

Preparation of Gallium Sulfide Nanosheets by Liquid Exfoliation and Their Application As Hydrogen Evolution Catalysts

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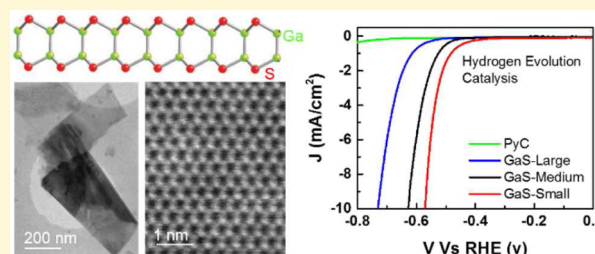
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Supporting Information

ABSTRACT: Here, we demonstrate the production of large quantities of gallium sulfide (GaS) nanosheets by liquid exfoliation of layered GaS powder. The exfoliation was achieved by sonication of the powder in suitable solvents. The variation of dispersed concentration with solvent was consistent with classical solution thermodynamics and showed successful solvents to be those with Hildebrand solubility parameters close to 21.5 MPa^{1/2}. In this way, nanosheets could be produced at concentrations of up to ~0.2 mg/mL with lateral sizes and thicknesses of 50–1000 nm and 3–80 layers, respectively. The nanosheets appeared to be relatively defect-free although oxygen was observed in the vicinity of the edges. Using controlled centrifugation techniques, it was possible to prepare dispersions containing size-selected nanosheets. Spectroscopic measurements showed the optical properties of the dispersions to vary strongly with nanosheet size, allowing the elucidation of spectroscopic metrics for in situ estimation of nanosheet size and thickness. These techniques allow the production of nanosheets with controlled sizes, which will be important for certain applications. To demonstrate this, we prepared films of GaS nanosheets of three different sizes for use as hydrogen evolution electrocatalysts. We found a clear correlation between performance and size, showing small nanosheets to be more effective. This is consistent with the catalytically active sites residing on the nanosheet edges.



INTRODUCTION

Over the past few years, liquid phase exfoliation (LPE) has become an increasingly important technique for the production of two-dimensional nanomaterials.^{1–3} This method involves the delamination of layered crystals, usually by exposure to ultrasonication^{4,5} or high shear rates,^{6–9} to form large quantities of two-dimensional nanosheets. These are then stabilized against aggregation through interaction with appropriate solvents^{1,10–20} or ionic liquids²¹ or by coating with surfactants^{4,22–32} or adsorbed polymer chains.^{33–36} The resultant nanosheets range in lateral size from <100 nm to >2000 nm, depending on the material in question.^{37,38} The nanosheet thickness tends to be broadly distributed between approximately 1 and 10 layers, with monolayer contents of tens of percent achievable.^{6,37} That nanosheets can be stably dispersed in liquids greatly facilitates further processing. For example, nanosheets can be size-selected by controlled centrifugation^{23,37–41} or chemically modified via functionalization protocols.^{42–44} In addition, the dispersions can easily be formed into thin films or mixed with other materials to form composites.^{4,6} This has enabled the use of liquid exfoliated

nanosheets in a range of applications from electrocatalysis,^{37,45,46} to composite reinforcement,^{47–51} to inkjet-printed devices.^{52,53} Additional advantages of this technique are that it is low-cost, relatively straightforward to implement, and can easily be scaled up to produce very large quantities of nanosheets.

Liquid phase exfoliation was first used to produce graphene from graphite in 2008.^{5,54} Subsequent work has seen the scale up and commercialization of this method as a graphene production technique.⁶ Possibly the greatest strength of this technique is its versatility. Liquid phase exfoliation has been used to produce nanosheets from a wide range of layered crystals beyond graphite including: boron nitride; transition metal dichalcogenides, from MoS₂ to WTe₂; MoO₃; and black phosphorus.^{2,4,5,11,23,39,40,50,55–58} However, because there are hundreds of different types of layered crystals,² this material set represents only the tip of the iceberg of materials that could be

