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LMECA2660 - Homework: simulating convection

Good Job!

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Friday, 18 march 2022

We consider the 1-D convection equation in the domain $x = [-L/2, L/2]$, with constant velocity c and periodic boundary conditions:

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0 \quad (1)$$

1 Stability of the temporal integration scheme

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As seen during the lectures, since we deal with a periodic function, we can decompose it into modes of different wavenumbers k and different amplitudes A_k . The mode k is given by:

$$u_k(x, t) = A_k(t) e^{ikx}$$

When we insert this expression of $u_k(x, t)$ into equation (1), we obtain the "eigenvalue" of the continuous mode k , $\lambda = \imath ck$:

$$\frac{dA_k(t)}{dt} + [\imath ck] A_k(t) = 0$$

When we discretize the partial derivative ∂_x of equation (1) with the mode $u_k(x, t)$, we find a modified version of λ :

$$\frac{dA_k(t)}{dt} + \left[\imath ck \frac{k^* h}{kh} \right] A_k(t) = 0$$

Hence, the eigenvalues of the schemes are given by:

$$\lambda_j \Delta t = \imath \frac{c \Delta t}{h} (k^* h) = \imath \cdot CFL \cdot (k^* h)$$

We want all the eigenvalues to be inside the marginal stability curve of RK4C. Since they are all purely imaginary, the condition is $|\lambda_j| < 2\sqrt{2}$ for $0 \leq k_j h \leq \pi$. We can then determine the CFL condition as: $CFL \max(k^* h) \leq 2\sqrt{2}$. Based on the results presented in table (1), I chose $CFL = 1.0$. ✓

Scheme	$k^* h$ as a function of kh (from [1])	$\max k^* h $	CFL number
E2	$\sin(kh)$	1.	2.828
E4	$[8 \sin(kh) - \sin(2kh)] / 6$	1.372	2.061
E6	$[45 \sin(kh) - 9 \sin(2kh) + \sin(3kh)] / 30$	1.586	1.783
I4	$[3 \sin(kh)] / [2 + \cos(kh)]$	1.732	1.632
I6	$[28 \sin(kh) + \sin(2kh)] / [18 + 12 \cos(kh)]$	1.989	1.421

Table 1. CFL number of various finite-difference schemes using RK4C. ✓

2 Validation of the periodic hypothesis

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The initial condition is a Gaussian with standard deviation σ . We also recall its Fourier transform.

$$u(x, 0) = U \exp\left(-\frac{x^2}{\sigma^2}\right)$$

$$\mathcal{F}(u(x, 0)) = U\sqrt{\pi}\sigma \exp\left(-\frac{k^2\sigma^2}{4}\right)$$

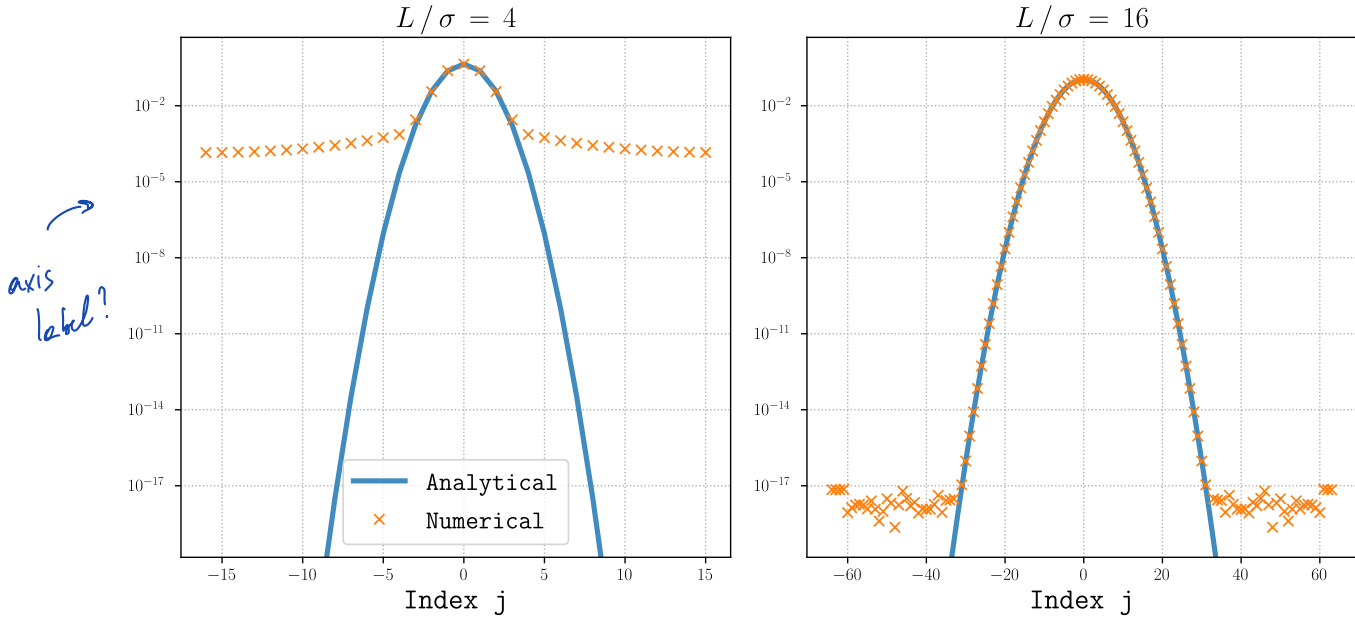


Figure 1. Amplitude of the Fourier transform and the FFT of the Gaussian initial function.

In both of the diagrams, the step size h is proportional to the width of the Gaussian: $h/\sigma = 1/8$. Hence, when $L/\sigma = 4$, we have $N = 32$ points, and when $L/\sigma = 16$, we have $N = 128$ points.

