

## A. Transits

*a. Ignition Condition.* Runaway fusion only occurs in the degenerate WD interior where thermal expansion is suppressed as a cooling mechanism. The outer layers of the WD, however, are composed of a non-degenerate gas and it is therefore essential that a DM candidate penetrate this layer in order to ignite a SN. We parameterize this by a DM stopping power  $(dE/dx)_{\text{SP}}$ , the kinetic energy lost by the DM per distance traveled in the non-degenerate layer, and demand that

$$\left(\frac{dE}{dx}\right)_{\text{SP}} \ll \frac{m_\chi v_{\text{esc}}^2}{R_{\text{envelope}}}, \quad (1)$$

where  $R_{\text{envelope}} \approx 50$  km is the width of a WD envelope [? ].

The energy deposited during a continuous heating event such as a DM transit is best described in terms of a linear energy transfer  $(dE/dx)_{\text{LET}}$ , the kinetic energy of SM particles produced per distance traveled by the DM. If these products have a heating length  $L_0$  then the relevant energy deposit must at minimum be taken as the energy transferred over the transit distance  $L_0$ . Of course, we can always choose to consider energy deposits over a longer segment of the DM trajectory. Importantly, as per the general condition (??) such a deposition is *less* explosive unless  $L_0$  is smaller than the trigger size  $\lambda_T$ . Thus, we consider the energy deposited in a transit over the larger of these two length scales. Assuming the energy of the DM is roughly constant over this heating event, the ignition condition for transit heating is:

$$\left(\frac{dE}{dx}\right)_{\text{LET}} \gtrsim \frac{\mathcal{E}_{\text{boom}}}{\lambda_T} \cdot \max\left\{\frac{L_0}{\lambda_T}, 1\right\}^2. \quad (2)$$

Note that the DM stopping power in the non-degenerate layer  $(dE/dx)_{\text{SP}}$  and the linear energy transfer in the degenerate interior  $(dE/dx)_{\text{LET}}$  are possibly controlled by different physics and may have very different numerical values. In addition, a transit heating event satisfying condition (1) will have negligible energy loss over the parametrically smaller trigger size or heating length  $L_0$ , validating (2).

The above argument sums the individual energy deposits along the DM trajectory as though they are all deposited simultaneously. This is possible if the DM moves sufficiently quickly so that this energy does not diffuse out of the region of interest before the DM has traversed the region. We therefore require that the diffusion time  $\tau_{\text{diff}} \approx 10^{-12}$  s across a heated region at temperature  $T_f$  be larger than the DM crossing-time:

$$\tau_{\text{diff}} \sim \frac{L^2}{\alpha(T_f)} \gg \frac{L}{v_{\text{esc}}}, \quad (3)$$

where  $\alpha(T)$  is the temperature-dependent diffusivity, and the DM transits at the stellar escape velocity  $v_{\text{esc}} \sim 10^{-2}$ . This condition is more stringent for smaller regions, so we focus on the smallest region of interest,  $L = \lambda_T$ . (3) is then equivalent to demanding that the escape speed is greater than the conductive speed of the fusion wave front,  $v_{\text{cond}} \sim \alpha(T_f)/\lambda_T$ . Numerical calculations of  $v_{\text{cond}}$  are tabulated in [? ], and indeed condition (3) is satisfied for all WD densities.

*b. Event Rate: Wind Scenario.* The rate of transit events is given by the flux of DM passing through a WD

$$\Gamma_{\text{trans}} \sim \frac{\rho_\chi}{m_\chi} R_{\text{WD}}^2 \left(\frac{v_{\text{esc}}}{v_{\text{halo}}}\right)^2 v_{\text{halo}}, \quad (4)$$

where  $m_\chi$  is the DM mass,  $\rho_\chi$  is the DM density in the region of the WD, and  $R_{\text{WD}}$  is the WD radius. Here  $v_{\text{halo}} \sim 10^{-3}$  is the virial velocity of our galactic halo, and the transit rate contains an  $(v_{\text{esc}}/v_{\text{halo}})^2 \sim 100$  enhancement due to gravitational focusing.

