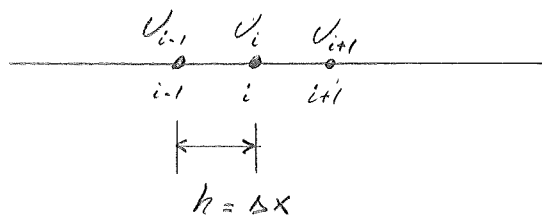


# Diffusion Equation

$$\frac{\partial v}{\partial t} = \epsilon \frac{\partial^2 v}{\partial x^2} \quad \text{PDE for } v(x, t)$$

↓ Finite differences (E2)

$$\frac{d}{dt} v_i = \epsilon \frac{(v_{i-1} - 2v_i + v_{i+1}))}{h^2}$$



$$\frac{\partial^2 v}{\partial x^2} \Big|_i + \frac{(\Delta x)^2}{12} \frac{\partial^4 v}{\partial x^4} \Big|_i + \dots$$

= system of ODEs for the  $v_i(t)$

Periodic problem  $\Leftrightarrow$  modal analysis (Fourier modes)

Continuous problem:  $v(x, t) = \sum_k A_k(t) e^{ikx}$

$$\left\{ \begin{array}{l} \lambda_k = -\epsilon k^2 \text{ real and negative} \\ \frac{dA_k}{dt} = -\epsilon k^2 A_k \Rightarrow A_k(t) = A_k(0) e^{-\epsilon k^2 t} \end{array} \right.$$

Discrete problem:  $v_i(t) = \sum_k A_k(t) e^{ikx_i}$  (still continuous in time)

$$\begin{aligned} \frac{d}{dt} v_i &= \sum_k e^{ikx_i} \frac{dA_k}{dt} = \sum_k \frac{\epsilon}{h^2} (e^{ikx_{i+1}} - 2e^{ikx_i} + e^{ikx_{i-1}}) A_k \\ &= \sum_k \frac{\epsilon}{h^2} e^{ikx_i} (e^{ikh} - 2 + e^{-ikh}) A_k \end{aligned}$$

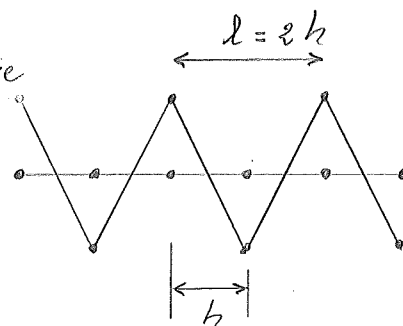
"von Neumann analysis"

$$\frac{d}{dt} A_k = -4 \frac{\epsilon}{h^2} \sin^2\left(\frac{kh}{2}\right) A_k \Rightarrow A_k(t) = A_k(0) e^{\lambda_k t}$$

$\lambda_k$  real and negative

•  $|\lambda_k|_{\max}$  if  $\frac{kh}{2} = \frac{\pi}{2}$

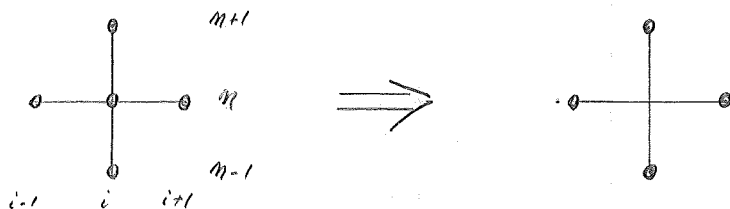
$kh = \pi \equiv$  "flip-flop mode"





Dugout - Frankel (1953)

Replace  $v_i^n$  by  $\frac{1}{2}(v_i^{n+1} + v_i^{n-1})$



$$v_i^{n+1} = v_i^{n-1} + 2\tau \left( v_{i+1}^n - (v_i^{n+1} + v_i^{n-1}) + v_{i-1}^n \right)$$

$$(1 + 2^m) \psi_i^{m+1} = (1 - 2^m) \psi_i^{m-1} + 2^m (\psi_{i+1}^m + \psi_{i-1}^m)$$

↳ Unconditionally stable for all  $\tau$  (exercise). One obtains:

- $r \leq 1/2$ :  $\beta_1$  and  $\beta_2$  are real
- $r > 1/2$ :  $\beta_1$  and  $\beta_2$  are complex for modes with  $\text{Re}(kh) > \frac{1}{2r} \equiv \text{bad.}$

↳ Problem! In fact, what we do is:

$$\frac{(V_i^{m+1} - V_i^{m-1})}{2 \Delta t} = \frac{e}{h^2} (V_{i+1}^m - 2V_i^m + V_{i-1}^m) + \frac{e}{h^2} (2V_i^m - (V_i^{m+1} + V_i^{m-1}))$$

$$\left( \frac{\partial \psi}{\partial t} + \frac{(\Delta t)^2}{6} \frac{\partial^3 \psi}{\partial t^3} + \dots \right) = \epsilon \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{h^2}{12} \frac{\partial^4 \psi}{\partial x^4} + \dots \right) - \epsilon \left( \frac{\Delta t}{h} \right)^2 \left( \frac{\partial^2 \psi}{\partial t^2} + \frac{(\Delta t)^2}{12} \frac{\partial^4 \psi}{\partial t^4} + \dots \right)$$

$$\frac{\partial \psi}{\partial t} = \epsilon \frac{\partial^2 \psi}{\partial x^2} - \epsilon \left( \frac{\Delta t}{h} \right)^2 \frac{\partial^2 \psi}{\partial t^2} + \frac{\epsilon h^2}{12} \frac{\partial^4 \psi}{\partial x^4} - \frac{(\Delta t)^2}{6} \frac{\partial^3 \psi}{\partial t^3} + \dots$$

"Modified equation" = equation approximated by the scheme

1. The truncation error of  $\Delta F$  is  $O((\Delta t)^2, h^2, (\frac{\Delta t}{h})^2)$

$\hookrightarrow \Delta F$  is consistent if  $\frac{(\Delta t)}{h} \rightarrow 0$  when  $\Delta t \rightarrow 0$  and  $h \rightarrow 0$

2. If  $\frac{\Delta t}{h} = \frac{1}{\text{constant}}$  when  $\Delta t \rightarrow 0$  and  $h \rightarrow 0$ ,  $\Delta F$  is inconsistent as it approximates:

$$\frac{\partial^2 U}{\partial t^2} = \epsilon \left( \frac{\partial^2 U}{\partial x^2} - \frac{1}{c_0^2} \frac{\partial^2 U}{\partial t^2} \right) \quad \text{wave-like behavior!}$$

$\equiv$  not physical

3. Also:  $\epsilon \left( \frac{\Delta t}{h} \right)^2 = \left( \frac{\epsilon \Delta t}{h^2} \right) \Delta t = \tau \Delta t$ . Hence DF modified equation

also is:  $\frac{\partial u}{\partial t} = \epsilon \frac{\partial^2 u}{\partial x^2} - \tau \Delta t \frac{\partial^2 u}{\partial t^2} + \frac{\epsilon h^2}{12} \frac{\partial^4 u}{\partial x^4} - \frac{(\Delta t)^2}{6} \frac{\partial^3 u}{\partial t^3} + \dots$

DF scheme is consistent when  $\Delta t \rightarrow 0$  and  $h \rightarrow 0$  with  $\tau$  constant.

