

CFD

**March 21: last course
Homework for Math students**

April 4: projects defense

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8. Boundary conditions

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Heat Equation in 1D

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

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Dirichlet:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$u = g$$

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Dirichlet:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$u = g$$

$$\frac{u_1^{n+1} - u_1^n}{\Delta t} = \kappa \frac{u_2^n - 2u_1^n + g}{\Delta x^2}$$

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Neumann:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = 0$$

$$\frac{u_1 - u_0}{\Delta x} = 0$$

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Neumann:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = 0$$

$$\frac{-3u_0 + 4u_1 - u_2}{2\Delta x} = 0$$

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Neumann:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = 0$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = \frac{u_1 - u_{-1}}{2\Delta x} + \mathcal{O}(\Delta x^2) = 0$$

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Neumann:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = 0$$

$$\begin{aligned} u_0^{n+1} &= u_0^n + \frac{\kappa \Delta t}{\Delta x^2} [u_{-1}^n - 2u_0^n + u_1^n] \\ &= u_0^n + \frac{\kappa \Delta t}{\Delta x^2} [-2u_0^n + 2u_1^n]. \end{aligned}$$

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Neumann:

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = 0$$

$$\left. \frac{\partial u}{\partial x} \right|_{1/2} = \frac{u_1 - u_0}{\Delta x} + \mathcal{O}(\Delta x^2) = 0$$

$$u_0 = u_1$$

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Neumann (non-homogeneous):

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2} \quad \text{with } x \in [0, L]$$

$$\left. \frac{\partial u}{\partial x} \right|_0 = \gamma \quad u_0 = u_1 - \gamma \Delta x$$

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Robin / Fourier / Mixed :

$$\frac{\partial u}{\partial x} = \alpha u$$

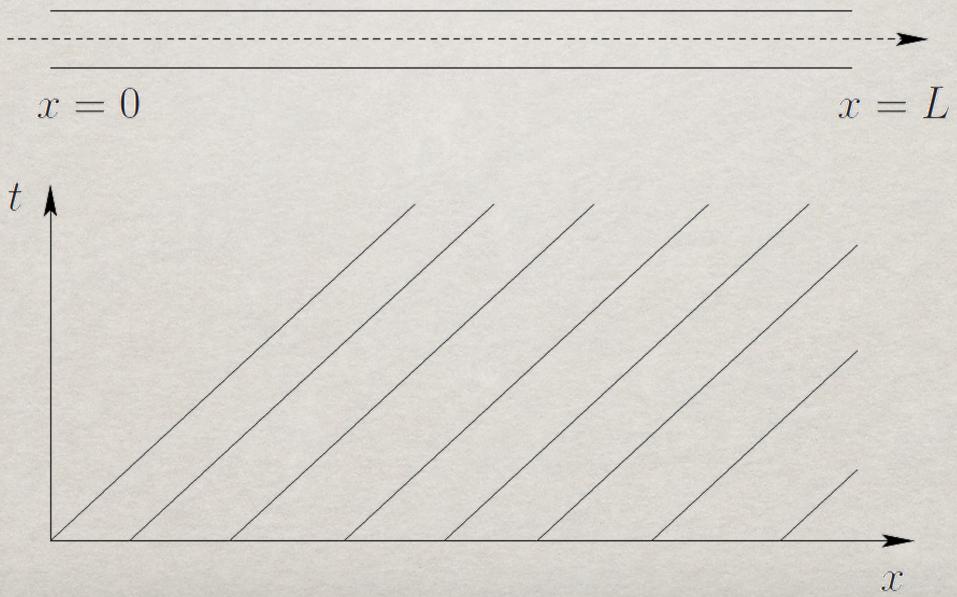
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Advection:

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0$$

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Advection:



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Sound waves

$$\frac{1}{\rho_0} \frac{\partial \rho'}{\partial t} + \nabla \cdot \mathbf{u} = 0$$

$$\rho_0 \frac{\partial \mathbf{u}}{\partial t} + c_0^2 \nabla \rho' = 0$$

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Sound waves

$$\frac{1}{\rho_0} \frac{\partial \rho'}{\partial t} + \nabla \cdot \mathbf{u} = 0$$

$$\rho_0 \frac{\partial \mathbf{u}}{\partial t} + c_0^2 \nabla \rho' = 0$$

$$\frac{\partial \tilde{\rho}}{\partial t} + \frac{\partial u}{\partial x} = 0$$

$$\frac{\partial \mathbf{u}}{\partial t} + c_0^2 \frac{\partial \tilde{\rho}}{\partial x} = 0.$$

