

# Blow-up of Solutions of Semilinear Parabolic Equations

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# 1 Sufficient conditions for blow-up

In this section, we consider blow-up of solutions of the problem

$$(F) \quad \begin{cases} u_t = \Delta u + f(u), & x \in \Omega, \\ u = 0, & x \in \partial\Omega, \\ u(x, 0) = u_0(x), & x \in \Omega. \end{cases} \quad (\text{if } \partial\Omega \neq \emptyset)$$

where  $\Omega$  is a bounded or unbounded domain in  $\mathbb{R}^N$ ,  $N \geq 1$ ,  $u_0$  is continuous and  $u_0 \in L^\infty(\Omega)$ . We begin with the fundamental result concerning the existence of solutions.

**Theorem 1.** *If  $f \in C^1$ , then (F) has a unique (classical) solution  $u$  defined for  $t \in [0, T)$ ,  $0 < T \leq \infty$ . If  $T < \infty$ , then  $\|u(\cdot, t)\|_\infty \rightarrow \infty$  as  $t \rightarrow T$  and  $u$  is said to blow up at time  $T$ .*

*Idea of proof.* Let  $e^{t\Delta}$  be the heat semigroup in  $\Omega$  with the Dirichlet boundary condition, that is,  $u = e^{t\Delta}u_0$  is the solution of

$$\begin{aligned} u_t &= \Delta u && \text{in } \Omega \times (0, \infty), \\ u &= 0 && \text{on } \partial\Omega \times (0, \infty), \\ u(x, 0) &= u_0(x) && \text{in } \Omega. \end{aligned}$$

In particular, if  $\Omega = \mathbb{R}^N$ , then

$$(e^{t\Delta}u_0)(x) = (4\pi t)^{-\frac{N}{2}} \int_{\mathbb{R}^N} u_0(y) \exp\left(-\frac{|x-y|^2}{4t}\right) dy.$$

By using  $e^{t\Delta}$ , the problem (F) is converted into the integral equation

$$u(\cdot, t) = e^{t\Delta}u_0 + \int_0^t e^{(t-s)\Delta} f(u(\cdot, s)) ds.$$

With the Banach fixed point theorem, a unique solution is obtained, and the solution can be continued as long as it is bounded. q.e.d.

Next, we give a sufficient condition for the blowup of solutions.

**Theorem 2** (Kaplan, 1963 [24]). *Let  $\Omega$  be bounded. Let  $\lambda_1, \varphi_1$  denote the first eigenvalue and eigenfunction of  $-\Delta$  with the Dirichlet boundary condition, that is,*

$$\begin{aligned} \Delta\varphi_1 + \lambda_1\varphi_1 &= 0 && \text{in } \Omega, \\ \varphi_1 &= 0 && \text{on } \partial\Omega, \\ \varphi_1 &> 0 && \text{in } \Omega, \quad \int_\Omega \varphi_1(x) dx = 1. \end{aligned}$$

Assume that  $f : [0, \infty) \rightarrow [0, \infty)$  is convex,  $f(u) > 0$  for  $u > 0$ , and  $\int_1^\infty \frac{du}{f(u)} < \infty$ . Define  $g(u) = f(u) - \lambda_1 u$  and let  $A$  be the larger root of  $g(u) = 0$  if it has a root. Suppose that  $u_0 \geq 0$  and that  $\int_\Omega \varphi(x) u_0(x) dx > A$  if  $g(u) = 0$  has a root, otherwise  $u_0 \not\equiv 0$ . Then  $u$  blows up in finite time.

*Proof.* Multiplying  $u_t = \Delta u + f(u)$  by  $\varphi_1$  and integrating over  $\Omega$ , we have

$$\begin{aligned} \frac{d}{dt} \int_\Omega u \varphi dx &= -\lambda_1 \int_\Omega u \varphi dx + \int_\Omega f(u) \varphi dx \\ &\geq -\lambda_1 \int_\Omega u \varphi dx + f\left(\int_\Omega u \varphi dx\right) \end{aligned}$$

by the Jensen inequality. Set

$$y(t) := \int_\Omega u \varphi_1 dx.$$

Then  $y(t)$  satisfies

$$y' \geq -\lambda_1 y + f(y) = g(y).$$

By assumption that  $f$  is convex and  $\int_1^\infty \frac{dy}{f(y)} < \infty$ , we obtain

$$\lim_{u \rightarrow \infty} \frac{f(u)}{u} = \infty.$$

Then there exists  $U > A$  such that for  $u \geq U$

$$f(u) > 2\lambda_1 u.$$

It follows that

$$\begin{aligned} t &\leq \int_{y(0)}^{y(t)} \frac{dy}{g(y)} \\ &\leq \int_{y(0)}^\infty \frac{dy}{g(y)} \\ &\leq \int_{y(0)}^U \frac{dy}{g(y)} + \int_U^\infty \frac{2}{f(y)} dy \\ &< \infty. \end{aligned}$$

This completes the proof. q.e.d.

**Remark 1.** The assumption  $\int_1^\infty \frac{du}{f(u)} < \infty$  is necessary. Indeed, if  $\int_1^\infty \frac{du}{f(u)} = \infty$ , then the solution of  $v_t = f(v)$ ,  $v(0) = \|u_0\|_\infty$  is a global supersolution for  $u$ .

**Remark 2.** The assumption  $\int_1^\infty \frac{du}{f(u)} < \infty$  is not sufficient for blow-up, if  $f$  is not convex. (See **Fila, Ninomiya and Vázquez [8]**).

**Remark 3.** The proof does not necessarily imply that  $y(t) \rightarrow \infty$  as  $t \rightarrow T$ .

When  $f(u) = u^p$ , there appears a critical exponent for the blowup and global existence of positive solutions.

**Theorem 3** (Fujita, 1966 [12]). Consider the equation

$$u_t = \Delta u + u^p, \quad x \in \mathbb{R}^N, \quad p > 1.$$

- (i) If  $p < 1 + \frac{2}{N}$ , all positive solutions blow up.
- (ii) If  $p > 1 + \frac{2}{N}$ , there exist both blowing up and global positive solutions.

**Hayakawa, 1973 [21], Kobayashi, Sirao, Tanaka, 1977 [25]:** The critical case  $p = 1 + \frac{2}{N}$  is as (i).

Why is  $p = 1 + \frac{2}{N}$  critical? The following is an intuitive explanation. The fundamental solution of  $u_t = \Delta u$  decays like  $t^{-\frac{N}{2}}$ , while the solution of  $u_t = u^p$  is given by  $C(T - t)^{-\frac{1}{p-1}}$ . They are balanced if

$$\frac{N}{2} = \frac{1}{p-1} \Leftrightarrow p = 1 + \frac{2}{N}.$$

Better explanation is as follows. Define  $v(y, s) = e^{\frac{s}{p-1}} u(e^{\frac{s}{2}} y, e^s - 1)$ . Then  $v$  solves

$$v_s = \Delta v + \frac{y}{2} \cdot \nabla v + \frac{1}{p-1} v + v^p, \quad y \in \mathbb{R}^N, \quad s > 0.$$

Here the principal eigenvalue of  $Lv := \Delta v + \frac{y}{2} \cdot \nabla v$  in  $L_\rho^2$  with  $\rho(y) = e^{\frac{|y|^2}{4}}$  is given by  $\lambda_1 = -\frac{N}{2}$ . This implies that the trivial solution  $v \equiv 0$  is stable if  $\frac{N}{2} > \frac{1}{p-1}$ , and is unstable if  $\frac{N}{2} < \frac{1}{p-1}$ . As a consequence of instability, the solution blows up.

