{n_1, n_2} \delta_{m_1, m_2} \delta_{i, j} \quad (23)$$

where $d\Omega = \sin \phi d\phi d\theta$ is the solid angle element, which measures the element of area on a sphere of radius 1.

Demostració. Let $N_{n_1, m_1}, N_{n_2, m_2}$ be the normalization factors of the spherical harmonics $Y_{n_1, m_1}, Y_{n_2, m_2}$ respectively. Note that we can separate the variables in the integral of Eq. (23). So if $i \neq j$, the integral over θ becomes $\int_0^{2\pi} \cos(m_1\theta) \sin(m_2\theta) d\theta$ which is equal to 0 regardless of the values of m_1 and m_2 . So from now on assume that $i = j$. Due to the symmetry between the cosine and the sine we can suppose that $i = c$. Thus:

$$\begin{aligned} \int_0^{2\pi} \int_0^\pi Y_{n_1, m_1}^i(\theta, \phi) Y_{n_2, m_2}^j(\theta, \phi) d\Omega &= \\ &= N_{n_1, m_1} N_{n_2, m_2} \int_0^\pi P_{n_1, m_1}(\cos \phi) P_{n_2, m_2}(\cos \phi) \sin \phi d\phi \int_0^{2\pi} \cos(m_1\theta) \cos(m_2\theta) d\theta \quad (24) \end{aligned}$$

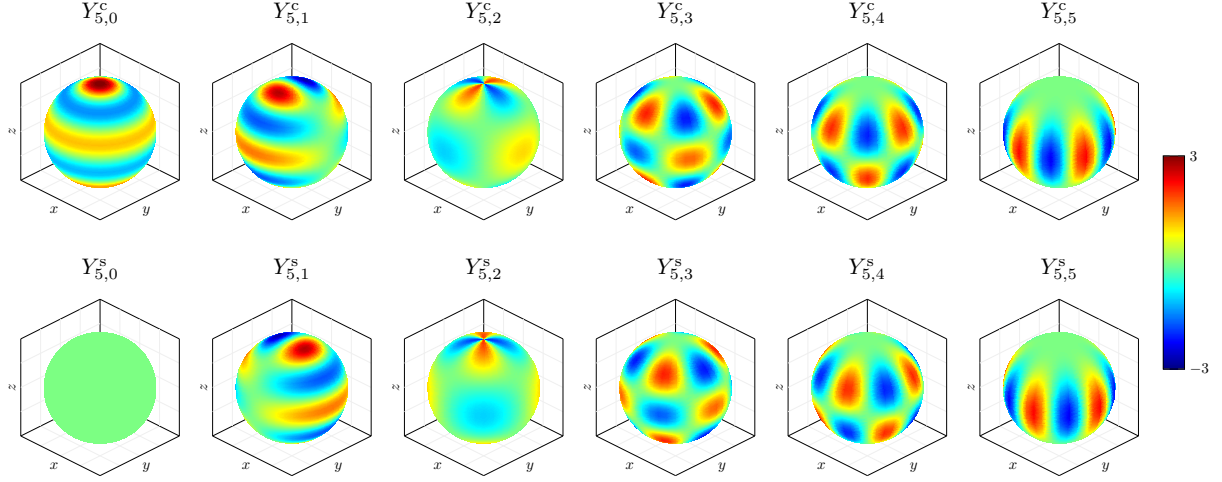


Figure 3: 3D color gradient representation of the spherical harmonics of degree $n = 5$. The first row correspond to the cosine spherical harmonics and the second row correspond to the sine spherical harmonics.