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$$\begin{aligned}\frac{\partial \tilde{\rho}}{\partial t} + \frac{\partial u}{\partial x} &= 0 \\ \frac{\partial \mathbf{u}}{\partial t} + c_0^2 \frac{\partial \tilde{\rho}}{\partial x} &= 0\end{aligned}$$

$$\begin{aligned}\frac{\partial}{\partial t}(c_0 \tilde{\rho} + u) + c_0 \frac{\partial}{\partial x}(c_0 \tilde{\rho} + u) &= 0 \\ \frac{\partial}{\partial t}(c_0 \tilde{\rho} - u) - c_0 \frac{\partial}{\partial x}(c_0 \tilde{\rho} - u) &= 0\end{aligned}$$

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$$\begin{aligned}\frac{\partial}{\partial t}(c_0 \tilde{\rho} + u) + c_0 \frac{\partial}{\partial x}(c_0 \tilde{\rho} + u) &= 0 \\ \frac{\partial}{\partial t}(c_0 \tilde{\rho} - u) - c_0 \frac{\partial}{\partial x}(c_0 \tilde{\rho} - u) &= 0\end{aligned}$$

$$X^\pm = c_0 \tilde{\rho} \pm u$$

$$\frac{\partial}{\partial t} X^\pm \pm c_0 \frac{\partial}{\partial x} X^\pm = 0$$

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Shallow water:

$$\frac{\partial}{\partial t} \begin{pmatrix} u \\ \eta \end{pmatrix} + \begin{pmatrix} U & 1 \\ c^2 & U \end{pmatrix} \frac{\partial}{\partial x} \begin{pmatrix} u \\ \eta \end{pmatrix} = 0$$

$$c = \sqrt{gH}$$

$$V_{\pm} = u \pm \eta/c$$

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Shallow water:

$$\frac{\partial}{\partial t} \begin{pmatrix} u \\ \eta \end{pmatrix} + \begin{pmatrix} U & 1 \\ c^2 & U \end{pmatrix} \frac{\partial}{\partial x} \begin{pmatrix} u \\ \eta \end{pmatrix} = 0$$

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Shallow water:

$$c = \sqrt{gH} \quad V_{\pm} = u \pm \eta/c$$

$$\frac{\partial}{\partial t} \begin{pmatrix} V_+ \\ V_- \end{pmatrix} + \begin{pmatrix} U + c & 0 \\ 0 & U - c \end{pmatrix} \frac{\partial}{\partial x} \begin{pmatrix} V_+ \\ V_- \end{pmatrix} = 0$$

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Boundary layers

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Fridtjof Nansen (1861-1930)



Fridtjof Nansen

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Vagn Ekman (1874-1954)

depth to $e^{-2\pi} = 1/535$ th part for each time its direction rotates four right angles. The direction and velocity of the

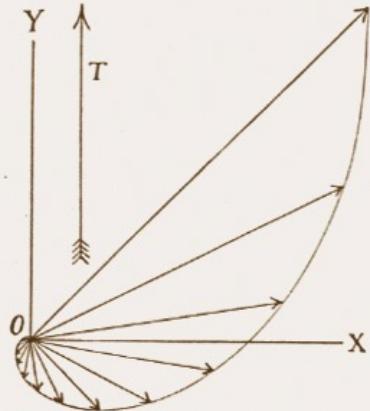


Fig. 1.

current at different depths are represented by the arrows in Fig. 1 above; the longest arrow refers to the surface, the (the water) does not however vary appreciably with the height, within



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Ekman spiral

$$\begin{aligned} -2\Omega u_y &= -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \frac{\partial^2 u_x}{\partial z^2}, \\ 2\Omega u_x &= -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \frac{\partial^2 u_y}{\partial z^2}. \end{aligned}$$

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Ekman spiral

$$2\Omega u_x^o = -\frac{1}{\rho} \frac{\partial p^0}{\partial y},$$
$$0 = -\frac{1}{\rho} \frac{\partial p^0}{\partial X},$$

$$-2\Omega u_y = \nu \frac{\partial^2 u_x}{\partial z^2},$$
$$2\Omega (u_x - u_x^0) = \nu \frac{\partial^2 u_y}{\partial z^2}.$$

