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Study of jet quenching with Z+jet correlations in PbPb and pp collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02\,\mathrm{TeV}$

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Abstract

The production of jets in association with Z bosons, reconstructed via the $\mu^+\mu^-$ and e^+e^- decay channels, is studied in pp and, for the first time, in PbPb collisions. Both data samples were collected by the CMS experiment at the LHC, at a nucleon-nucleon center-of-mass energy of 5.02 TeV. The PbPb collisions were analyzed in the 0–30% centrality range. The back-to-back azimuthal alignment was studied in both pp and PbPb collisions for Z bosons with transverse momentum $p_{\mathrm{T}}^{\mathrm{Z}} > 60\,\mathrm{GeV}/c$ and a recoiling jet with $p_{\mathrm{T}}^{\mathrm{jet}} > 30\,\mathrm{GeV}/c$. The p_{T} imbalance, $x_{\mathrm{jZ}} = p_{\mathrm{T}}^{\mathrm{jet}}/p_{\mathrm{T}}^{\mathrm{Z}}$, as well as the average number of jet partners per Z, R_{jZ}, were studied in intervals of $p_{\mathrm{T}}^{\mathrm{Z}}$. The R_{jZ} is found to be smaller in PbPb than in pp collisions, which suggests that in PbPb collisions a larger fraction of partons associated with the Z bosons fall below the $30\,\mathrm{GeV}/c$ $p_{\mathrm{T}}^{\mathrm{jet}}$ threshold because they lose energy.

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The correlated production of vector bosons and jets in hard parton scatterings occurring in ultrarelativistic heavy ion collisions provides an ideal probe of the quark-gluon plasma (QGP), a deconfined state of quarks and gluons [1, 2]. Final-state jets are created by the fragmentation of outgoing partons that interact strongly with the produced medium and lose energy [3–11], a phenomenon ("jet quenching") observed at RHIC [12, 13] and the LHC [14–18]. The transverse momentum (p_T) of the jet is highly correlated (through momentum conservation) with that of the associated Z boson, which is not affected by the medium [19–21] and reflects the initial energy of the parton. The lost energy can be related, via theoretical models, to the thermodynamical and transport properties of the medium [9–11, 22–24]. At LHC energies, Z+jet production is dominated by quark jets for $p_T^{jet} \gtrsim 30\,\text{GeV}/c$ [21], the primary subprocess being $q(\overline{q}) + g \to Z + q(\overline{q})$ [19], hence providing information on the parton flavor (quark or gluon) and kinematics, and allowing detailed studies of the energy loss with a well-defined production process. The Z-jet correlations are particularly well suited to perform tomographic studies of the QGP, given the minimal contributions from background channels [20, 25–27]. Correlations of jets with isolated photons are accessible at higher rates and carry similar information on parton energy loss [25–29], but suffer from an irreducible background of photons from jet fragmentation [17, 30], as well as larger uncertainties arising from the experimental selection of photon candidates.

This Letter describes the identification of Z+jet pairs in pp and PbPb collisions, and the first characterization of parton energy loss through angular and p_T correlations between the jet and the Z, reconstructed in dimuon or dielectron decays. The back-to-back azimuthal alignment of the Z and jets is studied through the difference $\Delta\phi_{jZ} = |\phi^{jet} - \phi^Z|$. The Z+jet momentum imbalance is studied using the $x_{jZ} = p_T^{jet}/p_T^Z$ ratio and the p_T^Z dependence of its mean value, $\langle x_{jZ} \rangle$. The average number of jet partners per Z boson, R_{jZ}, is also reported. The analysis exploits PbPb and pp data samples collected by CMS at a nucleon-nucleon center-of-mass energy of 5.02 TeV, corresponding to integrated luminosities of $404\,\mu\text{b}^{-1}$ and $27.4\,\text{pb}^{-1}$, respectively.

The central feature of CMS is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward hadron calorimeters extend the pseudorapidity (η) coverage and are used for PbPb event selection. Muons are measured in gas-ionization detectors located outside the solenoid. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [31].