An easy check shows that if $m_1 \neq m_2$ then the integral over θ is zero (and the same applies with sines). So suppose $m_1 = m_2 = m$. In that case, if $m \neq 0$ we have $\int_0^{2\pi} (\cos m\theta)^2 d\theta = \int_0^{2\pi} (\sin m\theta)^2 d\theta = \pi$ and if $m = 0$, the cosine integral evaluates to 2π whereas the sine integral is 0. We can omit this latter case because $Y_{n,0}^s$ is identically zero. Thus:

$$\frac{2\pi}{2 - \delta_{0,m}} N_{n_1,m} N_{n_2,m} \int_0^\pi P_{n_1,m}(\cos \phi) P_{n_2,m}(\cos \phi) \sin \phi d\phi = \frac{2\pi}{2 - \delta_{0,m}} N_{n_1,m} N_{n_2,m} \int_{-1}^1 P_{n_1,m}(x) P_{n_2,m}(x) dx \quad (25)$$

By [Lemma 7](#) this latter integral is $\frac{2}{2n_1+1} \frac{(n_1+m)!}{(n_1-m)!} \delta_{n_1,n_2}$. Finally, if $n_1 = n_2 = n$, putting all normalization factors together we get:

$$\frac{2\pi}{2 - \delta_{0,m}} N_n^m N_n^m \frac{2}{2n+1} \frac{(n+m)!}{(n-m)!} = 4\pi \quad (26)$$

□

Moreover, an important result in the Sturm-Liouville Theory of second order differential equations ([\[Wikb; Wan+09\]](#)) says that the family of spherical harmonics $\{Y_{n,m}^c(\theta, \phi), Y_{n,m}^s(\theta, \phi) : n \in \mathbb{N} \cup \{0\}, m \leq n\}$ form a complete set in the sense that any smooth function defined on the sphere $f : S^2 \rightarrow \mathbb{R}$ can be expanded in a series of spherical harmonics:

$$f(\theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n (c_{n,m} Y_{n,m}^c(\theta, \phi) + s_{n,m} Y_{n,m}^s(\theta, \phi)) \quad (27)$$

This will be useful in [Section 1.2.3](#) when expanding the gravitational potential created by the Earth at some arbitrary point in spherical harmonics.

1.2.2 Laplace's equation in spherical coordinates

We start first with this proposition that give us the laplacian of a function in spherical coordinates.

Proposition 9. Let $f : \mathbb{R}^3 \rightarrow \mathbb{R}$ be a twice-differentiable function. Then:

$$\Delta f = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \phi} \frac{\partial}{\partial \phi} \left(\sin \phi \frac{\partial f}{\partial \phi} \right) + \frac{1}{r^2 (\sin \phi)^2} \frac{\partial^2 f}{\partial \theta^2} \quad (28)$$

where $r \in [0, \infty)$ denotes the radial distance, $\theta \in [-\pi, \pi)$ denotes the longitude, and $\phi \in [0, \pi]$, the colatitude:

$$\begin{aligned} x &= r \sin \phi \cos \theta \\ y &= r \sin \phi \sin \theta \\ z &= r \cos \phi \end{aligned} \quad (29)$$

We are now interested in solving the Laplace equation. [Theorem 10](#) gives the solution of it as a function of the spherical harmonics.

Theorem 10. The regular solutions in a bounded region $\Omega \subseteq \mathbb{R}^3$ such that $0 \notin \overline{\Omega}$ of the Laplace equation in spherical coordinates are of the form

$$f(r, \theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n (a_n r^n + b_n r^{-n-1}) P_{n,m}(\cos \phi) (c_{n,m} \cos(m\theta) + s_{n,m} \sin(m\theta)) \quad (30)$$

$$= \sum_{n=0}^{\infty} \sum_{m=0}^n (a_n r^n + b_n r^{-n-1}) (\tilde{c}_{n,m} Y_{n,m}^c(\theta, \phi) + \tilde{s}_{n,m} Y_{n,m}^s(\theta, \phi)) \quad (31)$$

where $a_n, b_n, c_{n,m}, s_{n,m}, \tilde{c}_{n,m}, \tilde{s}_{n,m} \in \mathbb{R}$.