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Ekman spiral

$$-2\Omega u_y = \nu \frac{\partial^2 u_x}{\partial z^2},$$
$$2\Omega (u_x - u_x^0) = \nu \frac{\partial^2 u_y}{\partial z^2}.$$

$$u_x = u_x^0 \left[1 - e^{-z/\delta} \cos(z/\delta) \right],$$
$$u_y = u_x^0 e^{-z/\delta} \sin(z/\delta),$$

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Ekman spiral

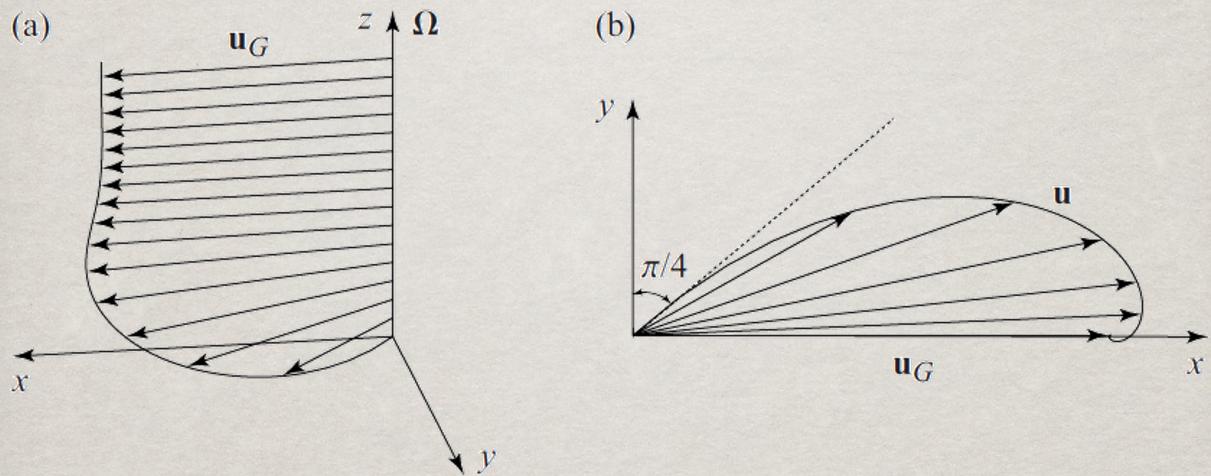
$$E \partial_{zz} \mathcal{V} = i \mathcal{V}$$

$$\zeta = \sqrt{2E} z$$

$$\frac{\partial^2 \mathcal{V}}{\partial \zeta^2} = 2i \mathcal{V}$$

$$\mathcal{V} \propto e^{-\zeta(1+i)}$$

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Vagn Ekman (1874-1954)

depth to $e^{-2\pi} = 1/535$ th part for each time its direction rotates four right angles. The direction and velocity of the

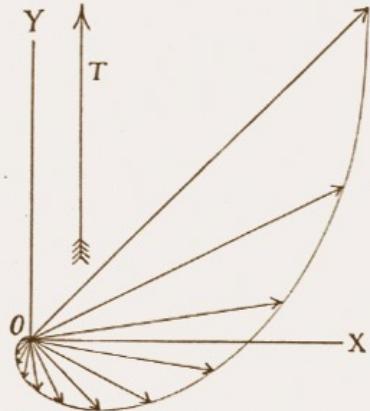
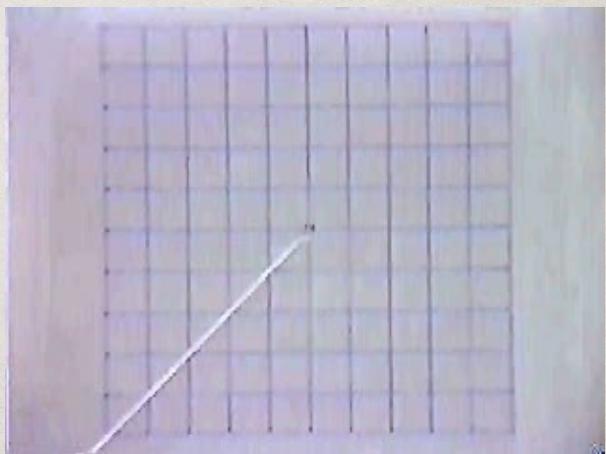
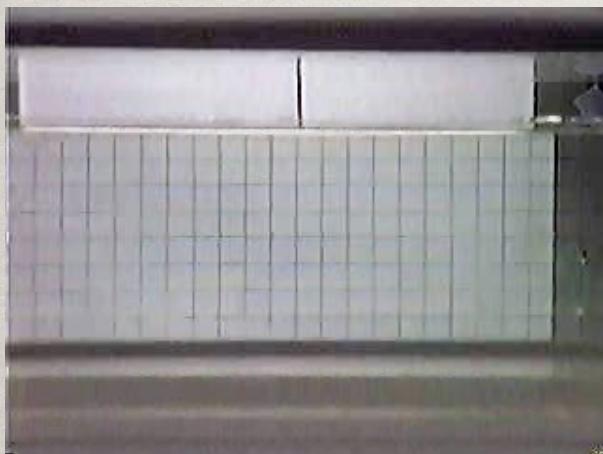


