

# 2D Turbulence spreading

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# Introduction

**Aim:** Understand how turbulence spreads in a 2D incompressible fluid when we constantly inject vortices in the domain.

The natural questions that arise are:

- Do the vortices injected spread until reaching the boundaries?
- If so, which profile does the energy and enstrophy distributions follow?

Additionally, we worked on a [point vortex model](#) to compare and observe similarities between the two models.



[Figure:](#) Turbulence in the atmosphere

# Problem setup (Navier-Stokes)

We consider the **2D incompressible Navier-Stokes equations** which, after nondimensionalizing, read:

$$\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p + \frac{1}{\text{Re}} \Delta \mathbf{u} + \mathbf{f}$$
$$\nabla \cdot \mathbf{u} = 0$$

The force is localized in the center of the domain.

- It injects vortices of size  $\sim 1/k_\ell$ .
- The injection region is a disk of size  $\pi/k_r$ .

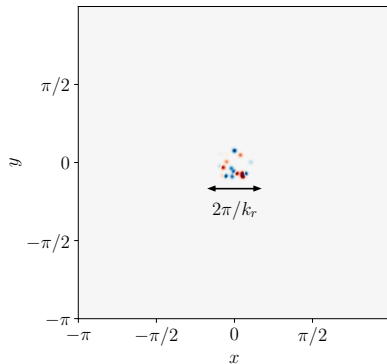


Figure: Forcing region for  $k_r = 8$

Thus, there are three parameters that control the dynamics of the system:  $k_\ell$ ,  $k_r$  and  $\text{Re}$ .

**Goal:** Understand the behavior of the system in the limit  $\text{Re} \rightarrow \infty$  and infinite domain ( $k_r \rightarrow \infty$ ).

We carried out two types of simulations:

- Fully parallel simulations: The domain is split among the processors.
- Embarrassingly parallel simulations: Several simulations are run simultaneously, each one in a different processor, in order to produce statistics.

The following table summarizes the values used in the simulations:

$k_r \backslash \text{Re}$	0.25	0.5	1	2	4	8	16	32	64	128
8	✓ <sub>512</sub>	✓ <sub>512</sub>	✓ <sub>512</sub>	✓✓ <sub>512</sub>	✓✓ <sub>1024</sub>	✓✓ <sub>1024</sub>	✓✓ <sub>1024</sub>	✓✓ <sub>2048</sub>	✓ <sub>2048</sub>	✓ <sub>4096</sub>
16				✓ <sub>1024</sub>	✓ <sub>2048</sub>	✓✓ <sub>2048</sub>	✓✓ <sub>2048</sub>	✓✓ <sub>2048</sub>	✓ <sub>4096</sub>	✓ <sub>4096</sub>
32				✓ <sub>2048</sub>	✓ <sub>4096</sub>	✓✓ <sub>4096</sub>	✓✓ <sub>4096</sub>	✓✓ <sub>4096</sub>	✓ <sub>8192</sub>	✓ <sub>8192</sub>
64						✓ <sub>8192</sub>	✓ <sub>8192</sub>	✓ <sub>8192</sub>		

**Table:** Simulations carried out during the study. ✓ = fully parallel simulations, ✓✓ = embarrassingly parallel simulations.

In all the simulations we set  $k_\ell = 4k_r$ .

The quantities monitored during the simulations are the following:

$$E_r = \sum_{r-\Delta r < \|\mathbf{x}\| \leq r} \|\mathbf{u}(\mathbf{x})\|^2$$

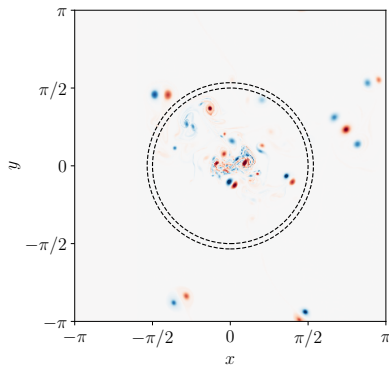
$$\Omega_r = \sum_{r-\Delta r < \|\mathbf{x}\| \leq r} |\omega(\mathbf{x})|^2$$

which account for the energy and enstrophy in annuli of radius  $r$  and width  $\Delta r$ .