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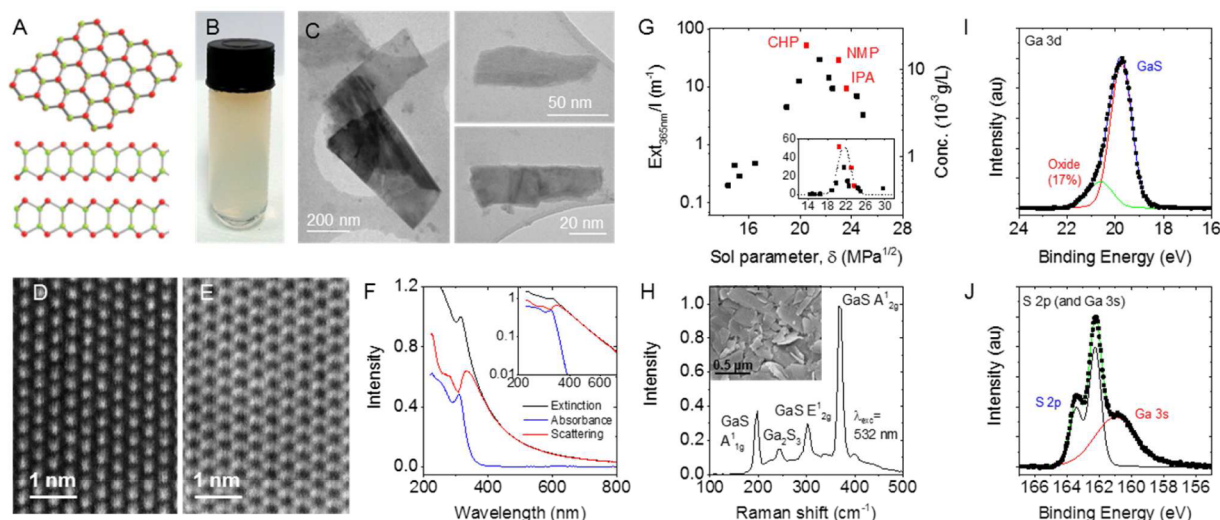


Figure 1. (A) Structure of gallium sulfide, GaS, in top and side view. Red, S; green, Ga. (B) Photograph of a typical dispersion (in 2-propanol). (C) Representative low-resolution bright field TEM images of GaS nanosheets exfoliated in IPA. (D) High angle annular dark field (HAADF) STEM image and (E) bright-field scanning transmission TEM image of GaS showing the intact lattice. (F) Optical extinction, absorbance and scattering spectra of a typical GaS nanosheet dispersion in IPA. Here intensity represents *Ext*, *Abs* or *Sca* as appropriate. (G) Dispersed concentration of GaS expressed as optical extinction at 365 nm divided by cell length plotted as a function of the Hildebrand solubility parameter δ_s of the solvent. The concentration (right axis) is estimated by the extinction coefficient at 365 nm. The dashed line in the inset is a fit to eq 1. The dispersions in panel G were prepared with initial GaS concentration $C_i = 1$ g/L and each was sonicated for $t_s = 6$ h and centrifuged at $f = 2$ krpm for $t_{CF} = 180$ min. (H) Raman spectrum (excitation wavelength 532 nm, mean of 20 spectra) of the filtered standard dispersion (IPA, $C_i = 45$ g/L, $t_s = 6$ h, $f = 2.5$ krpm, $t_{CF} = 180$ min) showing the characteristic phonons of GaS and a small impurity contribution which we associate with Ga_2S_3 . Inset: SEM image of the film. (I, J) Fitted XPS core level spectra of a filtered dispersion confirming the chemical nature of the nanosheets. (I) Ga 3d core level spectrum, (J) S 2p core level spectrum.

exfoliated by LPE. We believe it is critically important to use LPE to exfoliate yet untested layered crystals and to produce new types of nanosheets. These new materials will be of interest both for basic studies and because of their potential for use in a range of new applications.

The family of III–VI layered semiconductors is a good example of a largely untapped source of two-dimensional materials. These materials generally exist in the form MX where $M = \text{Ga}, \text{In}$ and $X = \text{S}, \text{Se}, \text{Te}$ (see Figure 1A), although other stoichiometries also exist.⁵⁹ These materials are wide bandgap semiconductors and are of interest for a range of applications from electrochemistry⁶⁰ to optoelectronics,^{61–63} gas sensing,⁶⁴ and nonlinear optics.⁶⁵ Over the last 30 years, a number of papers have studied the properties of such layered crystals with common examples being GaSe, InTe, and the related compound In_2Se_3 .^{65–67} However, work on exfoliation has only started recently with a number of papers reporting production of few-layer InSe and GaSe by mechanical cleavage or chemical exfoliation as well as vapor phase growth of thin layers.^{59,68–71} Such studies have found these exfoliated materials to be of interest for applications such as photo-detectors and in nonlinear optics.^{59,68–72}

A typical representative example of this family is gallium sulfide, GaS. Although found in a variety of structures such as nanobelts and tubes,⁷³ this material is most commonly encountered as a layered crystal and is particularly attractive due to its relatively low cost. Although a number of papers have studied the layered form of GaS,^{62,63,74} work on its exfoliation to give nanosheets is in the very early stages. Exfoliation has up till now only been achieved by micromechanical cleavage,^{61,64,75} which suffers from low throughput and can only produce material quantities suitable for fundamental studies. However, it is clear that the properties of exfoliated gallium sulfide are

interesting and differ from the bulk form.^{60,61,64,75} For example, exfoliated nanosheets of GaS have been used to fabricate sensitive photodetectors.^{61,64} In addition, it is likely that GaS is useful in applications beyond (opto)electronics. Gallium sulfide produced by atomic layer deposition has been combined with carbon nanotubes to produce high performance anodes in Li ion batteries.⁶⁰ However, for such materials to be competitive in applications such as battery electrodes, large quantities would be needed. Applications such as this highlight the need to develop a scalable method to produce nanosheets of gallium sulfide and related materials.