Fig. 1.

current at different depths are represented by the arrows in Fig. 1 above; the longest arrow refers to the surface, the (the water) does not however vary appreciably with the height, within

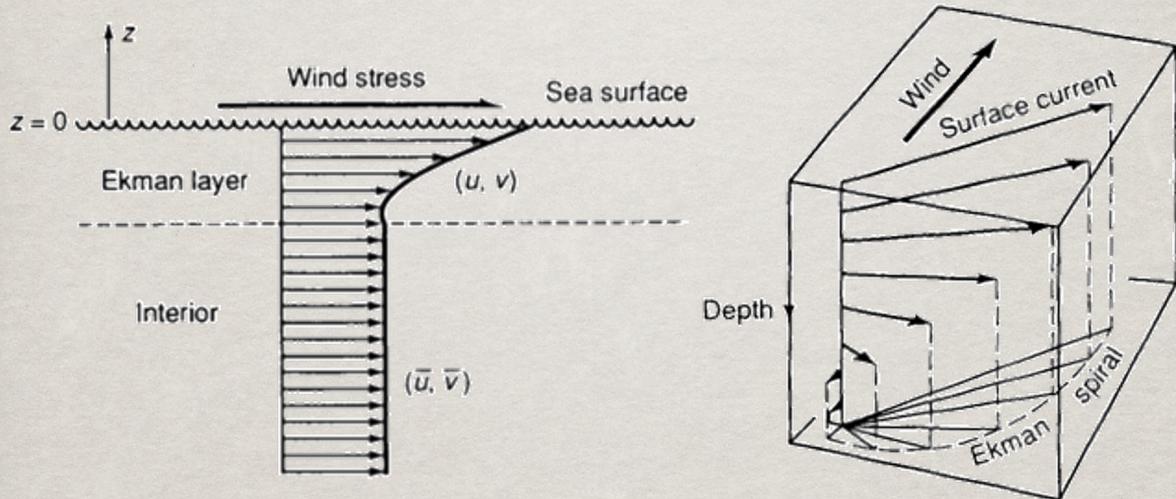


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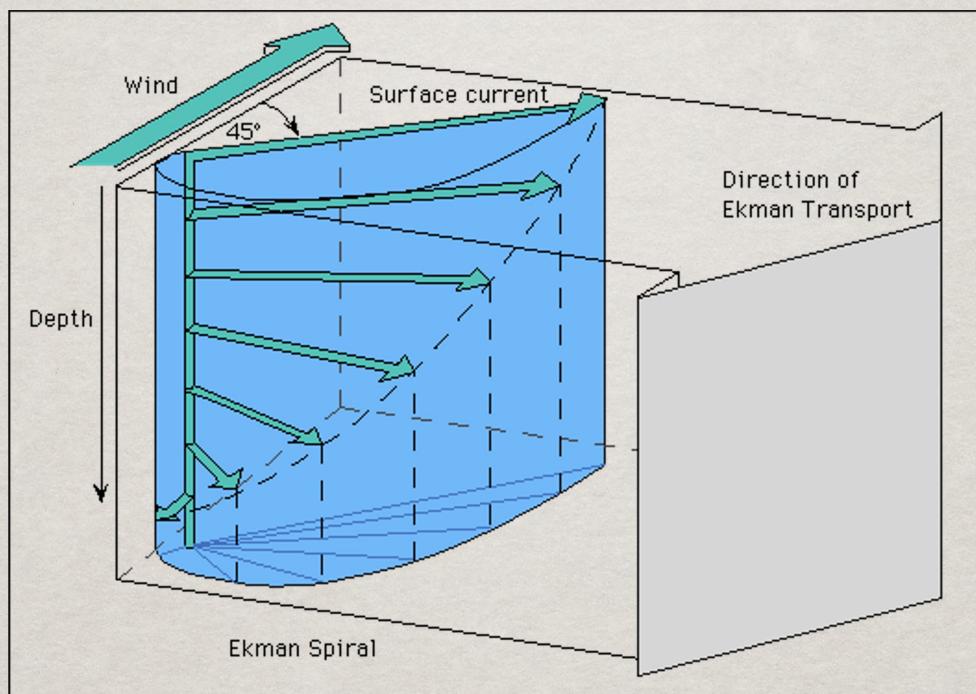
(GFD Dennou Japon)