The event samples are selected online with dedicated lepton triggers, and cleaned offline to remove noncollision events, such as beam-gas interactions or cosmic-ray muons [32]. In addition, events are required to have at least one reconstructed primary interaction vertex. The $Z \to e^+e^-$ events are triggered if two ECAL clusters [33] have transverse energy greater than 15 GeV and $|\eta| < 2.5$, while the $Z \to \mu^+\mu^-$ triggers require one muon of $p_T > 15\,\text{GeV}/c$ or two muons of $p_T > 10\,\text{GeV}/c$.

For the analysis of PbPb collisions, the "centrality" (overlap of the two colliding nuclei) is determined by the sum of the total energy deposited in both forward hadron calorimeters [15]. The results refer to the 30% most central collisions, to focus on the region of highest physics interest. After all the other analysis selections, 78% of the Z boson events fall in this centrality range.

The PYTHIA 8.212 [34] Monte Carlo (MC) event generator, with tune CUETP8M1 [35], is used

to simulate Z+jet signal events, with $p_T^Z > 30\,\text{GeV}/c$ and rapidity $|y^Z| < 2.5$. A sample with a Z boson without any kinematic selection was produced using a next-to-leading order (NLO) generator, MADGRAPH5_aMC@NLO [36]. In the PbPb case, a PYTHIA+HYDJET sample is created by embedding PYTHIA signal events in heavy ion events generated with HYDJET 1.9 [37] and tune HydroQJets. The generated events are propagated through the CMS apparatus using the GEANT4 [38] package. No unfolding is performed for the results presented. The recipe for applying a smearing of the jet p_T resolution is provided in Appendix A.

Electrons are identified as ECAL superclusters [39] matched in position and energy to tracks reconstructed in the tracker. They must have $p_T > 20\,\text{GeV}/c$, above the trigger threshold, and each supercluster must be within the acceptance of the tracker, $|\eta| < 2.5$. Electron candidates in the transition region between the barrel and endcap subdetectors ($1.44 < |\eta| < 1.57$) are excluded. In pp collisions, the electrons are selected via standard identification criteria [39]. A narrow transverse shape of showers in the ECAL and a low HCAL over ECAL energy ratio are required to reject misidentified electrons. Additional tracking information is used to distinguish electrons from charged hadrons [39]. For PbPb collisions, the identification criteria have been optimized to compensate for the higher background levels in the calorimeters. With these selections, the pp and PbPb electron reconstruction purities (efficiencies) are identical within 1% (10%).

Muons are selected by requiring segments in at least two muon detector planes and a good-quality fit when connecting them to tracker segments. This suppresses hadronic punch-through and muons from in-flight decays of hadrons. A minimum number of hits in the pixel and strip layers is required, and the reconstructed muon tracks must point to the primary vertex in the transverse and longitudinal directions [40]. The same selections are applied for both pp and PbPb data. In order to suppress the background continuum under the Z peak, mostly originating from uncorrelated simultaneous decays of heavy flavour mesons, the muons are required to have $p_T > 10\,\text{GeV}/c$. In addition, the muon tracks must fall in the acceptance of the muon detectors, $|\eta^\mu| < 2.4$.

Jet reconstruction uses the anti-k_T algorithm implemented in FASTJET [41], following the procedure of Ref. [16]. A small distance parameter, $R = 0.3$, minimizes the effects of fluctuations in the underlying event (UE), dominantly formed by soft processes in heavy ion collisions. The UE energy subtraction [42] is performed for PbPb as described in Refs. [15–17]. Closure tests, done on MC samples without medium induced jet energy loss, show no over subtraction of the UE in the PbPb sample. No subtraction is applied in the pp sample, where the UE contribution is negligible. The jet energy is calibrated applying η^{jet}- and p_T^{jet}-dependent correction factors derived with the PYTHIA signal sample [43]. Then, dijet and photon+jet balance techniques [44] are used to correct for the residual detector response differences between measured and simulated samples. In addition, a centrality-dependent correction obtained from simulation studies is applied to remove the residual effects from the UE in PbPb collisions. The UE from PbPb data and MC samples are compared using the p_T density [43, 45, 46], defined as the median of the ratio of the jet transverse momentum to the jet area, for all jets in the event. Given the coarse centrality range used in the analysis, the difference between the measured and simulated PbPb events has a negligible effect on jet reconstruction.

