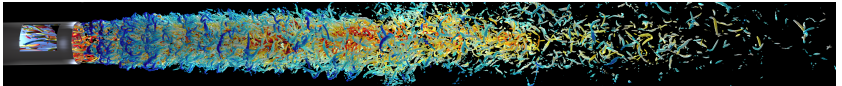


Turbulence

Navier Stokes and Symmetries



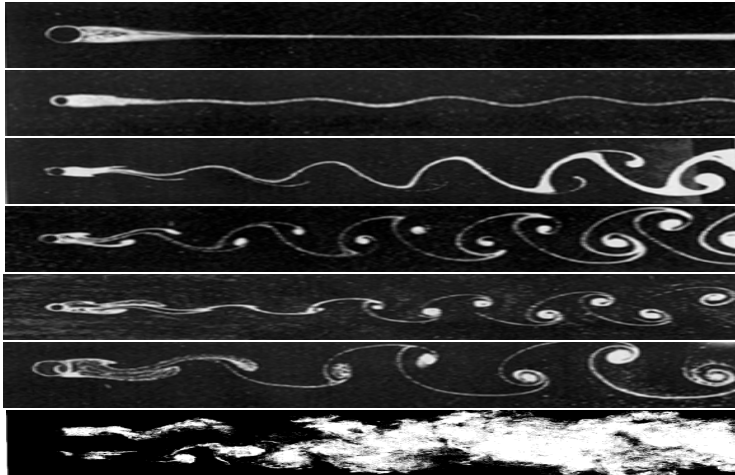
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Dep. Physique ENS Ulm

- Symmetry in fields and in equations
- Symmetries of the Navier Stokes
- Breaking of symmetries

Flow behind a cylinder



Re= 32

Re= 55

Re= 60

Re= 73

Re=102

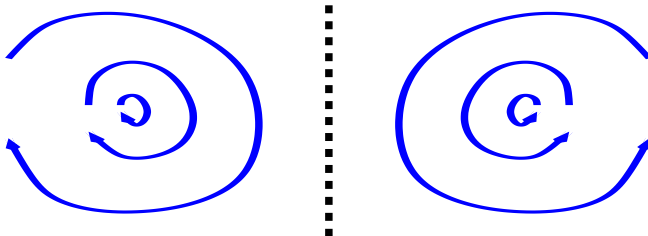
Re=161

Re=1800

Symmetries of Fields

We will say that a field $\mathbf{u}(\mathbf{x}, t)$ is invariant under a transformation \mathcal{T} or that it has a \mathcal{T} -symmetry if under the act of the transformation it remains the same:

$$\mathcal{T}[\mathbf{u}] = \mathbf{u}$$

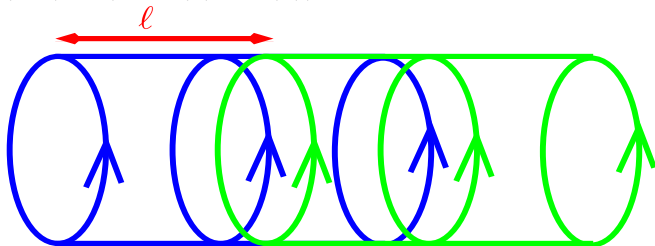


Space translations

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}(\mathbf{x} + \hat{\mathbf{e}}_i \ell, t)$$

Example:

- $\mathbf{u}(\mathbf{x}, t) = (0, \sin(z), \cos(y))$



- continuous symmetry in x : $\mathbf{u}(x, y, z) = \mathbf{u}(x + \ell_x, y, z, t)$
- discrete symmetry in y, z : $\mathbf{u}(x, y, z) = \mathbf{u}(x, y + 2n_y\pi, z + 2n_z\pi, t)$