Herein, we report that layered III–VI semiconductors can be exfoliated in solvents by bath sonication. The resultant dispersions contain nanosheets that are of high quality and appear to be defect-free except for an increased oxygen content in the vicinity of edges. In addition, spectroscopic properties strongly vary as a function of nanosheet size allowing us to establish quantitative metrics to determine mean length and thickness spectroscopically. This facilitates the production of dispersions with well-defined nanosheet sizes and so specific properties. These were subsequently used to fabricate GaS electrodes for electrocatalysis of hydrogen production. Strong size effects are found with smaller nanosheets performing better.

RESULTS AND DISCUSSION

Evidence of Exfoliation and Basic Characterization. In this work, we study the exfoliation of layered GaS in a number of solvents. We chose GaS as a representative member of the family of layered III–VI semiconductors (structure see Figure 1A) in part due to its commercial availability. To perform the exfoliation, we added GaS powder to a variety of solvents and agitated using an ultrasonic bath for a fixed period. The

dispersions were then centrifuged to remove unexfoliated material. In most cases, we obtained pale yellow colored liquids such as the one shown in Figure 1B.

In order to confirm the exfoliation of the layered powder to 2D nanosheets, we performed low-resolution transmission electron microscopy (TEM) imaging (Figure 1C). These measurements showed all dispersions tested to contain large quantities of electron-transparent 2D nanosheets. In addition, high angle annular dark field dark (HAADF) (Figure 1D) and bright field (Figure 1E) scanning transmission electron microscopy (STEM) imaging confirmed the 2D crystal lattice to be intact.

Typically, the dispersed concentration of stable, exfoliated nanomaterials has been estimated from measurements of the optical extinction at a given wavelength^{4–6,39} (the extinction, Ext , is defined via the optical transmittance: $T = 10^{-Ext}$). The measured extinction spectrum of a typical GaS dispersion is plotted as the black curve in Figure 1F (solvent isopropanol, initial GaS concentration $c_i = 45$ g/L, sonication time $t_s = 6$ h, centrifugation time $t_{cf} = 180$ min at 2.5 krpm equivalent to 665 g). This plot shows a near monotonic increase with decreasing wavelength with no features of note, bar a small peak at 315 nm. As with dispersions of other 2D materials, such curves are relatively featureless because extinction spectra of dispersed nano-objects contain a significant contribution from scattering (Sca) in addition to the actual absorbance (Abs) of the material (NB: extinction, absorbance and scattering are related by $Ext(\lambda) = Abs(\lambda) + Sca(\lambda)$).

The absorbance and scattering components can be differentiated from the overall extinction spectra using an integrating sphere³⁷ as shown in Figure 1F. The absorbance is very low in the high wavelength regime, becoming appreciable only for $\lambda < 400$ nm, consistent with the semiconducting nature of GaS. This data does indeed show a significant scattering component, particularly in the nonresonant regime ($\lambda > 400$ nm). The shape of the spectra will be discussed in more detail below.

Using extinction spectra to estimate dispersed concentration has been complicated by the recent realization that both absorbance and scattering coefficients are generally dependent on nanosheet size.^{37,76,77} The wavelength used has to be chosen with care and requires an understanding of the size-dependent extinction coefficients as discussed below. In the case of GaS, we found the extinction coefficient at 365 nm to be relatively nanosheet-size-independent (for nanosheet sizes of < 300 nm) with a value of $\epsilon_{365nm} \approx 3654 \text{ L g}^{-1} \text{ m}^{-1}$. This allows the estimation of the dispersed nanosheet concentration, C , using $Ext_{365nm} = \epsilon_{365nm} Cl$ (l is the cell length).

To identify the most appropriate solvents for the liquid exfoliation of GaS and to investigate the exfoliation/stabilization mechanism, we have sonicated and centrifuged the powder in 15 solvents to produce nanosheets under identical processing conditions (see Materials and Methods). We measured optical extinction spectra for each dispersion, using Ext_{365nm} as a measure of the dispersed concentration. In this way, we found large variations in dispersed concentration among the solvents studied, with the highest nanosheet contents found for amide solvents such as *N*-methyl-2-pyrrolidone (NMP) and *N*-cyclohexyl-2-pyrrolidone (CHP). To understand the exfoliation/stabilization mechanism, we plot Ext_{365nm}/l versus the solvent Hildebrand solubility parameter δ_s in Figure 1G. The concentrations of exfoliated GaS, estimated using $\epsilon_{365nm} \approx 3654 \text{ L g}^{-1} \text{ m}^{-1}$, are shown on the right axis. We find a clear and well-defined peak centered at around 21.5