On one hand, in the left diagram of figure (1) ($L/\sigma = 4$), the amplitude of the coefficients of the FFT and the Fourier transform in an infinite domain only coincide at low wavenumbers when $|j| \leq 3$. This represents approximately one fourth of the spectrum. At higher wavenumbers, they do not coincide while also being not negligible: $\approx 10^{-4}$. \rightarrow Periodic assumption is not valid.

On the other hand, in the right diagram of figure (1) ($L/\sigma = 16$), the coefficients of the FFT have the right amplitude until they reach machine precision $\epsilon \approx 10^{-16}$. They coincide for $|j| \leq 31$: approximately half of the spectrum. \rightarrow amplitude of the modes that are not correctly captured in the domain is negligible

In view of this, it is clear that in our further analysis we will take a large domain compared to the width of the Gaussian: $L/\sigma = 16$. \rightarrow periodic assumption is valid!

3 Results obtained with a uniform grid

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In this section, we will consider a grid with a constant Δx from $-L/2$ to $L/2$. In the following simulations, if U is not explicitly mentioned, it is supposed to be set to 1.

3.1 Comparison of analytical and numerical solutions

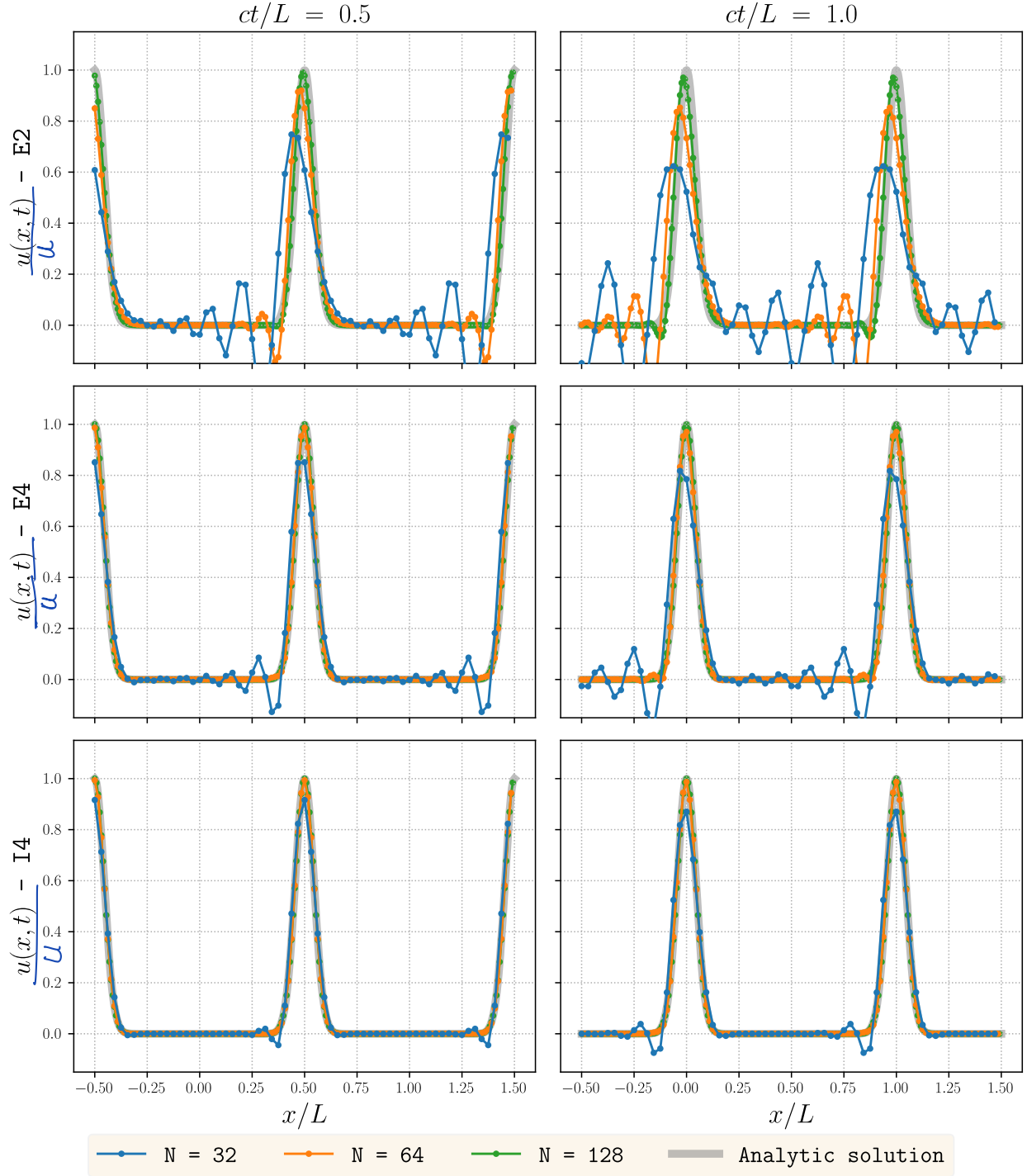


Figure 2. Solutions of the transport equation. Each row corresponds to one of the schemes E2, E4, I4. Each column shows the solution at a specific time. Then the colors correspond to the quality of the discretization.

As expected, in figure (2), the scheme E2 gives the worst results, followed by E4 and then I4 provides the best approximation. This follows from the leading term of the truncation error T of each of these schemes [1, p. 5]:

$$T_{E2} = -\frac{1}{6}(kh)^2 \quad T_{E4} = -\frac{1}{30}(kh)^4 \quad T_{I4} = -\frac{1}{180}(kh)^4$$

For all the schemes, the large step size $h = \sigma/2$ ($N = 32$) gives bad results, particularly at the left of the peak of the Gaussian. In the case of the E2 scheme, still with $N = 32$, the signal appears "destroyed" already at $ct/L = 0.5$.

