## Units and dB Conversion

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## FIELD & POWER QUANTITIES

 $\vec{E}$  is the electric field intensity [volt/meter or V/m] and  $\vec{D}$  is the electric flux density [coulomb/meter<sup>2</sup> or C/m<sup>2</sup>]:

$$\vec{D} = \varepsilon \vec{E}$$
 with  $\varepsilon = \varepsilon_r \varepsilon_0$ , with  $\varepsilon_0 = 8.854 \times 10^{-12}$  (1)

 $\vec{H}$  is the magnetic field intensity [ampere/meter or A/m] and  $\vec{B}$  is the magnetic flux density [tesla or weber/meter<sup>2</sup> or T, Wb/m<sup>2</sup>]:

$$\vec{B} = \mu \vec{H}$$
 with  $\mu = \mu_r \mu_0$ , with  $\mu_0 = 4\pi 10^{-7}$  (2)

and

$$\eta_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}} = 376.7\Omega \approx 120\pi$$

In the **far-field** and **free-space**:  $\vec{H} = \frac{1}{\eta} \hat{r} \times \vec{E} = \frac{1}{\eta_0} \hat{r} \times \vec{E}$ , and the time-average radiated power **density** (based on the Poyinting vector for peak-valued vector phasors/sinusoidal excitations/**permanent response**) reduces to:

$$P_{radiated} [\text{w/m}^2] = \text{Re} \{ \frac{1}{2} \vec{E} \times \vec{H}^* \} = \text{Re} \{ \frac{1}{2\eta_0} \vec{E} \times (\hat{r} \times \vec{E}^*) \} =$$

$$P_{radiated} = \operatorname{Re} \left\{ \frac{1}{2\eta_0} [\vec{E} \bullet \vec{E} * \hat{r} - \vec{E} \bullet \hat{r} \hat{r}] \right\} = \operatorname{Re} \left\{ \frac{1}{2\eta_0} \vec{E} \bullet \vec{E} * \right\} = \frac{1}{2\eta_0} \left| \vec{E} \right|^2$$

OBS.: If working with time-varying fields, the *rms* power density does **not** contain the factor  $\frac{1}{2}$  ( $\vec{S} = \vec{E} \times \vec{H}$ , quantities are not phasors). This is also referred to as the *instantaneous* power density. Following a similar derivation

$$P_{radiated} \left[ W/m^2 \right] = \vec{E} \times \vec{H} = \frac{1}{\eta_0} \left| \vec{E} \right|^2$$
 (3)

where  $\vec{E}$  contains explicit time-dependence (not a phasor) and is given in [V/m]. Equation (3) is the one that is going to be used in class given that rms values are employed. Note that the relationship between the peak-valued phasor and the rms quantity is  $|\vec{E}_{peak}| = \sqrt{2} |\vec{E}_{rms}|$ , and with that both equations yield the same result. "In this class  $|\vec{E}| = |\vec{E}_{rms}|$ , unless explicitly indicated".

Note that if an antenna with gain G is employed in the system (and remember that you may consider the receive case equal to the transmit due to reciprocity – isotropic materials), the radiated power density can be written as

$$P_{radiated} \left[ W/m^2 \right] = \frac{P_{in}G}{4\pi r^2} \tag{4}$$

Where  $P_{in}$  is the input power to the antenna in [W] and r is the distance in [m]. Equaling (3) and (4):

$$\left| \vec{E} \right| = \sqrt{\frac{P_{in} G \eta_0}{4\pi r^2}} \tag{5}$$

which yields the electric field intensity in [V/m] at a distance r. Note that if you measure  $\vec{E}$  with a field meter, you can use (5) to determine the antenna gain at that given direction (considering the probe truly isotropic). The ratio of two quantities in dB is given by

POWER [dB] = 
$$10\log_{10}\left(\frac{P_2}{P_1}\right)$$
 (6)

VOLTAGE [dB] = 
$$20\log_{10}\left(\frac{V_2}{V_1}\right)$$
 (7)

$$CURRENT [dB] = 20 \log_{10} \left( \frac{I_2}{I_1} \right)$$
 (8)

For example, a voltage of 1 V is 120 dBμV:

VOLTAGE [dB
$$\mu$$
V] = 20 log<sub>10</sub>  $\left(\frac{1V}{1\mu V = 10^{-6}V}\right)$  = 20 log<sub>10</sub> 10<sup>6</sup> = 120 dB $\mu$ V