Demostració. Let $f(r, \theta, \phi)$ be a solution of the Laplace equation:

$$\Delta f = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \phi} \frac{\partial}{\partial \phi} \left(\sin \phi \frac{\partial f}{\partial \phi} \right) + \frac{1}{r^2 (\sin \phi)^2} \frac{\partial^2 f}{\partial \theta^2} \quad (32)$$

Using separation variables $f(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi)$ we can write:

$$\frac{\Theta\Phi}{r^2} (r^2 R')' + \frac{R\Theta}{r^2 \sin \phi} (\sin \phi \Phi')' + \frac{R\Phi}{r^2 (\sin \phi)^2} \Theta'' = 0 \quad (33)$$

Here, we are making and abuse of notation denoting all the derivatives with a prime, but the reader should have no confusion with it. Isolating R from Θ and Φ yields:

$$\frac{(r^2 R')'}{R} = -\frac{1}{\sin \phi \Phi} (\sin \phi \Phi')' - \frac{1}{(\sin \phi)^2 \Theta} \Theta'' \quad (34)$$

Since the left-hand side depends entirely on r and the right-hand side does not, it follows that both sides must be constant. Therefore:

$$\frac{(r^2 R')'}{R} = \lambda \quad (35)$$

$$\frac{1}{\sin \phi \Phi} (\sin \phi \Phi')' + \frac{1}{(\sin \phi)^2 \Theta} \Theta'' = -\lambda \quad (36)$$

with $\lambda \in \mathbb{R}$. Similarly, separating variables from [Eq. \(36\)](#) we obtain that the equations

$$\frac{1}{\Theta} \Theta'' = -m^2 \quad (37)$$

$$\frac{\sin \phi}{\Phi} (\sin \phi \Phi')' + \lambda (\sin \phi)^2 = m^2 \quad (38)$$

must be constant with $m \in \mathbb{C}$ (a priori). The solution to the well-known [Eq. \(37\)](#) is a linear combination of the $\cos(m\theta)$ and $\sin(m\theta)$. Note, though, that since Θ must be a 2π -periodic function, that is satisfying $\Theta(\theta + 2\pi) = \Theta(\theta) \forall \theta \in \mathbb{R}$, m must be an integer. On the other hand making the change of variables $x = \cos \phi$ and $y = \Phi(\phi)$ in [Eq. \(38\)](#) and using the chain rule, that equation becomes:

$$(1-x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + \left(\lambda - \frac{m^2}{1-x^2} \right) y = 0 \quad (39)$$

which is the associate Legendre equation. We have argued in [Proposition 5](#) that we need $\lambda = n(n+1)$ and $m \leq n$ in order to obtain regular solutions at $x = \cos \phi = \pm 1$. Moreover, these solutions are $P_{n,m}(\cos \phi)$.

Finally, note that equation [Eq. \(35\)](#) is a Cauchy-Euler equation (check [\[Wika\]](#)) and so the general solution of it is given by

$$R(r) = c_1 r^n + c_2 r^{-n-1} \quad (40)$$

because $\lambda = n(n+1)$ (the reader may check that r^n and r^{-n-1} are indeed two independent solutions of [Eq. \(35\)](#)). So the general solution becomes a linear combination of each solution found varying $n \in \mathbb{N} \cup \{0\}$ and $m \in \{0, 1, \dots, n\}$:

$$f(r, \theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n (a_n r^n + b_n r^{-n-1}) P_{n,m}(\cos \phi) (c_{n,m} \cos(m\theta) + s_{n,m} \sin(m\theta)) \quad (41)$$

□

We ignore the singularity at $r = 0$ of Eq. (31) from now. In the next section we discuss it in more detail.

1.2.3 Expansion in spherical harmonics

We have just seen that if V satisfies the exterior Dirichlet problem for the Laplace equation, then, by uniqueness of solutions, it can be expressed as:

$$V(r, \theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n (a_n r^n + b_n r^{-n-1}) (\tilde{c}_{n,m} Y_{n,m}^c(\theta, \phi) + \tilde{s}_{n,m} Y_{n,m}^s(\theta, \phi)) \quad (42)$$

for some $a_n, b_n, \tilde{c}_{n,m}, \tilde{s}_{n,m} \in \mathbb{R}$. If we impose V to satisfy the third condition of Eq. (11), we must have $a_n = 0$. Finally, if we choose R_{\oplus} as a reference radius for a spherical model of the Earth, using the boundary condition on $\partial\Omega$

$$f(\theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{b_n}{R_{\oplus}^{n+1}} (\tilde{c}_{n,m} Y_{n,m}^c(\theta, \phi) + \tilde{s}_{n,m} Y_{n,m}^s(\theta, \phi)) \quad (43)$$