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$$\mathcal{V} \propto e^{-\zeta(1+i)}$$

Cushman-Roisin

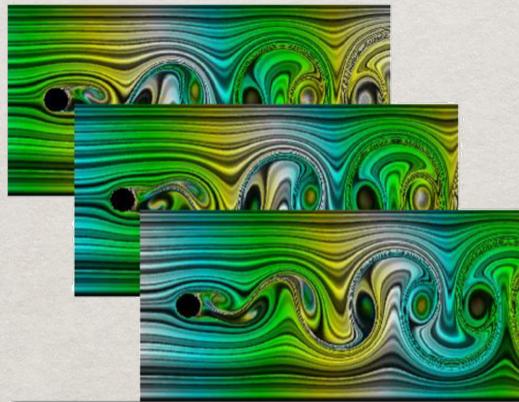
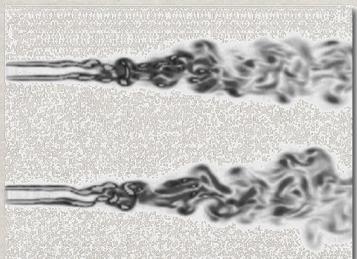


Adapted from Thurman, Harold V. **Essentials of Oceanography**, 5th ed.
Prentice-Hall, Inc., 1996.

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Opened domains:

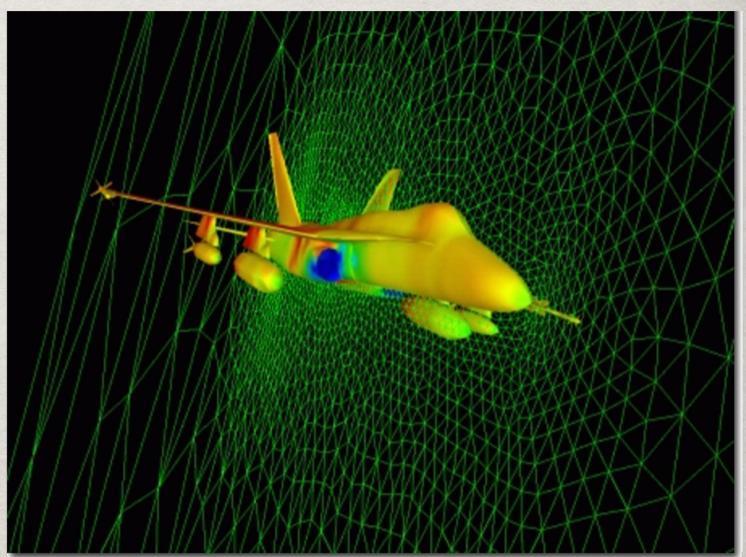
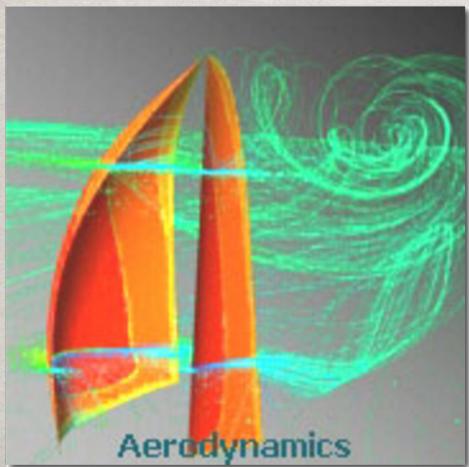
- **wakes, jets and plumes**
- aerodynamics
- acoustics



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Opened domains:

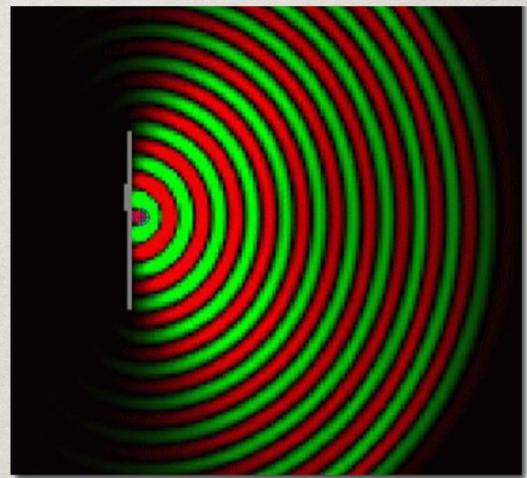
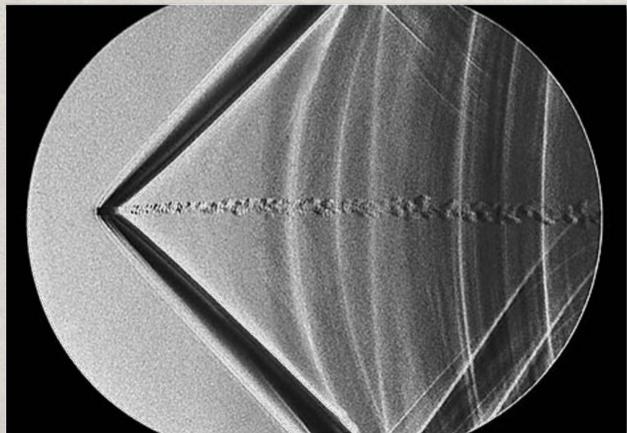
- wakes, jets and plumes
- **aerodynamics**
- acoustics



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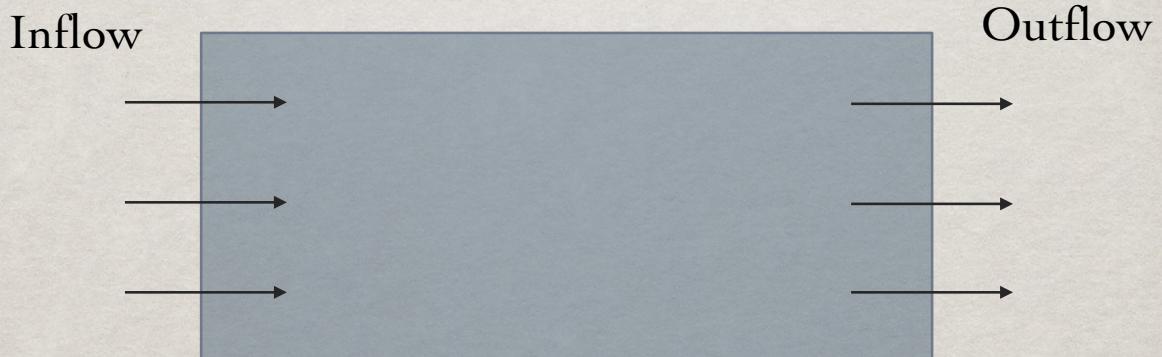
Opened domains:

- wakes, jets and plumes
- aerodynamics
- **acoustics**



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Opened domains:



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Opened domains for waves:

General idea:

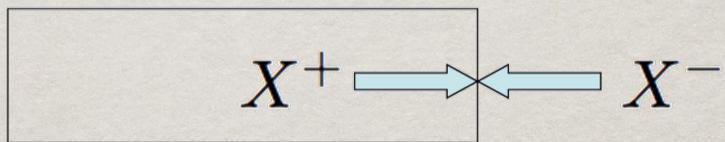
If the perturbation can be decomposed into wave components, one can manipulate the dispersion relation at the boundary to suppress the incoming waves.