$\Delta t = \tau \frac{h^2}{\epsilon}$

$\frac{\partial u}{\partial t} = \epsilon \frac{\partial^2 u}{\partial x^2} - \frac{\tau^2 h^2}{\epsilon} \frac{\partial^2 u}{\partial t^2} + \frac{\epsilon h^2}{12} \frac{\partial^4 u}{\partial x^4} - \frac{\tau^2 h^4}{6 \epsilon^2} \frac{\partial^3 u}{\partial t^3} + \dots$

$\frac{\partial^2 u}{\partial t^2} = \epsilon \frac{\partial^2}{\partial x^2} \left( \frac{\partial u}{\partial t} \right) - \dots = \epsilon^2 \frac{\partial^4 u}{\partial x^4} - \dots$

$\frac{\partial u}{\partial t} = \epsilon \frac{\partial^2 u}{\partial x^2} - \underbrace{\epsilon h^2 \left( \tau^2 - \frac{1}{12} \right)}_m \frac{\partial^4 u}{\partial x^4} + \dots$

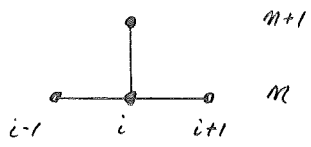
changes sign when  $\tau = \frac{1}{\sqrt{12}} \approx 0.2887$

$\tau = \frac{1}{\sqrt{12}} \Rightarrow$  hyper-diffusion term absent  $\equiv$  optimum?

$\tau > \frac{1}{\sqrt{12}} \Rightarrow$  hyper-diffusion of proper sign

$\tau < \frac{1}{\sqrt{12}} \Rightarrow$  hyper-diffusion of improper sign

Euler explicit



$\frac{(u_i^{n+1} - u_i^n)}{\Delta t} = \frac{\epsilon}{h^2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$

$O(\Delta t, h^2)$

$u_i^{n+1} = u_i^n + \tau (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$

simplified notation

$u_i^{n+1} = \tau u_{i+1}^n + (1-2\tau) u_i^n + \tau u_{i-1}^n$

$u_i^{n+1} = (I + \tau \delta^2) u_i^n$

Stability:

o.k. if all  $\lambda_k \Delta t$  are within the stability region of the scheme:

$-2 \leq -\frac{4\epsilon \Delta t}{h^2} \sin^2\left(\frac{kh}{2}\right) \leq 0$  for all  $kh$

$\tau \leq \frac{1}{2}$

• Other way  $U_i^m = \sum_k \beta_k^m e^{ikx_i} A_k(0)$

$$\beta_k^{m+1} e^{ikx_i} = r \beta_k^m e^{ik(x_i+h)} + (1-2r) \beta_k^m e^{ikx_i} + r \beta_k^m e^{ik(x_i-h)}$$

$$\beta_k = r (e^{ikh} + e^{-ikh}) + (1-2r)$$

$$= 2r \cos(kh) + (1-2r)$$

$$= 1 - 2r (1 - \cos(kh))$$

$$\beta_k = 1 - 4r \sin^2\left(\frac{kh}{2}\right) \quad \text{real}$$

Stable if  $|\beta_k| \leq 1$  for all  $kh$ .

$$\Leftrightarrow -1 \leq \beta_k \leq 1 \quad (\text{as } \beta_k \text{ is real})$$

$$-1 \leq 1 - 4r \sin^2\left(\frac{kh}{2}\right) \leq 1$$

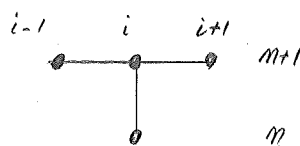
↓  $\hookrightarrow$  always true

$$2 \sin^2\left(\frac{kh}{2}\right) \leq \frac{1}{2}$$

$$\boxed{r \leq \frac{1}{2}}$$

O.K. Same result

Euler implicit



$$\frac{U_i^{m+1} - U_i^m}{\Delta t} = \frac{c}{h^2} (U_{i+1}^{m+1} - 2U_i^{m+1} + U_{i-1}^{m+1})$$