More generally, the larger the step size (the lower the N), the bigger is the dissipation (decrease of amplitude) and the dispersion (difference of velocity). Practically, in order to obtain results with sufficient quality, we can use the E2 scheme with at least $N = 128$, and the E4 and I4 schemes with at least $N = 64$.

3.2 Evolution of the diagnostics with time

There is almost no dissipation in this case, but the modes are not transported at the right velocity. The decrease of the amplitude is only due to dispersion error!

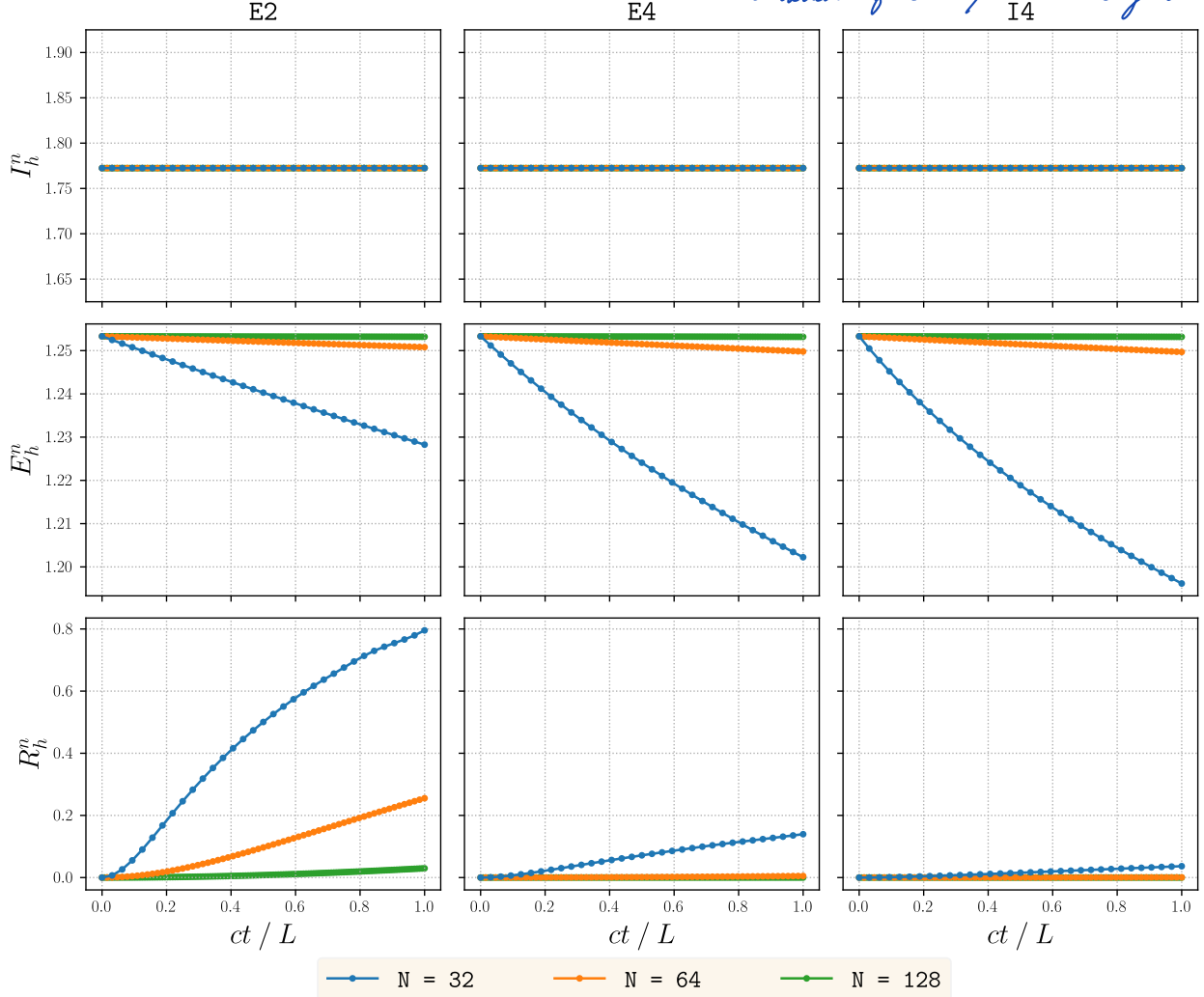


Figure 3. Presentation of the three diagnostic I_h , E_h and R_h as functions of a dimensionless time ct/L . The three schemes are arranged in column and the three diagnostic in rows.

We recall the diagnostics: I for the integral, E for the energy, and R for the error. In the following, u_i^n indicates the discrete approximation of $u(x, t)$ at position $x_i = ih$ and time $t = n\Delta t$:

$$I_h^n = \frac{h}{\sigma U} \sum_{i=0}^{N-1} u_i^n \quad E_h^n = \frac{h}{\sigma U^2} \sum_{i=0}^{N-1} (u_i^n)^2 \quad R_h^n = \frac{h}{\sigma U^2} \sum_{i=0}^{N-1} [u_i^n - u(x_i, t^n)]^2$$

The first global diagnostic is the discrete version of the integral of the function $u(x, t)$ from $x = -b = -L/2$ to $x = b = L/2$. We see in figure (3) that it is constant. We can also show that it is constant and find its value in the continuous case:

$$\begin{aligned} \frac{d}{dt} \int_{-b}^b u \, dx &= \int_{-b}^b \frac{\partial u}{\partial t} \, dx = -c \int_{-b}^b \frac{\partial u}{\partial x} \, dx = -c \left[u \left(\frac{L}{2} \right) - u \left(-\frac{L}{2} \right) \right] = 0 \\ \frac{1}{\sigma U} \int_{-b}^b u \, dx &= 2 \int_0^{\frac{L}{2}} \exp(-x^2) \, dx = \sqrt{\pi} \operatorname{erf}(8) \approx 1.772 \end{aligned}$$

The second diagnostic concerns the energy contained in u . For all schemes, the worse the quality of the grid (the smaller the N), the faster is the decrease of energy. Even if it may appear counter-intuitive, the order of the scheme has the opposite effect. For all N , the lower the truncation error (the better the scheme), the faster is the decrease of energy.