and the orthogonality of the spherical harmonics, we can deduce that the coefficients $b_n \tilde{c}_{n,m}$ and $b_n \tilde{s}_{n,m}$ are given by:

$$b_n \tilde{c}_{n,m} = \frac{R_{\oplus}^{n+1}}{4\pi} \int_0^{2\pi} \int_0^{\pi} f(\theta, \phi) Y_{n,m}^c(\theta, \phi) \sin \phi \, d\phi \, d\theta \quad (44)$$

$$b_n \tilde{s}_{n,m} = \frac{R_{\oplus}^{n+1}}{4\pi} \int_0^{2\pi} \int_0^{\pi} f(\theta, \phi) Y_{n,m}^s(\theta, \phi) \sin \phi \, d\phi \, d\theta \quad (45)$$

Hence, introducing the gravitational constant G and the Earth's mass M_{\oplus} into the equation, our final expression for the gravitational potential is

$$V(r, \theta, \phi) = \frac{GM_{\oplus}}{R_{\oplus}} \sum_{n=0}^{\infty} \sum_{m=0}^n \left(\frac{R_{\oplus}}{r} \right)^{n+1} (\bar{C}_{n,m} Y_{n,m}^c(\theta, \phi) + \bar{S}_{n,m} Y_{n,m}^s(\theta, \phi)) \quad (46)$$

where the coefficients $\bar{C}_{n,m}, \bar{S}_{n,m} \in \mathbb{R}$ are given by the formulas:

$$\bar{C}_{n,m} = \frac{1}{4\pi} \frac{R_{\oplus}}{GM_{\oplus}} \int_0^{2\pi} \int_0^{\pi} f(\theta, \phi) Y_{n,m}^c(\theta, \phi) \sin \phi \, d\phi \, d\theta \quad (47)$$

$$\bar{S}_{n,m} = \frac{1}{4\pi} \frac{R_{\oplus}}{GM_{\oplus}} \int_0^{2\pi} \int_0^{\pi} f(\theta, \phi) Y_{n,m}^s(\theta, \phi) \sin \phi \, d\phi \, d\theta \quad (48)$$

The coefficients $\bar{C}_{n,m}, \bar{S}_{n,m}$ are called *geopotential coefficients*, and they describe the dependence on the Earth's internal structure. They are obtained from observation of the perturbations seen in the orbits of other satellites [MG05], because it is not possible to measure the Earth's density directly. Other methods for obtaining such data are through surface gravimetry, which provides precise local and regional information about the gravity field, or through altimeter data, which can be used to provide a model for the geoid of the Earth (which is the shape that the ocean surface takes under the influence of the gravity of Earth) which in turn can be used to obtain the geopotential coefficients.

1.3 Numerical computation

1.3.1 Gravitational acceleration

Up to this point, we have only studied the gravitational potential exerted by the non-homogeneous Earth on a satellite. But, in order to integrate the equations of motion of the satellite, we need to compute

the gravitational acceleration $\mathbf{g} = -\nabla V$ instead. In order to do this efficiently, we will make use of the following formulas given in [MG05; Cun70]. First, let

$$V_{n,m}(\theta, \phi) = \left(\frac{R_{\oplus}}{r}\right)^{n+1} P_{n,m}(\cos \phi) \cos(m\theta) \quad W_{n,m}(\theta, \phi) = \left(\frac{R_{\oplus}}{r}\right)^{n+1} P_{n,m}(\cos \phi) \sin(m\theta) \quad (49)$$

Thus, we can write:

$$V = \frac{GM_{\oplus}}{R_{\oplus}} \sum_{n=0}^{\infty} \sum_{m=0}^n (\bar{C}_{n,m} N_{n,m} V_{n,m} + \bar{S}_{n,m} N_{n,m} W_{n,m}) \quad (50)$$