$O(\Delta t, h^2)$

$$U_i^{m+1} - r (U_{i+1}^{m+1} - 2U_i^{m+1} + U_{i-1}^{m+1}) = U_i^m$$

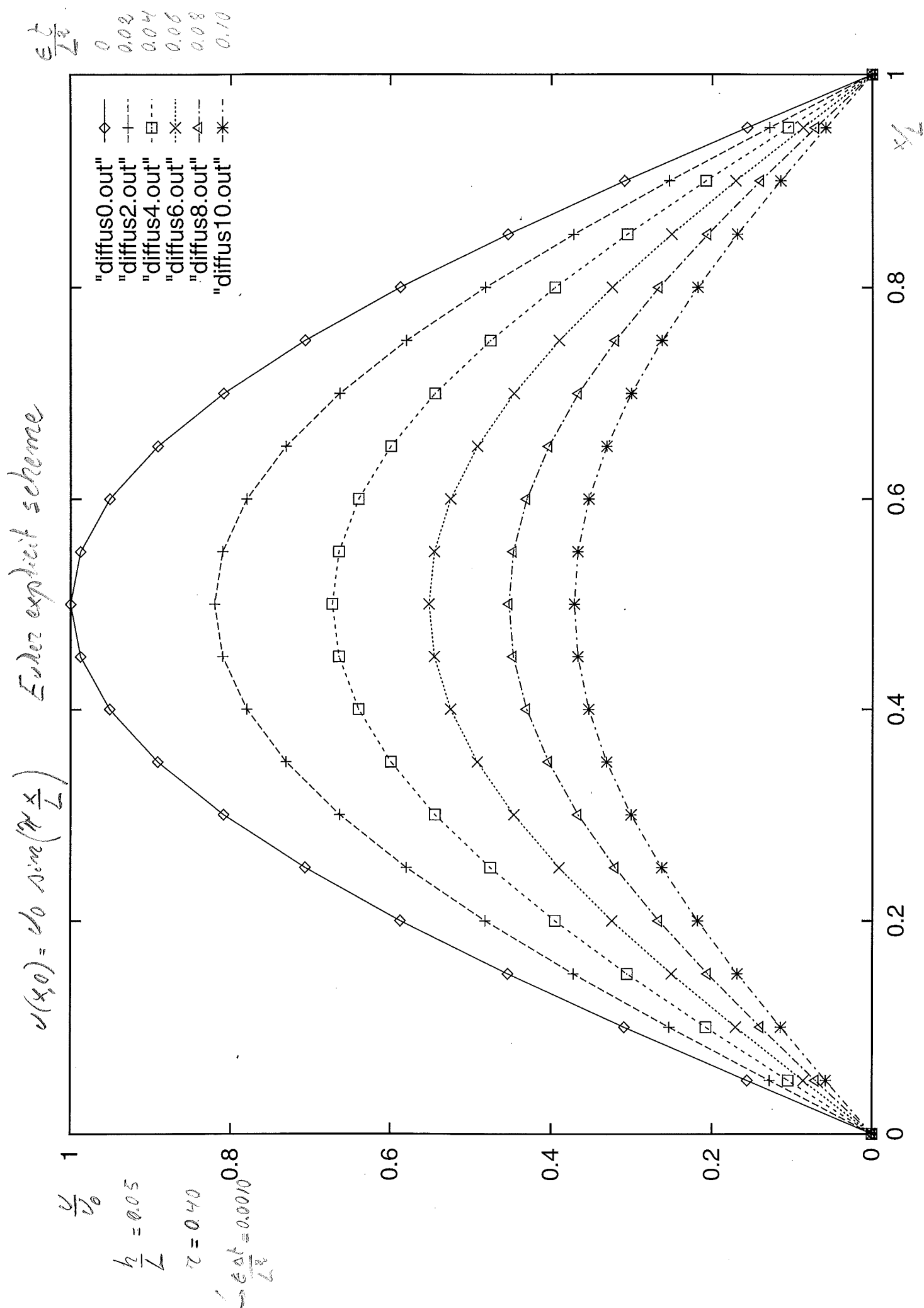
unconditionally stable

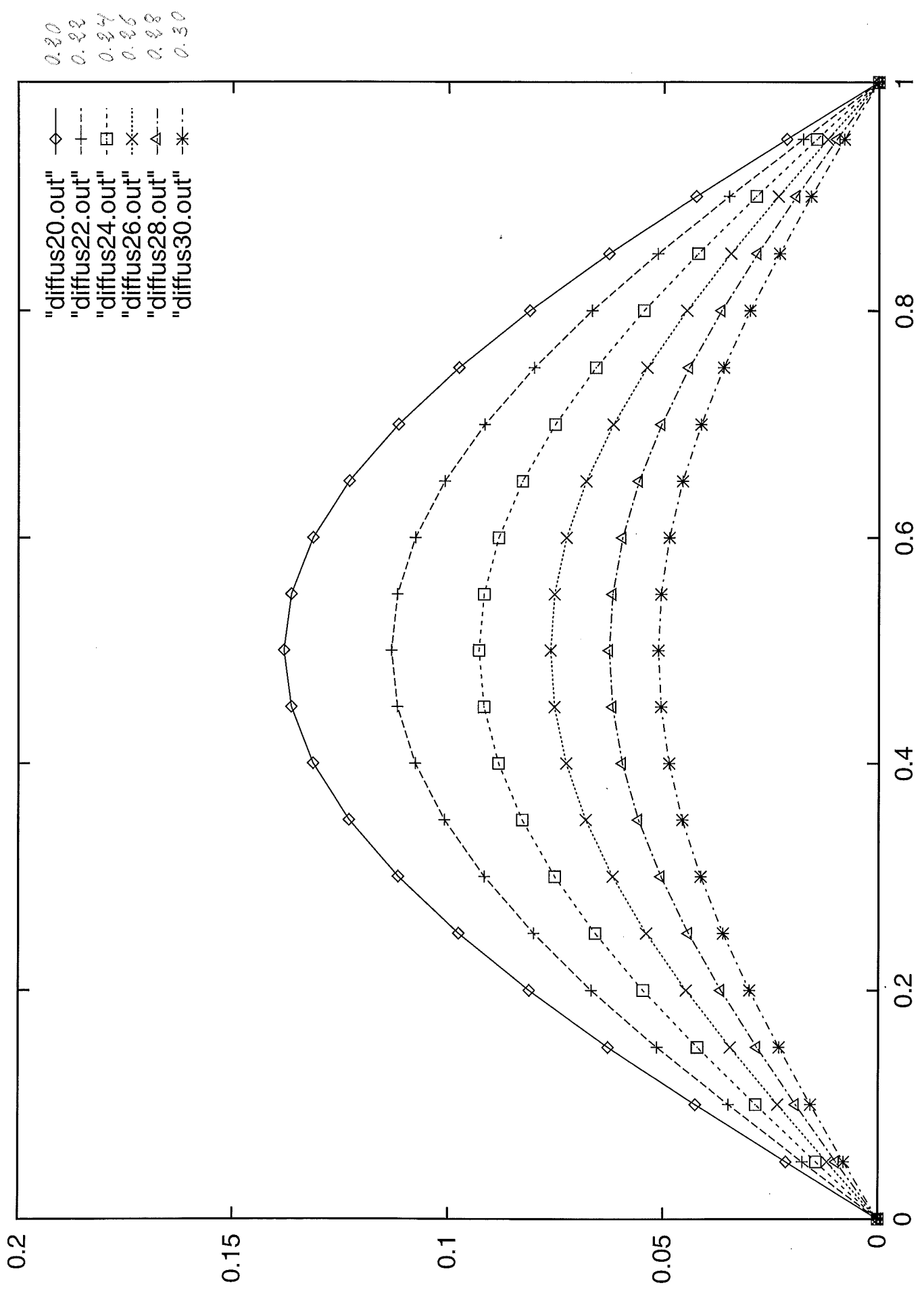
$$-2U_{i+1}^{m+1} + (1+2r)U_i^{m+1} - rU_{i-1}^{m+1} = U_i^m$$

notation

$$\boxed{(I - rS^2) U_i^{m+1} = U_i^m}$$

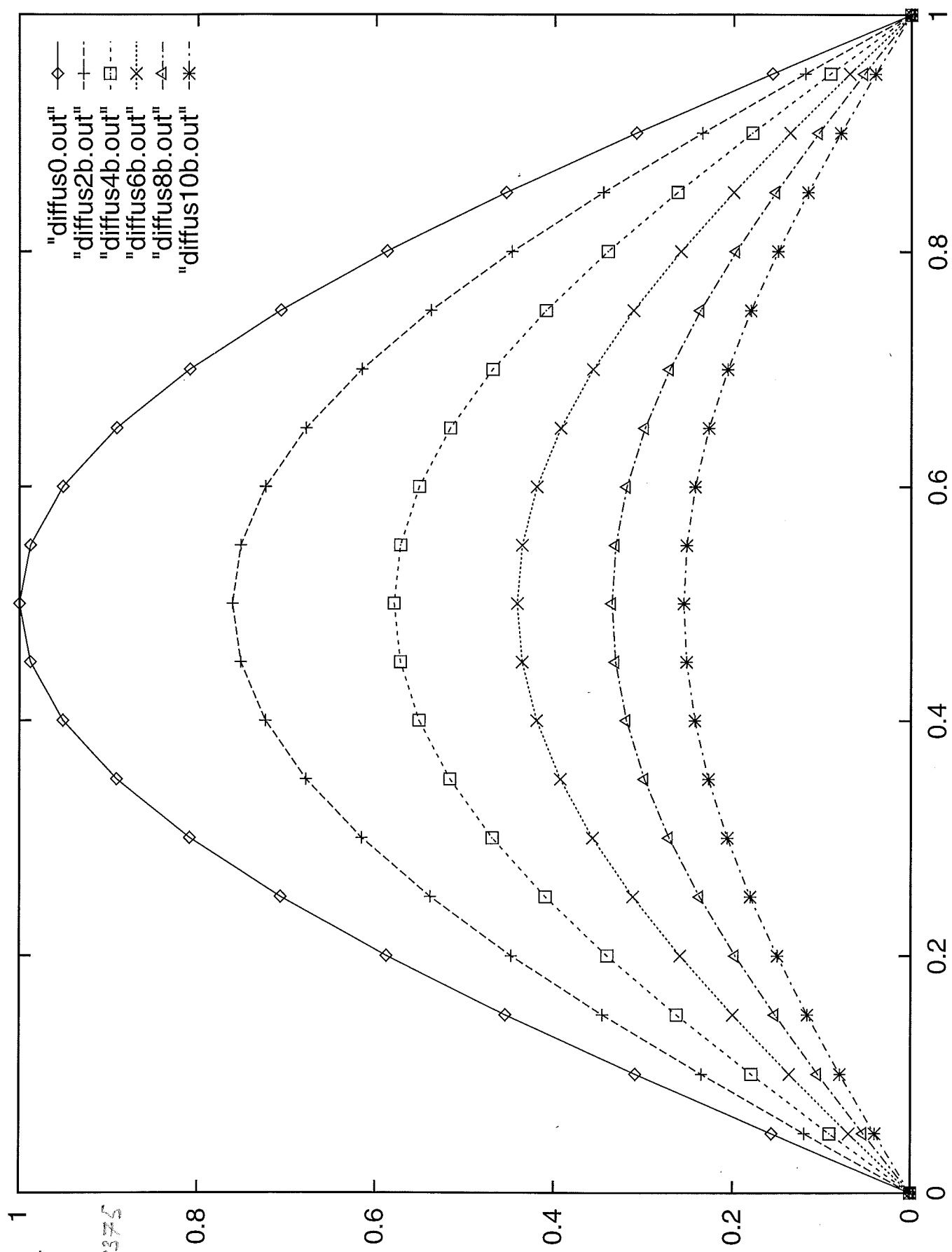
$\hookrightarrow$  Need to solve a tridiagonal system