Time translations

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}(\mathbf{x}, t + T)$$

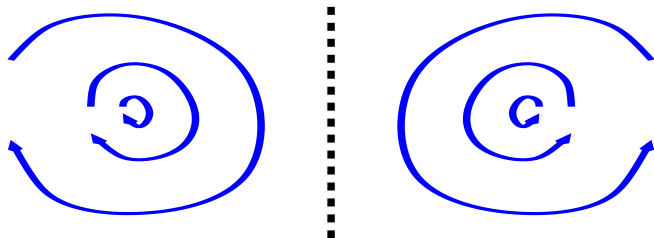
Example:

- Time periodic flows have discrete symmetries
- Constant in time flows have continuous symmetries

$$\mathcal{T}[(u_x(x, y, z, t), u_y(x, y, z, t), u_z(x, y, z, t))] = \\ (-u_x(-x, y, z, t), u_y(-x, y, z, t), u_z(-x, y, z, t))$$

Example:

- $\mathbf{u}(\mathbf{x}, t) = (\sin(x) \cos(y), -\cos(x) \sin(y), 0)$

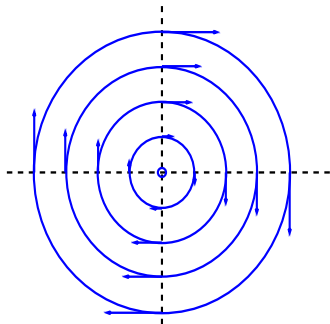


Rotations

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathcal{R}\mathbf{u}(\mathcal{R}^{-1}\mathbf{x}, t) + \mathbf{c}$$

Where \mathcal{R} is the rotation matrix **Example:**

- $\mathbf{u}(\mathbf{x}, t) = (y, -x, 0)$



Galilean Transformations

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}(\mathbf{x} + \mathbf{c}t, t) + \mathbf{c}$$

Example:

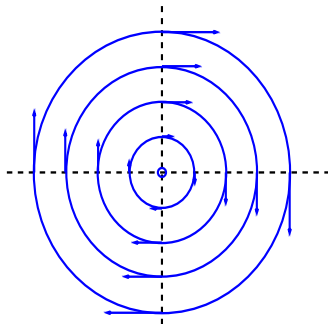
- $\mathbf{u}(\mathbf{x}, t) = -x/t$

Scaling Transformations

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \lambda^\alpha \mathbf{u}(\lambda \mathbf{x}, t)$$

Example:

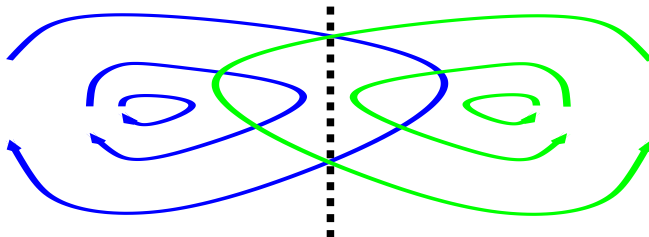
- $\mathbf{u}(\mathbf{x}, t) = (y, -x, 0)(x^2 + y^2)^{\alpha/2-1}$



Symmetries of equations

We will say that an equation (e.g. the Navier-Stokes) is invariant under a transformation \mathcal{T} or that it has a \mathcal{T} -symmetry if for any solution $\mathbf{u}(\mathbf{x}, t)$ of this equation $\mathcal{T}[\mathbf{u}(\mathbf{x}, t)]$ is also a solution.

Note that $\mathbf{u}(\mathbf{x}, t)$ does not have to be a symmetric field.



Symmetries of equations

Lemma: If the initial conditions $\mathbf{u}(\mathbf{x}, 0)$ has one of the spatial symmetries of the Navier-Stokes

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, 0)] = \mathbf{u}(\mathbf{x}, 0),$$

then if $\mathbf{u}(\mathbf{x}, t)$ remains smooth it retains this symmetry for all times.