Let $C_{n,m} := \bar{C}_{n,m} N_{n,m}$ and $S_{n,m} := \bar{S}_{n,m} N_{n,m}$. If $\mathbf{g} = (\ddot{x}, \ddot{y}, \ddot{z})$, then:

$$\ddot{x} = \sum_{n=0}^{\infty} \sum_{m=0}^n \ddot{x}_{n,m} \quad \ddot{y} = \sum_{n=0}^{\infty} \sum_{m=0}^n \ddot{y}_{n,m} \quad \ddot{z} = \sum_{n=0}^{\infty} \sum_{m=0}^n \ddot{z}_{n,m} \quad (51)$$

where the *partial* accelerations $\ddot{x}_{n,m}$, $\ddot{y}_{n,m}$, $\ddot{z}_{n,m}$ are given by:

$$\ddot{x}_{n,m} = \begin{cases} -\frac{GM_{\oplus}}{R_{\oplus}^2} C_{n,0} V_{n+1,1} & \text{if } m = 0 \\ -\frac{GM_{\oplus}}{R_{\oplus}^2} \cdot \frac{1}{2} \left[C_{n,m} V_{n+1,m+1} + S_{n,m} W_{n+1,m+1} - \frac{(n-m+2)!}{(n-m)!} (C_{n,m} V_{n+1,m-1} + S_{n,m} W_{n+1,m-1}) \right] & \text{if } m > 0 \end{cases} \quad (52)$$

$$\ddot{y}_{n,m} = \begin{cases} -\frac{GM_{\oplus}}{R_{\oplus}^2} C_{n,0} W_{n+1,1} & \text{if } m = 0 \\ -\frac{GM_{\oplus}}{R_{\oplus}^2} \cdot \frac{1}{2} \left[C_{n,m} W_{n+1,m+1} - S_{n,m} V_{n+1,m+1} - \frac{(n-m+2)!}{(n-m)!} (C_{n,m} W_{n+1,m-1} - S_{n,m} V_{n+1,m-1}) \right] & \text{if } m > 0 \end{cases} \quad (53)$$

$$\ddot{z}_{n,m} = -\frac{GM_{\oplus}}{R_{\oplus}^2} (n-m+1) (C_{n,m} V_{n+1,m} + S_{n,m} W_{n+1,m}) \quad (54)$$

and the functions $V_{n,m}$, $W_{n,m}$ are calculated using the following recurrence relations:

$$\left\{ \begin{array}{l} V_{n,m} = \frac{2n-1}{n-m} \frac{R_{\oplus}}{r} \cos \phi V_{n-1,m} - \frac{n+m-1}{n-m} \frac{R_{\oplus}^2}{r^2} V_{n-2,m} \\ W_{n,m} = \frac{2n-1}{n-m} \frac{R_{\oplus}}{r} \cos \phi W_{n-1,m} - \frac{n+m-1}{n-m} \frac{R_{\oplus}^2}{r^2} W_{n-2,m} \end{array} \right. \quad \text{if } 0 \leq m \leq n-2$$

$$\left\{ \begin{array}{l} V_{n,n-1} = (2n-1) \frac{R_{\oplus}}{r} \cos \phi V_{n-1,n-1} \\ W_{n,n-1} = (2n-1) \frac{R_{\oplus}}{r} \cos \phi W_{n-1,n-1} \end{array} \right. \quad \text{if } m = n-1 \quad (55)$$

$$\left\{ \begin{array}{l} V_{n,n} = (2m-1) \frac{R_{\oplus}}{r} \sin \phi [\cos \theta V_{n-1,n-1} - \sin \theta W_{n-1,n-1}] \\ W_{n,n} = (2m-1) \frac{R_{\oplus}}{r} \sin \phi [\cos \theta W_{n-1,n-1} + \sin \theta V_{n-1,n-1}] \end{array} \right. \quad \text{if } m = n$$

starting from the initial quantities $V_{00} = \frac{R_{\oplus}}{r}$ and $W_{00} = 0$ and using the scheme of Fig. 4 [MG05].

1.4 Other perturbations

So far we have only considered one force acting on the satellite: the gravitational force exerted by the Earth. However, there are other important forces worth-considering, and they relative importance varies depending on the satellite's altitude.

