

Homework 5. Problem 2 (Plücker 2.4)

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The energy (neglecting translation as it does not couple to rotation)

is:

$$H = \frac{1}{2} \pm \underbrace{(\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta)}_{\text{kinetic energy of rotation}} - \underbrace{\mu E \cos \theta}_{\text{simple dipolar interaction with the applied field (along z)}}$$

The question

is formulated in terms of canonical momenta.

If \mathcal{L} is the Lagrangian $\mathcal{L} = \mathcal{L}(q_i, \dot{q}_i, t)$, canonically conjugate variables (Q_i, P_i)

are introduced by the transformation:

$$q_i \rightarrow Q_i \\ P_i = \frac{\partial \mathcal{L}}{\partial \dot{q}_i}$$

In this case, the Jacobian of the transformation is 1:

$$\int d\vec{q} d\vec{p} = \int d\vec{Q} d\vec{P}$$

In our case, the kinetic energy $\mathcal{L}_K = \frac{1}{2} I (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta)$. The interaction term does not depend on the velocity. Hence: in polar coordinates:

$$P_\theta = \frac{\partial \mathcal{L}_K}{\partial \dot{\theta}} = \frac{1}{2} I \dot{\theta} = I \dot{\theta}$$

$$P_\phi = \frac{\partial \mathcal{L}_K}{\partial \dot{\phi}} = \frac{1}{2} I \sin^2 \theta \dot{\phi} = I \sin^2 \theta \dot{\phi}$$

In terms of the canonically conjugate momenta:

$$H = \frac{1}{2} I \left(\left(\frac{P_\theta}{I} \right)^2 + \left(\frac{P_\phi}{I \sin^2 \theta} \right)^2 \sin^2 \theta \right) - \mu E \cos \theta \\ = \frac{P_\theta^2}{2I} + \frac{P_\phi^2}{2I \sin^2 \theta} - \mu E \cos \theta$$

The partition function is (single particle), in canonical coordinates:

$$Z = \int \frac{d\theta d\phi dP_\theta dP_\phi}{h^2} e^{-\beta H} = \frac{1}{h^2} \int_0^\pi d\theta \int_0^{2\pi} d\phi \int_{-\infty}^{\infty} dP_\theta \int_{-\infty}^{\infty} dP_\phi e^{-\beta \left[\frac{P_\theta^2}{2I} + \frac{P_\phi^2}{2I \sin^2 \theta} \right] - \mu E \cos \theta}$$

The Gaussian integral: $\int_{-\infty}^{\infty} dx e^{-ax} = \sqrt{\frac{\pi}{a}}$

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Hence:

$$Z = \frac{1}{h^2} \int_0^\pi \int_0^\pi \frac{1}{2\pi} \sqrt{\frac{2\pi I}{\beta}} \sqrt{\frac{2\pi I \sin^2 \theta}{\beta}} d\theta e^{\beta \mu E \cos \theta}$$

$$Z = \frac{4\pi^2 I}{h^2 \beta} \int_0^\pi \sin \theta e^{\beta \mu E \cos \theta} d\theta \quad \text{Call } x = \beta \mu E \cos \theta$$

$$dx = -\beta \mu E \sin \theta d\theta$$

$$\Rightarrow Z = \frac{4\pi^2 I}{h^2 \beta} \int_{\beta \mu E}^{-\beta \mu E} \left(-\frac{dx}{\beta \mu E \sin \theta} \right) e^x =$$

$$= \frac{4\pi^2 I}{h^2 \beta \mu E} \int_{\beta \mu E}^{-\beta \mu E} dx e^x = \frac{4\pi^2 I}{h^2 \beta \mu E} (e^{\beta \mu E} - e^{-\beta \mu E}) \rightarrow$$

$$\boxed{Z = \frac{8\pi^2 I}{h^2 \beta \mu E} \sinh(\beta \mu E) = \frac{2I}{h^2 \beta \mu E} \sinh(\beta \mu E)}$$

b) Polarization is defined.

$$P = \frac{N}{V} \langle \mu \cos \theta \rangle$$

Given that $Z = \int \frac{d\theta dp_\theta dp_\phi}{h^2} e^{-\beta K} e^{\beta \mu E \cos \theta}$

$$= \int \frac{d\theta dp_\theta dp_\phi}{h^2} e^{\beta \mu E \cos \theta}$$

$$P = \frac{N}{V} \frac{1}{Z} \frac{1}{h^2} \int d\theta dp_\theta dp_\phi (\mu \cos \theta) e^{-\beta K} e^{\beta \mu E \cos \theta}$$

$$= \frac{N}{V} \frac{1}{Z} \frac{\partial}{\partial \beta E} \int d\theta dp_\theta dp_\phi e^{-\beta K} e^{\beta \mu E \cos \theta} = \frac{N}{V} \frac{\partial Z}{\partial (\beta E)} = \left[\frac{N}{V} \frac{\partial \ln Z}{\partial (\beta E)} \right] = P$$

$$\begin{aligned}
 \therefore c) \quad P &= \frac{Nk_B T}{V} \frac{\partial}{\partial E} \ln \left[\frac{2I}{\hbar^2 \rho \mu E} \sinh(\beta \mu E) \right] = \\
 &= \frac{Nk_B T}{V} \frac{\partial}{\partial E} \ln \left(\frac{1}{E} \sinh(\beta \mu E) \right) = \frac{Nk_B T}{V} \left(-\frac{1}{E} + \frac{\partial}{\partial E} \sinh(\beta \mu E) \right) \\
 &= \frac{Nk_B T}{V} \left[-\frac{1}{E} + \frac{\cosh(\beta \mu E) \beta \mu}{\sinh(\beta \mu E)} \right] = \boxed{\frac{N}{V} \left[\mu \coth(\beta \mu E) - \frac{k_B T}{E} \right] = P}
 \end{aligned}$$

d) Now:

$$\epsilon E = \epsilon_0 E + P = \epsilon_0 E + \frac{N}{V} \left[\mu \coth(\beta \mu E) - \frac{k_B T}{E} \right]$$

For small field: $\coth x = \frac{1}{x} + \frac{x}{3} + \dots$

$$\Rightarrow \epsilon E \simeq \epsilon_0 E + \frac{N}{V} \left[\cancel{\mu} \frac{1}{\cancel{\beta \mu E}} + \frac{1}{3} \beta \mu E - \frac{\cancel{k_B T}}{E} \right] =$$

$$= \left(\epsilon_0 + \frac{N}{V} \frac{1}{3} \beta \mu \right) E \quad \Rightarrow \quad \boxed{\epsilon = \epsilon_0 + \frac{N}{V} \frac{1}{3} \beta \mu}$$