Except in Fig. 4, the resolutions of the measured jet energy and azimuthal angle in the pp samples are smeared to match those of the PbPb sample. The jet energy resolution can be quantified using the Gaussian standard deviation σ of the $p_T^{\text{reco}}/p_T^{\text{gen}}$ ratio, where p_T^{reco} is the UE-subtracted, detector-level jet p_T and p_T^{gen} is the generator-level jet p_T without any contributions from the UE in PbPb. It is determined using PYTHIA+HYDJET (for PbPb) and PYTHIA

(for pp) samples and parametrized as a function of p_T^{gen} using the expression $\sigma\left(p_T^{gen}\right) = C \oplus (S/\sqrt{p_T^{gen}}) \oplus (N/p_T^{gen})$, where $\oplus$ stands for the sum in quadrature and the parameters C, S, and N are determined from simulation studies. The same parametrization is used to determine the jet azimuthal angle resolution, quantified by the Gaussian standard deviation σ_ϕ of the $|\phi^{reco} - \phi^{gen}|$ difference.
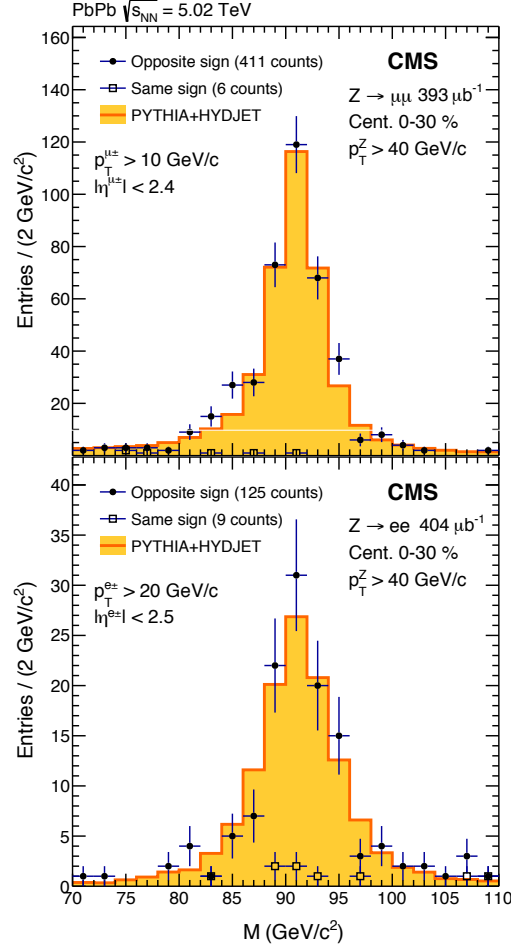

Figure 1: Invariant mass distributions of the selected dimuons (top) and dielectrons (bottom), compared to PYTHIA+HYDJET $Z(\ell\ell)$+jet events. The MC histogram is normalized to the number of events in the data.

The Z candidates are defined as opposite-charge electron or muon pairs, with a reconstructed invariant mass ($M^{\ell\ell}$) in the interval 70–110 GeV/c^2 and $p_T > 40$ GeV/c. The invariant mass distributions of all the dileptons used in the PbPb analysis are shown in Fig. 1. Each Z candidate is paired with all jets in the same event that pass the $p_T^{jet} > 30$ GeV/c and $|\eta^{jet}| < 1.6$ selection. Simulation studies show that the jet selection efficiency and the energy resolution are well understood for this kinematic range. Additional energy corrections are applied to the jet p_T, to account for residual performance degradations observed in simulation studies. Jets reconstructed within $\Delta R < 0.4$ from a lepton are rejected, to eliminate jet energy contamination by leptons from Z decays.

For the analysis of PbPb collisions, the background contribution from jets not produced in the same parton-parton interaction as the Z boson needs to be considered. This contribution arises from misidentified jets reconstructed from regional energy fluctuations in the high-multiplicity

heavy ion UE, or from additional initial hard interactions not related to the primary Z+jet production. The background jet contributions are estimated constructing a mixed-event jet background by correlating the Z boson from each candidate Z+jet event with jets reconstructed in subsets of 40 minimum bias events. All events must pass the offline event selection and have the same centrality and interaction vertex as the Z+jet candidate event. The resulting background jet spectrum is subtracted from the raw jet spectrum, eliminating coincidental Z+jet pairs and ensuring that the final Z+jet observables reflect the correlations of Z bosons and associated jets.

