Module 6 : Wave Guides

Lecture 42 : Phase Velocity and Dispersion

Objectives

In this course you will learn the following

Cut-off Frequency of a Mode.

Lecture 42: Phase Velocity and Dispersion

Cut-off Frequency of a Mode

The propagation constant of the modal fields in the Z-direction is given as

$$\beta = \beta_1 \sqrt{1 - (\frac{m\lambda_1}{2d})^2} - \dots (6.17)$$

$$= \sqrt{\beta_1^2 - (\frac{m\pi}{d})^2} - \dots (6.18)$$

- If the propagation constant β is real the mode will be at travelling mode whereas, if β becomes imaginary the wave exponentially decays in the z-direction and the fields do not represent a wave. These fields are then called 'EVANASENT FIELDS'. The evanasent fields do not carry any power. The power is carried only by the travelling modes.
- For the travelling mode therefore we need β to be real which implies

$$\beta = real$$

$$\Rightarrow \qquad \beta_1 \ge \frac{m\pi}{d} \qquad (6.19)$$

Since

$$\beta_1 = \frac{2\pi}{\lambda_1} = \frac{2\pi f}{v_1},$$
 (6.20)

where v_1 is the velocity of the uniform plane wave in medium 1, we get

$$f \ge \frac{mv_1}{2d} - \dots (6.21)$$

$$\Rightarrow \lambda_1 \le \frac{2d}{m} - \dots (6.22)$$

- We can note the following important things at this stage:
- For a given waveguide height d, the frequency has to be higher than certain threshold frequency for propagation of a particular mode. The threshold frequency is called **CUT-OFF** freuquency at the mode and given by

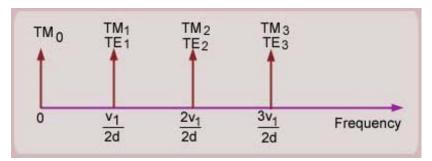
$$f_{cm} = \frac{mv_1}{2d}$$
 ----- (6.23)

The corresponding cut-off wavelength is

$$\lambda_{cm} = \frac{2d}{m} - \cdots - (6.24)$$

- For a given waveguide height, d, and frequency, f, only those modes propagate for which $m \le \frac{2df}{v_1} = \frac{2d}{\lambda}$. This means inside a wave guide there is a possibility of only finite number of modes at a given frequency.
- As the mode number (m) increases and the cut-off frequency also increases meaning higher order mode get excited only at higher frequencies.

The cut-off frequencies for different modes are shown in the following figure:



If a mode has the cut-off frequency less than the frequency of operation the mode propagates otherwise it does not propagate.

NOTE:

The $\overline{TM_0}$ mode which is also the TEM mode has no cut-off frequency. This is the mode which can propagate at any frequencies starting from dc.

Lecture 42: Phase Velocity and Dispersion

The phase velocity of a mode is

$$v_p = \frac{\omega}{\beta} = \frac{\omega}{\beta_1 \sqrt{1 - (\frac{m\lambda_1}{2d})^2}} = \frac{v_1}{\sqrt{1 - (\frac{m\lambda_1}{2d})^2}}$$

where v_1 is the velocity of the uniform plane wave in the medium filling the region between the wave guide walls.

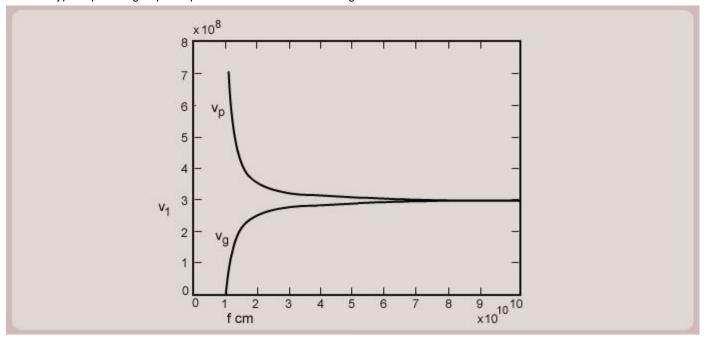
Noting that $\frac{2d}{m} = \lambda_{cm}$ the phase velocity can be written as

$$v_p = \frac{v_1}{\sqrt{1 - (\frac{\lambda_1}{\lambda_{cm}})^2}} = \frac{v_1}{\sqrt{1 - (\frac{f_{cm}}{f})^2}}$$
 -----(6.26)

- It is clear from the expression that in general the phase velocity of a mode is a function of frequency except when the cutoff frequency is zero. This phenomena is called ' **WAVE DISPERSION**'.
- In general, one can then say that the modal propagation on a wave guide is dispersive in nature.
- The group velocity of the mode is

$$v_g = \frac{{v_1}^2}{v_p} = v_1 \sqrt{1 - \left(\frac{fcm}{f}\right)^2}$$
 ----- (6.27)

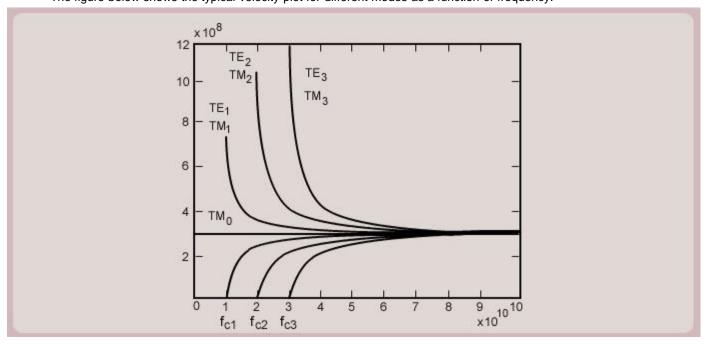
Typical plot for group and phase velocities is shown in figure below



At the cut-off frequency the phase velocity approaches infinity whereas the group velocity approaches zero. That means

the energy propagation seizes as the mode approaches cut-off. As the frequency increases both group and phase velocities asymptotically approach the velocity of the plane wave in the media.

The figure below shows the typical velocity plot for different modes as a function of frequency.



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Recap

In this course you have learnt the following

Cut-off Frequency of a Mode.