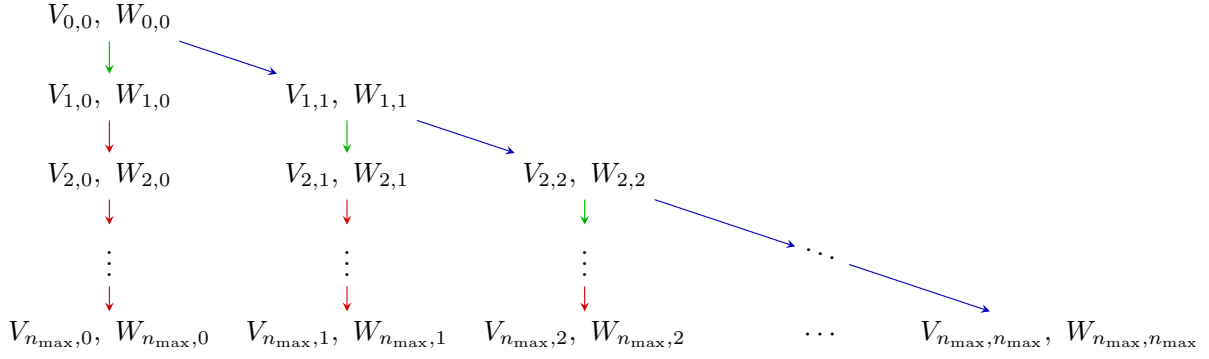


Figure 4: Scheme for the calculation of the coefficients $V_{n,m}$ and $W_{n,m}$ for $0 \leq m \leq n \leq n_{\max}$. The red arrows indicate that the coefficients are calculated using the first of the above recursions; the green arrows indicate that they are calculated using the second recursion; and the blue arrows indicate that they are calculated using the third recursion.

1.4.1 Atmospheric drag

For *LEO* (*Low Earth Orbit*) satellites, that is, satellites with an altitude of less than 1 000 km, the most important perturbation is the atmospheric drag. Indeed, as these satellites travel in the upper layers of the atmosphere, they are subject to a drag force caused by the interaction with air particles. The acceleration due to drag can be written as:

$$\mathbf{a}_D = -\frac{1}{2}C_D \frac{A}{m} \rho v_{\text{rel}} \mathbf{v}_{\text{rel}} \quad (56)$$

where A is the cross-sectional area of the satellite, m is its mass, ρ is the atmospheric density, \mathbf{v}_{rel} is the relative velocity between the satellite and the air particles, $v_{\text{rel}} = \|\mathbf{v}_{\text{rel}}\|$, and C_D is the drag coefficient, which depends on the shape of the satellite and the properties of the air particles. In order to determine \mathbf{v}_{rel} , we will suppose that the atmospheric particles rotate with the Earth, and thus the relative velocity is given by:

$$\mathbf{v}_{\text{rel}} = \dot{\mathbf{r}} - \boldsymbol{\omega}_{\oplus} \times \mathbf{r} \quad (57)$$

where $\boldsymbol{\omega}_{\oplus}$ is the angular velocity of the Earth.

The complexity of modeling this force arises from the challenge of representing accurately the atmospheric density as a function of the altitude. We will not delve into this topic, but in the simulation we will use the density model of Harris-Priester, which is valid for altitudes between 100 and 1 000 km (see [MG05] for more details).

1.4.2 Sun and Moon gravitational perturbations

For *MEO* (*Medium Earth Orbit*) satellites, within an altitude between 1 000 and 35 780 km [Val13]; *GEO* (*Geostationary Earth Orbit* or *Geosynchronous Earth Orbit*) satellites, with an altitude of around 35 780 km; and *HEO* (*High Earth Orbit*) satellites, with an altitude of more than 35 780 km, the most important perturbations are the gravitational perturbations caused by the Moon and the Sun, or even other celestial bodies.