Next, consider

$$\begin{aligned} u_t &= \Delta u + f(u), & x &\in \Omega, \\ u &= 0, & x &\in \partial\Omega, \\ u(x, 0) &= u_0(x), \end{aligned}$$

where  $\Omega$  is a bounded domain. For  $v \in C^1(\overline{\Omega})$ , define

$$E(v) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx - \int_{\Omega} F(v) dx, \quad F(v) = \int_0^v f(u) du.$$

**Lemma 1.** *If  $u$  is a solution for  $t \in [0, T)$ , then*

$$\frac{d}{dt}E(u(\cdot, t)) = - \int_{\Omega} u_t^2 dx \leq 0.$$

*Proof.* We calculate

$$\begin{aligned} \frac{d}{dt}E(u(\cdot, t)) &= \int_{\Omega} \nabla u \cdot \nabla u_t dx - \int_{\Omega} f(u)u_t dx \\ &= - \int_{\Omega} (\Delta u + f(u))u_t dx = - \int_{\Omega} u_t^2 dx \leq 0. \end{aligned}$$

q.e.d.

**Theorem 4** (Levine, 1973 [28]). *Let  $\Omega$  be bounded and  $f(u) = |u|^{p-1}u$ ,  $p > 1$ . If  $u_0 \in C^1(\overline{\Omega})$  satisfies  $E(u_0) < 0$ , then  $u$  blows up in finite time .*

*Proof.* Multiplying  $u_t = \Delta u + |u|^{p-1}u$  by  $u$  and integrating over  $\Omega$ , we have

$$\int_{\Omega} uu_t dx = \int_{\Omega} u \Delta u dx + \int_{\Omega} |u|^{p+1} dx,$$

and hence

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\Omega} u^2 dx &= - \int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} |u|^{p+1} dx \\ &= -2E(u) + \frac{p-1}{p+1} \int_{\Omega} |u|^{p+1} dx. \end{aligned}$$

From the Hölder inequality, it follows that

$$\int_{\Omega} |u|^{p+1} dx = \int_{\Omega} (u^2)^{\frac{p+1}{2}} dx \geq C(\Omega) \left( \int_{\Omega} u^2 dx \right)^{\frac{p+1}{2}}.$$

Define  $y(t) := \int_{\Omega} u^2 dx$ . Then the function  $y(t)$  satisfies

$$\frac{1}{2} y' > C(\Omega)^{-1} \frac{p-1}{p+1} y^{\frac{p+1}{2}}.$$

Since  $y$  cannot exist globally, the solution  $u$  cannot exist globally either.  
q.e.d.

Consider the stationary problem

$$(S) \quad \begin{cases} v_{xx} + v^p = 0, & x \in (-L, L), \quad L > 0, \quad p > 1, \\ v(\pm L) = 0. \end{cases}$$

**Theorem 5.** *For every  $L > 0$  and  $p > 1$ , there is a unique positive solution of (S).*

*Proof.* At first, we prove that  $x = 0$  is the only maximum point of a positive solution of (S). Assume  $x = x_0$  is a maximum point of a positive solution of (S),  $v(x_0) = M > 0$ . Since a positive solution of (S) is concave by  $v_{xx} < 0$  on  $(-1, 1)$ , this maximum point is unique. Multiplying  $v_{xx} + v^p = 0$  by  $v_x$ , we have

$$\begin{aligned} v_x v_{xx} + v_x v^p &= 0, \\ \left( \frac{1}{2} v_x^2 + \frac{1}{p+1} v^{p+1} \right)_x &= 0. \end{aligned}$$

Integrating this over  $[x_0, x]$ ,  $x > x_0$ , by  $v_x(x_0) = 0$ , we obtain

$$\frac{1}{2} v_x^2 + \frac{1}{p+1} v^{p+1} = \frac{1}{p+1} M^{p+1}.$$

By  $v_x \leq 0$  on  $[x_0, L]$ , we have

$$\begin{aligned} v_x^2 &= \frac{2}{p+1} (M^{p+1} - v^{p+1}), \\ v_x &= -\sqrt{\frac{2}{p+1}} \sqrt{M^{p+1} - v^{p+1}}. \end{aligned}$$

By separation of variables, it follows that for  $x_0 \leq x \leq L$

$$\begin{aligned} \frac{dv}{\sqrt{M^{p+1} - v^{p+1}}} &= -\sqrt{\frac{2}{p+1}} dx, \\ \int_{v(x)}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} &= \sqrt{\frac{2}{p+1}} (x - x_0). \end{aligned}$$

By  $v(L) = 0$ , we obtain

$$\int_0^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} = \sqrt{\frac{2}{p+1}} (L - x_0).$$

Similarly, by  $v_x \geq 0$  on  $[L, x_0]$ , we obtain

$$\int_0^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} = \sqrt{\frac{2}{p+1}} (L + x_0).$$

It follows that  $x_0 = 0$ .

Next, we prove that a positive solution of (S) is even. As above, we have

$$\begin{aligned}\int_{v(x)}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} &= \sqrt{\frac{2}{p+1}}x \text{ for } 0 \leq x \leq L, \\ \int_{v(x)}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} &= -\sqrt{\frac{2}{p+1}}x \text{ for } 0 \leq x \leq L.\end{aligned}$$

It follows that

$$\int_{v(x)}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} = \int_{v(-x)}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}}.$$

Since

$$\int_{\alpha}^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}}$$

is monotone decreasing on  $[0, M]$  with  $\alpha$ , it follows that  $v(x) = v(-x)$ , and so a positive solution of (S) is even.

Finally, as a solution of (S) is even, it is sufficient to prove that there is a unique positive solution of

$$(SR) \quad \begin{cases} v_{xx} + v^p = 0, & x \in (0, L), \\ v_x(0) = v(L) = 0. \end{cases}$$

A positive solution of (SR) can be obtained as a solution

$$(SR)' \quad \begin{cases} v_{xx} + v^p = 0, & x > 0, \\ v_x(0) = 0, & v(0) = M > 0. \end{cases}$$

Define

$$\varphi(M) := \int_0^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}}.$$

Then it is sufficient to show for every  $L > 0$  that  $\varphi(M) = \sqrt{\frac{2}{p+1}}L$  has unique positive solution. We calculate  $\varphi(M)$  as

$$\varphi(M) = \int_0^M \frac{dw}{\sqrt{M^{p+1} - w^{p+1}}} = M^{-\frac{p+1}{2}} \int_0^M \frac{dw}{\sqrt{1 - (\frac{w}{M})^{p+1}}},$$

and we change the variable  $\frac{w}{M} = y$ ,  $dw = Mdy$  to obtain

$$\varphi(M) = M^{-\frac{p+1}{2}} \int_0^1 \frac{dy}{\sqrt{1 - y^{p+1}}},$$

where  $\int_0^1 \frac{dy}{\sqrt{1-y^{p+1}}}$  is a constant dependent on  $p$ . Define  $C_p := \int_0^1 \frac{dy}{\sqrt{1-y^{p+1}}}$ , then it follows that

$$\varphi(M) = C_p M^{-\frac{p-1}{2}}.$$

Since  $\varphi$  is monotone decreasing and the range of  $\varphi$  is  $(0, \infty)$ ,  $\varphi(M) = \sqrt{\frac{2}{p+1}}L$  has a unique positive solution for every  $L > 0$ .

q.e.d.

**Theorem 6.** *Consider the problem*

$$(P) \quad \begin{cases} u_t = u_{xx} + u^p, & |x| < 1, \ p > 1, \\ u = 0, & |x| = 1, \\ u(x, 0) = u_0(x), & |x| < 1. \end{cases}$$

*If  $u_0(x) \geq v(x)$ ,  $u_0 \not\equiv v(x)$ ,  $v$  is the solution of (S), then  $u$  blows up in finite time.*

*Proof.* By the maximum principle, if  $\tau > 0$ , then  $u(x, \tau) > v(x)$  for  $|x| < 1$ . Moreover  $u_x(1, \tau) < v_x(1)$  and  $u_x(-1, \tau) > v_x(-1)$ . Hence there is  $k > 1$  such that  $u(x, \tau) > kv(x)$ .

Thus, it is sufficient to consider the case where  $u_0(x) = kv(x)$ ,  $k > 1$ . In this case, we have

$$\begin{aligned} (u_t)_t &= (u_t)_{xx} + pu^{p-2}u_t, & |x| < 1, \\ u_t &= 0, & |x| = 1, \\ u_t(x, 0) &= kv_{xx} + k^p v^p > k(v_{xx} + v^p) = 0. \end{aligned}$$

By the maximum principle, we obtain  $u_t \geq 0$  as long as  $u$  exists.