## B. Collisions and Decays

*a. Ignition Condition.* For a point-like DM-DM collision or DM decay event releasing particles of heating length  $L_0$ , ignition will occur if the total energy in SM products satisfies condition (??). Such an event will likely result in both SM and dark sector products, so we parameterize the resulting energy in SM particles as a fraction  $f_{\text{SM}}$  of the DM mass. For non-relativistic DM, the DM mass is the dominant source of energy and therefore  $f_{\text{SM}} \lesssim 1$  regardless of the interaction details. With this parameterization, a single DM-DM collision or DM decay has an ignition condition:

$$m_\chi f_{\text{SM}} \gtrsim \mathcal{E}_{\text{boom}} \cdot \max\left\{\frac{L_0}{\lambda_T}, 1\right\}^3. \quad (5)$$

We are thus sensitive to DM masses  $m_\chi \gtrsim 10^{16}$  GeV.

However, there is the possibility if DM is captured in the WD that allows collisions or decays of lower mass DM to ignite the star. Multiple DM-DM collisions or decays in a sufficiently small region can occur rapidly enough to be counted as a single heating event. This is similar in nature to a transit heating event, where multiple scatters across a transit length  $\lambda_T$  can release an energy  $\mathcal{E}_{\text{boom}}$  and satisfy (2) even if any individual scatter is not explosive by itself. If a single DM-DM collision is unable to ignite the star, the sum total of the energy released in many collisions can still result in a SN if

$$m_\chi f_{\text{SM}} \gtrsim \frac{\mathcal{E}_{\text{boom}}}{N_{\text{mult}}} \cdot \max\left\{\frac{L_0}{\lambda_T}, 1\right\}^3, \quad N_{\text{mult}} \gtrsim 1, \quad (6)$$

We define  $N_{\text{mult}}$  as the number of collisions within a region of size  $\max\{\lambda_T, L_0\}^3$  (or smaller) during a diffusion time  $\tau_{\text{diff}}$ . This necessarily depends on additional DM parameters and the evolution of the captured DM in the star. These are discussed in detail below.

*b. Event Rate: DM Wind.* For the remainder of this section, all numerical quantities are evaluated assuming a WD lifetime  $\tau_{\text{WD}} \sim 5$  Gyr and central WD density  $n_{\text{ion}} \sim 10^{31} \text{ cm}^{-3}$ . At this density, the relevant WD parameters are approximately:

$$M_{\text{WD}} \approx 1.25 M_\odot, \quad R_{\text{WD}} \approx 4000 \text{ km}, \quad v_{\text{esc}} \approx 2 \times 10^{-2}. \quad (7)$$

We also assume a typical WD temperature  $T \sim \text{keV}$ . DM that is not captured traverses the WD in  $R_{\text{WD}}/v_{\text{esc}} \approx 0.1$  s, and the rate of DM-DM collisions within the WD parameterized by cross-section  $\sigma_{\chi\chi}$  is:

$$\Gamma_{\text{ann}} \sim \left(\frac{\rho_\chi}{m_\chi}\right)^2 \sigma_{\chi\chi} \left(\frac{v_{\text{esc}}}{v_{\text{halo}}}\right)^3 v_{\text{halo}} R_{\text{WD}}^3. \quad (8)$$

Similarly the net DM decay rate inside the WD parameterized by a lifetime  $\tau_\chi$  is:

$$\Gamma_{\text{decay}} \sim \frac{1}{\tau_\chi} \frac{\rho_\chi}{m_\chi} \left(\frac{v_{\text{esc}}}{v_{\text{halo}}}\right) R_{\text{WD}}^3. \quad (9)$$

*c. Event Rate: DM Capture.* For the DM to be captured in a WD, it must lose energy  $\sim m_\chi v^2$ , where  $v$  is the relative DM velocity (in the rest frame of the WD) asymptotically far away. Properly, this DM velocity is described by a (boosted) Maxwell distribution peaked at the galactic virial velocity  $v_{\text{halo}} \sim 10^{-3}$ . Since typically  $v \ll v_{\text{esc}}$ , the DM has initial velocity  $v_{\text{esc}}$  in the star and must lose a fraction  $(v/v_{\text{esc}})^2$  of its energy to become captured.