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}(\mathbf{x}, t)$$

Proof: If $\mathbf{u}(\mathbf{x}, t)$ is a solution then (since \mathcal{T} is one of the symmetries of the Navier-Stokes) $\mathcal{T}[\mathbf{u}(\mathbf{x}, t)]$ is also a solution that has the same initial conditions as $\mathbf{u}(\mathbf{x}, t)$ ($\mathcal{T}[\mathbf{u}(\mathbf{x}, 0)] = \mathbf{u}(\mathbf{x}, 0)$). Thus either there is non-uniqueness of solutions or

$$\mathcal{T}[\mathbf{u}(\mathbf{x}, t)] = \mathbf{u}(\mathbf{x}, t).$$

(Non-uniqueness occurs only for non-smooth $\mathbf{u}(\mathbf{x}, t)$).

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $\mathbf{u}(\mathbf{x} + \ell, t)$ is also a solution.”
 - $\partial_x \mathbf{u}(\mathbf{x} + \ell, t) = \frac{\partial x'}{\partial x} \partial_{x'} \mathbf{u}(\mathbf{x}', t) = \partial_{x'} \mathbf{u}(\mathbf{x}', t)$

$$\partial_t \mathbf{u}(x + \ell, t) + \mathbf{u}(x + \ell, t) \cdot \nabla \mathbf{u}(x + \ell, t) = -\nabla P + \nu \nabla^2 \mathbf{u}(x + \ell, t) + \mathbf{f}(\mathbf{x}, t)$$

$$\partial_t \mathbf{u}(x', t) + \mathbf{u}(x', t) \cdot \nabla' \mathbf{u}(x', t) = -\nabla' P + \nu \nabla'^2 \mathbf{u}(x', t) + \mathbf{f}(\mathbf{x}' - \ell, t)$$

which is the original Navier Stokes (if $\mathbf{f}(\mathbf{x}, t) = \mathbf{f}(\mathbf{x} + \ell, t)$)

ie, $\mathbf{u}(x + \ell, t)$ satisfies the same equations as $\mathbf{u}(x, t)$

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
- Time translations: $t' = t + T$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $\mathbf{u}(\mathbf{x}, t + T)$ is also a solution.”
 - $\partial_t \mathbf{u}(\mathbf{x}, t + T) = \frac{\partial t'}{\partial t} \partial_{t'} \mathbf{u}(\mathbf{x}, t') = \partial_{t'} \mathbf{u}(\mathbf{x}, t')$

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
- Time translations: $t' = t + T$
- Galilean transformations: $\mathbf{x}' = \mathbf{x} - \mathbf{c}t$, $t' = t$, $\mathbf{u}' = \mathbf{u} + \mathbf{c}$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $\mathbf{u}(\mathbf{x} - \mathbf{c}t, t) + \mathbf{c}$ is also a solution.”
 - $\partial_t(\mathbf{u}(\mathbf{x} - \mathbf{c}t, t) + \mathbf{c}) = \left(\frac{\partial t'}{\partial t}\right) \partial_{t'} \mathbf{u}(\mathbf{x}', t') + \left(\frac{\partial \mathbf{x}'}{\partial t}\right) \nabla_{\mathbf{x}'} \mathbf{u}(\mathbf{x}', t')$
 - $\partial_t \mathbf{u}(\mathbf{x}', t') = \partial_{t'} \mathbf{u}(\mathbf{x}', t') - \mathbf{c} \nabla' \mathbf{u}(\mathbf{x}', t')$
 - $(\mathbf{u} + \mathbf{c}) \cdot \nabla(\mathbf{u} + \mathbf{c}) = \mathbf{u} \cdot \nabla \mathbf{u} + \mathbf{c} \cdot \nabla \mathbf{u}$

$$\partial_t \mathbf{u}' + \mathbf{u}' \cdot \nabla \mathbf{u}' = -\nabla P + \nu \nabla^2 \mathbf{u}' + \mathbf{f}$$

$$\cancel{\partial_t' \mathbf{u}' - \mathbf{c} \cdot \nabla' \mathbf{u}'} + \mathbf{u} \cdot \nabla' \mathbf{u} + \cancel{\mathbf{c} \cdot \nabla' \mathbf{u}'} = -\nabla' P + \nu \nabla'^2 \mathbf{u} + \mathbf{f}$$

$$\text{if } \mathbf{f}(\mathbf{x}, t) = \mathbf{f}(\mathbf{x} - \mathbf{c}t, t)$$

