Turbulence Equilibrium Dynamics

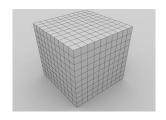


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Real space

Discretising space:

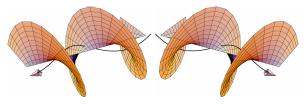


$$\mathbf{u}(\mathbf{x},t) o \mathbf{u}(\mathbf{x_n},t)$$
 where $\mathbf{x_n} = [x_{n_x},y_{n_y},z_{n_z}] = [n_x,n_y,n_z]\,\delta x,$ with $\mathbf{n} \in \mathbb{Z}^3$
$$N = L/\delta x, \quad N^3\text{-points, 3-vector components} +$$
 Incompressibility $\nabla \cdot \mathbf{u} = 0$

Fourier space

$$\tilde{\mathbf{u}}_{\mathbf{k}} = \frac{1}{L^3} \int e^{i\mathbf{k}\mathbf{x}} \mathbf{u} d\mathbf{x}^3, \quad \mathbf{u}(x) = \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{x}} \tilde{\mathbf{u}}_{\mathbf{k}}$$

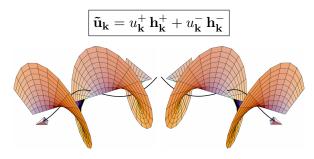
$$\mathbf{k}\cdot\mathbf{\tilde{u}}_{\mathbf{k}}=\mathbf{0}$$



$$\mathbf{\tilde{u}_k} = u_\mathbf{k}^+ \mathbf{h}_\mathbf{k}^+ + u_\mathbf{k}^- \mathbf{h}_\mathbf{k}^-$$

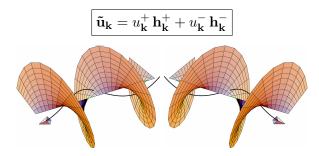
$$\mathbf{h}_{\mathbf{k}}^{\pm} = \frac{\mathbf{k} \times (\mathbf{k} \times \hat{\mathbf{e}})}{\sqrt{2} |\mathbf{k} \times (\mathbf{k} \times \hat{\mathbf{e}})|} \pm i \frac{\mathbf{k} \times \hat{\mathbf{e}}}{\sqrt{2} |\mathbf{k} \times \hat{\mathbf{e}}|}$$





$$\mathbf{h}_{\mathbf{k}}^{\pm} = \frac{\mathbf{k} \times (\mathbf{k} \times \hat{\mathbf{e}})}{\sqrt{2} |\mathbf{k} \times (\mathbf{k} \times \hat{\mathbf{e}})|} \pm i \frac{\mathbf{k} \times \hat{\mathbf{e}}}{\sqrt{2} |\mathbf{k} \times \hat{\mathbf{e}}|}$$
$$i \mathbf{k} \times \mathbf{h}_{\mathbf{k}}^{\pm} = \pm k \mathbf{h}_{\mathbf{k}}^{\pm}, \quad i \mathbf{k} \cdot \mathbf{h}_{\mathbf{k}}^{\pm} = 0$$

$$\mathbf{h}_{\mathbf{k}}^{s}\cdot(\mathbf{h}_{\mathbf{k}}^{s'})^{*}=\mathbf{h}_{\mathbf{k}}^{s}\cdot\mathbf{h}_{\mathbf{k}}^{-s'}=\mathbf{h}_{\mathbf{k}}^{s}\cdot\mathbf{h}_{-\mathbf{k}}^{s'}=\delta_{s,s'}$$



Energy:

$$\mathcal{E} = \frac{1}{2} \sum_{\mathbf{k}} \left[|u_{\mathbf{k}}^+|^2 + |u_{\mathbf{k}}^-|^2 \right]$$

Helicity:

$$\mathcal{H} = \frac{1}{2} \sum_{\mathbf{k}} k \left[|u_{\mathbf{k}}^{+}|^{2} - |u_{\mathbf{k}}^{-}|^{2} \right]$$



Navier Stokes in helical modes

Let $s_{\mathbf{k}}=\pm 1$, $s_{\mathbf{q}}=\pm 1$, $s_{\mathbf{p}}=\pm 1$. Then

$$\partial_t \left\langle \mathbf{h}_{-\mathbf{k}}^{\mathbf{s}_{\mathbf{k}}} e^{-i\mathbf{k}\cdot\mathbf{x}} \cdot \mathbf{w} \right\rangle = \left\langle \mathbf{h}_{-\mathbf{k}}^{\mathbf{s}_{\mathbf{k}}} e^{-i\mathbf{k}\cdot\mathbf{x}} \cdot \left(\nabla \times (\mathbf{u} \times \mathbf{w}) + \nu \nabla^2 \mathbf{w} + \nabla \times \mathbf{f} \right) \right\rangle$$