Suppose  $u$  is global. If  $u$  is bounded, there exists a solution  $w(x)$  of (S) such that

$$u(x, t) \rightarrow w(x) \text{ as } t \rightarrow \infty, \ w(x) > v(x), \ |x| < 1,$$

this is a contradiction with Theorem 5.

If  $u$  is unbounded, then

$$u(x, t) \equiv u(-x, t), \ u_x(x, t) < 0 \text{ for } x \in [0, 1),$$

and either

$$(i) \quad \lim_{t \rightarrow \infty} u(x, t) < \infty, \ x \in (a, 1], \ a \in (0, 1),$$

or

$$(ii) \quad \lim_{t \rightarrow \infty} u(x, t) = \infty, \ |x| < 1,$$

occurs. In case (i), there exists  $w$  such that

$$\lim_{t \rightarrow \infty} u(x, t) = w(x), \quad x \in (a, 1],$$

$$(S^*) \begin{cases} w_{xx} + w^p = 0, \\ w(a) = \infty, \quad w(1) = 0, \end{cases}$$

but  $(S^*)$  does not have a solution, because  $w$  is concave. In case (ii), there is  $t_0 > 0$  such that

$$\int_{-1}^1 u(x, t_0) \varphi(x) dx > \lambda_1^{\frac{1}{p-1}} \quad (= A \text{ in Theorem 2}),$$

this is a contradiction with Theorem 2.

q.e.d.

## 2 Blow-up set

**Definition 1.** Let  $u$  blow up in finite time  $T > 0$ . Then  $x_0 \in \Omega$  is a **blow-up point** if  $u(x_n, t_n) \rightarrow \infty$  for some  $\{(x_n, t_n)\}_{n=1}^\infty \subset \Omega \times (0, T)$  such that  $(x_n, t_n) \rightarrow (x_0, T)$  as  $n \rightarrow \infty$ . **The blow-up set**  $B$  is the set of all blow-up points.

**Weissler, 1984 [37]:** If  $u$  is the solution of (P) with  $u_0(x) = kv(x)$ , where  $v(x)$  is the unique stationary solution and  $k > 1$ , then  $B = \{0\}$ .

**Theorem 7** (Friedman and McLeod, 1985 [11]). Consider the problem

$$\begin{aligned} u_t &= u_{rr} + \frac{N-1}{r} u_r + u^p, & (r, t) &\in (0, R) \times (0, \infty), \quad p > 1, \\ u(R, t) &= u_r(0, t) = 0, & t &\in (0, \infty), \\ u(r, 0) &= u_0(r), & r &\in [0, R]. \end{aligned}$$

Let  $u_0 \in C^2([0, R])$ ,  $(u_0)_r < 0$  for  $r \in (0, R]$ ,  $(u_0)_{rr} < 0$ . Then for any  $q \in (1, p)$ , there is  $K = K(u_0, p, q, R)$  such that

$$u(r, t) \leq K r^{-\frac{2}{q-1}} \quad \text{for } t \in (0, T).$$

(If  $T < \infty$ , then  $B = \{0\}$ .)

*Proof.* Define

$$J(r, t) = r^{N-1} u_r + c(r) F(u).$$

Assuming

$$\sup_{r>0} \frac{c(r)}{r^N} < \infty,$$

we claim that

$$\begin{aligned}
J_t + \frac{N-1}{r}J_r - J_{rr} + b(r, t)J &= -c(r)A(r, t), \\
b(r, t) &\text{ bounded for } 0 \leq r \leq R, \\
A(r, t) &= pu^{p-1}F - u^pF' + \frac{c^2}{r^{2N-2}}F'F^2 - \frac{2c'}{r^{N-1}}F'F + \frac{2(N-1)}{r^N}cF'F \\
&\quad + \left(c'' - \frac{N-1}{r}c'\right)\frac{F}{c},
\end{aligned}$$

and if we choose  $c(r) = \epsilon r^{N+\delta}$ ,  $\epsilon, \delta > 0$ ,  $F(u) = u^\gamma$ ,  $1 < \gamma < p$ , then for any  $\gamma \in (1, p)$ ,  $\delta > 0$ , there is  $\epsilon > 0$  such that  $A > 0$ .

*Proof of claim.* We compute

$$\begin{aligned}
&J_t + \frac{N-1}{r}J_r - J_{rr} + c(r)A(r, t) \\
&= \left\{ pu^{p-1} + 2(N-1)cF'r^{-N} - \frac{c^3}{r^{2N-2}}(r^{N-1}u_r + cF) \right\} J.
\end{aligned}$$

The function

$$b(r, t) = pu^{p-1} + 2(N-1)cF'r^{-N} - (r^{N-1}u_r + cF)c^3/r^{2N-2}$$

is bounded for  $0 < r < R$  by the assumption that  $\sup_{r>0} c(r)/r^N < \infty$ . We choose  $c(r) = \epsilon r^{N+\delta}$ ,  $0 < \epsilon < 1$ ,  $\delta > 0$ ,  $F(u) = u^\gamma$ ,  $1 < \gamma < p$ . Then we have

$$\begin{aligned}
A(r, t) &= u^\gamma \{ (p - \gamma)u^{p-1} - 2\epsilon(\delta + 1)\gamma r^\delta u^{\gamma-1} \\
&\quad + \epsilon^2\gamma(\gamma - 1)r^{2\delta+2}u^{3\gamma-2} + (N + \delta)\delta r^{-2} \}.
\end{aligned}$$

By  $p > \gamma$ , there exists  $M > 0$  independent of  $\epsilon$  such that if  $u > M$ , then  $A > 0$ . With  $u(R, t) = 0$ , there is  $R_0 > 0$  independent of  $\epsilon$  such that for  $R_0 < r < R$ , we have

$$(p - \gamma)u^{p-1} - 2\epsilon(\delta + 1)\gamma r^\delta u^{\gamma-1} + \epsilon^2\gamma(\gamma - 1)r^{2\delta+2}u^{3\gamma-2} + (N + \delta)\delta r^{-2} > 0,$$

and  $A(r, t) > 0$  for  $R_0 < r < R$ . There is  $m > 0$  such that  $u(r, t) > m$  for  $0 < r < R_0$ . If  $u \leq M$ ,  $0 < r < R_0$ , then we obtain

$$\begin{aligned}
A(r, t) &\geq u^\gamma \{ (p - \gamma)m^{p-1} - 2\epsilon(\delta + 1)\gamma r^\delta M^{\gamma-1} \\
&\quad + \epsilon^2\gamma(\gamma - 1)r^{2\delta+2}m^{3\gamma-2} + (N + \delta)\delta r^{-2} \}.
\end{aligned}$$

If  $\epsilon$  is small, then it follows that  $A(r, t) > 0$ . Consequently, the proof is complete. q.e.d.

Moreover, we have

$$J(0, t) = 0, \quad J(R, t) \leq 0.$$

If  $\epsilon$  is small, we obtain  $J(r, 0) \leq 0$ . By the maximum principle, it follows that  $J \leq 0$  for  $(r, t) \in [0, R] \times [0, T)$ . Then we have

$$u_r \leq -\epsilon r^{1+\delta} u^\gamma, \quad \frac{1}{1-\gamma} (u^{1-\gamma})_r \leq -\epsilon r^{1+\delta}.$$

Integrating this over  $[0, r]$ , we obtain

$$u^{1-\gamma}(r, t) \geq u^{1-\gamma}(r, 0) - u^{1-\gamma}(0, t) \geq \epsilon(\gamma - 1) \frac{1}{2+\gamma} r^{2+\gamma},$$

$$u(r, t) \leq K r^{-\frac{2+\delta}{\gamma-1}},$$

and so we choose  $\gamma, \delta$  such that  $\frac{2+\delta}{\gamma-1} = \frac{2}{q-1}$ . q.e.d.