We also defined quantities that measures where most of the energy or enstrophy is located:

$$\mathcal{R}_E^2 = \frac{\sum_{\Delta r < r \leq \pi} r^2 E_r}{\sum_{\Delta r < r \leq \pi} E_r} \quad \mathcal{R}_\Omega^2 = \frac{\sum_{\Delta r < r \leq \pi} r^2 \Omega_r}{\sum_{\Delta r < r \leq \pi} \Omega_r}$$

We call them *mean energy radius* and *mean enstrophy radius*, respectively.



## Problem setup (Point vortex)

We consider the simplest model of point vortices in 2D, which is based on the evolution by advection of a set of points where the vorticity  $\omega$  is singular. At all other points in  $\mathbb{R}^2$ , the vorticity is zero.

At all instants of time  $\omega$  can be thought as a sum of Dirac deltas:

$$\omega(\mathbf{x}, t) = \sum_{i=1}^N \Gamma_i \delta(\mathbf{x} - \mathbf{x}_i(t))$$

where  $\Gamma_i$  is the circulation of the  $i$ -th vortex, and  $\mathbf{x}_i(t)$  is its position.

A set of ODEs can be derived for  $\mathbf{x}_i(t)$  by imposing that they satisfy the 2D incompressible Euler equations:

$$\begin{aligned}\partial_t \omega + \mathbf{u} \cdot \nabla \omega &= 0 \\ \nabla \cdot \mathbf{u} &= 0\end{aligned}$$

In the point vortex simulations, we input vortices in the center and remove them when they reach the boundaries.

For this model, we monitor the number of vortices  $N_r$  in annuli of radius  $r$  and width  $\Delta r$ , in order to later on give an evolution of the density of number of vortices in rings of radius  $r$ :

$$\rho_r = \lim_{\Delta r \rightarrow 0} \frac{N_r}{2\pi r \Delta r}$$

We also monitor an equivalent *mean radius*  $\mathcal{R}_N$  weighted with the number of vortices in each annulus:

$$\mathcal{R}_N^2 = \frac{\sum_{\Delta r < r \leq \pi} r^2 N_r}{\sum_{\Delta r < r \leq \pi} N_r}$$



# Results

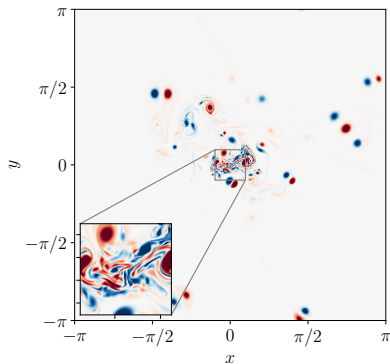


Figure: Navier-Stokes simulation with  $k_r = 16$  and  $\text{Re} = 128$ .

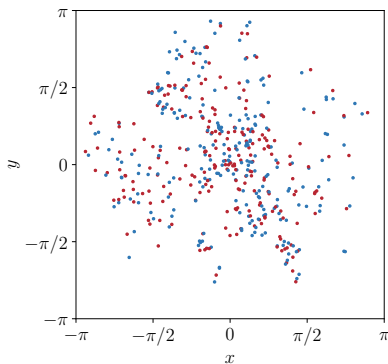
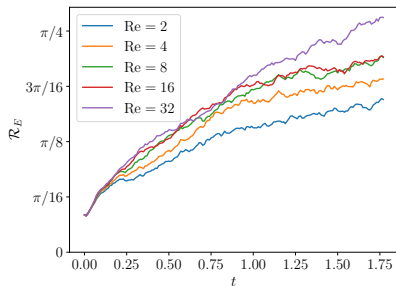
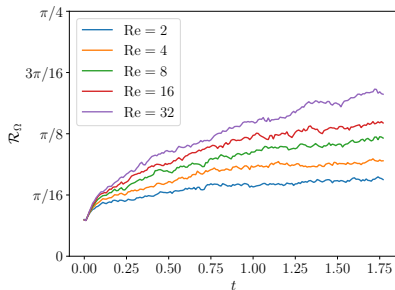


Figure: Point vortex simulation with  $k_r = 16$ .