Navier Stokes:

$$\partial_t u_{\mathbf{k}}^{s_{\mathbf{k}}} = \left(\sum_{\mathbf{p} + \mathbf{q} = \mathbf{k}, s_{\mathbf{q}}, s_{\mathbf{p}}} C_{\mathbf{k}, \mathbf{q}, \mathbf{p}}^{s_{\mathbf{k}}, s_{\mathbf{q}}, s_{\mathbf{p}}} u_{\mathbf{q}}^{s_{\mathbf{q}}} u_{\mathbf{p}}^{s_{\mathbf{p}}} \right) - \nu |\mathbf{k}|^2 u_{\mathbf{k}}^{s_{\mathbf{k}}} + f_{\mathbf{k}}^{s_{\mathbf{k}}}$$

where

$$C_{\mathbf{k},\mathbf{q},\mathbf{p}}^{s_{\mathbf{k}},s_{\mathbf{q}},s_{\mathbf{p}}}u^{s_{\mathbf{q}}}u^{s_{\mathbf{p}}} = \frac{1}{2}(s_{\mathbf{q}}q - s_{\mathbf{p}}p) \langle \mathbf{h}_{-\mathbf{k}}^{\mathbf{s}_{\mathbf{k}}} \cdot \mathbf{h}_{\mathbf{q}}^{\mathbf{s}_{\mathbf{q}}} \times \mathbf{h}_{\mathbf{p}}^{\mathbf{s}_{\mathbf{p}}} \rangle$$

Number of Degrees of Freedom of N wavenumber modes

$$\tilde{\mathbf{u}}_{\mathbf{k}} = u_{\mathbf{k}}^{+} \, \mathbf{h}_{\mathbf{k}}^{+} + u_{\mathbf{k}}^{-} \, \mathbf{h}_{\mathbf{k}}^{-}$$

$$\partial_t u_{\mathbf{k}}^{s_{\mathbf{k}}} = \left(\sum_{\mathbf{p} + \mathbf{q} = \mathbf{k}, s_{\mathbf{q}}, s_{\mathbf{p}}} C_{\mathbf{k}, \mathbf{q}, \mathbf{p}}^{s_{\mathbf{k}}, s_{\mathbf{q}}, s_{\mathbf{p}}} u_{\mathbf{q}}^{s_{\mathbf{q}}} u_{\mathbf{p}}^{s_{\mathbf{p}}} \right) - \nu |\mathbf{k}|^2 u_{\mathbf{k}}^{s_{\mathbf{k}}} + f_{\mathbf{k}}^{s_{\mathbf{k}}}$$

If N wavenumbers are sufficient to resolve a given problem then the number of degrees of freedom N_F are

- (2) $u_{\mathbf{k}}^+, u_{\mathbf{k}}^-$ two modes
- (2) $u_{\mathbf{k}}^{\pm}$ is complex
- (1/2) $u_{-\mathbf{k}}^{\pm} = (u_{\mathbf{k}}^{\pm})^*$ realizability condition

$$N_F = 2 \times 2 \times \frac{1}{2} \times N = 2N$$

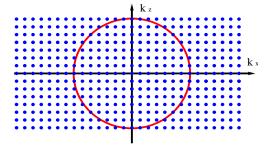


What if we forget dissipation?

Euler-equations

If $\nu=0$ dissipation can still occur by the formation of singularities. The Galerkin Truncated Euler-equations

$$\frac{d}{dt}u_{\mathbf{k}}^{s_{\mathbf{k}}} = \sum_{\mathbf{p}+\mathbf{q}=\mathbf{k}, s_{\mathbf{q}}, s_{\mathbf{p}}}^{|\mathbf{k}|, |\mathbf{q}|, |\mathbf{p}| \le k_{\max}} C_{\mathbf{k}, \mathbf{q}, \mathbf{p}}^{s_{\mathbf{k}}, s_{\mathbf{q}}, s_{\mathbf{p}}} u_{\mathbf{q}}^{s_{\mathbf{q}}} u_{\mathbf{p}}^{s_{\mathbf{p}}}$$