The systematic uncertainties related to Z boson reconstruction are sizable (negligible) in the dielectron (dimuon) channel. Comparing the measured and simulated dielectron invariant mass peaks shows that the average deviation between electron p_T^{reco} and p_T^{gen} is 0.5%. A systematic uncertainty is evaluated by shifting the electron p_T by $\pm 0.5\%$, resulting in changes of $\langle x_{jZ} \rangle$ and R_{jZ} for PbPb (pp) by 0.5% (0.3%) and 3% (0.8%), respectively. The simulated Z dielectrons reconstructed in central PbPb collisions have a p_T resolution of 5% for $p_T^Z > 40\,\text{GeV}/c$. In PbPb simulated events, p_T^Z is smeared by 5%, resulting in variations of $\langle x_{jZ} \rangle$ and R_{jZ} by 1.5% and 0.8%, respectively. When combining the two lepton results, a weighting is applied to the electron sample, to compensate for the different centrality dependencies of the Z boson reconstruction in the electron and muon channels. The difference between the corrected and uncorrected $\langle x_{jZ} \rangle$ and R_{jZ} values, 0.3% and 5.8% respectively, is taken as systematic uncertainty.

Simulation studies show that the jet energy scale, $\langle p_T^{reco} / p_T^{gen} \rangle$, can deviate from unity by up to 2%. Additional deviations can arise from differences between the fragmentation pattern of jets in measured and simulated events. To evaluate the corresponding systematic uncertainty, the jet energy scale is shifted for PbPb (pp) upward by 6% (2%) and downward by 4% (2%). The higher upward variation reflects the relatively high energy scale of quark jets, which contribute more to the Z+jet events than the gluon jets. The relative change in $\langle x_{jZ} \rangle$ and R_{jZ} for PbPb (pp) is 5.4% (2.4%) and 4.6% (2.4%), respectively. Finally, differences between the measured and simulated samples suggest that the jet energy resolution is up to 15% worse in data. The related systematic uncertainty is evaluated smearing p_T^{jet} by 15% in the PbPb MC. The pp data are smeared to simulate the poor resolution due to the UE fluctuations in PbPb data. The smearing is performed with the relative resolution, $\sigma_{rel} = \sqrt{\sigma_{PbPb}^2 - \sigma_{pp}^2}$, where σ_{PbPb} and σ_{pp} correspond to the parametrizations described above. A systematic uncertainty is assigned by varying the relative resolution by $\pm 15\%$. The PbPb (pp) relative change in $\langle x_{jZ} \rangle$ and R_{jZ} due to jet energy resolution is 2.5% and 3.7% (0.5% and 0.7%), respectively. The jet angular resolution correction implies an additional uncertainty on the pp sample, of 0.1% for $\langle x_{jZ} \rangle$ and 0.2% for R_{jZ}.

The total systematic uncertainties for PbPb (pp) amount to 6.2% (2.5%) and 8.9% (2.6%) for the $\langle x_{jZ} \rangle$ and R_{jZ} results, respectively, of which 5.7% and 8.0% are uncorrelated between the pp and PbPb results; the uncorrelated uncertainties do not reflect possible differences between p_T^{reco} and p_T^{gen}.