0.20  
0.22  
0.24  
0.26  
0.28  
0.30

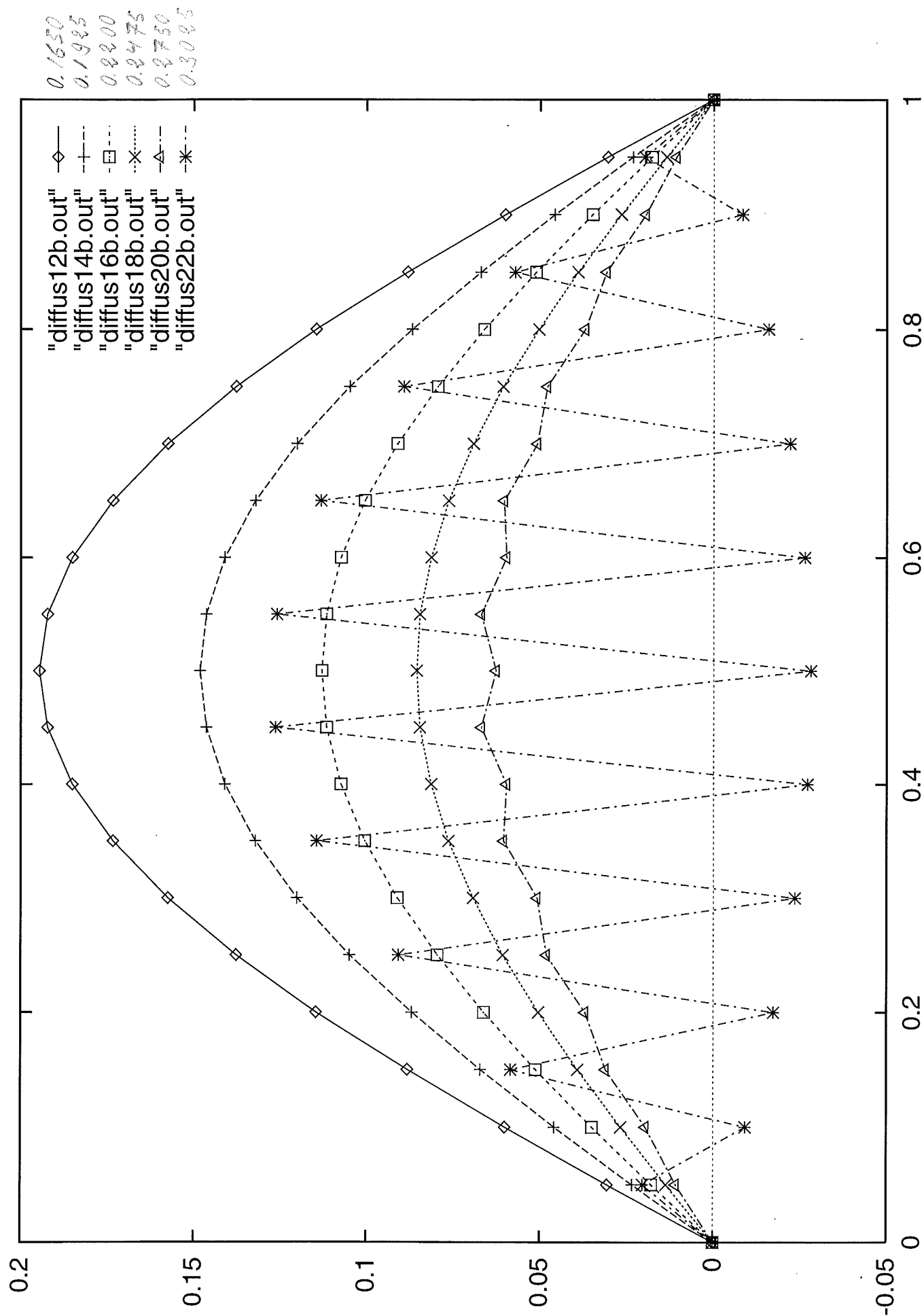
$\frac{e t}{L^2}$   
 0  
 0.0275  
 0.0550  
 0.0825  
 0.0110  
 0.1375



$\alpha = 0.55$

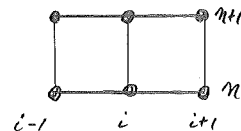
$\rightarrow \frac{e \Delta t}{L^2} = 0.001375$   
 $\rightarrow \frac{e \Delta t}{L^2} = 0.001375$





# Crank-Nicolson Generalized

$$\frac{(u_i^{n+1} - u_i^n)}{\Delta t} = \frac{\epsilon}{h^2} \left( \theta (u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}) + (1-\theta) (u_{i+1}^n - 2u_i^n + u_{i-1}^n) \right)$$



$$-\theta \tau u_{i+1}^{n+1} + (1+2\theta \tau) u_i^{n+1} - \theta \tau u_{i-1}^{n+1} = 2(1-\theta) u_{i+1}^n + (1-2\tau(1-\theta)) u_i^n + 2(1-\theta) u_{i-1}^n$$

$\theta = 1$ : Euler implicit,  $\mathcal{O}(\Delta t, h^2)$

$\theta = 0$ : Euler explicit,  $\mathcal{O}(\Delta t, h^2)$

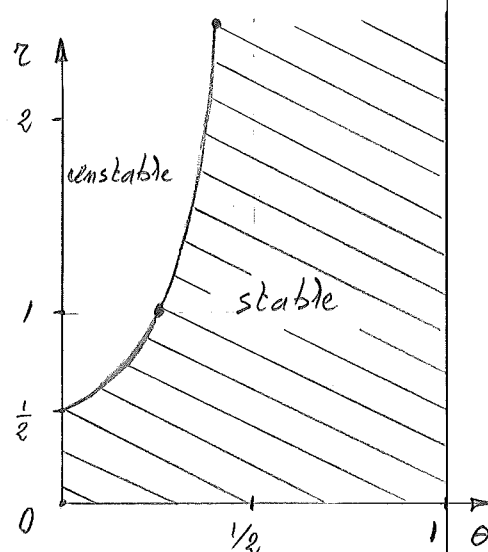
$\theta = \frac{1}{2}$ : { Crank-Nicolson,  $\mathcal{O}((\Delta t)^2, h^2)$   
Trapezoid rule

$\theta$  general: (exercise)

$$\begin{cases} 0 \leq \theta < \frac{1}{2} & \text{stable if } \tau \leq \frac{1}{2(1-2\theta)} \\ \frac{1}{2} \leq \theta \leq 1 & \text{uncond. stable} \end{cases}$$

notation:

$$(I - \theta \tau S^2) u_i^{n+1} = (I + (1-\theta) \tau S^2) u_i^n$$



Other schemes? OK as long as all the  $(\lambda_k \Delta t)$  are within the stability region of the scheme used:

RK2  $\tau \leq \frac{1}{2}$

RK3  $\tau \leq \frac{5}{8}$

AB2  $\tau \leq \frac{1}{4}$

AM3  $\tau \leq \frac{3}{2}$  etc.

## Scheme of Douglas

$$\frac{(v_i^{n+1} - v_i^n)}{\Delta t} = \epsilon \frac{1}{2} \left( \frac{\partial^2 v}{\partial x^2} \Big|_i^{n+1} + \frac{\partial^2 v}{\partial x^2} \Big|_i^n \right) = \mathcal{O}((\Delta t)^2)$$

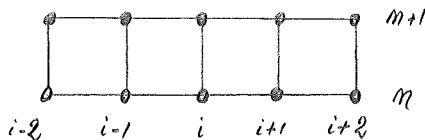
$$= \frac{\epsilon}{2 \hbar^2} \left( I - \frac{S^2}{12} + \frac{S^4}{90} \dots \right) (\psi_i^{n+1} + \psi_i^n)$$

$$(v_i^{m+1} - v_i^m) = \frac{\tau}{2} S^2 \left( I - \frac{S^2}{12} + \frac{S^4}{90} \dots \right) (v_i^{m+1} + v_i^m)$$

- If we only use  $\delta^2$ , we obtain the Crank-Nicolson scheme  $\Rightarrow \mathcal{O}(h^2)$
- How to obtain  $\mathcal{O}(h^4)$ ?

(1) use  $S^2(I - \frac{S^2}{I_2^2})_i (v_i^{m+1} + v_i^m) \Rightarrow$  Widerspruch!