Equilibrium Dynamics

(Lee 1952; Hopf 1952; Kraichnan 1967, 1973; Orszag 1977)

$$\frac{d}{dt}u_{\mathbf{k}}^{s_{\mathbf{k}}} = \sum_{\mathbf{p}+\mathbf{q}=\mathbf{k}, s_{\mathbf{q}}, s_{\mathbf{p}}}^{|\mathbf{k}|, |\mathbf{q}|, |\mathbf{p}| \le k_{\max}} C_{\mathbf{k}, \mathbf{q}, \mathbf{p}}^{s_{\mathbf{k}}, s_{\mathbf{q}}, s_{\mathbf{p}}} u_{\mathbf{q}}^{s_{\mathbf{q}}} u_{\mathbf{p}}^{s_{\mathbf{p}}}$$

Liouville's condition

$$\frac{\partial}{\partial u_{\mathbf{k}}^{s_{\mathbf{k}}}} \left(\sum_{\mathbf{p}+\mathbf{q}=\mathbf{k}, s_{\mathbf{q}}, s_{\mathbf{p}}}^{|\mathbf{k}|, |\mathbf{p}| \leq k_{\max}} C_{\mathbf{k}, \mathbf{q}, \mathbf{p}}^{s_{\mathbf{k}}, s_{\mathbf{q}}, s_{\mathbf{p}}} u_{\mathbf{q}}^{s_{\mathbf{q}}} u_{\mathbf{p}}^{s_{\mathbf{p}}} \right) = 0$$

Density of states constant along trajectories

+Ergodicity assumption ⇒

 $\mathcal{P}[\mathbf{u}]$ is determined by the invariants of the system $(\mathcal{E},\mathcal{H})$

Equilibrium Dynamics

Micro-Canonical Ensamble

$$\mathcal{P}[\mathbf{u}] = \frac{1}{Z} \delta \left[\mathcal{E} - \frac{1}{2} \sum_{s_{\mathbf{k}, |\mathbf{k}| \le k_{\max}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^2 \right] \delta \left[\mathcal{H} - \frac{1}{2} \sum_{s_{\mathbf{k}, |\mathbf{k}| \le k_{\max}} s_{\mathbf{k}} k |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^2 \right]$$

Canonical Ensemble

$$\mathcal{P}[\mathbf{u}] = \frac{1}{Z} \exp \left[\frac{1}{2} \beta \sum_{s_{\mathbf{k}}, |\mathbf{k}| \le k_{\text{max}}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^2 + \frac{1}{2} \gamma \sum_{s_{\mathbf{k}}, |\mathbf{k}| \le k_{\text{max}}} s_{\mathbf{k}} k |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^2 \right]$$



Micro-Canonical Ensamble (Neglecting Helicity)

$$\begin{split} \mathcal{P}[\mathbf{u}] &= \frac{1}{Z} \delta \left[\mathcal{E} - \frac{1}{2} \sum_{s_{\mathbf{k}}, |\mathbf{k}| \leq k_{\text{max}}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right] \\ \left\langle |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right\rangle &= \frac{1}{Z} \int |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \delta \left[\mathcal{E} - \frac{1}{2} \sum_{s_{\mathbf{q}}, |\mathbf{q}| \leq k_{\text{max}}} |u_{\mathbf{q}}^{s_{\mathbf{q}}}|^{2} \right] \prod_{\mathbf{q}} du_{\mathbf{q}}^{s_{\mathbf{q}}} \\ &= \frac{1}{NZ} \int \left(\sum_{\mathbf{k}|, s_{\mathbf{k}}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right) \delta \left[\mathcal{E} - \frac{1}{2} \sum_{s_{\mathbf{k}}, |\mathbf{k}| \leq k_{\text{max}}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right] \prod_{\mathbf{q}, s_{\mathbf{q}}} du_{\mathbf{q}}^{s_{\mathbf{q}}} \\ &= \frac{2\mathcal{E}}{NZ} \int \delta \left[\mathcal{E} - \frac{1}{2} \sum_{s_{\mathbf{k}}, |\mathbf{k}| \leq k_{\text{max}}} |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right] \prod_{\mathbf{q}} du_{\mathbf{q}}^{s_{\mathbf{q}}} = 2\mathcal{E}/N \\ &E(k) = \frac{1}{2} \sum_{\mathbf{k} \leq |\mathbf{q}| \leq k+1} \left\langle |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \right\rangle \propto k^{2} \end{split}$$