Figure 2-top shows the $\Delta\phi_{jZ}$ distribution of Z+jet pairs that pass all the selections; only Z+jet pairs with $p_T^Z > 60\,\text{GeV}/c$ were included to reduce the fraction of events where energy loss effects cause the jet partner to fall below the $p_T^{jet} > 30\,\text{GeV}/c$ threshold. There are 678 and 232 events that pass the $p_T^Z > 60\,\text{GeV}/c$ selection in pp and in the 30% most central PbPb collisions, respectively. To study if the angular distribution of jets with respect to the Z boson is affected by interactions of the parton with the medium, a Kolmogorov–Smirnov (KS) test was performed using pseudo data generated from identical underlying shapes. This test is useful to quantify shape differences since it is sensitive to adjacent bins fluctuating in the same direction but not
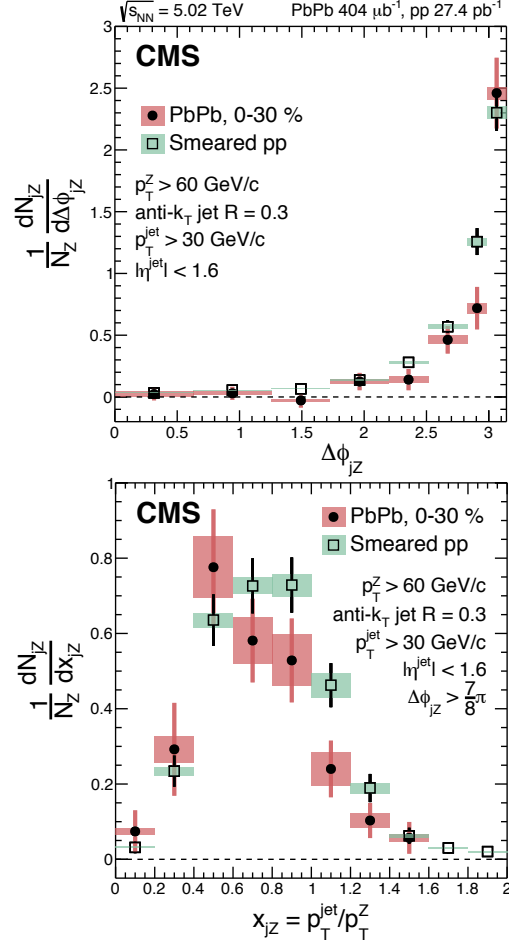
Figure 2: Distributions of the azimuthal angle difference $\Delta\phi_{jZ}$ between the Z boson and the jet (top), and of the transverse momentum ratio x_{jZ} between the jet and the Z boson with $\Delta\phi_{jZ} > 7\pi/8$ (bottom). The distributions are normalized by the number of Z events, N_Z. Vertical lines (bands) indicate statistical (systematic) uncertainties.

to the overall normalization. No significant difference is seen between the pp and PbPb $\Delta\phi_{jZ}$ distributions; the probability to obtain a KS value larger than that observed in the data, p-value, is greater than 0.40, even if systematic uncertainties are excluded.

For the x_{jZ} and R_{jZ} results, shown in Figs. 2 and 3, only events with $\Delta\phi_{jZ} > 7\pi/8$ are used, to select mostly back-to-back Z+jet pairs; it keeps 63% and 73% of the pp and PbPb events, respectively. Figure 2-bottom shows the x_{jZ} distributions for PbPb and pp collisions. Jet energy loss is expected to manifest itself both as a shift in the x_{jZ} distribution and an overall decrease in the number of Z+jet pairs as jets fall below the p_T^{jet} threshold. Therefore, the KS test was applied to the x_{jZ} distribution and a separate overall normalization χ^2 test was applied to the total number of Z+jet pairs per Z leading to p-values of $p_1 = 0.07$ and $p_2 = 0.01$, respectively. The systematic uncertainties and their correlations were included in these calculations. The combined p-value [47] is $p_1 p_2 (1 - \ln(p_1 p_2)) = 0.0064$ when including Z+jet pairs with $p_T^Z > 40\,\text{GeV}/c$, indicating that the two x_{jZ} distributions are significantly different.