Engquist-Majda (1977)

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Characteristics:

In one dimension, we want to construct an outflow boundary condition at the right boundary



outflow boundary condition $X^- = 0$

Engquist-Majda (1977)

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Final Projects

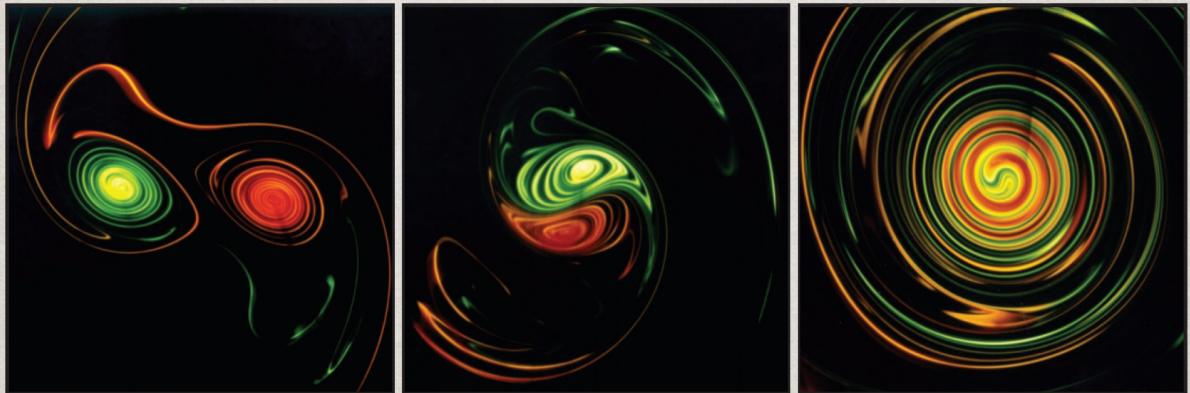
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Kelvin wake

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Vortices interactions

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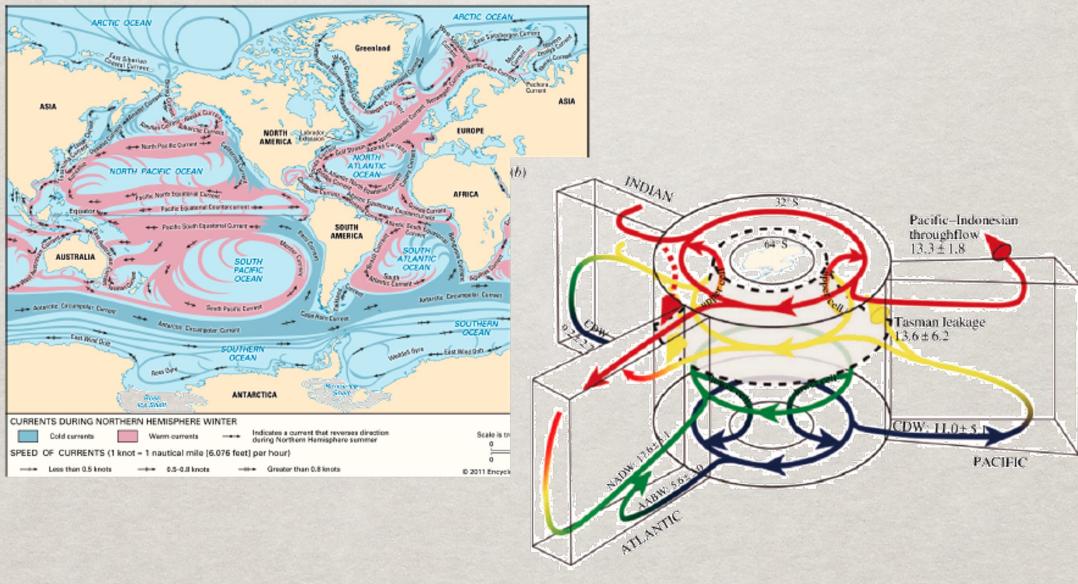
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Thermal convection