The physics of DM capture can be made more precise for a specific interaction. Consider a spin-independent, elastic scattering off ions with cross section  $\sigma_{\chi A}$ . Assuming  $m_{\text{ion}} \ll m_\chi$ , the typical momentum transfer in an elastic scatter is  $q \sim \mu_A v_{\text{esc}} \approx 200$  MeV, where  $\mu_A \sim m_{\text{ion}}$  is the reduced mass of the DM-nuclei system. This corresponds to an energy transfer  $q^2/m_{\text{ion}} \sim m_{\text{ion}} v_{\text{esc}}^2 \approx 10$  MeV. The average number of DM scatters during a full transit of the WD is simply a ratio of the mean free path to the size of the WD

$$N_{\text{scat}} \sim n_{\text{ion}} \sigma_{\chi A} R_{\text{WD}}. \quad (10)$$

If  $N_{\text{scat}} < 1$ , then  $N_{\text{scat}}$  is the probability for a *single* scatter to occur during the transit. Thus, DM with initial velocities less than

$$v_{\text{cap}}^2 \sim v_{\text{esc}}^2 \left(\frac{m_{\text{ion}}}{m_\chi}\right) \max\{N_{\text{scat}}, 1\}. \quad (11)$$

will be captured in the WD. A detailed calculation of the rate of DM capture [?] yields

$$\Gamma_{\text{cap}} \sim \Gamma_{\text{trans}} \cdot \min\{N_{\text{scat}}, 1\} \left(\frac{v_{\text{cap}}}{v_{\text{halo}}}\right)^2. \quad (12)$$

Here we assume  $v_{\text{cap}} \lesssim v_{\text{halo}}$ ; otherwise, the capture rate is simply  $\Gamma_{\text{trans}}$ . Since the momentum transfer  $q$  is roughly of order the inverse nuclear size, it is reasonable to expect the DM coherently scatters off all nucleons in the nucleus. Indeed, the average per-nucleon cross section (spin-independent) is

$$\sigma_{\chi A} = A^2 \left(\frac{\mu_A}{\mu_n}\right)^2 F^2(q) \sigma_{\chi n}, \quad (13)$$

where  $F^2(q) \approx 0.1$  is the Helm form factor [? ]. We can compare the cross section sufficient for capture (11) to the limits from direct detection experiments. Currently, the bound on spin-independent DM nuclear elastic scatters from XENON 1T is

$$\sigma_{\chi n} < 10^{-45} \text{ cm}^2 \left( \frac{m_\chi}{10^3 \text{ GeV}} \right). \quad (14)$$

It is interesting that any DM candidate whose scattering cross section barely avoids the direct detection constraint can be efficiently captured in a WD,  $v_{\text{cap}} \approx 0.5 v_{\text{halo}}$ .

If the DM is captured, it eventually thermalizes to an average velocity

$$v_{\text{th}} \sim \sqrt{\frac{T}{m_\chi}} \approx 10^{-12} \left( \frac{m_\chi}{10^{16} \text{ GeV}} \right)^{-1/2} \quad (15)$$

and settles at the thermal radius

$$R_{\text{th}} \sim \left( \frac{T}{G m_\chi \rho_{\text{WD}}} \right)^{1/2} \approx 0.1 \text{ cm} \left( \frac{m_\chi}{10^{16} \text{ GeV}} \right)^{-1/2}, \quad (16)$$