The relative shift between the pp and PbPb x_{jZ} distributions is studied using their means, $\langle x_{jZ}\rangle$, shown in Fig. 3-top as a function of p_T^Z. The minimum p_T of the partner jet imposes a lower limit on the value of x_{jZ}. As p_T^Z increases relative to the p_T^{jet} cutoff, the kinematic phase space
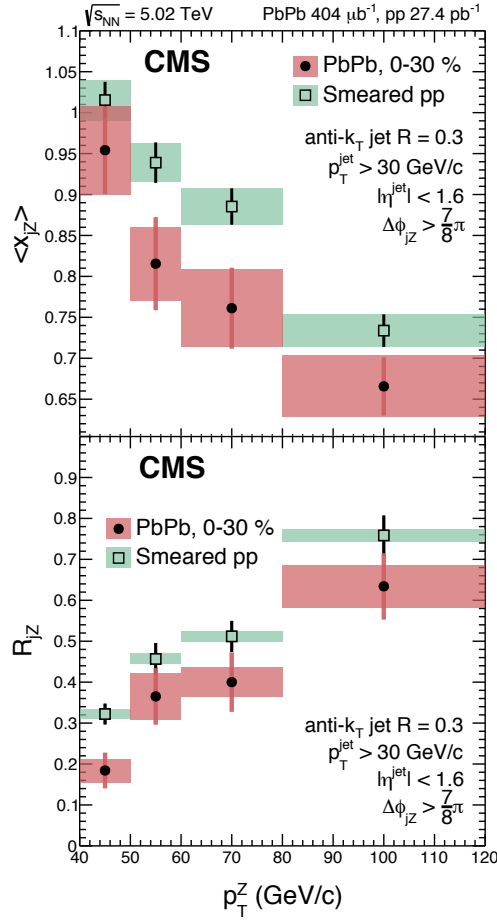
Figure 3: The mean value of the x_{jZ} distribution (top) and the average number of jet partners per Z boson R_{jZ} (bottom), as a function of p_T^Z. Vertical lines (bands) indicate statistical (systematic) uncertainties.

for lower x_{jZ} opens up, resulting in a shift towards lower x_{jZ} for higher p_T^Z. For all ranges, $\langle x_{jZ} \rangle$ is found to be lower in PbPb collisions than in pp collisions, as expected from energy loss models of partons traversing the medium. Also R_{jZ} is expected to increase as a function of p_T^Z, as the $p_T^{jet} > 30\,\text{GeV}/c$ threshold restricts the phase space of jets counted for a given p_T^Z selection. Figure 3-bottom shows the dependence of R_{jZ} on p_T^Z. The R_{jZ} values are found to be smaller in PbPb than in pp. As their difference is approximately constant as a function of p_T^Z, a relatively smaller fraction of jets is lost in PbPb collisions for larger initial (before traversing the medium) parton energies.

Figure 4 compares the x_{jZ} results to several theoretical calculations, using the same kinematic selections as the data. The PbPb results are compared to three models that incorporate the phenomenon of jet quenching: JEWEL [26], Hybrid [25], and GLV [27]. The x_{jZ} pp measured results are compared to several non-quenching scenarios: the pp references used as inputs to the PbPb models (different tunes of the PYTHIA LO event generator) and the MADGRAPH5_aMC@NLO generator [36], which includes matrix elements for Z plus 0, 1, and 2 jets at NLO, and Z+3 jets at LO. The pp calculations were smeared to reflect the detector resolution affecting the pp data. The JEWEL model is a dynamical, perturbative framework for jet quenching, which has been extended to simulate boson-jet events [26]. This PbPb x_{jZ} calculation is consistent with the data within the current precision, despite the poor agreement of its baseline with the pp measurement. The baseline PYTHIA8 tunes used by the Hybrid [25] and GLV [27] models, as
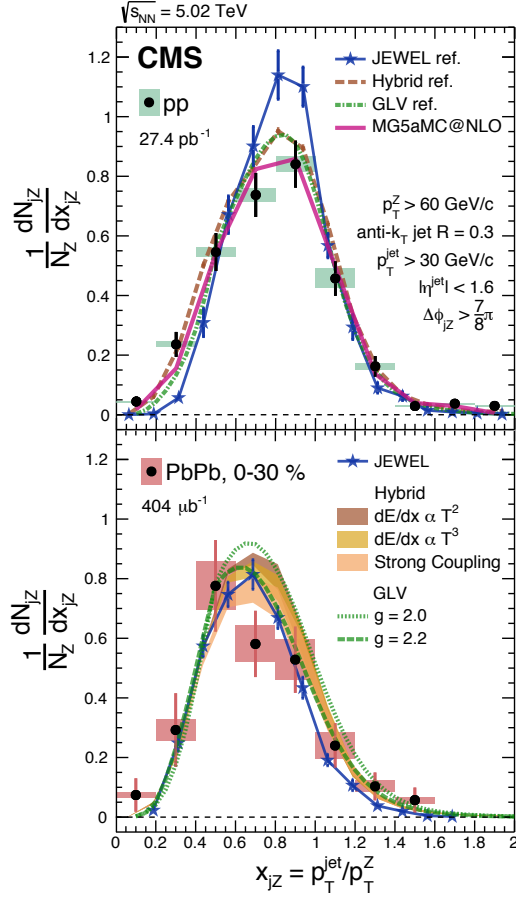
Figure 4: Comparison of the measured pp (top) and PbPb (bottom) x_{jZ} distributions with several theoretical models, smeared by the respective jet energy resolution: JEWEL [26], Hybrid [25], and GLV [27]. The JEWEL error bars represent statistical uncertainties while the widths of the Hybrid bands represent systematic variations. A MADGRAPH5_aMC@NLO calculation [36] is also shown.