MPa^{1/2} very similar to studies on other 2D nanomaterials.^{4,5,10,39}

According to classical solution thermodynamics, in the simplest case,^{55,78} the saturated concentration, C , of 2D solutes, such as the GaS nanosheets studied here is given by

$$C \propto \exp \left[-\frac{\nu_{NS}}{3kT} (\delta_s - \delta_{NS})^2 \right] \quad (1)$$

where δ_{NS} represents the Hildebrand parameter of the 2-dimensional solute, ν_{NS} represents the unit volume of the solute, and the factor of 3 stems from the solute dimensionality.⁷⁸ The dashed line in Figure 1G (inset) is a fit to eq 1 and shows very good agreement with the experimental data, further confirming the validity of solution thermodynamics in the case of nanomaterial dispersions. The data implies the Hildebrand solubility parameter of GaS nanosheets to be $\sim 21.5 \text{ MPa}^{1/2}$, close to values of 23, ~ 22.5 , 22 and 21 reported recently for graphene,¹⁰ BN,⁴ MoS₂ (ref 55) and MoO₃ (ref 39), respectively. We note that analysis in terms of Hildebrand parameters has been shown to be equivalent to analysis in terms of surface energies.⁷⁸ Essentially, this means successful solvents are those with surface tensions close to 40 mJ/m². Taken together, this means that, as has been observed previously for the materials mentioned above, liquid phase exfoliation of GaS is most favorable in solvents whose interaction with the nanosheets is such that the energy of exfoliation is minimized.⁵

Even though we found empirically that NMP and CHP are the most suitable solvents in terms of dispersed concentration, we have used 2-propanol (IPA) for the remainder of the study (these three solvents are marked in red in Figure 1G). Although the concentration of GaS that can be exfoliated in IPA is only a quarter of that achievable in NMP, use of IPA brings significant advantages in processing and analysis due to its low toxicity and low boiling point. We fully optimized the exfoliation conditions in IPA in terms of centrifugation rate and time, sonication time, and initial GaS concentration as presented in the Supporting Information Figure S1. We found the following parameters to be optimized for the production of standard dispersions of exfoliated GaS in IPA: initial GaS concentration $c_i = 45$ g/L, sonication time $t_s = 6$ h, centrifugation rate $f = 2.5$ krpm (equivalent to 665 g), centrifugation time $t_{CF} = 180$ min. Under these circumstances, we could produce dispersions with concentrations of roughly $C = 0.22$ g/L of dispersed nanosheets.

It is important to confirm that the exfoliated GaS is pristine and defect-free. For this purpose, samples were subjected to Raman and X-ray photoelectron spectroscopies (XPS). Figure 1H shows the Raman spectrum (excitation wavelength 532 nm, mean of 20 spectra) of a filtered film of the standard dispersion. The films are homogeneous in appearance as evidenced by scanning electron microscopy (SEM) (inset in Figure 1H). The characteristic lattice vibrations of GaS are detected at 195, 300, and 370 cm⁻¹ as assigned in the figure.^{61,74,75} Only minor contributions from other materials are observed such as the $\nu_1(A_1)$ mode of the GaS₄ molecular unit as typically found in Ga₂S₃ which we see at $\sim 240 \text{ cm}^{-1}$.^{79,80} The fitted XPS core level spectra (Ga 3d in Figure 1I and S 2p in Figure 1J) further confirm that the exfoliated material is GaS with minor contributions from Ga_xO_y. We note that such oxide species were also found in the starting powder (Supporting Information Figure S2). In addition, all XPS peaks measured for the starting powder are significantly broadened in terms of

full width at half-maximum of the fit components, suggesting a lower degree of order or purity in the starting powder compared to the exfoliated material. This shows that the sonication and centrifugation partially purifies the starting material by removal of unwanted components, such as nonlayered impurities. This can be illustrated by the presence of significant amounts of Ga_2S_3 in the starting powder as shown by X-ray diffraction (XRD) (Supporting Information Figure S3) and Raman spectroscopy (Supporting Information Figure S4). Because Ga_2S_3 is not layered, it is not exfoliated and, therefore, is largely removed during centrifugation.

Electron Energy Loss Spectroscopy. Because XPS shows that oxides are present in both GaS powder and exfoliated nanosheets (Figure 1I and Supporting Information Figure S2) and virtually nothing is known about the long-term stability of GaS in the exfoliated state, it is important to track where oxides species reside. This is particularly crucial in light of recent investigations on exfoliated black phosphorus nanosheets, which have shown significant degradation under exposure to ambient conditions.^{81–83} To gain insights into potential oxidation of GaS, we have analyzed liquid exfoliated nanosheets by STEM imaging and electron energy loss spectroscopy (EELS; see Figure 2). A representative STEM image of a nanosheet at the edge region is shown in Figure 2A. EEL spectra were recorded from the same sample region to form a map, with the intensity of each pixel in the map (ranging from black to yellow/white) corresponding to the integrated intensity of the oxygen K-edge in the pixel location. From this data, an oxygen content map could be constructed from the