**Remark 4.** *Theorem 7 holds for  $f(u)$  more general than  $u^p$ . It is essential that  $f(u) \geq C u \log^q u$ ,  $u \geq U > 1$ ,  $q > 2$ .*

**Remark 5.** *If  $f(u) = (u + a) \log^q(u + a)$ ,  $a > 1$ , then blow-up occurs if and only if  $q > 1$ . This follows from Theorem 2 and Remark 1.*

**Lacey, 1986 [25]:** Consider the problem

$$\begin{aligned} u_t &= \Delta u + f(u), \quad x \in \Omega - \text{bounded}, \\ u &= 0, \quad x \in \partial\Omega. \end{aligned}$$

- (i) If  $\Omega$  is large and  $f(u) = (u + 2) \log^2(u + 2)$ , then there is  $\Omega_0 \subset \Omega$ ,  $|\Omega_0| > 0$ ,  $\overline{\Omega_0} \subset \Omega$ , such that  $B = \{\overline{\Omega_0}\}$  (regional blow-up), if  $\Omega$  is small, then  $B = \overline{\Omega}$ .
- (ii) If  $f(u) = (u + 2) \log^q(u + 2)$ ,  $1 < q < 2$ , then  $B = \overline{\Omega}$ .

**Chen and Matano, 1989 [5]:** Consider the problem

$$\begin{aligned} u_t &= u_{xx} + f(u), \quad |x| < 1, \\ u &= 0, \quad |x| = 1, \end{aligned}$$

under some assumption on  $f$ . Assume  $T < \infty$ , then  $B$  consists of finite number of points. This number is smaller than or equal to the number of local maxima of  $u_0$ .

**Merle, 1992 [30]:** Consider the problem

$$\begin{aligned} u_t &= u_{xx} + u^p, & |x| < 1, \quad p > 1, \\ u &= 0, & |x| = 1, \\ u(x, 0) &= u_0(x), & |x| \leq 1. \end{aligned}$$

Given any positive integer  $k$  and  $-1 < x_1 < \dots < x_k < 1$ , there is  $u_0$  such that  $u$  blows up at  $t = T < \infty$  and  $B = \{x_1, \dots, x_k\}$ .

**Giga and Kohn, 1989 [18]:** Consider the problem

$$\begin{aligned} u_t &= \Delta u + u^p, & x \in \Omega, \\ u &= 0, & x \in \partial\Omega, \\ u(x, 0) &= u_0(|x|) \geq 0, & x \in \Omega, \end{aligned}$$

where  $\Omega = \{|x| < R\} \subset \mathbb{R}^N$ ,  $p > 1$  and  $(N - 2)p < N + 2$ . Then there is  $u_0$  such that  $u$  blows up on a sphere  $B = \{|x| = R_0\}$ ,  $0 < R_0 < R$ .

**Velázquez, 1993 [36]:** Consider the equation

$$u_t = \Delta u + u^p, \quad x \in \mathbb{R}^N, \quad p > 1, \quad (N - 2)p < N + 2.$$

Since  $(N - 1)$ -dimensional Hausdorff measure is finite in every compact sets, the dimension of  $B$  is at most  $(N - 1)$  in every compact sets if  $u \not\equiv C(T - t)^{-\frac{1}{p-1}}$ .

### 3 Blow-up rate

In this section, we consider the blow-up rate of solutions for

$$u_t = \Delta u + |u|^{p-1}u.$$

Define  $M(t) := \max_{\overline{\Omega}} u(\cdot, t)$ , then

$$M'(t) \leq M^p(t), \quad M(t) \geq ((p - 1)(T - t))^{-\frac{1}{p-1}}.$$

**Definition 2.** *Blow-up is of **type** I if  $(T - t)^{-\frac{1}{p-1}} \|u(\cdot, t)\|_\infty \leq C$ , otherwise blow-up is of **type** II.*

**Weissler, 1985 [38]:** Assume the condition

$$(N - 2)p < N + 2, \quad \Omega = \{|x| < R\} \subset \mathbb{R}^N, \quad u - \text{radial}, \quad u_r, \quad u_{rr} \leq 0, \quad u_t \geq 0.$$

Then blow-up is type I.

**Friedman and McLeod, 1985 [11]:** Blow-up is of type I, if  $\Omega$  is bounded,  $u, u_t \geq 0$  ( $u_0 \geq 0, \Delta u_0 + u_0^p \geq 0$  in  $\Omega$ ).

*Proof.* Define  $J(r, t) = u_t - \epsilon u^p$  for  $\epsilon > 0$ . Then we have

$$\begin{aligned} J_t - \Delta J - p u^{p-1} J &= \epsilon p(p-1) u^{p-2} |\nabla u|^2 \geq 0, \\ J &= 0 \text{ on } \partial\Omega, \quad J \geq 0 \text{ if } t = 0. \end{aligned}$$

By the maximum principle, we obtain  $J \geq 0$  in  $\Omega \times (0, T)$ . It follows that

$$u_t \geq \epsilon u^p,$$

and so we obtain

$$u(x, t) \leq ((p-1)\epsilon(T-t))^{-\frac{1}{p-1}}.$$

q.e.d.

The results on type I blow-up are listed as follows:

- **Galaktionov and Posashkov, 1986 [13]:**  $N=1$ .
- **Giga and Kohn, 1985 [16], 1989, [18]:**  $\Omega = \mathbb{R}^N$  or  $\Omega$  is bounded, convex and either (i)  $u \geq 0, (N-2)p < N+2$  or (ii)  $(3N-4)p < 3N+8$ .
- **Giga, Matsui and Sasayama, 2004 [19], [20]:**  $\Omega = \mathbb{R}^N$  or  $\Omega$  is bounded, convex,  $(N-2)p < N+2$ .
- **Filippas, Herrero and Velázquez, 2000 [10]:**  $\Omega = \{|x| < R\}$ ,  $u$  is radial,  $u_r \geq 0$ ,  $p = \frac{N+2}{N-2}$ .
- **Matano and Merle, 2004 [29]:**  $\Omega = \{|x| < R\}$ ,  $u$  is radial,  $p > \frac{N+2}{N-2}$ ,  $N > 2$ ,  $p < p_{JL} = 1 + \frac{4}{N-4-2\sqrt{N-1}}$  if  $N > 10$ .

The results on type II blowup are listed as follows:

- **Herrero and Velázquez, 1993 [22]:** There are solutions with type II blow-up if  $p > p_{JL}$ .
- **Mizoguchi, 2004 [32]:** The same result as [22] with a similar but shorter proof.

In order to explain why type II blow-up occurs, we define

$$w(y, s) = (T - t)^{\frac{1}{p-1}} u(x, t), \quad y = \frac{x}{\sqrt{T-t}}, \quad s = -\log(T-t).$$

Then  $w$  satisfies

$$(W) \quad w_s = \Delta w - \frac{y}{2} \nabla w + w^p - \frac{1}{p-1} w, \quad y \in \mathbb{R}^N, \quad s > -\log T.$$

There is a solution  $w$  such that

$$w(y, s) \rightarrow w^*(y) \text{ as } s \rightarrow \infty,$$

where  $w^*(y)$  is the singular stationary solution  $C|y|^{-\frac{2}{p-1}}$ .

## 4 Blow-up profile

In this section, we consider behavior at the blow-up time of solution for

$$u_t = \Delta u + u^p, \quad p > 1.$$

We note that (W) has a stationary solution  $k = (p-1)^{-\frac{1}{p-1}}$ .

**Giga and Kohn, 1987 [17]:** The constant  $k$  is the only positive bounded stationary solution of (W). Moreover  $w(y, s) \rightarrow k$  as  $s \rightarrow \infty$  uniformly for  $|y| \leq C$ ,  $C > 0$ .