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$$\mathcal{P}[\mathbf{u}] = \frac{1}{Z} \exp\left[\frac{1}{2} \left[\sum_{s_{\mathbf{k}},|\mathbf{k}| \le k_{\max}} (\beta + s_{\mathbf{k}} \gamma k) |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2}\right]\right]$$

$$\langle |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \rangle = \frac{1}{Z} \int |u_{\mathbf{k}}^{s_{\mathbf{k}}}|^{2} \mathcal{P}[\mathbf{u}] \prod_{\mathbf{q}} du_{\mathbf{q}}^{s_{\mathbf{q}}}$$

$$= \frac{1}{(\beta + s_{\mathbf{k}} \gamma k)}, \qquad \left(\int_{-\infty}^{+\infty} e^{-ax^{2}} dx = \sqrt{\frac{\pi}{a}}\right)$$

$$\langle u_{-\mathbf{k}}^{s_{\mathbf{k}}} \cdot w_{\mathbf{k}}^{s_{\mathbf{k}}} \rangle = \frac{s_{\mathbf{k}} k}{(\beta + s_{\mathbf{k}} \gamma k)}, \qquad \left(\int_{-\infty}^{+\infty} x^{2} e^{-ax^{2}} dx = \sqrt{\frac{\pi}{4a^{3}}}\right)$$

$$\langle e_{k} \rangle = \frac{1}{2} \langle |u_{\mathbf{k}}^{+}|^{2} + |u_{\mathbf{k}}^{-}|^{2} \rangle = \frac{\beta}{(\beta^{2} - \gamma k^{2})},$$

$$\langle h_{k} \rangle = \frac{k}{2} \langle |u_{\mathbf{k}}^{+}|^{2} - |u_{\mathbf{k}}^{-}|^{2} \rangle = \frac{\gamma k^{2}}{(\beta^{2} - \gamma^{2} k^{2})}$$

$$\langle e_k \rangle = \frac{1}{2} \langle |u_{\mathbf{k}}^+|^2 + |u_{\mathbf{k}}^-|^2 \rangle = \frac{\beta}{(\beta^2 - \gamma k^2)},$$

$$\langle h_k \rangle = \frac{k}{2} \langle |u_{\mathbf{k}}^+|^2 - |u_{\mathbf{k}}^-|^2 \rangle = \frac{\gamma k^2}{(\beta^2 - \gamma^2 k^2)}$$

$$E(k) = \frac{4\pi k^2}{(\beta^2 - \gamma^2 k^2)}, \quad H(k) = \frac{4\pi \gamma k^4}{(\beta^2 - \gamma^2 k^2)}$$

$$\mathcal{E} = \sum_{\mathbf{k}} \frac{\beta}{(\beta^2 - \gamma k^2)}, \quad \mathcal{H} = \sum_{\mathbf{k}} \frac{\gamma k^2}{(\beta^2 - \gamma k^2)}$$

$$\langle e_k \rangle = \frac{1}{2} \langle |u_{\mathbf{k}}^+|^2 + |u_{\mathbf{k}}^-|^2 \rangle = \frac{\beta}{(\beta^2 - \gamma k^2)},$$

$$\langle h_k \rangle = \frac{k}{2} \langle |u_{\mathbf{k}}^+|^2 - |u_{\mathbf{k}}^-|^2 \rangle = \frac{\gamma k^2}{(\beta^2 - \gamma^2 k^2)}$$

$$E(k) = \frac{4\pi k^2}{(\beta^2 - \gamma^2 k^2)}, \quad H(k) = \frac{4\pi \gamma k^4}{(\beta^2 - \gamma^2 k^2)}$$

$$k_{\text{max}} < |\beta/\gamma|$$

$$E(k) = \frac{4\pi k^2}{(\beta^2 - \gamma k^2)}, \quad H(k) = \frac{4\pi \gamma k^4}{(\beta^2 - \gamma k^2)}$$

 $Thermal\ equilibrium\ of\ large\ scales$

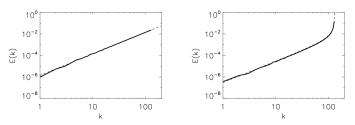


Figure 1: The energy spectra from two simulations of the truncated Euler equations with $k_{max} = 128$ and zero helicity (left) and $\mathcal{H}/\mathcal{E}k_{max} = 0.82$ (right). The dashed lines show the theoretical predictions given in eq. 1.2.