where its kinetic energy balances against the gravitational potential energy of the (enclosed) WD mass. For simplicity we take a constant WD density  $\rho_{\text{WD}} \sim n_{\text{ion}} m_{\text{ion}}$  within  $R_{\text{th}}$ . Of course, the timescale to reach thermalization depends on the nature of the DM-SM interaction. This has been explicitly calculated in the case that the DM loses energy via elastic nuclear scatters, see [? ]. First, the DM passes through the WD many times before the size of its orbit becomes fully contained within the star. This occurs after a time

$$t_1 \sim \left( \frac{m_\chi}{m_{\text{ion}}} \right)^{3/2} \frac{R_{\text{WD}}}{v_{\text{esc}}} \frac{1}{N_{\text{scat}}} \frac{1}{\max\{N_{\text{scat}}, 1\}^{1/2}} \approx 2 \times 10^3 \text{ yr} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right)^{3/2} \left( \frac{\sigma_{\chi A}}{10^{-38} \text{ cm}^2} \right)^{-3/2}. \quad (17)$$

Note that the time to complete a single orbit is simply the free-fall timescale:

$$t_{\text{ff}} \sim \sqrt{\frac{1}{G \rho_{\text{WD}}}} \approx 0.5 \text{ s}. \quad (18)$$

This stage is relevant only if the energy loss after a single transit is does not exceed  $\sim m_\chi v_{\text{esc}}^2$ :

$$\left( \frac{m_{\text{ion}}}{m_\chi} \right) \max\{N_{\text{scat}}, 1\} < 1. \quad (19)$$

This is the case for any cross sections which satisfy the XENON bound (14). Subsequently, the DM completes many orbits within the star until dissipation from elastic scatters reduces the orbital size to the thermal radius. This occurs after a characteristic time

$$t_2 \sim \left( \frac{m_\chi}{m_{\text{ion}}} \right) \frac{1}{n_{\text{ion}} \sigma_{\chi A}} \frac{1}{v_{\text{ion}}} \approx 30 \text{ yr} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right) \left( \frac{\sigma_{\chi A}}{10^{-38} \text{ cm}^2} \right)^{-1}. \quad (20)$$

where  $v_{\text{ion}} \sim \sqrt{\frac{T}{m_{\text{ion}}}}$ . Thus, DM will only settle at the thermal radius if the total thermalization time is shorter than the age of the WD:

$$t_1 + t_2 < \tau_{\text{WD}}. \quad (21)$$

We now turn towards the rate of DM-DM collisions for captured DM. Ultimately, the most interesting constraint in the capture scenario will be from the focusing of annihilations during self-gravitational collapse of a DM sphere at the thermal radius. However, we first carefully examine the effect of annihilations on the evolution of captured DM. To begin, the settling DM constitutes a number density of DM throughout the WD volume as well as outside the star. We can compare the total rate of annihilations of in-falling DM to the rate of DM capture. This annihilation rate is dominated by the DM density inside the star with orbits near the thermal radius, and the depletion of in-falling DM can be ignored as long as: bluecondition

After a settling time has passed, DM will begin steadily accumulating at  $R_{\text{th}}$ . If (??) is satisfied, the accumulation rate is roughly the same as the capture rate. However, this density of accumulating DM is also depleting due to annihilations. Eventually, these two rates become comparable and there is an equilibrium number of DM particles

$$N_{\text{eq}} \sim \left( \frac{\Gamma_{\text{cap}} R_{\text{th}}^3}{\sigma_{\chi\chi} v_{\text{th}}} \right)^{1/2} \approx 5 \times 10^{27} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right)^{-3/2} \left( \frac{\sigma_{\chi\chi}}{10^{-40} \text{ cm}^2} \right)^{-1/2} \left( \frac{\sigma_{\chi A}}{10^{-38} \text{ cm}^2} \right)^{1/2} \left( \frac{\rho_\chi}{0.4 \text{ GeV/cm}^3} \right)^{1/2}. \quad (22)$$