**Herrero and Velázquez, 1992 [23]:** Consider the problem

$$\begin{aligned} u_t &= u_{xx} + u^p, \quad x \in \mathbb{R}, \\ u(x, 0) &= u_0(x) \geq 0, \end{aligned}$$

where  $u_0$  is continuous, bounded,  $u_0 \not\equiv C$  and such that  $u$  blows up at  $T < \infty$  with the blowup set  $B = \{0\}$ . Then, one of the following alternatives holds:

$$(a) \quad \lim_{x \rightarrow 0} \left( \frac{|x|^2}{|\log|x||} \right)^{\frac{1}{p-1}} u(x, T) = \left( \frac{8p}{(p-1)^2} \right)^{\frac{1}{p-1}},$$

(b) There exist a constant  $C > 0$  and even integer  $m \geq 4$  such that

$$\lim_{x \rightarrow 0} |x|^{\frac{m}{p-1}} u(x, T) = C.$$

The case (a) occurs if  $u_0$  has a single maximum. The case (b) with  $m = 4$  occurs if  $u_0$  has two local maxima for  $t < T$  which merge at  $T = t$ .

**Bricmont and Kupiainen, 1994 [3]:** All profiles in (b) do occur.

## 5 Complete blow-up

In this section, we consider

$$(P)' \quad \begin{cases} u_t = \Delta u + u^p, & x \in \Omega : \text{bounded}, p > 1, \\ u = 0, & x \in \partial\Omega, \\ u(\cdot, 0) = u_0 \geq 0, u_0 \in L^\infty(\Omega). \end{cases}$$

**Definition 3** (Baras and Cohen, 1987 [2]). *Let  $f_n(u) = \min\{u^p, n\}$ , let  $u_n$  be the solution of  $(P)'$  with  $f_n(u)$  instead of  $u^p$ , let  $u$  be a solution of  $(P)'$  which blows up at  $t = T < \infty$ . Then blow-up is **complete** if*

$$\lim_{n \rightarrow \infty} u_n(x, t) = \infty \text{ for } (x, t) \in \Omega \times (T, \infty).$$

**Baras and Cohen, 1987 [2]:** Let  $u_0$  be such that  $u$  blows up. If  $(N-2)p < N+2$  or if  $u_0 \in C^2$ ,  $\Delta u_0 + u_0^p \geq 0$ , then blow-up is complete.

## 6 Continuation after blow-up

**Definition 4.** *The function  $u$  is a **weak solution** of  $(P)'$  on  $[0, T]$  if*

(a)  $u \in C([0, T]; L^1(\Omega)),$

(b)  $u^p \in L^1(\Omega \times (0, T)),$

(c) 
$$\int_{\Omega} u(x, t_2) \psi(x, t_2) dx - \int_{\Omega} u(x, t_1) \psi(x, t_1) dx - \int_{t_1}^{t_2} \int_{\Omega} u \psi_t dx ds$$

$$= \int_{t_1}^{t_2} \int_{\Omega} (u \Delta \psi + u^p \psi) dx ds$$

for every  $\psi \in C^2(\bar{\Omega} \times [0, T])$  with  $\psi = 0$  on  $\partial\Omega$ ,  $0 \leq t_1 < t_2 \leq T$ .

*The function  $u$  is a **global weak solution** if it is a weak solution on  $[0, T]$  for every  $T > 0$ . If  $\Omega = \mathbb{R}^N$ , then we require that (c) holds for every  $\psi \in C^2(\mathbb{R}^N \times [0, T])$  with compact support and (a), (b) hold for  $L^1_{loc}$ .*

**Theorem 8** (Ni, Sacks and Tavantzis, 1984 [34]). *Let  $p \geq \frac{N+2}{N-2}$ ,  $N > 2$ , and  $\Omega$  be convex. Then for every  $f \in L^\infty(\Omega)$  with  $f \geq 0$ ,  $f \not\equiv 0$ , there is  $\lambda^* > 0$  such that the problem  $(P)'$  with  $u_0 = \lambda^* f$  has an unbounded global weak solution.*

*Idea of Proof.* Define

$$\lambda^* = \sup \{ \lambda > 0; \text{ the solution } u \text{ of } (P)' \text{ with } u_0 = \lambda f \\ \text{ is a global classical solution, } u(\cdot, t) \rightarrow 0 \text{ as } t \rightarrow \infty \}.$$

Then  $0 < \lambda^* < \infty$ . The function  $u(x, t; \lambda^* f)$  is not a global classical solution converging to zero. If  $\Omega$  is convex, there is no positive stationary solution, and hence  $u(x, t; \lambda^* f)$  cannot be uniformly bounded. Because it is a monotone limit of global classical solutions converging to zero,  $u(x, t; \lambda^* f)$  must be a global weak solution. q.e.d.

**Galaktionov and Vázquez, 1997 [15]:** Let  $p \geq \frac{N+2}{N-2}$ ,  $N > 2$ , and let  $p < 1 + \frac{6}{N-10}$  if  $N > 10$ . Suppose that  $\Omega = \{|x| < R\}$  and  $u_0$  is radial. Then the solution  $u$  from Theorem 8 blows up in finite time.

**Mizoguchi, 2005 [31]:** The restriction  $p < 1 + \frac{6}{N-10}$  if  $N > 10$  can be removed.

Consider the problem

$$\begin{aligned} u_t &= u_{xx} + a(x, t)u_x + b(x, t)u, \\ u(x_0, t_0) &= 0, \quad u_x(x_0, t_0) = 0, \quad u_{xx}(x_0, t_0) = 0, \end{aligned}$$

where  $a(x, t)$  and  $b(x, t)$  are bounded. Define

$$z(u(\cdot, t)) = \# \{ r; u(r, t) = 0 \}.$$

With the fact that  $z(\cdot, t)$  is nonincreasing in time (Sturm, 1836) and  $z(\cdot, t)$  decreases when a multiple zero occurs (Angenent, 1986 [1]), we obtain the following results.

**Fila, Matano and Poláčik, 2005 [7]:** Let  $p \geq \frac{N+2}{N-2}$ ,  $N > 2$ , and let

$$p < p_{JL} = 1 + \frac{4}{N - 4 - 2\sqrt{N - 1}}$$

if  $N > 10$ ,  $\Omega = \{|x| < R\}$  Suppose that  $u_0$  is radial and such that  $u$  blows up but it is a minimal global weak solution. Then there are  $t_1 = T < t_2 < t_3 < \dots < t_k < \infty$  such that  $\|u(\cdot, t_j)\| = \infty$  and  $\|u(\cdot, t_j)\| < \infty$  if  $t \neq t_j$ ,  $j = 1, \dots, k$ .

**Mizoguchi, 2005 [31]:** If  $N > 10$ ,  $p > p_{JL}$ ,  $\Omega = \mathbb{R}^N$ , then there is  $u_0$  such that  $k = 2$ .

**Example 1.** The function  $u(r, t) = \{\varphi(t) + r^2\}^{-1}$  satisfies

$$u_t = u_{rr} + \frac{N-1}{r}u_r + g(r, t)u^2,$$

with

$$g(r, t) = 2N - \varphi'(t) - \frac{8r^2}{\varphi(t) + r^2}.$$

Assume  $\varphi(t) \geq 0$ . If  $\varphi(T) = 0$ , then  $u$  blows up. If  $\varphi(t) = 1 - e^{-(t-1)^2}$ , then  $u$  blows up at  $t = 1$ . If  $N > 4$ , then there are  $0 < c_1 < c_2 < \infty$  such that  $c_1 \leq g \leq c_2$ . The function  $u$  is a global weak solution if and only if  $N > 4$ .