well as MADGRAPH5_aMC@NLO, describe the pp data reasonably well. For PbPb collisions, the Hybrid model calculation labeled 'Strong Coupling' combines a perturbative description of the weakly coupled physics of jet production and evolution, with a gauge/gravity duality description of the strongly coupled dynamics of the medium, and of the soft exchanges between the jet and the medium. Two weak coupling benchmark calculations are also shown, where the energy loss has a quadratic temperature dependence (collisional energy loss), or a cubic dependence (radiative energy loss). Given the large experimental and theoretical uncertainties, all three scenarios describe the PbPb data reasonably well and cannot be distinguished. Nevertheless, the Strong Coupling curve appears closest to the data. The GLV model [27] generates the energy loss via out-of-cone radiation and collisional energy dissipation. Two curves are shown, for different coupling strengths between the jet and the medium, $g = 2.0$ and 2.2, reflecting previous analyses of jet quenching measurements at 2.76 TeV [48, 49]; the $g = 2.2$ curve seems favored by the data.

In summary, correlations of $p_T^Z > 40$ GeV/c Z bosons with $p_T^{jet} > 30$ GeV/c jets have been studied in pp and, for the first time, in PbPb collisions. The data were collected with the CMS experiment during the 2015 data taking period, at $\sqrt{s_{NN}} = 5.02$ TeV. No significant difference was found between the distributions of the azimuthal angle difference of the Z and the jet in pp and

PbPb collisions. The x_{jZ} distributions indicate that the PbPb values tend to be lower than those measured in pp collisions. Correspondingly, the average value of the transverse momentum ratio $\langle x_{jZ} \rangle$ is smaller in PbPb than in pp collisions, for all p_T^Z intervals. The average number of jet partners per Z, R_{jZ}, is lower in PbPb than in pp collisions, for all p_T^Z intervals, which suggests that in PbPb collisions a larger fraction of partons associated with Z bosons lose energy and fall below the 30 GeV/c p_T^{jet} threshold. These measurements provide new input for the determination of jet quenching parameters using a selection of partons with well-defined flavor and initial kinematics.

Acknowledgments

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A   Supplemental Material

A proper comparison between the data presented in this Letter and theory can be done when the theory calculations smear the generator-level jet p_T to the resolution observed in detector-level jets. The procedure is given below.

The p_T^{gen} is smeared into p_T^{reco} by multiplying the p_T^{gen} with a number sampled from a Gaussian distribution with mean 1 and variance $\sigma^2(p_T^{\text{gen}})$. The functional form is the same for pp and PbPb :

$$\sigma\left(p_T^{\text{gen}}\right) = \sqrt{C^2 + \frac{S^2}{p_T^{\text{gen}}} + \frac{N^2}{\left(p_T^{\text{gen}}\right)^2}}$$

The parameter S describes the p_T dependence of the jet energy resolution, C represents the high-p_T limit of the resolution, and N reflects the effect of UE fluctuations on the energy resolution. The parameters for $\sigma\left(p_T^{\text{gen}}\right)$ are determined using PYTHIA+HYDJET (for PbPb) and PYTHIA (for pp) samples. The parameter C is determined only from the PYTHIA sample, and has the same value for pp and PbPb, 0.061 ± 0.001. The PbPb (pp) numerical values for S and N are 1.24 ± 0.04 (0.95 ± 0.01) and 8.08 ± 0.15 (0.001 ± 0.001), respectively.

Arguably a more accurate comparison would include the smearing of generator-level jet ϕ as well. However, simulation studies showed that, compared to jet p_T smearing, jet ϕ smearing has negligible effect on the observables presented and therefore is omitted in this guide.

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