↳ Need to solve a pentadiagonal system



(2) Idea: operate on both sides using  $(I + \frac{\delta^2}{12})$ :

$$\begin{aligned} \left(I + \frac{\delta^2}{12}\right) (v_i^{n+1} - v_i^n) &= \frac{\tau}{2} \left(I + \frac{\delta^2}{12}\right) \delta^2 \left(I - \frac{\delta^2}{12} + \frac{\delta^4}{90} \dots\right) (v_i^{n+1} + v_i^n) \\ &= \frac{\tau}{2} \delta^2 \left(I - \cancel{\frac{\delta^2}{12}} + \frac{\delta^4}{90} + \cancel{\frac{\delta^2}{12}} - \frac{\delta^4}{144} \dots\right) (v_i^{n+1} + v_i^n) \\ &= \frac{\tau}{2} \delta^2 \left(I + \left(\frac{1}{90} - \frac{1}{144}\right) \delta^4 \dots\right) (v_i^{n+1} + v_i^n) \end{aligned}$$

✓✓ scheme

$$(I + \frac{\delta^2}{\beta^2})(v_i^{n+1} - v_i^n) = \frac{2}{\delta} \delta^2 (v_i^{n+1} + v_i^n)$$

$$\left(I - \frac{1}{2} \left(7 - \frac{1}{6}\right) S^2\right) \psi_i^{n+1} = \left(I + \frac{1}{2} \left(7 + \frac{1}{6}\right) S^2\right) \psi_i^n$$

$$12 v_i^{m+1} - (6n-1) s^2 v_i^{m+1} = 12 v_i^m + (6n+1) s^2 v_i^m$$

$$\begin{aligned}
 (1-6\tau) v_{i+1}^{m+1} + (10+12\tau) v_i^{m+1} + (1-6\tau) v_{i-1}^{m+1} \\
 = (1+6\tau) v_{i+1}^m + (10-12\tau) v_i^m + (1+6\tau) v_{i-1}^m
 \end{aligned}$$

Error is  $\mathcal{O}((\Delta t)^2, h^4, (\Delta t)h^2)$   
 and we still have a tridiagonal system

---



•  $\lambda_{kl}/\max$  if  $\frac{k \Delta x}{2} = \frac{l \Delta y}{2} = \frac{\tau}{2}$

$$\hookrightarrow \lambda_{crit} = -4\epsilon \left( \frac{1}{(\Delta x)^2} + \frac{1}{(\Delta y)^2} \right)$$

$$\lambda_{crit} \Delta t = -4\epsilon \left( \frac{1}{(\Delta x)^2} + \frac{1}{(\Delta y)^2} \right) \Delta t = -4(\tau_x + \tau_y)$$

$$\begin{cases} \tau_x = \frac{\epsilon \Delta t}{(\Delta x)^2} \\ \tau_y = \frac{\epsilon \Delta t}{(\Delta y)^2} \end{cases}$$

notation

Example: Euler explicit:

$$4(\tau_x + \tau_y) \leq 2$$

$$\boxed{\tau_x + \tau_y \leq \frac{1}{2}}$$

Case  $\Delta x = \Delta y = h \Rightarrow \tau_x = \tau_y = \tau \Rightarrow \text{need } \boxed{\tau \leq \frac{1}{4}}$

$$u_{ij}^{n+1} = u_{ij}^n + \epsilon \Delta t \left( \frac{S_x^2}{(\Delta x)^2} + \frac{S_y^2}{(\Delta y)^2} \right) u_{ij}^n$$

$$u_{ij}^{n+1} = u_{ij}^n + \frac{\epsilon \Delta t}{(\Delta x)^2} S_x^2 u_{ij}^n + \frac{\epsilon \Delta t}{(\Delta y)^2} S_y^2 u_{ij}^n$$

$$u_{ij}^{n+1} = (I + (\tau_x S_x^2 + \tau_y S_y^2)) u_{ij}^n$$

$$\boxed{u^{n+1} = (I + (\tau_x S_x^2 + \tau_y S_y^2)) u^n}$$

Another simplification  
of notation

Generalized

Crank Nicolson (Implicit)

$$\frac{u^{n+1} - u^n}{\Delta t} = \epsilon \left( \frac{\delta_x^2}{(\Delta x)^2} + \frac{\delta_y^2}{(\Delta y)^2} \right) (\theta u^{n+1} + (1-\theta) u^n)$$

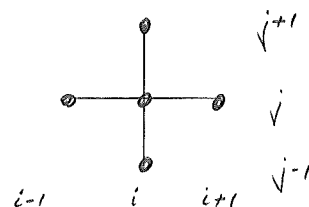
$$u^{n+1} = u^n + (\eta_x \delta_x^2 + \eta_y \delta_y^2) (\theta u^{n+1} + (1-\theta) u^n)$$

$$(I - \theta (\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^{n+1} = (I + (1-\theta) (\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^n$$

$O((\Delta t)^2)$  if  $\theta = 1/2$  (= Crank Nicolson)

$$(I - \frac{1}{2} (\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^{n+1} = (I + \frac{1}{2} (\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^n$$

coupling of  $u_j^{n+1}$  with the 4 neighbors



But the  $u_{ij}$  matrix is stored in the computer as a vector! Hence it is not convenient to solve the above numerically.

→ We need to seek alternatives

→ "Operator splitting" methods

ADI methods ("Alternate Direction Implicit")

Peaceman - Rachford (1955), Douglas (1955)

Step 1:

$$\frac{(u^* - u^n)}{\Delta t/2} = \epsilon \left( \frac{\delta_x^2}{(\Delta x)^2} u^* + \frac{\delta_y^2}{(\Delta y)^2} u^n \right)$$

one tridiagonal system  
to solve for each  $j$   
implicit in  $x$   
explicit in  $y$

$u^*$  is essentially an approximation of  $u^{n+1/2}$

$$\frac{(v^{n+1} - v^*)}{\Delta t/2} = c \left( \frac{\Delta x^2}{(\Delta t)^2} v^* + \frac{\Delta x^2}{(\Delta t)^2} v^{n+1} \right) \quad \begin{array}{l} \text{explicit in } x \\ \text{implicit in } t \end{array}$$

Using the simplified notation:

implicit in  $y$   
 → one tri-diagonal systems  
 to solve for each  $i$

$$\left( I - \frac{1}{2} \gamma_x \delta x^2 \right) \psi^* = \left( I + \frac{1}{2} \gamma_y \delta y^2 \right) \psi^m \quad (1)$$

$$\left(I - \frac{1}{2} \gamma_Y S_Y^2\right) \psi^{n+1} = \left(I + \frac{1}{2} \gamma_X S_X^2\right) \psi^* \quad (9)$$

### Composite form of the scheme

$$\begin{aligned} \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \left( I - \frac{1}{2} \gamma_y S_y^2 \right) \psi^{n+1} &= \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \left( I + \frac{1}{2} \gamma_x S_x^2 \right) \psi^* \\ &\quad \swarrow \searrow S_x^2 \text{ commutes with itself} \\ &= \left( I + \frac{1}{2} \gamma_x S_x^2 \right) \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \psi^* \\ &= \left( I + \frac{1}{2} \gamma_x S_x^2 \right) \underbrace{\left( I + \frac{1}{2} \gamma_y S_y^2 \right)}_{III} \psi^n \end{aligned}$$

$$\hookrightarrow \left( I - \frac{1}{2} (\gamma_x \delta_x^2 + \gamma_y \delta_y^2) + \frac{1}{4} \gamma_x \gamma_y \delta_x^2 \delta_y^2 \right) \psi^{n+1} = \left( I + \frac{1}{2} (\gamma_x \delta_x^2 + \gamma_y \delta_y^2) + \frac{1}{4} \gamma_x \gamma_y \delta_x^2 \delta_y^2 \right) \psi^n$$

This is basically the Crank-Nicolson scheme, with a higher order

error term:  $\frac{1}{4} r_x r_y S_x^2 S_y^2 (v^{n+1} - v^n) \approx \frac{1}{4} r_x r_y \Delta t S_x^2 S_y^2 \left( \frac{\partial v}{\partial t} \right)^{n+1/2}$

$\hookrightarrow$  scheme is  $O((\Delta t)^2)$  still.