Finally, the third diagnostic measures the global squared error. This time, the results are not surprising. The error (1) increases over time, (2) decreases as the truncation error of the scheme decreases, and (3) decreases as we refine the grid (N increasing).

3.3 Order of convergence

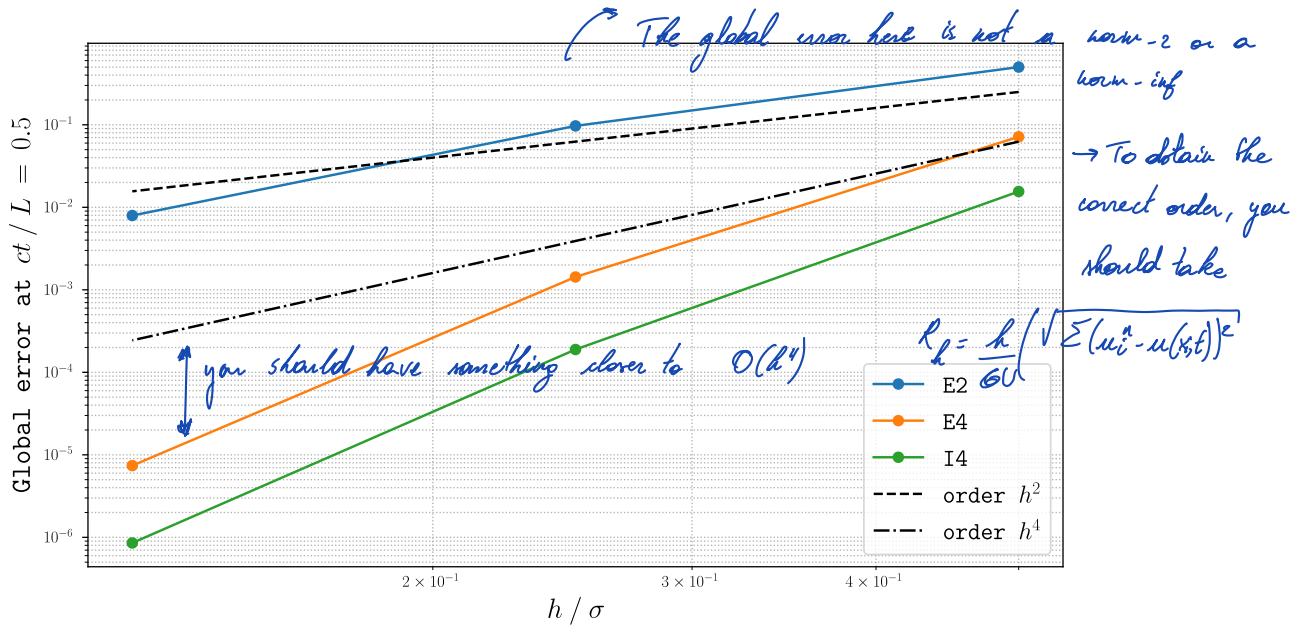


Figure 4. Error R_h at time $ct/L = 0.5$ as a function of h/σ for the different schemes.

The global error shown in figure (4) is decreasing slightly faster than $O(h^2)$ using the E2 scheme, and slightly faster than $O(h^4)$ using the E4 and I4 schemes. As mentioned a few times before, the error made using I4 is smaller than the error made with E4, since their truncation errors differ by a factor of 6. What is your exact convergence order for each scheme?

4 Results obtained with a non-uniform grid

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We now consider a mapping, constant in time, from the "numerical space" $-\frac{L}{2} \leq \xi \leq \frac{L}{2}$ with a uniform grid to the "physical space" $-\frac{L}{2} \leq x \leq \frac{L}{2}$ with a non-uniform grid:

$$x = g(\xi) = \xi - a \frac{L}{2\pi} \sin\left(2\pi \frac{\xi}{L}\right) \quad \text{and} \quad \frac{dx}{d\xi} = g'(\xi) = 1 - a \cos\left(2\pi \frac{\xi}{L}\right)$$

We can use the chain rule to modify the transport equation (1).

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial \xi} \frac{d\xi}{dx} = 0$$

$$\frac{\partial}{\partial t} \left(u \frac{dx}{d\xi} \right) + \frac{\partial(cu)}{\partial \xi} = 0$$