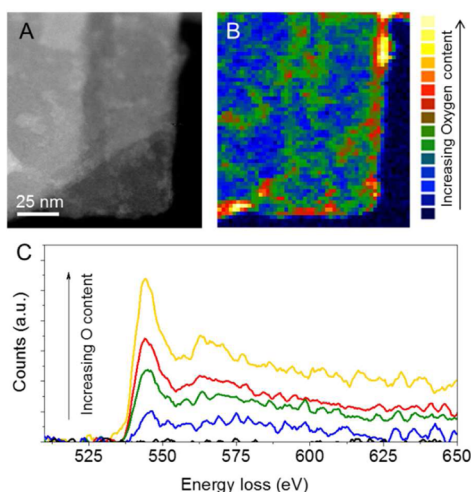


Figure 2. (A) Scanning transmission electron microscopic image of a GaS nanosheet. (B) Color-coded electron energy loss (EEL) map of the same region. Each pixel in the map corresponds to the integrated intensity of the energy loss signal of the oxygen K-edge after background removal to the same edge. Hence, the map shows the variation in the oxygen content across the region where increasing oxygen content is shown from black to green to red to yellow to white. The oxygen signal is most prominent on the edges of the nanosheet and at step edges in the center of the nanosheet as indicated in the figure (yellow/white regions). (C) Representative EEL spectra extracted at different positions of the nanosheet, with the colors of the spectra corresponding to the colors in the map in panel B. The spectra shown in yellow represent the highest O content, as found near the edges of the nanosheets, and the spectra shown in black and blue representing the lowest O content, as found on the surfaces inside the flake and over vacuum.

same sample region (Figure 2B) as shown in STEM image (Figure 2A). Hence, the map is color-coded and shows increasing oxygen content from black/blue (no/low oxygen) to green to red to yellow to white. It is clear that the oxides reside mostly near the edges of the nanosheets and at step edges throughout the nanosheets. The EEL spectra corresponding to regions with different oxygen content are displayed in Figure 2C.

Size Selection. A great advantage of liquid exfoliation is not only that nanosheets are processable from liquids but also that they can be size-selected using well-established techniques by controlled centrifugation.^{37,38,40} This is extremely important considering that potential applications often require control over both thickness and lateral dimensions. To demonstrate this, we have performed size-selection by controlled centrifugation with increasing centrifugation velocity in consecutive steps (see Materials and Methods). This enables the production of nanosheet dispersions with varying mean sizes from the same stock dispersion in large quantities. By this procedure, five different sizes were produced and subjected to detailed characterization. For the main manuscript we focused on three different dispersions. Further data can be found in the Supporting Information. Displayed in Figure 3A–C are TEM length histograms and representative images of three different size-selected GaS dispersions in IPA (data for all five sizes is given in the Supporting Information Figure S5). The mean lateral dimensions were determined from the statistical TEM analysis yielding TEM mean lengths (L) ranging from 130–405 nm.

Because it has previously been reported that optical extinction, absorbance, and scattering spectra change as a function of size, the optical response was measured for each dispersion.^{37,40} Any systematic spectral changes can be valuable because, once calibrated, they can be used to establish metrics to potentially quantify both lateral dimensions and thickness of the liquid exfoliated nanosheets.³⁷ Optical extinction (normalized to the local minimum), absorbance (normalized to the local minimum in extinction) spectra of the size-selected GaS are displayed in Figure 3D–F (plots showing the un-normalized coefficient spectra can be found in the Supporting Information, Figure S12). All exhibit well-defined and systematic changes as a function of size, which will be analyzed below. Of greatest interest are the absorbance spectra (Figure 3E), particularly in comparison to the extinction spectra (Figure 3D). As mentioned above, the extinction spectra are almost featureless. However, GaS is known to display excitonic transitions at ~ 410 nm (A-exciton, direct transition at the Γ -point) and ~ 315 nm (B-exciton, transition at the M-point).⁸⁴ Although the B-exciton is discernible in the extinction spectra as the local maximum, the A-excitonic transitions are invisible due to masking by the scattering background. However, in the absorbance spectra, both excitons can be clearly resolved. In addition, very weak features at higher wavelength are discernible in the absorbance spectra (inset Figure 3E), which may stem from defects.⁸⁵ However, these are widely invariant with size so that we suggest they are not edge or exfoliation induced.

In addition to quantifying lateral dimensions by statistical TEM, we have determined the thickness of the exfoliated GaS nanosheets using atomic force microscopy (AFM) after deposition onto Si/SiO₂ wafers (see Materials and Methods). We find reasonably thin nanosheets similar in appearance to the TEM images (Figure 3G inset). Because the direct measure-