## Stability

$$\psi_i^m = \rho^m e^{i(kx_i + l y_i)} A(0)$$

$\equiv$  simplified notation ... it stands for  $\beta_{kl}$  and  $A_{kl}(0)$   
and there should also be  $\sum_{kl} \dots$



$$u_{ij}^{m+1} = \underbrace{t_2 \cdot t_1 \cdot \rho^m}_{\rho} e^{i(k \cdot x_{i+1} + y_j)} A(0) \text{ after step 2}$$

$$\lambda_1 = \frac{(1 - \frac{v}{c} \sin^2(\frac{\lambda \Delta y}{\lambda}))}{(1 + \frac{v}{c} \sin^2(\frac{\lambda \Delta x}{\lambda}))} = \frac{(1 - \frac{1}{2} \alpha_y)}{(1 + \frac{1}{2} \alpha_x)} = \frac{(1 + \frac{1}{2} \lambda_y \Delta t)}{(1 - \frac{1}{2} \lambda_x \Delta t)}$$

and

$$\beta_z = \frac{(1 - 2\gamma_x \sin^2(\frac{k \Delta x}{2}))}{(1 + 2\gamma_y \sin^2(\frac{\lambda \Delta y}{2}))} = \frac{(1 - \frac{1}{2} \alpha_x)}{(1 + \frac{1}{2} \alpha_y)} = \frac{(1 + \frac{1}{2} \lambda_x \Delta t)}{(1 - \frac{1}{2} \lambda_y \Delta t)}$$

Finally:

$$\rho = \xi_2 \xi_1 = \left( \frac{1 - \frac{1}{2} dx}{1 + \frac{1}{2} dy} \right) \left( \frac{1 - \frac{1}{2} dy}{1 + \frac{1}{2} dx} \right) = \underbrace{\left( \frac{1 - \frac{1}{2} dx}{1 + \frac{1}{2} dx} \right)}_{\rho_x} \underbrace{\left( \frac{1 - \frac{1}{2} dy}{1 + \frac{1}{2} dy} \right)}_{\rho_y} = \rho_x \cdot \rho_y$$

$\beta_x \equiv$  same as when applying CN scheme on  $x \Rightarrow |\beta_x| \leq 1$

$\rho_y \equiv$  name as when applying CW scheme on  $y \Rightarrow |\rho_y| \leq 1$

$\hookrightarrow |A| \leq 1 \Rightarrow$  scheme is unconditionally stable.

### Treatment of boundary conditions

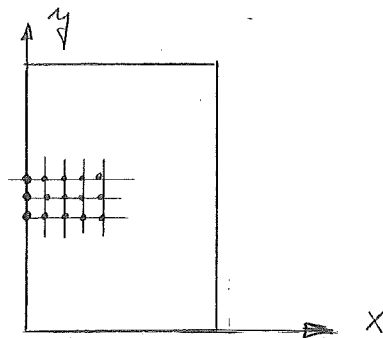
↳ We need to derive what is the proper condition to use for  $v^*$  so as to preserve the order of the scheme.

Example: case with  $v$  given on the boundaries

$$\left( \mathbf{I} - \frac{1}{\rho} \gamma_x \gamma_x^T \right) \mathcal{V}^* = \left( \mathbf{I} + \frac{1}{\rho} \gamma_y \gamma_y^T \right) \mathcal{V}^m \quad (11)$$

$$\left( \mathbb{I} - \frac{1}{\ell} \gamma \gamma^T \right) \psi^{n+1} = \left( \mathbb{I} - \frac{1}{\ell} \gamma \gamma^T \right) \psi^* \quad (2)$$

$$(1) \Rightarrow \frac{1}{\rho} \gamma_x \delta x^2 v^* = v^* - \left( I^* \frac{1}{\rho} \gamma_y \delta y^2 \right) v^m \quad (3)$$



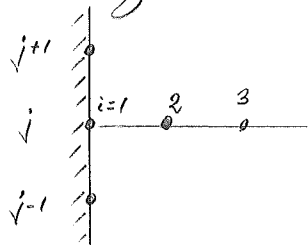
Insert (3) into (2):

$$\begin{aligned} \left( I - \frac{1}{2} \gamma \gamma^2 \right) \psi^{n+1} &= \psi^* + \left( \psi^* - \left( I + \frac{1}{2} \gamma \gamma^2 \right) \psi^n \right) \\ &= \psi^* - \left( I + \frac{1}{2} \gamma \gamma^2 \right) \psi^n \end{aligned}$$

$$\begin{aligned} \psi^* &= \frac{1}{2} \left[ \left( I - \frac{1}{2} r_y^2 S_y^2 \right) \psi^{n+1} + \left( I + \frac{1}{2} r_y^2 S_y^2 \right) \psi^n \right] \\ &= \frac{1}{2} (\psi^{n+1} + \psi^n) - \frac{1}{2} r_y^2 S_y^2 (\psi^{n+1} - \psi^n) \equiv \text{proper BC for } \psi^* \end{aligned}$$

Example: left boundary on the vertical boundaries

Example: left boundary



$v_{i,j}^n$  is given for all  $j$  and  $n$   
as  $v(x_i, y_j, t^n)$  is given

$$v_{ij}^* = \frac{1}{2} \left[ \left( v_{ij}^{n+1} - \frac{1}{2} \gamma \left( v_{i,j+1}^{n+1} - 2 v_{ij}^{n+1} + v_{i,j-1}^{n+1} \right) \right) \right. \\ \left. + \left( v_{ij}^n + \frac{1}{2} \gamma \left( v_{i,j+1}^n - 2 v_{ij}^n + v_{i,j-1}^n \right) \right) \right]$$

Note: when BC is not function of  $t$ , we have:  $V_{ij}^* = V_{ij}$   
when BC is not function of  $y$ , we have:  $V_{ij}^* = \frac{1}{2}(V_i^{m+1} + V_i^m)$

ADI Douglas - Rachford

$$(I - \eta_x S_x^2) \cup^* = (I + \eta_y S_y^2) \cup^n \quad (1)$$

$$(I - \gamma \delta \gamma^2) \psi^{n+1} = \psi^* - \gamma \delta \gamma^2 \psi^n \quad (2)$$

Composite form:

$$\begin{aligned} (I - \gamma_x S_x^2) (I - \gamma_y S_y^2) \psi^{n+1} &= \underbrace{(I - \gamma_x S_x^2) \psi^n}_{III} - (I - \gamma_x S_x^2) \gamma_y S_y^2 \psi^n \\ &= (I + \gamma_y S_y^2) \psi^n - \gamma_y S_y^2 \psi^n + \gamma_x \gamma_y S_x^2 S_y^2 \psi^n \\ &= (I + \gamma_x \gamma_y S_x^2 S_y^2) \psi^n \end{aligned}$$

$$(I - (r_x S_x^2 + r_y S_y^2) + r_x r_y S_x^2 S_y^2) v^{n+1} = (I + r_x r_y S_x^2 S_y^2) v^n$$

This is basically the Euler scheme, with a higher order error term:

$$\tau_x \tau_y \delta x^2 \delta y^2 (v^{m+1} - v^n) \approx \tau_x \tau_y \Delta t \delta x^2 \delta y^2 \left( \frac{\partial v}{\partial t} \right)^{m+1/2}$$

$\hookrightarrow$  scheme is  $\mathcal{O}(\Delta t)$  stiff.