# Results (Navier-Stokes)



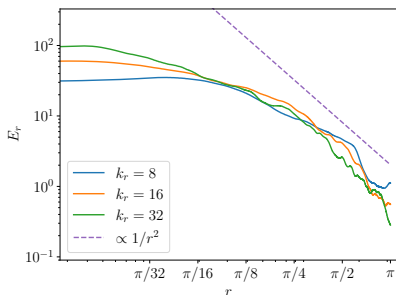
(a) Energy mean radius



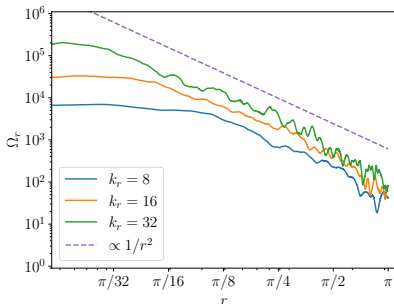
(b) Enstrophy mean radius

**Figure:** Mean energy radius and mean enstrophy radius for several runs varying  $Re$  at  $k_r = 16$ .

- $\mathcal{R}_E$  and  $\mathcal{R}_\Omega$  increase with time, which means that vortices spread far from the perturbation region.
- The higher the Reynolds number, the faster the spreading.



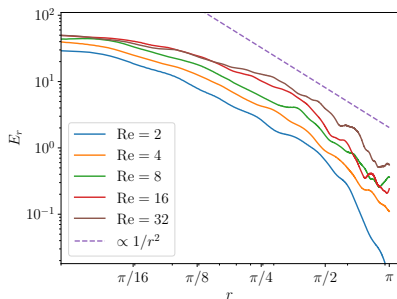
(a) Energy profile



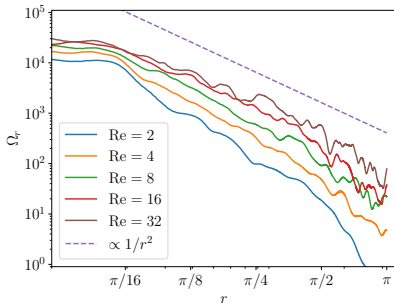
(b) Enstrophy profile

**Figure:** Energy and enstrophy profiles varying  $k_r$  at  $t = 1.66$  and  $\text{Re} = 32$ .

- Energy plots are smoother, while enstrophy plots are more spiky.
- As  $k_r$  increases, the apparent power law  $A/r^2$  for the enstrophy extends.



(a) Energy profile



(b) Enstrophy profile

**Figure:** Energy and enstrophy profiles varying  $Re$  at  $t = 1.75$  and  $k_r = 16$ .

- Again, energy plots are smoother than enstrophy plots.
- As  $Re$  increases, the magnitudes of both quantities increase due to less dissipation acting on the system.

# Results (Point vortex)

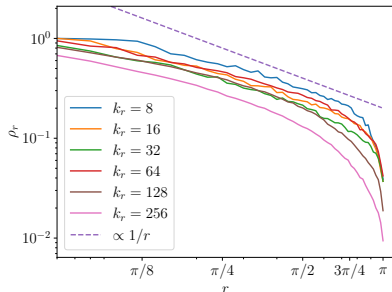


Figure: Density of the number of vortices as a function of the radius.

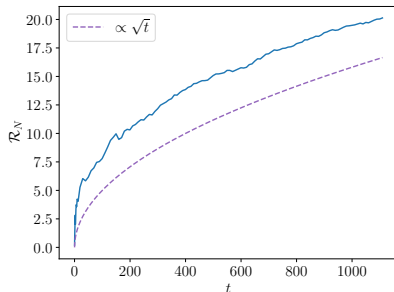


Figure: Mean radius weighted with the number of vortices.

- The density of vortices  $\rho_r$  following a power law  $A/r$  in the middle range of  $r$  implies that the flux of vortices on that region is almost constant.
- We observe a monotonous increase of  $\mathcal{R}_N$  with time, which resembles the Navier-Stokes case.

# Conclusions

- Energy injected in the center seems to spread across the domain, unlike the 3D case.
- The enstrophy seems to follow a power law  $A/r^2$ , that for large  $k_r$ , extends to the whole domain.
- When a stationary state is reached for the point vortex model, we observe a constant flux of vortices in the middle range of  $r$ .

## Future improvements:

- Extend the integration time of the Navier-Stokes simulations.
- Simulate cases for  $k_r = 64$  and higher Reynolds numbers.
- Replicate results adding a drag term  $-\alpha \mathbf{u}$  to the equations. Can we conclude the same?