*Proof.* At first, we prove that there are  $0 < c_1 < c_2 < \infty$  such that  $c_1 \leq g \leq c_2$ , if  $N > 4$ . To consider maximum and minimum of  $g(\cdot, t)$ , we calculate

$$g(r, t) = 2N - (t-1)e^{-(t-1)^2} + \frac{8r^2}{1 - e^{-(t-1)^2} + r^2},$$

$$2N - 8 - 2(t-1)e^{-(t-1)^2} \leq \min_r g \leq \max_r g \leq 2N - 2(t-1)e^{-(t-1)^2}.$$

We consider  $f(t) = (t-1)e^{-(t-1)^2}$ , we calculate

$$f'(t) = \{1 - 2(t-1)^2\}e^{-(t-1)^2}$$

and

$$f(0) = -e^{-1}, \quad f(1 - \frac{1}{\sqrt{2}}) = -\frac{1}{\sqrt{2}}e^{-\frac{1}{2}}, \quad f(1 + \frac{1}{\sqrt{2}}) = \frac{1}{\sqrt{2}}e^{-\frac{1}{2}}, \quad \lim_{t \rightarrow \infty} f(t) = 0.$$

By this calculation, we have

$$-\frac{1}{\sqrt{2}}e^{-\frac{1}{2}} \leq f(t) \leq \frac{1}{\sqrt{2}}e^{-\frac{1}{2}}.$$

Consequently, we obtain

$$0 < 2N - 8 - \sqrt{2}e^{-\frac{1}{2}} \leq g \leq 2N + \sqrt{2}e^{-\frac{1}{2}} < \infty.$$

Next, we find a necessary and sufficient condition on  $N$  such that  $u$  is a global weak solution. Assume  $N > 4$ . Fix  $T > 0$ . By  $0 \leq u(r, t) \leq 1/r^2$ , the Lebesgue convergence theorem and  $c_1 \leq g \leq c_2$ , we have

$$u \in C([0, T]; L_{loc}^1(\mathbb{R}^N)), \quad g(r, t)u \in L_{loc}^1(\mathbb{R}^N \times (0, T)).$$

We take  $\psi(x, t) \in C_0^2(\mathbb{R}^N \times [0, T])$ ,  $0 \leq t_1 < t_2 \leq T$  and  $R > 0$ . Then we have

$$\begin{aligned} & \int_{\mathbb{R}^N} u(x, t_2)\psi(x, t_2)dx - \int_{\mathbb{R}^N} u(x, t_1)\psi(x, t_1)dx - \int_{t_1}^{t_2} \int_{\mathbb{R}^N} u\psi_t dx ds \\ &= \int_{|x| \leq R} u(x, t_2)\psi(x, t_2)dx - \int_{|x| \leq R} u(x, t_1)\psi(x, t_1)dx + \int_{t_1}^{t_2} \int_{|x| > R} u_t \psi dx ds \\ & \quad - \int_{t_1}^{t_2} \int_{|x| \leq R} u\psi_t dx ds. \end{aligned}$$

Moreover, we have

$$\begin{aligned}
\int_{t_1}^{t_2} \int_{|x|>R} u_t \psi dx ds &= \int_{t_1}^{t_2} \int_{|x|>R} (\Delta u \psi + g(r, t) u \psi) dx ds \\
&= \int_{t_1}^{t_2} \int_{|x|>R} (u \Delta \psi + g(r, t) u \psi) dx ds - \int_{t_1}^{t_2} \int_{|x|=R} \frac{\partial u}{\partial r} \psi dS dt \\
&\quad + \int_{t_1}^{t_2} \int_{|x|=R} u \frac{\partial \psi}{\partial r} dS dt.
\end{aligned}$$

With  $0 \leq u \leq 1/r^2$ , we obtain

$$\begin{aligned}
\left| \int_{|x| \leq R} u(x, t_2) \psi(x, t_2) dx \right|, \quad \left| \int_{|x| \leq R} u(x, t_1) \psi(x, t_1) dx \right| \\
\leq \int_{|x| \leq R} \frac{1}{r^2} \|\psi\|_\infty dx = C \|\psi\|_\infty R^{N-2} \rightarrow 0 \text{ as } R \rightarrow 0,
\end{aligned}$$

$$\begin{aligned}
\left| \int_{t_1}^{t_2} \int_{|x| \leq R} u \psi_t dx ds \right| &\leq \int_{t_1}^{t_2} \int_{|x| \leq R} \frac{1}{r^2} \|\psi_t\|_\infty dx ds \\
&= C(t_2 - t_1) \|\psi_t\|_\infty R^{N-2} \rightarrow 0 \text{ as } R \rightarrow 0,
\end{aligned}$$

and by  $0 \geq \frac{\partial u}{\partial r} \geq -\frac{2}{r^3}$ , we obtain

$$\begin{aligned}
\left| \int_{t_1}^{t_2} \int_{|x|=R} \frac{\partial u}{\partial r} \psi dS dt \right| &\leq \int_{t_1}^{t_2} \int_{|x|=R} \frac{2}{r^3} \|\psi\|_\infty dS dt \\
&= C(t_2 - t_1) \|\psi\|_\infty R^{N-4} \rightarrow 0 \text{ as } R \rightarrow 0, \\
\left| \int_{t_1}^{t_2} \int_{|x|=R} u \frac{\partial \psi}{\partial r} dS dt \right| &\leq \int_{t_1}^{t_2} \int_{|x|=R} \frac{1}{r^2} \|\nabla \psi\|_\infty dS dt \\
&= C(t_2 - t_1) \|\nabla \psi\|_\infty R^{N-3} \rightarrow 0 \text{ as } R \rightarrow 0.
\end{aligned}$$

By this calculation, it follows that as  $R \rightarrow \infty$

$$\begin{aligned}
\int_{\mathbb{R}^N} u(x, t_2) \psi(x, t_2) dx - \int_{\mathbb{R}^N} u(x, t_1) \psi(x, t_1) dx - \int_{t_1}^{t_2} \int_{\mathbb{R}^N} u \psi_t dx ds \\
= \int_{t_1}^{t_2} \int_{\mathbb{R}^N} (u \Delta \psi + g(r, t) u \psi).
\end{aligned}$$

This means that  $u$  is a weak solution on  $[0, T]$ . Because  $0 < T < \infty$  is arbitrary,  $u$  is a global weak solution.

Finally, assume  $u$  is a global weak solution. This needs  $u \in C([0, T]; L_{loc}^1(\mathbb{R}^N))$  and  $g(r, t)u \in L_{loc}^1(\mathbb{R}^N \times (0, T))$  for  $0 < T < \infty$ . By  $u \in C([0, T]; L_{loc}^1(\mathbb{R}^N))$

for  $0 < T < \infty$ , we obtain  $u(x, 1) = \frac{1}{r^2} \in L_{loc}^1(\mathbb{R}^N)$ . Then we need  $N > 2$ . We consider  $g(r, t)u \in L_{loc}^1(\mathbb{R}^N \times (0, T))$  for  $0 < T < \infty$ . With  $1 - e^{-t} \leq t$ , we have

$$g(r, t) \geq 2N - 2(t - 1)^2 - \frac{8r^2}{(t - 1)^2 + r^2}.$$

If  $r \leq C|t - 1|$ ,  $|t - 1| \leq C$ , we have

$$g(r, t) \geq 2N - 2C^2 - \frac{8C^2}{1 + C^2}.$$

If  $C$  is sufficiently small, obtain  $g(r, t) \geq N$ . By this, we obtain

$$\begin{aligned} \int_{1-C}^{1+C} \int_{r \leq C|t-1|} g(r, t) u(r, t)^2 dx dt &\geq \int_{1-C}^{1+C} \int_{r \leq C|t-1|} N \frac{1}{\{(t-1)^2 + r^2\}^2} \omega_N r^{N-1} dr dt \\ &\geq A_N \int_{1-C}^{1+C} |t-1|^{N-5} \int_{r \leq C} \frac{r^{N-1}}{(1+r^2)^2} dr dt \\ &= \infty, \text{ if } N \leq 4. \end{aligned}$$

This is a contradiction. It follows that  $N > 4$ .

Consequently, a necessary and sufficient condition on  $N$  such that  $u$  is a global weak solution is  $N > 4$ . q.e.d.

## 7 Boundedness of global solutions

In this section, we consider the boundedness of global solutions for

$$\begin{cases} u_t = \Delta u + f(u), & x \in \Omega \\ u = 0, & x \in \partial\Omega, \\ u(x, 0) = u_0(x), \quad u_0 \in L^\infty. \end{cases}$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^N$ . We address the following

Question: Does blow-up in infinite time occur?