Of course, there is no guarantee that this equilibrium is achieved within the age of the WD. In that case, annihilations can be ignored and the total number of DM particles accumulated is simply

$$N_{\text{life}} \sim \Gamma_{\text{cap}} \tau_{\text{WD}} \approx 10^{30} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right)^{-2} \left( \frac{\sigma_{\chi A}}{10^{-38} \text{ cm}^2} \right) \left( \frac{\rho_\chi}{0.4 \text{ GeV/cm}^3} \right) \quad (23)$$

However, if the collected mass of DM at the thermal radius ever exceeds the WD mass within this volume, then there is the possibility of self-gravitational collapse of the DM. The critical number of DM particles needed for collapse is given by

$$N_{\text{crit}} \sim \frac{\rho_{\text{WD}} R_{\text{th}}^3}{m_\chi} \approx 10^{27} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right)^{-5/2}. \quad (24)$$

This can only be achieved if the time to collect a critical mass of DM is shorter than the time for annihilations to deplete this mass sufficiently *and* shorter than the WD lifetime. Thus the condition for collapse is:

$$N_{\text{crit}} < N_{\text{eq}}, \quad N_{\text{crit}} < N_{\text{life}}. \quad (25)$$

Once the DM sphere becomes self-gravitating, it will start to collapse. This When the DM sphere collapses, its initial evolution is determined by the dynamical free-fall time  $t_{\text{ff}}$ . Soon after, further collapse will proceed as a result of the DM dissipating its energy through elastic scatters, thereby reducing its orbital radius. The timescale for collapse is initially independent of DM velocity but decreases once the DM velocity exceeds the velocity of ions  $v_{\text{ion}}$ :

$$t_{\text{col}} \sim \frac{m_\chi}{\rho_{\text{WD}} \sigma_{\chi A}} \frac{1}{\max\{v_{\text{ion}}, v\}}, \quad v \sim \sqrt{\frac{GN_{\text{crit}} m_\chi}{r}}. \quad (26)$$

Annihilations in the collapsing DM density can be ignored until the timescale for annihilations becomes of order the collapse time. This occurs at a characteristic radius

$$R_{\chi\chi} \sim R_{\text{th}} \min \left\{ \left( \frac{\sigma_{\chi\chi}}{\sigma_{\chi A}} \right)^{1/3}, \left( \frac{\sigma_{\chi\chi}}{\sigma_{\chi A}} \right)^{2/7} \left( \frac{m_{\text{ion}}}{m_\chi} \right)^{1/7} \right\}. \quad (27)$$

Essentially, the number of collapsing DM particles is depleting by an  $\mathcal{O}(1)$  fraction once the DM sphere has focused to within  $R_{\chi\chi}$ . Of course, such a collapse of the accumulated DM in the WD is only sensible if

$$R_{\chi\chi} < R_{\text{th}}. \quad (28)$$

This condition will ultimately be satisfied for any DM mass and annihilation cross section which satisfy the ignition condition (6) during a collapse.

We also briefly mention the possibility that the number of DM particles initially collapsing can be much larger than  $N_{\text{crit}}$ . Once the accumulated DM reaches  $N_{\text{crit}}$ , evolution of the DM profile can either be collapse or further collection. The later scenario occurs if the collapse time is greater than the time to collect a critical number of DM particles:

$$N_{\text{crit}} < \Gamma_{\text{cap}} t_{\text{col}}. \quad (29)$$

This is true for DM masses above  $\sim 10^{14}$  GeV, independent of scattering cross section  $\sigma_{\chi A}$ .