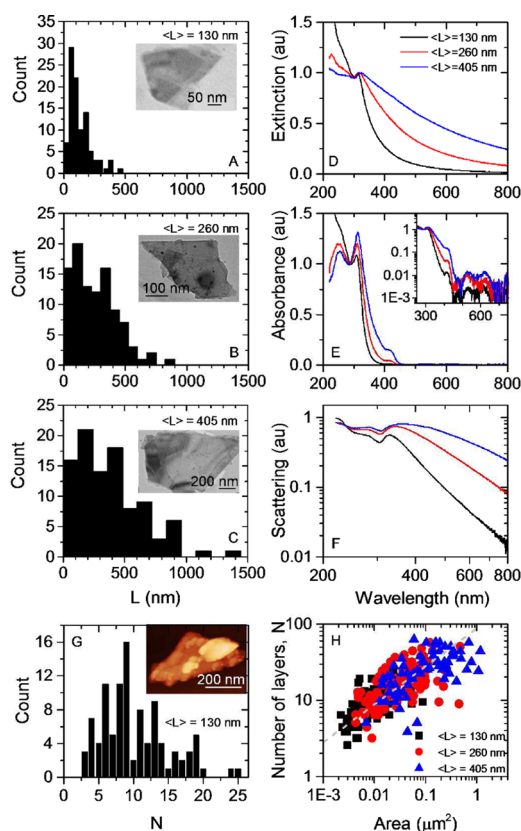


Figure 3. (A–C) Length histograms and representative images of size-selected GaS in IPA as insets. (A) Smallest-size GaS, (B) medium-sized GaS, (C) large GaS. The mean length $\langle L \rangle$ is shown as figure legend. Data for additional sizes are given in the Supporting Information. (D–F) Optical extinction, absorbance and scattering spectra of dispersions of different sized nanosheets showing systematic changes in spectral shape. (D) Extinction, (E) absorbance and (F) scattering. Extinction and absorbance spectra are normalized to their local minimum while scattering spectra are normalized to the minimum in the extinction spectra. The inset in (E) shows the absorbance on a log scale to highlight the A-exciton (~ 420 nm) and the gap features. (G) Typical AFM number of layers, N , histogram of smallest-size GaS in IPA drop cast onto Si/SiO₂ wafers. The mean number of layers was determined from step height analysis (see Supporting Information). Inset: representative image of a nanosheet. (H) Plot of number of layers per nanosheet as a function of flake area determined from AFM. The dashed line indicates square root behavior.

ment of the number of layers, N , in the case of liquid exfoliated nanomaterials is complicated by both solvent and nanosheet contributing to the apparent AFM thickness, we have used previously elaborated step height analysis^{6,37,40} to convert the thickness to the number of layers (Supporting Information Figure S6). The resultant N histogram of the GaS sample with mean length of 130 nm is displayed in Figure 3G showing that the nanosheets are reasonably well exfoliated with mean N of 10 layers. Histograms and representative images of the other sizes are shown in the Supporting Information Figure S7.

However, before being certain about conversion of apparent thickness to number of layers, it needs to be ensured that only individually deposited nanosheets are taken into account. To test this, we plot the mean AFM length of the counted nanosheets as a function of the TEM length in the Supporting Information Figure S8A, finding very good agreement. Furthermore, we can use our previous knowledge on other

liquid exfoliated 2D materials that has shown a square root dependence of N as a function of nanosheet area.^{37,39} As demonstrated by Figure 3H (and Supporting Information Figure S8B), the same behavior is observed in the case of GaS strongly supporting the accuracy of the AFM number of layer determination.

We furthermore note that the structural integrity of the size-selected nanosheets was confirmed by Raman and XPS (see Supporting Information Figure S9–11), which gave similar results for both non-size-selected standard dispersions and the size-selected nanosheets described in Figure 3. This shows the nanosheet spectroscopic properties other than the optical response to be generally independent of size in the regime our nanosheets are produced. For example, we do not observe layer-number induced shifts in the GaS Raman spectra as previously reported.⁷⁵ However, this is not surprising, as the Raman spectrum of GaS is only sensitive to layer numbers of <5 and the majority of our GaS nanosheets have thicknesses >5 layers.

Spectroscopic Metrics to Determine Size and Concentration. The precise quantification of both N and L as described above is important, as it provides the foundation to establish spectroscopic metrics based on extinction, absorbance, and scattering spectra. In the following, we show a number of these metrics, which can be used to determine information regarding GaS nanosheet size and thickness from optical measurements.

As mentioned above, the scattering spectra are sensitive to the nanosheet lateral size. Previous work has shown that, in the high wavelength regime where the absorbance is negligible, the scattering exponent scales as a power law with wavelength: $\sigma \propto \lambda^{-n}$, where σ is the scattering coefficient ($Sca = \sigma Cl$) and n is the size-dependent scattering exponent that can also be used as a potential metric to quantify nanosheet length.^{4,37} To test whether such a metric can be established for liquid-exfoliated GaS, we plot the long wavelength scattering exponent, n , as a function of mean nanosheet length (as measured by TEM) in Figure 4A. Importantly, because absorbance of GaS is negligible at high wavelength, we obtain the same result extracting n from either extinction or scattering spectra, implying that measurement in an integrating sphere is not essential to determine $\langle L \rangle$. We find a roughly linear relationship (valid in this size range only) between $\langle L \rangle$ and n that allows us to use the scattering coefficient to determine the mean length of the GaS nanosheets in the dispersion according to eq 2