It is easy to see that the answer is “Yes” if  $f(u) = \lambda u$ ,  $\lambda > 1$ .

A function  $f$  is said to be superlinear, if there exists  $\epsilon$ ,  $A > 0$  such that

$$uf(u) \geq (2 + \epsilon) \int_0^u f(v) dv > 0 \text{ for } \forall u \geq A.$$

What if  $f$  is superlinear? The answer is “No” if

- (i)  $N = 1$   $f$  is locally Lipschitz (Fila, 1992 [6]),
- (ii)  $N = 2, |f'(u)| \leq Ke^{|u|^q}, 0 < q < 2$  (Fila, 1992 [6]),
- (iii)  $N > 3$ ,  $f$  is locally Lipschitz,  $|f(u)| \leq C(1 + |u|^p)$ ,  $(N - 2)p < N + 2$  (Cazenave and Lions, 1984 [4]),
- (iv)  $N > 2$ ,  $f(u) = u^p$ ,  $u \geq 0$ ,  $p > \frac{N+2}{N-2}$ ,  $p < 1 + \frac{6}{N-10}$  if  $N > 10$ ,  $\Omega = B_R = \{|x| < R\}$ ,  $u$  is radial (Galaktionov and Vázquez, 1994 [14]),
- (v)  $3 \leq N \leq 9$ ,  $f(u) = e^u$ ,  $\Omega = B_R$ ,  $u$  is radial (Fila and Poláčik, 1999 [9]),
- (vi)  $N > 2$ ,  $f \in C^1$ ,  $f \geq 0$ ,  $\lim_{u \rightarrow \infty} \frac{f(u)}{u^p} = 1$ ,  $p > \frac{N+2}{N-2}$ ,  $\Omega = B_R$ ,  $u$  is radial,  $p < p_{JL}$  if  $N > 10$  (Chen, Fila and Guo).

The answer is “Yes”, if

- (vii)  $f(u) = u^p$ ,  $N > 2$ ,  $p = \frac{N+2}{N-2}$ ,  $\Omega = B_R$ ,  $u$  is radial, (Ni, Sacks and Tavantzis, 1984 [34], Galaktionov and Vázquez, 1994 [14]),
- (viii)  $f(u) = 2(N - 2)e^u$ ,  $N > 9$ ,  $\Omega = B_1$  (Lacey and Tzanetis, 1987 [27]).

### Explanations

(vii) Define

$$\lambda^* := \sup\{\lambda > 0; u(\cdot, t; \lambda\varphi) \rightarrow 0 \text{ as } t \rightarrow \infty\}$$

for  $\varphi \in L^\infty(\Omega)$ ,  $\varphi \geq 0$ ,  $\varphi \not\equiv 0$ . Then  $u(\cdot, t; \lambda^*\varphi)$  is a global weak solution. Galaktionov and Vázquez, 1997 [15], showed that blow-up in finite time is complete. This implies that  $u(\cdot, t; \lambda^*\varphi)$  is a global classical solution. If  $u(\cdot, t; \lambda^*\varphi)$  is bounded,  $u(\cdot, t; \lambda^*\varphi) \rightarrow 0$  as  $t \rightarrow \infty$  by Corollary 1, but 0 is stable, which contradicts the definition of  $\lambda^*$ . It follows that  $u(\cdot, t; \lambda^*\varphi)$  is a classical unbounded solution.

### Pohožaev identity, 1965 [35] (for a ball)

If  $u$  is a nontrivial solution of

$$(S) \quad \begin{cases} u_{rr} + \frac{N-1}{r}u_r + |u|^{p-1}u = 0, & 0 < r < 1, \\ u'(0) = u(1) = 0, \end{cases}$$

then

$$\left(N - 2 - \frac{2N}{p+1}\right) \int_0^1 r^{N-1} |u|^{p+1} dr = -u_r^2(1).$$

**Corollary 1.** *If  $p \geq \frac{N+2}{N-2}$ ,  $N \geq 2$ , then there is no nontrivial solution of (S).*

*Proof for Pohožaev identity.* This equation is converted into

$$(1) \quad (r^{N-1}u_r)_r + r^{N-1}|u|^{p-1}u = 0.$$

Multiplying (1) by  $u$  and integrating it over  $(0, 1)$  by parts, we obtain

$$(2) \quad \int_0^1 r^{N-1}u_r^2 dr = \int_0^1 r^{N-1}|u|^{p+1} dr.$$

Multiplying (1) by  $ru_r$  and integrating it over  $(0, 1)$  by parts, we have

$$u_r^2(1) - \int_0^1 r^{N-1}u_r^2 dr - \int_0^1 r^N u_r u_{rr} dr + \int_0^1 r^N |u|^{p-1} u u_r dr = 0.$$

In view of  $u_{rr} + \frac{N-1}{r}u_r + |u|^{p-1}u = 0$ , this yields

$$(3) \quad u_r^2(1) - (N-2) \int_0^1 r^{N-1}u_r^2 dr + 2 \int_0^1 r^N |u|^{p-1} u u_r dr = 0.$$

Here we have

$$(4) \quad \int_0^1 r^N |u|^{p-1} u u_r dr = -\frac{N}{p+1} \int_0^1 r^{N-1} |u|^{p+1} dr.$$

By combining (2),(3),(4), we obtain

$$\left(N-2-\frac{2N}{p+1}\right) \int_0^1 r^{N-1} |u|^{p+1} dr = -u_r^2(1).$$

q.e.d.

(b) The function  $\varphi^*(r) = -2 \log r$  satisfies

$$\varphi_{rr} + \frac{N-1}{r}\varphi_r + 2(N-2)e^\varphi = 0, \quad r > 0,$$

$$\varphi(1) = 0.$$

If  $u_0 < \varphi^*$ , then a solution  $u$  is global and

$$u \rightarrow \varphi \text{ as } t \rightarrow \infty.$$

Define

$$u_\lambda(x, t) = \begin{cases} \lambda^{p-1}u(\lambda x, \lambda^2 t) & \text{if } f(u) = u^p, \\ 2\lambda + u(e^\lambda x, e^{2\lambda}) & \text{if } f(u) = e^u. \end{cases}$$

The function  $u_\lambda$  is a solution if  $u$  is a solution. In (iv) and (v), the scaling invariance is used.

**Lemma 2.** *There is  $v^*$  satisfying*

$$v_{rr} + \frac{N-1}{r}v_r + f(v) = 0, \quad v_r < 0 < v, \quad r \in (0, \epsilon), \quad \epsilon > 0,$$

$\lim_{r \rightarrow 0} r^{\frac{2}{p-1}} v^*(r) = K = K(p, N)$  is the constant for which  $Kr^{-\frac{2}{p-1}}$  satisfies

$$\varphi_{rr} + \frac{N-1}{r}\varphi_r + \varphi^p = 0, \quad p > \frac{N}{N-2}.$$

**Ni and Sacks, 1985 [33]:** Every global classical solution is radially decreasing after some  $t_0 > 0$ .

*Proof of (6).* Suppose  $u$  is a global classical solution,  $u_r \leq 0$ ,  $\|u(\cdot, t_i)\|_\infty \rightarrow \infty$  as  $t_i \rightarrow \infty$ . Set

$$M_i = u(0, t_i), \quad R_i = M_i^{-\frac{p-1}{2}}, \quad w_i(\rho, \tau) = R_i^m u(R_i \rho, t_i + R_i^2 \tau), \quad m = \frac{2}{p-1}.$$

Then  $w$  satisfies

$$w_{i,\tau} - \Delta w_i = w_i^p g_i, \quad g_i = \frac{f(M_i w_i)}{(M_i w_i)^p} \rightarrow 1 \text{ as } i \rightarrow \infty, \quad w_i(\cdot, 0) \rightarrow \varphi_1,$$

where  $\varphi_1$  satisfies

$$\begin{aligned} \varphi_{rr} + \frac{N-1}{r}\varphi_r + \varphi^p &= 0, \quad r > 0, \\ \varphi_r(0) &= 0, \quad \varphi(0) = 1. \end{aligned}$$

If  $p < p_{JL}$ , then

$$z(\varphi^* - \varphi_1) = \infty,$$

and so  $z(\varphi_* - w_i(\cdot, 0))$  can be made arbitrarily large on  $[0, \rho^*]$ , if  $\rho^*$ ,  $i$  is large. Then  $z(w_i(\cdot, 0) - v_i^*)$  is large on  $[0, \rho^*]$ , where  $v_i^*(\rho) = R_i^m v^*(R_i \rho)$ ,  $z(u(\cdot, t_i) - v^*)$  is large on  $[0, \rho^* R_i]$ . But  $z(u(\cdot, t) - v^*)$  is nonincreasing in  $t$ , it is bounded by  $z(u(\cdot, 0) - v^*)$ , this is a contradiction. q.e.d.