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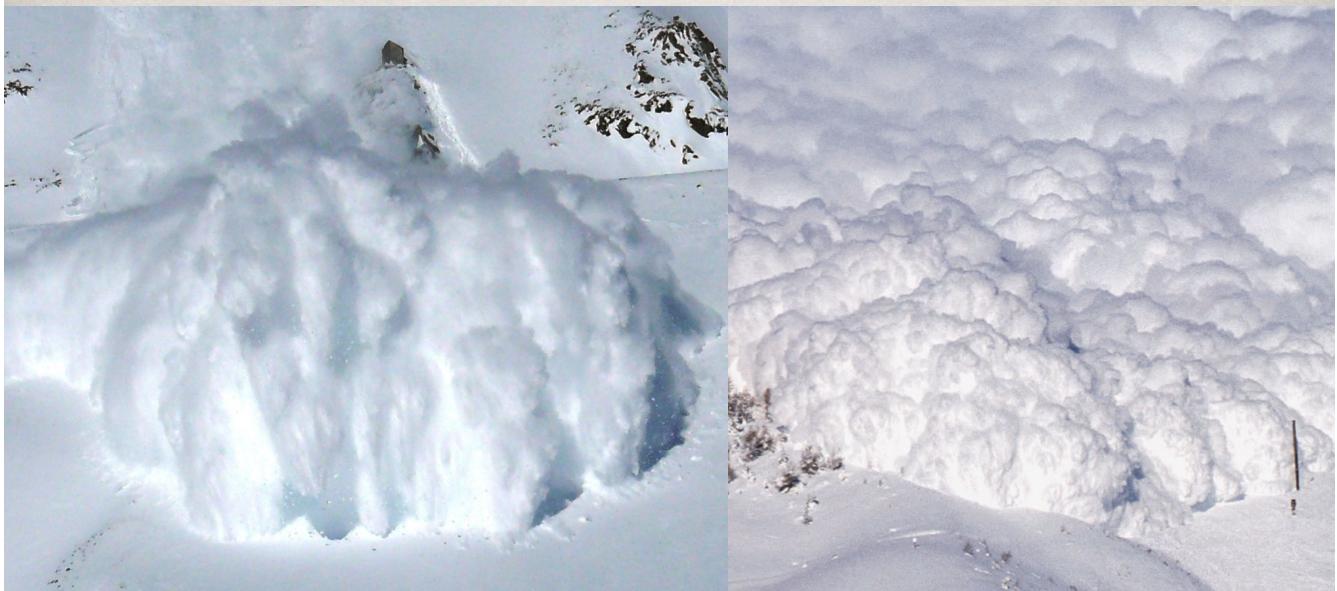
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Thermo-haline convection

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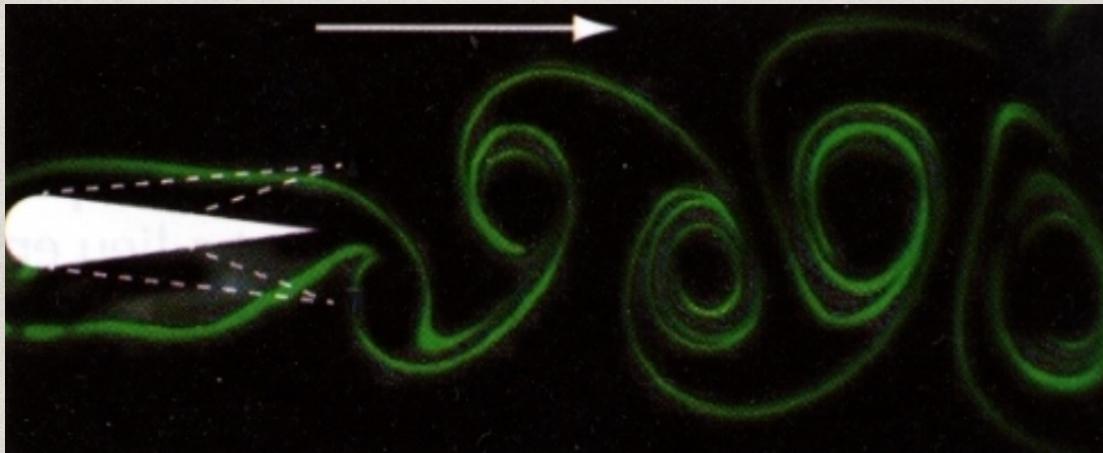


Avalanche model



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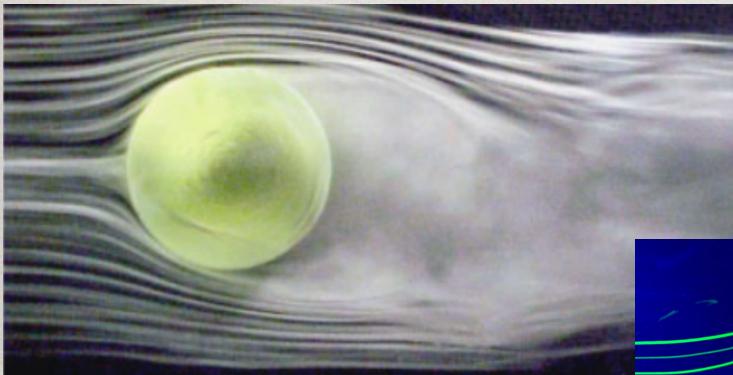


Propulsion



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Tennis/Football lift, Magnus effect



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Final Projects

Your own idea...