There are two potential evolutions of the captured DM: either the DM collapses or it does not. In the later case, the DM has reached its equilibrium number at the thermal radius or is still continuing to accumulate, not yet having the critical mass necessary for collapse within its lifetime. We have checked that this scenario does not yield any constraints on DM parameters which ignite the star. Thus we turn our attention to collapsing DM, characterized by (25). The number of collisions  $N_{\text{mult}}$  that can be counted as a single heating event is given by an integral of the annihilation rate over a diffusion time  $\tau_{\text{diff}}$ . This depends on the radius of the collapsing DM sphere, but there exists an optimum radius  $R_*$  at which the annihilation rate is maximized. Assuming the diffusion time is the shortest timescale during the collapse, this number is of order:

$$N_{\text{mult}} \sim \left( \frac{N_{\text{crit}}}{R_*^3} \right)^2 \sigma_{\chi\chi} v_{\text{ff}}(R_*) \min\{L_{\text{heat}}, R_*\}^3 \tau_{\text{diff}}, \quad (30)$$

Naively, we would suspect that the optimum radius is just set by the scale at which annihilations during the collapse become relevant  $R_* \sim R_{\chi\chi}$ . However, there may be some stabilizing pressure which prevents the DM from collapsing

multicollision.pdf

FIG. 1: Constraints on DM-DM annihilation cross-section into SM particles which deposit their energy compactly within a trigger size  $\lambda_T$  during self-gravitational collapse in a WD. Bounds come from observation of a single  $1.25 M_\odot$  WD assuming efficient capture of the DM and settling time (see text). We also take the DM collapse to be stabilized by formation of a BH.

below a certain radius. This is to be expected for heavy, composite DM, although such a stable radius would depend on unknown physics. Famously, gravity itself provides such a “pressure”, arresting collapses below the Schwarzschild radius by the formation of a black hole:

$$R_{\text{BH}} \sim G N_{\text{crit}} m_\chi \approx 5 \times 10^{-15} \text{ cm} \left( \frac{m_\chi}{10^{10} \text{ GeV}} \right)^{-3/2}. \quad (31)$$

If the DM stabilizes into a BH, the condition that not even a *single* collision occurs during the collapse to  $R_{\text{BH}}$  is simply

$$\left( \frac{N_{\text{crit}}}{R_{\text{BH}}^3} \right) \sigma_{\chi\chi} \frac{m_\chi}{\rho_{\text{WD}} \sigma_{\chi A}} \lesssim 1. \quad (32)$$

This is the most stringent bound that can be placed on DM masses greater than  $\mathcal{E}_{\text{boom}}$ , for which a single collision is capable of igniting the star.

Finally, we examine the decay rate of captured DM in the WD. Of course, this rate is proportional to the number of DM particles within the WD at any given instance. In the wind scenario (9), this number is given by  $\sim (\Gamma_{\text{trans}} \times t_{\text{ff}})$ . In the capture scenario, the number of DM particles available for decay is instead determined by the settling time inside the WD. Note that for DM masses above  $\mathcal{E}_{\text{boom}}$ , it is the case that gravitational collapse will occur within the age of the WD, i.e.  $N_{\text{crit}} < N_{\text{life}}$ . If the DM is captured via elastic scatters, the rate of decays is given by

$$\Gamma_{\text{decay}} \sim \frac{1}{\tau_\chi} \Gamma_{\text{cap}} t_2 \quad (33)$$

Since, (33) is independent of  $\sigma_{\chi A}$  and scales inversely with  $m_\chi$ , the rate of captured DM decays in the WD is simply enhanced by a factor  $\approx 10^5$  compared to the rate of transiting “wind” decays (9). Lastly, we have checked that the enhancement of the DM number in the star is never sufficient to yield novel bounds on the DM lifetime if multiple decays are required.

### C. Constraints

We now turn towards constraints on DM-DM collisions in the capture scenario. In Figure 1, we derive the constraints on  $\sigma_{\chi\chi}$  due to the observation of a single  $1.25 M_\odot$  in our local DM density. We distinguish between whether a single collision can blow up the WD or multiple collisions are required, although both bounds are derived from the focusing of annihilations during gravitational core collapse of DM in the star. We also take the DM collapse to be stabilized by formation of a BH. The results of Figure 1 are valid for any SM annihilation products which deposit their energy compactly upon release within the trigger size  $\lambda_T$  see discussion in section 3.