$$\langle L(\mu\text{m}) \rangle = 0.67 - 0.14n \quad (2)$$

Recent work has also shown that for nanosheet dispersions, the shape of the absorption spectra is sensitive to nanosheet lateral size due to edge effects.³⁷ These spectral changes in liquid-exfoliated 2D nanomaterials can be expressed as ratios of absorbance (or extinction) at two wavelengths. These intensity ratios contain nanosheet length information due to differences in electronic properties at nanosheet edge and center.³⁷ This effect can clearly be seen in the case of GaS when plotting the ratio of absorbance at the A-exciton/local minimum at 290 nm ($A_{420\text{nm}}/A_{290\text{nm}}$) as a function of $\langle L \rangle$ (Figure 4B). This plot shows a well-defined relationship allowing us to link mean nanosheet length to $A_{420\text{nm}}/A_{290\text{nm}}$ according to eq 3

$$\langle L(\mu\text{m}) \rangle = 0.93(A_{420\text{nm}}/A_{290\text{nm}})^{0.44} \quad (3)$$

We note that although this shows the validity of the approach, $\langle L \rangle$ determination from the scattering exponent in optical

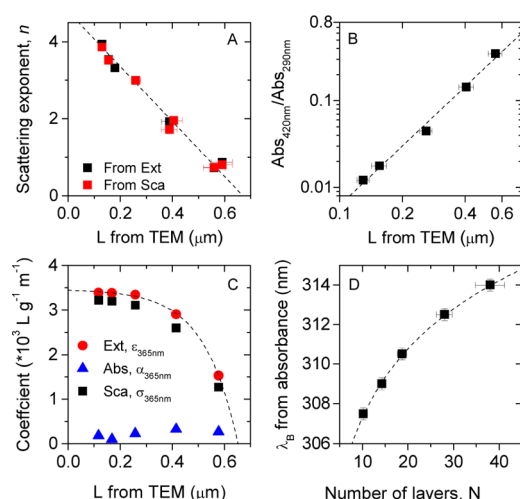


Figure 4. (A) Long wavelength scattering exponent, n , measured from both extinction and scattering spectra, plotted versus TEM mean length, L , for size-selected GaS in IPA. (B) Ratio of absorbance at 420 nm to that at 290 nm, $\text{Abs}_{420\text{nm}}/\text{Abs}_{290\text{nm}}$ plotted versus TEM length, L , for size selected GaS in IPA. The fit lines plotted in A and B can be used as metrics to determine L according to eqs 2 and 3, respectively. (C) Extinction, absorbance, and scattering coefficients as a function of L . The fit line can be used to find the extinction coefficient for a mean nanosheet length which in turn can be used to determine the dispersed concentration from the optical extinction spectra according to eq 4. (D) Position of the B-exciton, λ_B (measured from the absorbance spectra), plotted versus number of monolayers per nanosheet, N (determined from AFM). The relation can be used as metric to determine N according to eq 5.

extinction spectra is more straightforward, especially because the integrating sphere is not required.

Although the determination of $\langle L \rangle$ is useful, for most applications, it is critical to control concentration/mass of the nanosheets. An accurate determination of the extinction/absorbance coefficient, therefore, is required (N.B. the extinction, absorption and scattering coefficients are related to Ext , Abs and Sca by: $Ext = \epsilon Cl$, $Abs = \alpha Cl$ and $Sca = \sigma Cl$). Because the spectra change as a function of size, extinction and absorbance coefficients are expected to change as well (as observed for MoS_2).³⁷ Therefore, we also analyze extinction ϵ , absorbance α , and scattering σ coefficient spectra (Supporting Information Figure S12). The GaS nanosheet concentration in each case was determined by filtration and weighing. Depending on the wavelength, the coefficients vary significantly as a function of nanosheet size (Supporting Information Figure S12–13). However, we observe only minor changes in the extinction coefficient at 365 nm over a wide range of size from 100–300 nm (Figure 4C), making this an ideal wavelength for estimation of the concentration from the extinction spectra using $\epsilon_{365\text{nm}} = 3465 \text{ L g}^{-1} \text{ m}^{-1}$. Because mean nanosheet lengths from a single centrifugation step are typically below 300 nm, we have used the approximate extinction coefficient for our process optimization (as described above). If the nanosheet size is $>300 \text{ nm}$ (where $\langle L \rangle$ can be determined by eq 2 or 3), then the extinction coefficient can be found from an empirical fit of the data

$$\epsilon_{365\text{nm}} (\text{L g}^{-1} \text{ m}^{-1}) = 3465 - 21 \times e^{(L(\mu\text{m}))/0.128} \quad (4)$$

In addition to determination of length and concentration from the optical spectra, it would also be useful to have a metric to

assess the mean number of layers. In the case of MoS_2 , the energy/wavelength of the A-exciton provided such a metric due to confinement effects.³⁷ In the case of GaS, we correlate the mean nanosheet thickness (as measured from AFM) with the peak position of the B-exciton from the absorbance spectra, λ_B . We used the B-exciton rather than the A-exciton in this case, as the intensity of the A-exciton is comparatively weak for all GaS sizes. Shown in Figure 4D is a plot of λ_B versus $\langle N \rangle$ with the correlation being clear. By applying an empirical fit, we find the mean number of layers of the liquid exfoliated GaS nanosheets can be determined from the peak position of the B-exciton in the absorbance spectra, λ_B , according to eq 5

$$\langle N \rangle = 2.24 \times 10^{-26} \times e^{\lambda_B(\text{nm})/11.5} \quad (5)$$

A similar metric from the extinction spectra is presented in the Supporting Information (Figure S14). However, we note that the position of the B-exciton does not tend to a constant value that might be associated with bulk for nanosheets thicker than ~ 10 – 20 layers as might be expected. The reasons for this are unclear and will be the subject of a future study.