Adding a third body into the equations requires some carefulness. We will do the general deduction to make it clear for the reader. Suppose we have $N + 2$ bodies of masses M, M_1, \dots, M_N and $m \ll M, M_1, \dots, M_N$ at positions $\mathbf{s}, \mathbf{s}_1, \dots, \mathbf{s}_N$ and \mathbf{r} , respectively, in an inertial reference frame (we omit the dependence in time). The dynamics of this system are governed by the following system of differential equations:

$$\begin{cases} \ddot{\mathbf{s}} = \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{s}\|^3} (\mathbf{s}_i - \mathbf{s}) + \frac{Gm}{\|\mathbf{r} - \mathbf{s}\|^3} (\mathbf{r} - \mathbf{s}) \\ \ddot{\mathbf{s}}_i = \frac{GM}{\|\mathbf{s} - \mathbf{s}_i\|^3} (\mathbf{s} - \mathbf{s}_i) + \sum_{\substack{j=1 \\ j \neq i}}^N \frac{GM_j}{\|\mathbf{s}_j - \mathbf{s}_i\|^3} (\mathbf{s}_j - \mathbf{s}_i) + \frac{Gm}{\|\mathbf{r} - \mathbf{s}_i\|^3} (\mathbf{r} - \mathbf{s}_i) \quad \text{for } i = 1, \dots, N \\ \ddot{\mathbf{r}} = \frac{GM}{\|\mathbf{s} - \mathbf{r}\|^3} (\mathbf{s} - \mathbf{r}) + \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{r}\|^3} (\mathbf{s}_i - \mathbf{r}) \end{cases} \quad (58)$$

Here G is the gravitational constant. Since $m \ll M, M_1, \dots, M_N$, we can assume that the motion of the large bodies is not affected by the small body, i.e. we can make m tend to zero and assume:

$$\begin{cases} \ddot{\mathbf{s}} = \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{s}\|^3} (\mathbf{s}_i - \mathbf{s}) \\ \ddot{\mathbf{s}}_i = \frac{GM}{\|\mathbf{s} - \mathbf{s}_i\|^3} (\mathbf{s} - \mathbf{s}_i) + \sum_{\substack{j=1 \\ j \neq i}}^N \frac{GM_j}{\|\mathbf{s}_j - \mathbf{s}_i\|^3} (\mathbf{s}_j - \mathbf{s}_i) & \text{for } i = 1, \dots, N \\ \ddot{\mathbf{r}} = \frac{GM}{\|\mathbf{s} - \mathbf{r}\|^3} (\mathbf{s} - \mathbf{r}) + \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{r}\|^3} (\mathbf{s}_i - \mathbf{r}) \end{cases} \quad (59)$$

For our purposes, if the body of mass M is the Earth and the body of mass m is the satellite, we are interested in the position of the satellite relative to the Earth, that is $\mathbf{R} = \mathbf{r} - \mathbf{s}$. The dynamics of \mathbf{R} is governed by the following equation:

$$\ddot{\mathbf{R}} = -\frac{GM}{\|\mathbf{R}\|^3} \mathbf{R} + \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{s} - \mathbf{R}\|^3} (\mathbf{s}_i - \mathbf{s} - \mathbf{R}) - \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{s}\|^3} (\mathbf{s}_i - \mathbf{s}) \quad (60)$$

Taking the origin of the inertial reference frame at the center of mass of the Earth, i.e. $\mathbf{s} = 0$, the latter equation finally becomes:

$$\ddot{\mathbf{R}} = -\frac{GM}{\|\mathbf{R}\|^3} \mathbf{R} + \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i - \mathbf{R}\|^3} (\mathbf{s}_i - \mathbf{R}) - \sum_{i=1}^N \frac{GM_i}{\|\mathbf{s}_i\|^3} \mathbf{s}_i = -\frac{GM}{\|\mathbf{R}\|^3} \mathbf{R} + \sum_{i=1}^N \mathbf{a}_i \quad (61)$$

where $\mathbf{a}_i = \frac{GM_i}{\|\mathbf{s}_i - \mathbf{R}\|^3} (\mathbf{s}_i - \mathbf{R}) - \frac{GM_i}{\|\mathbf{s}_i\|^3} \mathbf{s}_i$ is the contribution of the i -th body to the acceleration of the satellite. In the simulation, we will replace the first term of the right-hand side of the equation by the expansion of the acceleration in spherical harmonics.