## References

- [1] S.B. Angenent, *The Morse-Smale property for a semilinear parabolic equation*, J. Differential Equations 62 (1986), no. 3, 427–442
- [2] P. Baras and L. Cohen, *Complete blow-up after  $T_{\max}$  for the solution of a semilinear heat equation*, J. Funct. Anal. 71 (1987), no. 1, 142–174

- [3] J. Bricmont and A. Kupiainen, *Universality in blow-up for nonlinear heat equations*, Nonlinearity 7 (1994), no. 2, 539–575
- [4] T. Cazenave and P.-L. Lions, *Solutions globales d'équations de la chaleur semi linéaires*, Comm. Partial Differential Equations 9 (1984), no. 10, 955–978
- [5] X.-Y. Chen and H. Matano, *Convergence, asymptotic periodicity, and finite-point blow-up in one-dimensional semilinear heat equations*, J. Differential Equations 78 (1989), no. 1, 160–190
- [6] M. Fila, *Boundedness of global solutions of nonlinear diffusion equations*, J. Differential Equations 98 (1992), no. 2, 226–240
- [7] M. Fila, H. Matano and P. Poláčik, *Immediate regularization after blow-up*, SIAM J. Math. Anal. 37 (2005), no. 3, 752–776
- [8] M. Fila, H. Ninomiya and J.L. Vázquez, *Dirichlet boundary conditions can prevent blow-up reaction-diffusion equations and systems*, Discrete Contin. Dyn. Syst. 14 (2006), no. 1, 63–74
- [9] M. Fila and P. Poláčik, *Global solutions of a semilinear parabolic equation*, Adv. Differential Equations 4 (1999), no. 2, 163–196
- [10] S. Filippas, M.A. Herrero and J.J.L. Velázquez, *Fast blow-up mechanisms for sign-changing solutions of a semilinear parabolic equation with critical nonlinearity*, R. Soc. Lond. Proc. Ser. A Math. Phys. Eng. Sci. 456 (2000), no. 2004, 2957–2982
- [11] A. Friedman and B. McLeod, *Blow-up of positive solutions of semilinear heat equations*, Indiana Univ. Math. J. 34 (1985), no. 2, 425–447
- [12] H. Fujita, *On the blowing up of solutions of the Cauchy problem for  $u_t = \Delta u + u^{1+\alpha}$* , J. Fac. Sci. Univ. Tokyo Sect. I 13 (1966), 109–124
- [13] V.A. Galaktionov and S.A. Posashkov, *Application of new comparison theorems to the investigation of unbounded solutions of nonlinear parabolic equations*, Differential cprime nye Uravneniya 22 (1986), no. 7, 1165–1173, 1285
- [14] V.A. Galaktionov and J.L. Vazquez, *Asymptotic behavior for an equation of superslow diffusion*, Asymptotic Anal. 8 (1994), no. 2, 145–159

- [15] V.A. Galaktionov and J.L. Vazquez, *Continuation of blowup solutions of nonlinear heat equations in several space dimensions*, Comm. Pure Appl. Math. 50 (1997), no. 1, 1–67
- [16] Y. Giga and R.V. Kohn, *Asymptotically self-similar blow-up of semilinear heat equations*, Comm. Pure Appl. Math. 38 (1985), no. 3, 297–319
- [17] Y. Giga and R.V. Kohn, *Characterizing blowup using similarity variables*, Indiana Univ. Math. J. 36 (1987), no. 1, 1–40
- [18] Y. Giga and R.V. Kohn, *Nondegeneracy of blowup for semilinear heat equations*, Comm. Pure Appl. Math. 42 (1989), no. 6, 845–884
- [19] Y. Giga, S. Matsui and S. Sasayama, *Blow up rate for semilinear heat equations with subcritical nonlinearity*, Indiana Univ. Math. J. 53 (2004), no. 2, 483–514
- [20] Y. Giga, S. Matsui and S. Sasayama, *On blow-up rate for sign-changing solutions in a convex domain*, Math. Methods Appl. Sci. 27 (2004), no. 15, 1771–1782
- [21] K. Hayakawa, *On nonexistence of global solutions of some semilinear parabolic differential equation*, Proc. Japan Acad. 49 (1973), 503–505
- [22] M.A. Herrero and J.J.L. Velázquez, *Some results on blow up for semilinear parabolic problems*, IMA Vol. Math. Appl. 47 1993
- [23] M.A. Herrero and J.J.L. Velázquez, *Blow-up of solutions of supercritical parabolic equations*, preprint
- [24] S. Kaplan, *On the growth of solutions of quasi-linear parabolic equations*, Comm. Pure Appl. Math. 16 (1963), 305–330
- [25] K. Kobayashi, T. Sirao and H. Tanaka, *On the growing up problem for semilinear heat equations*, J. Math. Soc. Japan 29 (1977), no. 3, 407–424
- [26] A.A. Lacey, *Global blow-up of a nonlinear heat equation*, Proc. Roy. Soc. Edinburgh Sect. A 104 (1986), no. 1-2, 161–167
- [27] A.A. Lacey and D. Tzanetis, *Global existence and convergence to a singular steady state for a semilinear heat equation*, Proc. Roy. Soc. Edinburgh Sect. A 105 (1987), 289–305
- [28] H.A. Levine, *Some nonexistence and instability theorems for solutions of formally parabolic equations of the form  $P(t)u_t = -Au + F(u)$* , Arch. Rational Mech. Anal. 51 (1973), 371–386

- [29] H. Matano and F. Merle, *On nonexistence of type II blowup for a supercritical nonlinear heat equation*, Comm. Pure Appl. Math. 57 (2004), no. 11, 1494–1541
- [30] F. Merle, *Solution of a nonlinear heat equation with arbitrarily given blow-up points*, Comm. Pure Appl. Math. 45 (1992), no. 3, 263–300
- [31] N. Mizoguchi, *Multiple blowup of solutions for a semilinear heat equation*, Math. Ann. 331 (2005), no. 2, 461–473
- [32] N. Mizoguchi, *Type-II blowup for a semilinear heat equation*, Adv. Differential Equations 9 (2004), no. 11-12, 1279–1316
- [33] W.-M. Ni and P.E. Sacks, *The number of peaks of positive solutions of semilinear parabolic equations*, SIAM J. Math. Anal. 16 (1985), no. 3, 460–471
- [34] W.-M. Ni, P.E. Sacks and J. Tavantzis, *On the asymptotic behavior of solutions of certain quasilinear parabolic equations*, J. Differential Equations 54 (1984), no. 1, 97–120
- [35] S.I. Pohožaev, *On the eigenfunctions of the equation  $\Delta u + \lambda f(u) = 0$* , Dokl. Akad. Nauk SSSR 165 1965 36–39
- [36] J.J.L. Velázquez, *Estimate on the  $(n-1)$ -dimensional Hausdorff measure of the blow-up set for semilinear heat equation*, Indiana Univ. Math. J. 42 (1993), no. 2, 445–476
- [37] F.B. Weissler, *Single point blow-up for a semilinear initial value problem*, J. Differential Equations 55 (1984), no. 2, 204–224
- [38] F.B. Weissler, *An  $L^\infty$  blow-up estimate for a nonlinear heat equation*, Comm. Pure Appl. Math. 38 (1985), no. 3, 291–295