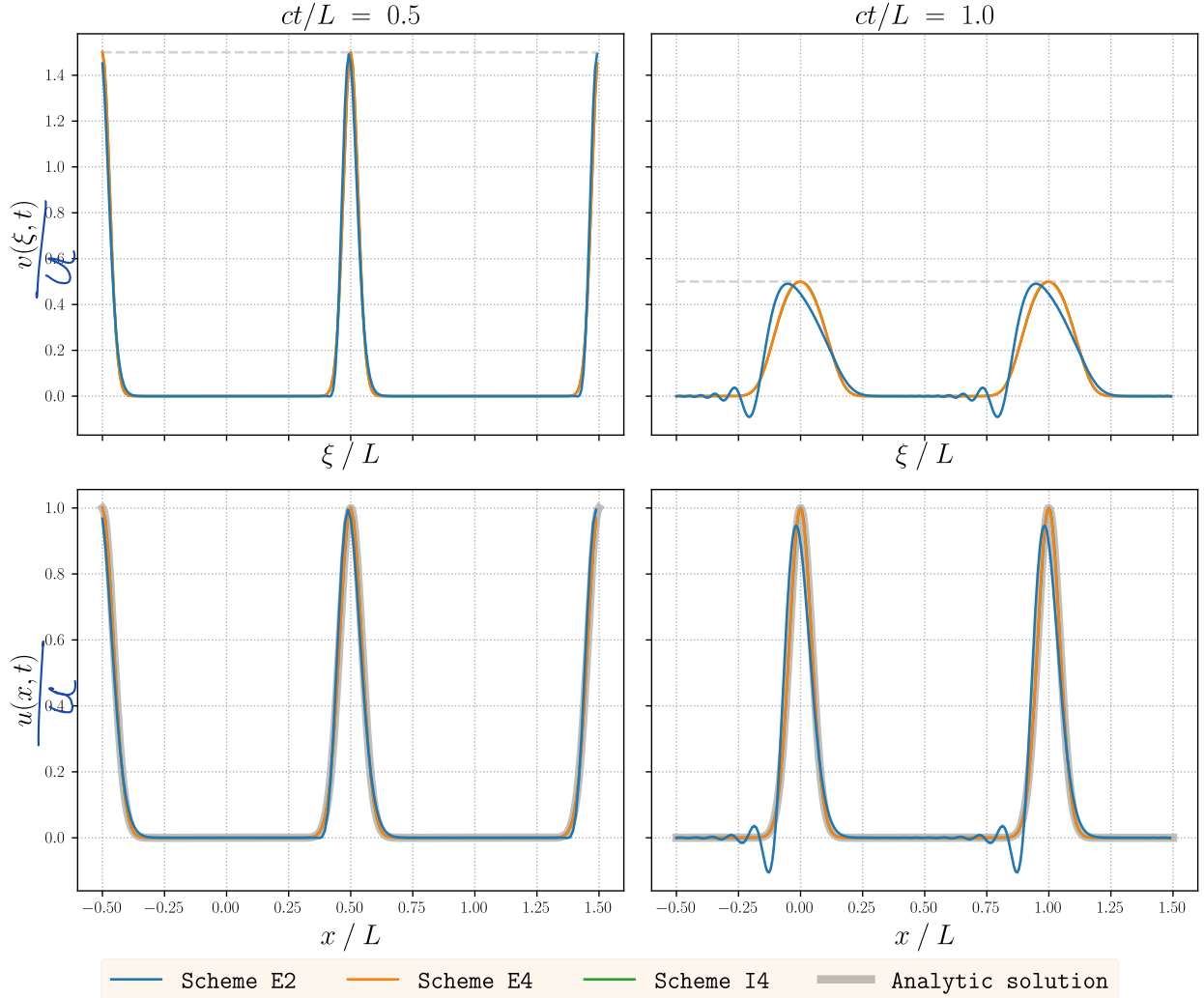
since both $x = g(\xi)$ and c do not depend on t

✓

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$$\frac{\partial v}{\partial t} + \frac{\partial(bv)}{\partial \xi} = 0$$

where $b(\xi) = \frac{c}{g'(\xi)}$ and $v(\xi, t) = g'(\xi) u(x(\xi), t)$



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Figure 5. The first row shows the solution $v(\xi, t)$ in the numerical domain while the second row shows the solution $u(x, t)$ in the physical domain. Each column shows the solution at a specific time.

When the peak of the Gaussian is located where the density of points ξ_i is high, it makes $v(\xi, t)$ more distributed and gives it a lower peak. This is what we observe in the right part of figure (5) at $ct/L = 1$ when the Gaussian is located at $x = 0 \bmod L$. The peak is enlarged since there are more points ξ under it, and it is flattened since $g'(\xi) < 1$ around $\xi = 0$.

The opposite happens at $ct/L = 0.5$, when the Gaussian is located at $L/2 \bmod L$.

Once again, E2 gives worse results than E4 and I4, which are almost indistinguishable.

→ Advantage of using a stretched grid?
Drawbacks?

5 Wave packet

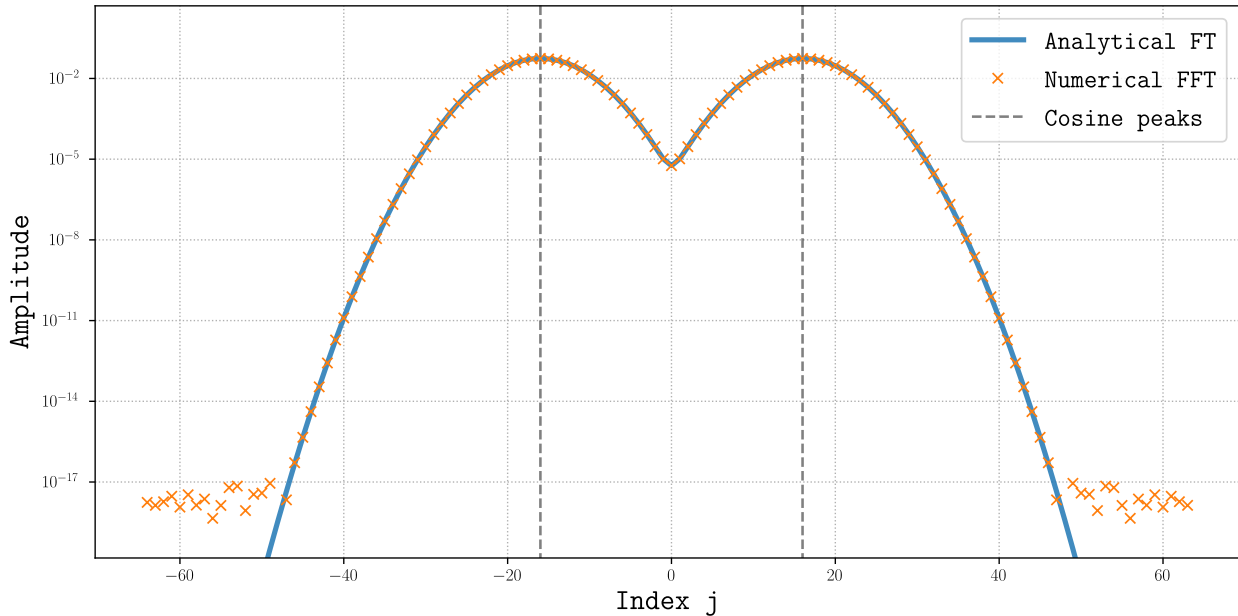
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Now, we consider a different initial condition, with $k_p = \frac{2\pi}{L} \cdot 16$:

$$\tilde{u}(x, 0) = U \cos(k_p x) \exp\left(-\frac{x^2}{\sigma^2}\right) = \cos(k_p x) u(x, 0)$$

Using the properties of the Fourier Transform, we know that a product in the time domain is a convolution in the frequency domain, up to a scaling factor $\frac{1}{2\pi}$:

$$\begin{aligned} \mathcal{F}(\tilde{u}(x, 0)) &= \frac{1}{2\pi} \mathcal{F}(\cos(k_p x)) * \mathcal{F}(u(x, 0)) \\ &= \frac{1}{2} \left[\delta(k - k_p) + \delta(k + k_p) \right] * \left[U \sqrt{\pi} \sigma \exp\left(-\frac{k^2 \sigma^2}{4}\right) \right] \\ &= \frac{1}{2} U \sqrt{\pi} \sigma \left[\exp\left(-\frac{(k - k_p)^2 \sigma^2}{4}\right) + \exp\left(-\frac{(k + k_p)^2 \sigma^2}{4}\right) \right] \end{aligned}$$



0,5/0,5

Figure 6. Amplitude of the Fourier transform and the FFT of the "wave packet" initial function.

As a result of the convolution, the Gaussian in the Fourier space is divided in two, as shown in figure (6). Its maximum located at the wavenumber k_p of the cosine function. The signal is shifted towards higher wavenumbers compared to the first initial condition $u(x, 0)$ we considered.

→ Is the periodic approximation still valid?

Before showing the results of the simulation, we compute the group velocities c_g^* for each scheme as a function of the wavelength. As reminded in the homework statement, the group velocity satisfies the relation $c_g^* = c \frac{dk^*}{dk}$ and is equal to zero at $kh = k_m h$.

Therefore, the modes with wavelengths exceeding k_m have a negative group velocity. In order to keep a numerical method of quality, we need to ensure that h is small enough such that the wavenumbers stay below k_m . ✓