Stability:

$$f_1 = \frac{(1 - a_1)}{(1 + a_1)}$$

$$\beta = \frac{1}{(1+a_y)} (\xi_1 + a_y) = \frac{1}{(1+a_y)} \left( \left( \frac{1-a_y}{1+a_x} \right) + a_y \right) = \frac{1+a_y a_x}{(1+a_y)(1+a_x)}$$

$|p| \leq 1 \Rightarrow$  uncond. stable

### Boundary conditions

$$(2) \Rightarrow v^* = (I - \gamma_y \delta_y^2) v^{n+1} + \gamma_y \delta_y^2 v^n$$

$$= v^{n+1} - \gamma_y \delta_y^2 (v^{n+1} - v^n)$$

ADI Mitchell - Fairweather

$$\left(I - \frac{1}{2} \left(\tau_x - \frac{1}{6}\right) \delta_x^2\right) u^* = \left(I + \frac{1}{2} \left(\tau_y + \frac{1}{6}\right) \delta_y^2\right) u^n \quad (1)$$

$$\left(I - \frac{1}{2} \left(\tau_y - \frac{1}{6}\right) \delta_y^2\right) u^{n+1} = \left(I + \frac{1}{2} \left(\tau_x + \frac{1}{6}\right) \delta_x^2\right) u^* \quad (2)$$

Composite Form

$$\left(I - \frac{1}{2} \left(\tau_x - \frac{1}{6}\right) \delta_x^2\right) \left(I - \frac{1}{2} \left(\tau_y - \frac{1}{6}\right) \delta_y^2\right) u^{n+1} = \left(I + \frac{1}{2} \left(\tau_x + \frac{1}{6}\right) \delta_x^2\right) \left(I + \frac{1}{2} \left(\tau_y + \frac{1}{6}\right) \delta_y^2\right) u^n$$

$$\hookrightarrow O(\Delta t)^2, (\Delta x)^4, (\Delta y)^4$$

Stability: uncond. stableBoundary conditions (exercise)Methods "Locally One-Dimensional" (LOD)

Yanenko (1971)

LOD Euler

$$\begin{cases} \frac{u^* - u^n}{\Delta t} = \epsilon \frac{\delta_x^2}{(\Delta x)^2} u^* & (1) \\ \frac{u^{n+1} - u^*}{\Delta t} = \epsilon \frac{\delta_y^2}{(\Delta y)^2} u^{n+1} & (2) \end{cases}$$

$$(I - \tau_x \delta_x^2) u^* = u^n \quad (1)$$

$$\zeta_1 = \frac{1}{1 + \alpha_x}$$

uncond. stable

$$(I - \tau_y \delta_y^2) u^{n+1} = u^* \quad (2)$$

$$\zeta_2 = \frac{1}{1 + \alpha_y}$$

uncond. stable

Composite form

$$(I - \tau_x \delta_x^2) (I - \tau_y \delta_y^2) u^{n+1} = u^n$$

 $\hookrightarrow$  This is basically the Euler scheme, with higher order errorterm:  $\tau_x \tau_y \delta_x^2 \delta_y^2 u^{n+1}$ . This is not as good as  $\tau_x \tau_y \delta_x^2 \delta_y^2 (u^{n+1} - u^n)$ !

$$\beta = \frac{1}{(1 + \alpha_y)} \cdot \frac{1}{(1 + \alpha_x)}$$

uncond. stable

$$\text{BC: (2)} \Rightarrow u^* = (I - \tau_y \delta_y^2) u^{n+1} \\ = u^{n+1} - \tau_y \delta_y^2 u^{n+1}$$

LOD Crank Nicolson

$$\left( I - \frac{1}{2} \gamma \delta \gamma^2 \right) \psi^{n+1} = \left( I + \frac{1}{2} \gamma \delta \gamma^2 \right) \psi^* \quad (2) \quad \xi_2 = \left( \frac{1 - \frac{1}{2} a_\gamma}{1 + \frac{1}{2} a_\gamma} \right) \quad \text{uncond. stable}$$

Composite form:

$$\begin{aligned} \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \left( I - \frac{1}{2} \gamma_y S_y^2 \right) \psi^{n+1} &= \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \left( I + \frac{1}{2} \gamma_y S_y^2 \right) \psi^* \\ &\quad \swarrow \searrow \text{as } S_x^2 \text{ and } S_y^2 \text{ commute} \\ &= \left( I + \frac{1}{2} \gamma_y S_y^2 \right) \left( I - \frac{1}{2} \gamma_x S_x^2 \right) \psi^* \\ &= \left( I + \frac{1}{2} \gamma_y S_y^2 \right) \left( I + \frac{1}{2} \gamma_x S_x^2 \right) \psi^n \end{aligned}$$

scheme is  $\mathcal{O}((\Delta t)^2)$ , as Crank-Nicolson.

$$\rho = \begin{pmatrix} \frac{1 - \frac{1}{2} a_y}{1 + \frac{1}{2} a_y} & \frac{1 - \frac{1}{2} a_x}{1 + \frac{1}{2} a_x} \end{pmatrix} \quad \text{uncond. stable.}$$

BC:

$$(1) \Rightarrow \frac{1}{2} \partial_X S_X^2 \psi^* = \psi^* - \left( I + \frac{1}{2} \partial_X S_X^2 \right) \psi^m$$

↳ we cannot insert that into (2) !

↳ Impossible to find a proper condition on  $v^*$  so as to preserve the order of the scheme at the boundary?

Def's:  $a_x = 4 \eta_x \sin^2\left(\frac{k \Delta x}{2}\right) = -\lambda_x \Delta t$

$a_y = 4 \eta_y \sin^2\left(\frac{l \Delta y}{2}\right) = -\lambda_y \Delta t$

Stability

• Euler implicit:  $\mathcal{O}(\Delta t)$

$(I - (\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^{n+1} = u^n$

$\beta = \frac{1}{1 + (a_x + a_y)}$

uncond.

• Crank-Nicolson:  $\mathcal{O}((\Delta t)^2)$

$(I - \frac{1}{2}(\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^{n+1} = (I + \frac{1}{2}(\eta_x \delta_x^2 + \eta_y \delta_y^2)) u^n$

$\beta = \frac{(1 - \frac{1}{2}(a_x + a_y))}{(1 + \frac{1}{2}(a_x + a_y))}$

uncond.

• ADI (Peaceman-Rachford):  $\mathcal{O}((\Delta t)^2)$

$(I - \frac{1}{2} \eta_x \delta_x^2) u^* = (I + \frac{1}{2} \eta_y \delta_y^2) u^n$

$\zeta_1 = \frac{(1 - \frac{1}{2} a_y)}{(1 + \frac{1}{2} a_x)}$

$(I - \frac{1}{2} \eta_y \delta_y^2) u^{n+1} = (I + \frac{1}{2} \eta_x \delta_x^2) u^*$

$\zeta_2 = \frac{(1 - \frac{1}{2} a_x)}{(1 + \frac{1}{2} a_y)}$

Composite:

$(I - \frac{1}{2} \eta_x \delta_x^2)(I - \frac{1}{2} \eta_y \delta_y^2) u^{n+1} = (I + \frac{1}{2} \eta_x \delta_x^2)(I + \frac{1}{2} \eta_y \delta_y^2) u^n$

$\beta = \zeta_2 \cdot \zeta_1 = \frac{(1 - \frac{1}{2} a_x)}{(1 + \frac{1}{2} a_y)} \cdot \frac{(1 - \frac{1}{2} a_y)}{(1 + \frac{1}{2} a_x)} = \frac{(1 - \frac{1}{2} a_x)}{(1 + \frac{1}{2} a_x)} \cdot \frac{(1 - \frac{1}{2} a_y)}{(1 + \frac{1}{2} a_y)}$

uncond.

• ADI (Douglas-Rachford):  $\mathcal{O}(\Delta t)$

$(I - \eta_x \delta_x^2) u^* = (I + \eta_y \delta_y^2) u^n$

$\zeta_1 = \frac{(1 - a_y)}{(1 + a_x)}$

$(I - \eta_y \delta_y^2) u^{n+1} = u^* - \eta_y \delta_y^2 u^n$

Composite:

$(I - \eta_x \delta_x^2)(I - \eta_y \delta_y^2) u^{n+1} = (I + \eta_x \eta_y \delta_x^2 \delta_y^2) u^n$

$\beta = \frac{1}{(1 + a_y)} (\zeta_1 + a_y) = \frac{(1 + a_y a_x)}{(1 + a_y)(1 + a_x)}$

uncond.