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
- Time translations: $t' = t + T$
- Galilean transformations: $\mathbf{x}' = \mathbf{x} - \mathbf{c}t$, $t' = t$, $\mathbf{u}' = \mathbf{u} + \mathbf{c}$
- Rotations $\mathbf{u}' = \mathcal{R}\mathbf{u}$, $\mathbf{x}' = \mathcal{R}^{-1}\mathbf{x}$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $\mathcal{R}\mathbf{u}(\mathcal{R}^{-1}\mathbf{x}, t)$ is also a solution.”

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
- Time translations: $t' = t + T$
- Galilean transformations: $\mathbf{x}' = \mathbf{x} - \mathbf{c}t$, $t' = t$, $\mathbf{u}' = \mathbf{u} + \mathbf{c}$
- Rotations $\mathbf{u}' = \mathcal{R}\mathbf{u}$, $\mathbf{x}' = \mathcal{R}^{-1}\mathbf{x}$
- Parity (reflections): $\mathbf{u}' = -\mathbf{u}$, $\mathbf{x}' = -\mathbf{x}$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $-\mathbf{u}(-\mathbf{x}, t)$ is also a solution.”

Symmetries of the Navier-Stokes

- Space translations: $x' = x + \ell$
- Time translations: $t' = t + T$
- Galilean transformations: $\mathbf{x}' = \mathbf{x} - \mathbf{c}t$, $t' = t$, $\mathbf{u}' = \mathbf{u} + \mathbf{c}$
- Rotations $\mathbf{u}' = \mathcal{R}\mathbf{u}$, $\mathbf{x}' = \mathcal{R}^{-1}\mathbf{x}$
- Parity (reflections): $\mathbf{u}' = -\mathbf{u}$, $\mathbf{x}' = -\mathbf{x}$
- Scaling $\mathbf{x}' = \mathbf{x}/\lambda$, $t' = t/\lambda^\alpha$, $\mathbf{u}' = \lambda^\beta \mathbf{u}$
 - “if $\mathbf{u}(\mathbf{x}, t)$ is a solution $\lambda^\beta \mathbf{u}(\mathbf{x}/\lambda, t/\lambda^\alpha)$ is also a solution.”

$$\partial_t \mathbf{u}' + \mathbf{u}' \cdot \nabla \mathbf{u}' = -\nabla P + \nu \nabla^2 \mathbf{u}' + \mathbf{f}$$

$$\lambda^{\beta-\alpha} \partial_t \mathbf{u} + \lambda^{2\beta-1} \mathbf{u} \cdot \nabla' \mathbf{u} = -\lambda^{2\beta-1} \nabla P + \nu \lambda^{\beta-2} \nabla^2 \mathbf{u} + \mathbf{f}$$

Is a solution if $\beta - \alpha = 2\beta - 1$ and $2\beta - 1 = \beta - 2$

$\beta = -1$ and $\alpha = 2$

In the absence of viscosity there is a scaling symmetry for any β
and $\alpha = 1 - \beta$

Symmetries of the Navier-Stokes

- If $\mathbf{u}(\mathbf{x}, 0)$ has one of the spatial symmetries of the Navier Stokes equations then $\mathbf{u}(\mathbf{x}, t)$ will retain this symmetry at all times
- Most of the symmetries of the Navier Stokes equations break down when non-constant forcing is considered
- Most of the symmetries of the Navier Stokes equations break down when non-infinite domain sizes are considered



Thank you
for your attention!