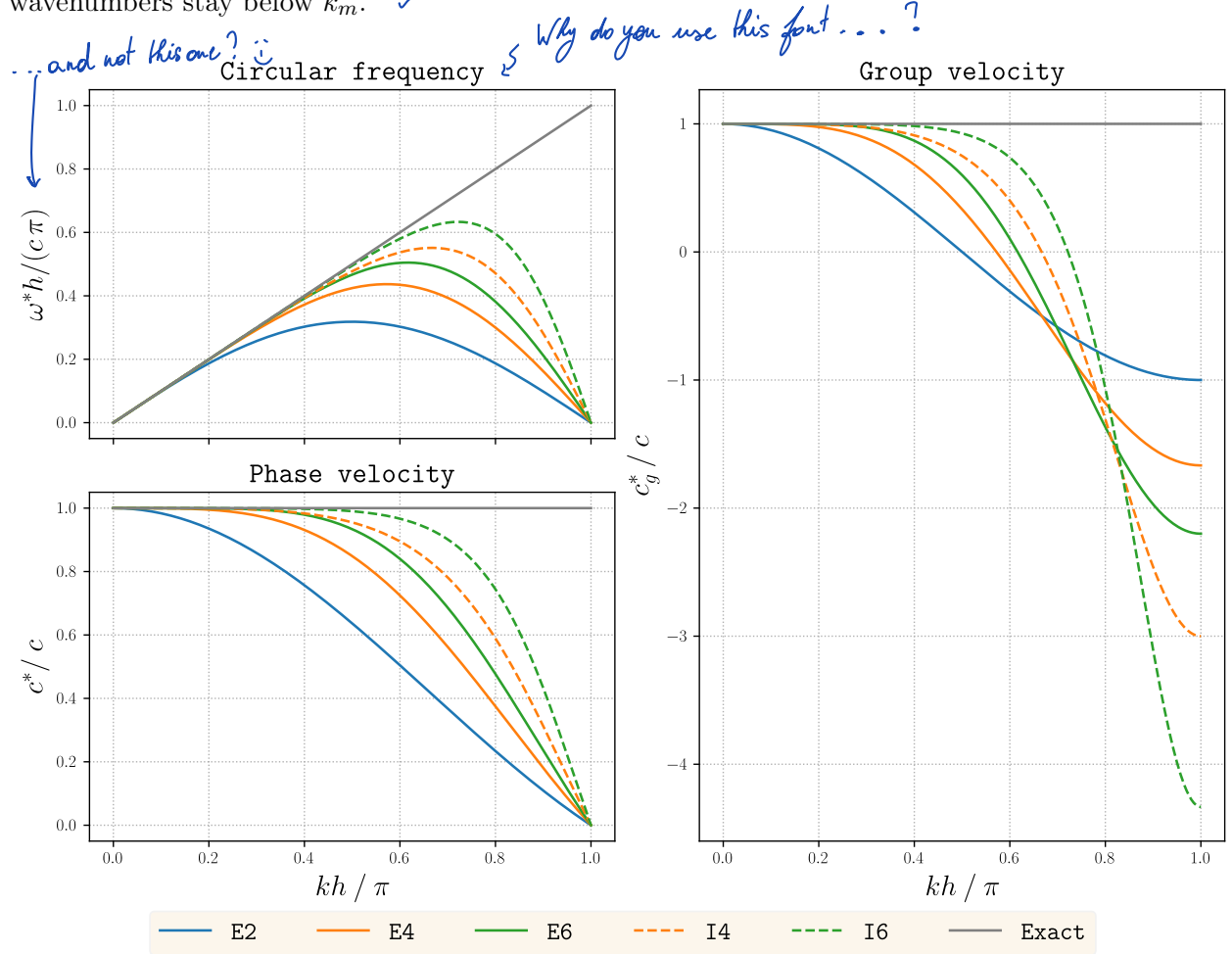


Figure 7. Circular frequency, phase and group velocities as functions of the wavenumber.

Scheme	c_g^* as a function of kh	$k_m h / \pi$	minimum value
E2	$\cos(kh)$	0.5	-1
E4	$[4 \cos(kh) - \cos(2kh)] / 3$	0.572	-1.666
E6	$[15 \cos(kh) - 6 \cos(2kh) + \cos(3kh)] / 10$	0.616	-2.2
I4	$3 [1 + 2 \cos(kh)] / [2 + \cos(kh)]^2$	0.666	-3
I6	$\frac{[28 \sin(kh) + \sin(2kh)] \sin(kh) + [3 + 2 \cos(kh)][14 \cos(kh) + \cos(2kh)]}{3 [3 + 2 \cos(kh)]^2}$	0.721	-4.333

Table 2. Group velocity of the different schemes.

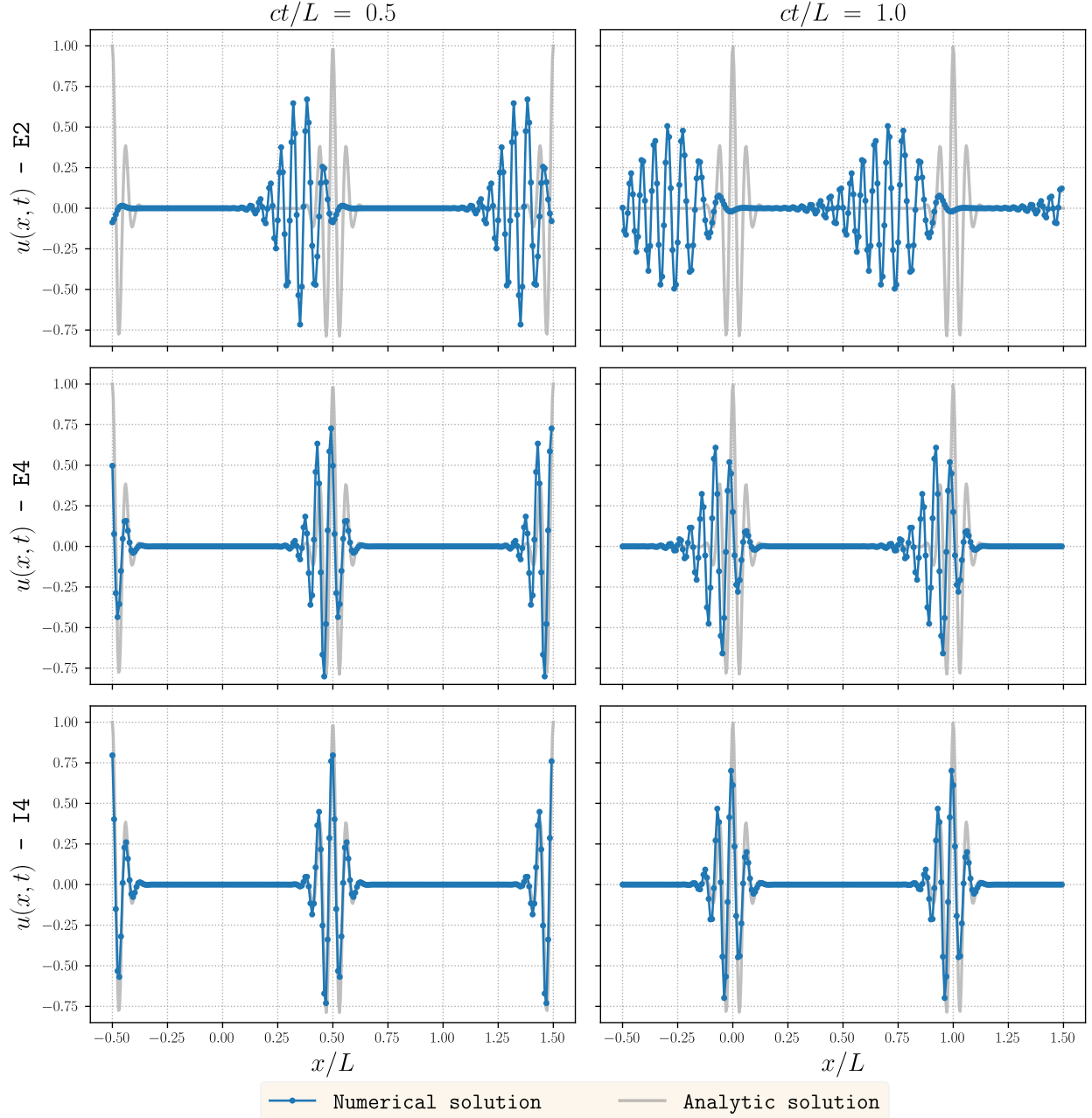


Figure 8. Solution of the transport equation with a wave packet as initial condition.