1.4.3 Solar radiation pressure

Satellites in medium and high-altitude orbits may also experience a force that arises from the absorption or reflection of photons emitted by the Sun [MG05]. We will expose a brief overview on how this solar radiation pressure is modeled and how it is included in the equations of motion. The acceleration due to solar radiation pressure is given by:

$$\mathbf{a}_R = -P_\odot C_R \frac{A_\odot}{m} \frac{\mathbf{s}_\odot - \mathbf{r}}{\|\mathbf{s}_\odot - \mathbf{r}\|} \quad (62)$$

Here $P_\odot = 4.57 \times 10^{-6} \text{ N/m}^2$ is the solar radiation pressure at around the distance of the Earth, $C_R = 1 + \varepsilon \in [0, 2]$ is the radiation pressure coefficient and ε is the reflectivity coefficient, being $C_R = 0$ for a perfectly translucent body, $C_R = 1$ for a perfectly absorbing body, and $C_R = 2$ for a perfectly reflecting body. \mathbf{r} denotes the position of the satellite relative to the Earth, \mathbf{s}_\odot is the position of the Sun relative to the Earth, m the mass of the satellite and A_\odot its fraction of area exposed to the Sun.

One should note that A_\odot and A (the cross-sectional area of the satellite used to compute the drag force) are in general not the same. We will see in the next section, though, that due to the difficulties on determining them, we will assume that they are equal and constant throughout the integration process.

In reality, Eq. (62) does not describe entirely well the perturbation due to the solar radiation. Indeed, an illumination factor $\nu \in [0, 1]$ should be added into the equation to account for the fraction of solar rays that actually reach the satellite. This cause the distinction of three regions: the illumination zone ($\nu = 1$), where the satellite is fully exposed to the Sun; the penumbral zone ($0 < \nu < 1$), where the satellite is partially exposed to the Sun, and the umbral zone ($\nu = 0$), where the satellite receives no solar radiation due to the Earth obstructing the light. Thus, Eq. (62) should be corrected to:

$$\mathbf{a}_R = -\nu P_\odot C_R \frac{A_\odot}{m} \frac{\mathbf{s}_\odot - \mathbf{r}}{\|\mathbf{s}_\odot - \mathbf{r}\|} \quad (63)$$

Fig. 5 shows a 2D schematic representation of the penumbra and umbral zones of the Earth. To determine whether the satellite is or not inside on of these zones, we must compute the dot product between \mathbf{r} and \mathbf{s}_\odot and if it is negative, then compare the relative “altitude” of the satellite with respect to the “altitudes” of the penumbral and umbral zones (y_{sat} , y_{pen} and y_{umb} respectively in the figure).

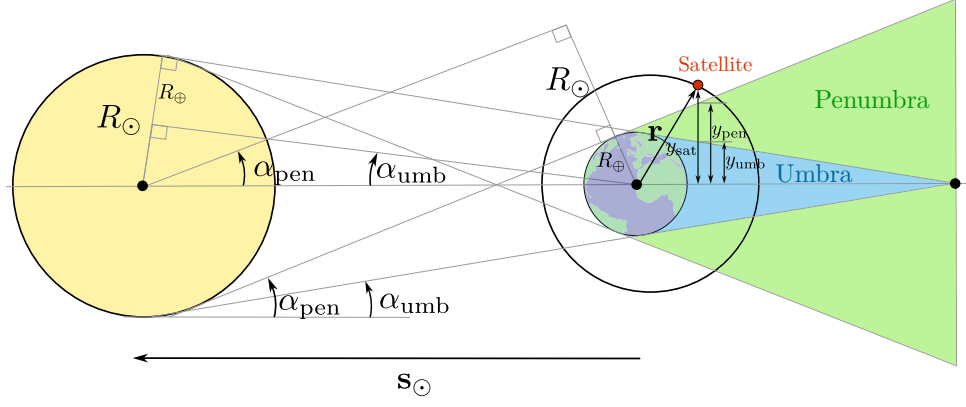


Figure 5: Schematic representation of the penumbra and umbral zones. Based on [Val13].

1.4.4 Minor perturbations

In addition to the major perturbations described above, there are other minor perturbations that we will not consider in this work, but they are worth-considering if accurate results are required. Some of these perturbations include the perturbations caused by other planets, relativistic effects, the Earth tide effects, or radiation pressure caused by the sunlight reflected by the Earth (albedo) [MG05].