Stability

• ADI (Mitchell - Fairweather):  $\mathcal{O}(\Delta t)^2, (\Delta x)^4, (\Delta y)^4$

$$\left( I - \frac{1}{2} \left( \eta_x - \frac{1}{\delta} \right) S_x^2 \right) \psi^* = \left( I + \frac{1}{2} \left( \eta_y + \frac{1}{\delta} \right) S_y^2 \right) \psi^n$$

$$\left( I - \frac{1}{2} \left( \eta_y - \frac{1}{\delta} \right) S_y^2 \right) \psi^{n+1} = \left( I + \frac{1}{2} \left( \eta_x + \frac{1}{\delta} \right) S_x^2 \right) \psi^*$$

Composite:

$$\left( I - \frac{1}{2} \left( \eta_x - \frac{1}{\delta} \right) S_x^2 \right) \left( I - \frac{1}{2} \left( \eta_y - \frac{1}{\delta} \right) S_y^2 \right) \psi^{n+1} = \left( I + \frac{1}{2} \left( \eta_x + \frac{1}{\delta} \right) S_x^2 \right) \left( I + \frac{1}{2} \left( \eta_y + \frac{1}{\delta} \right) S_y^2 \right) \psi^n$$

uncond.

• LOD Euler:  $\mathcal{O}(\Delta t)$

$$\left( I - \eta_x S_x^2 \right) \psi^* = \psi^n$$

$$\zeta_1 = \frac{1}{(1 + a_x)}$$

$$\left( I - \eta_y S_y^2 \right) \psi^{n+1} = \psi^*$$

$$\zeta_2 = \frac{1}{(1 + a_y)}$$

Composite:

$$\left( I - \eta_x S_x^2 \right) \left( I - \eta_y S_y^2 \right) \psi^{n+1} = \psi^n$$

$$\rho = \zeta_2 \cdot \zeta_1 = \frac{1}{(1 + a_y)} \cdot \frac{1}{(1 + a_x)}$$

uncond

• LOD Crank-Nicolson:  $\mathcal{O}(\Delta t)^2$

$$\left( I - \frac{1}{2} \eta_x S_x^2 \right) \psi^* = \left( I + \frac{1}{2} \eta_x S_x^2 \right) \psi^n$$

$$\zeta_1 = \frac{(1 - \frac{1}{2} a_x)}{(1 + \frac{1}{2} a_x)}$$

$$\left( I - \frac{1}{2} \eta_y S_y^2 \right) \psi^{n+1} = \left( I + \frac{1}{2} \eta_y S_y^2 \right) \psi^*$$

$$\zeta_2 = \frac{(1 - \frac{1}{2} a_y)}{(1 + \frac{1}{2} a_y)}$$

Composite: (Seulement si  $S_x^2$  et  $S_y^2$  commutent)

$$\left( 1 - \frac{1}{2} \eta_x S_x^2 \right) \left( 1 - \frac{1}{2} \eta_y S_y^2 \right) \psi^{n+1} = \left( 1 + \frac{1}{2} \eta_y S_y^2 \right) \left( 1 + \frac{1}{2} \eta_x S_x^2 \right) \psi^n$$

$$\rho = \zeta_2 \cdot \zeta_1 = \frac{(1 - \frac{1}{2} a_y)}{(1 + \frac{1}{2} a_y)} \cdot \frac{(1 - \frac{1}{2} a_x)}{(1 + \frac{1}{2} a_x)}$$

uncond

• Trivial extension of ADI Peaceman-Rachford to 3-D :  $\mathcal{O}(\Delta t)$ !

Stability

$$\begin{cases} (I - \frac{1}{3} \eta_x S_x^2) \psi^* = (I + \frac{1}{3} (\eta_y S_y^2 + \eta_z S_z^2)) \psi^n \\ (I - \frac{1}{3} \eta_y S_y^2) \psi^{**} = (I + \frac{1}{3} (\eta_z S_z^2 + \eta_x S_x^2)) \psi^* \\ (I - \frac{1}{3} \eta_z S_z^2) \psi^{***} = (I + \frac{1}{3} (\eta_x S_x^2 + \eta_y S_y^2)) \psi^{**} \end{cases}$$

Bad!

$$\rho = \zeta_3 \cdot \zeta_2 \cdot \zeta_1 = \frac{(1 - \frac{1}{3}(a_x + a_y))}{(1 + \frac{1}{3} a_z)} \cdot \frac{(1 - \frac{1}{3}(a_y + a_z))}{(1 + \frac{1}{3} a_x)} \cdot \frac{(1 - \frac{1}{3}(a_z + a_x))}{(1 + \frac{1}{3} a_y)}$$

No!

• <sup>ADI</sup> Douglas-Rachford :  $\mathcal{O}(\Delta t)$

$$\begin{cases} (I - \eta_x S_x^2) \psi^* = (I + (\eta_y S_y^2 + \eta_z S_z^2)) \psi^n \\ (I - \eta_y S_y^2) \psi^{**} = \psi^* - \eta_y S_y^2 \psi^n \\ (I - \eta_z S_z^2) \psi^{***} = \psi^{**} - \eta_z S_z^2 \psi^n \end{cases}$$

$$\zeta_1 = \frac{1 - (a_y + a_z)}{(1 + a_x)}$$

Composite:

$$(I - \eta_x S_x^2)(I - \eta_y S_y^2)(I - \eta_z S_z^2) \psi^{***} = (I + (\eta_x \eta_y S_x^2 S_y^2 + \eta_y \eta_z S_y^2 S_z^2 + \eta_z \eta_x S_z^2 S_x^2) - \eta_x \eta_y \eta_z S_x^2 S_y^2 S_z^2) \psi^n$$

$$\rho = \frac{1}{(1 + a_z)} \left( \frac{1}{(1 + a_y)} (\zeta_1 + a_y) + a_z \right) = \frac{1 + (a_x a_y + a_y a_z + a_z a_x) + a_x a_y a_z}{(1 + a_z)(1 + a_y)(1 + a_x)}$$

uncond

• <sup>ADI</sup> Douglas-Bryan :  $\mathcal{O}((\Delta t)^2)$

$$\begin{cases} (I - \frac{1}{2} \eta_x S_x^2) \psi^* = (I + \frac{1}{2} \eta_x S_x^2 + (\eta_y S_y^2 + \eta_z S_z^2)) \psi^n \\ (I - \frac{1}{2} \eta_y S_y^2) \psi^{**} = (I + \frac{1}{2} (\eta_x S_x^2 + \eta_y S_y^2) + \eta_z S_z^2) \psi^* + \frac{1}{2} \eta_x S_x^2 \psi^* \\ (I - \frac{1}{2} \eta_z S_z^2) \psi^{***} = (I + \frac{1}{2} (\eta_x S_x^2 + \eta_y S_y^2 + \eta_z S_z^2)) \psi^{**} + \frac{1}{2} \eta_x S_x^2 \psi^* + \frac{1}{2} \eta_y S_y^2 \psi^{**} \end{cases}$$

Composite:

$$(I - \frac{1}{2} \eta_x S_x^2)(I - \frac{1}{2} \eta_y S_y^2)(I - \frac{1}{2} \eta_z S_z^2) \psi^{***} = (I + \frac{1}{2} (\eta_x S_x^2 + \eta_y S_y^2 + \eta_z S_z^2) + \frac{1}{4} (\eta_x \eta_y S_x^2 S_y^2 + \eta_y \eta_z S_y^2 S_z^2 + \eta_z \eta_x S_z^2 S_x^2) - \frac{1}{8} \eta_x \eta_y \eta_z S_x^2 S_y^2 S_z^2) \psi^n$$

$$\rho = \frac{1 - \frac{1}{2}(a_x + a_y + a_z) + \frac{1}{4}(a_x a_y + a_y a_z + a_z a_x) + \frac{1}{8} a_x a_y a_z}{(1 + \frac{1}{2} a_x)(1 + \frac{1}{2} a_y)(1 + \frac{1}{2} a_z)} = \frac{((1 - \frac{1}{2} a_x)(1 - \frac{1}{2} a_y)(1 - \frac{1}{2} a_z) + \frac{1}{4} a_x a_y a_z)}{(1 + \frac{1}{2} a_x)(1 + \frac{1}{2} a_y)(1 + \frac{1}{2} a_z)}$$

uncond.

This scheme is the proper extension of the ADI Peaceman-Rachford to 3D.