In the case of the E2 scheme, the principal mode with $j = 16$ has a group velocity of only 71% of the velocity c since $\cos(2\pi \frac{16}{128}) = 1/\sqrt{2} \approx 0.71$. This is the reason why it has only traveled until $x/L \approx 0.71$ after $ct/L = 1$ in the upper right diagram of figure (8) while the true wave packet has traveled until $x = L$.

This effect is of course smaller for the other schemes, since they behave much better at this wavelength: $v_g^* = 94\%$ of c for E4 and $v_g^* = 98.8\%$ of c for I4.

6 Negative group velocity with the E6 scheme

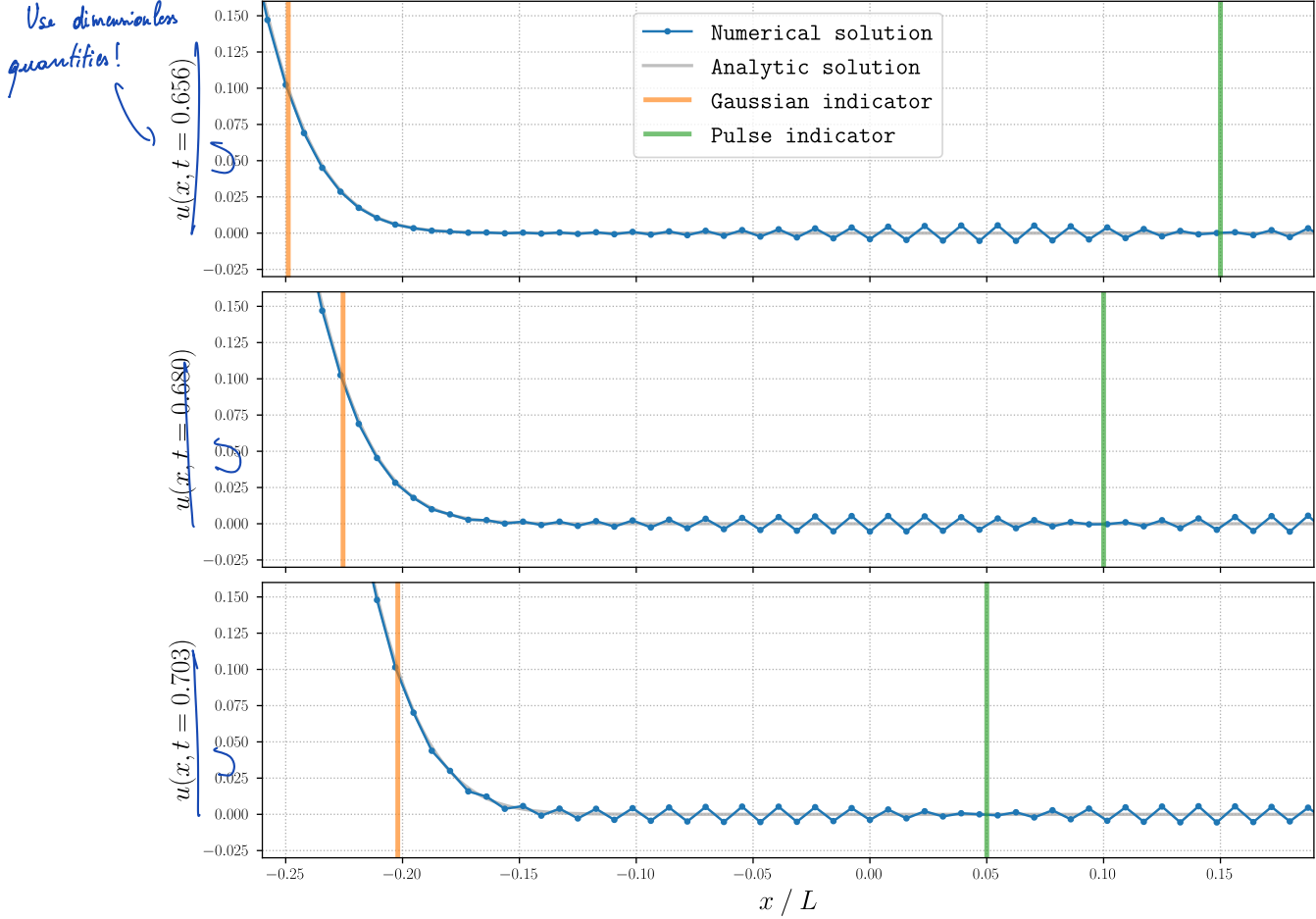


Figure 9. Solution of the transport equation of a Gaussian at three equally spaced times. The parameters of this simulation are $L = 1$ m, $c = 1$ m/s, $U = 1$, $N = 128$, $\Delta t = 7.812 \times 10^{-3}$ s

In figure (9), we observe that the solution contains a noise at the flip-flop mode. Since its wavenumber is maximal, its group velocity is minimal: $v_g^* = -2.2$ is given in table (2) and in figure (7) for the E6 scheme.

On one hand, the noise (indicated by the green line) travels a distance $\Delta x/L \approx -0.1 = 0.05 - 0.15$. On the other hand, the Gaussian has only travelled a distance $\Delta x/L \approx 0.047 = -0.202 - (-0.249)$. The ratio between both distances is ≈ 2.13 which is close to the 2.2 expected.

Without full certainty, we can safely assume that this phenomenon appears with the sixth-order schemes (E6 and I6) because we only use a fourth-order temporal scheme. In this case, the error is in $\mathcal{O}(\Delta t^4) = \mathcal{O}(\Delta x^4)$ since $\Delta x = c\Delta t/CFL$. In light of this, sixth-order schemes in space should not be used with RK4C.

References

- [1] G. Winckelmans. *LMECA2660 - Finite differences convection diffusion [Course notes]*. 2018.
- [2] G. Winckelmans. *LMECA2660 - Time integration schemes [Course notes]*. 2022.