The presented quantitative in situ spectroscopic metrics to determine $\langle L \rangle$, $\langle N \rangle$, and concentration of liquid exfoliated GaS underline the strength of LPE, as these will be extremely useful to prepare dispersions with known dimensions and concentrations to test in applications.

Hydrogen Evolution Catalysis. The size metrics and precise concentration control described above facilitates the use of liquid-exfoliated GaS nanosheets in applications where lateral nanosheet size is important. Such an application is as catalysts in the hydrogen evolution reaction (HER). It is known that edge sulfur atoms in MoS_2 nanosheets are involved in the electrocatalysis of hydrogen production.^{86,87} We hypothesized that GaS may also have catalytically active edge sites. If this were to be the case, liquid exfoliated GaS nanosheets should catalyze H_2 production with the rate of production increasing as the nanosheet size decreases. To test whether this was the case, we size-selected a GaS stock dispersion according to the procedure used throughout this manuscript to give three size-selected dispersions. The mean nanosheet lengths were determined from the scattering exponent to be 450, 280, and 180 nm, respectively. The dispersions were vacuum-filtered to give thin films ($\sim 0.65 \text{ mg/cm}^2$), which were transferred onto pyrolytic carbon (PyC) coated Si/SiO₂ substrates (Supporting Information Figure S15) and characterized for hydrogen evolution catalysis in a three-electrode electrochemical workstation (see Materials and Methods).

The current density versus potential relative to reversible hydrogen electrode (polarization curves) for small, medium, and large nanosheets are shown in Figure 5A. Clear size-dependence of hydrogen production is observed with small nanosheets performing much better than larger ones. This suggests that catalytic sites active in GaS also reside at edges as found for other 2D materials.^{86,87} For example, as the nanosheet size is reduced from 450 to 280 to 180 nm, the onset potential (the potential where $J = 1 \text{ mA/cm}^2$) decreases from 0.62 to 0.53 to 0.48 V, whereas the current density at 0.6 V increases from 0.6 to 6 to 22 mA/cm^2 (see Supporting Information Figure S16 and Table S2). We note that although we observe an increased content of oxides in the vicinity of the edges according to EELS (Figure 2), this does not necessarily affect the catalytic reaction, as it was proposed that only the very outer rim of atoms are catalytically active.^{86,87} In addition, we do not see a decrease in catalytic performance after repeated

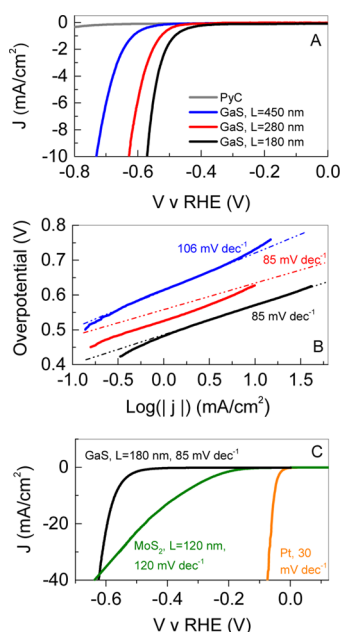


Figure 5. (A) Linear sweep voltammograms (5 mV s^{-1}) of vacuum-filtered GaS nanosheets transferred onto pyrolytic carbon with a reversible hydrogen (RHE) reference electrode. The electrocatalytic response toward the hydrogen-evolution reaction is compared for size-selected GaS with three different mean sizes. A clear size-dependence is evident. Measured potentials were subjected to iR correction. Supporting electrolyte was $0.5 \text{ M H}_2\text{SO}_4$. (B) Overpotential versus current density plots showing Tafel slopes expressed as millivolts per decade (mV dec^{-1}). (C) Linear sweep voltammograms of s-GaS compared to liquid-exfoliated 2H-MoS₂ of similar lateral dimensions and mass per area and Pt as benchmark. Tafel slopes are indicated in the figure. Despite of the high onset potential, s-GaS reaches the same current density as MoS₂ at $\sim 0.6 \text{ eV}$ due to its smaller Tafel slope.

measurement of the polarization curves (Supporting Information Figure S17) and do not observe a significantly increased oxide content according to XPS after this procedure (Supporting Information Figure S18).

Tafel plots of overpotential versus current density (Figure 5B) show the expected behavior in the linear regime with Tafel slopes between 106 and 85 mV/dec . These slopes are smaller than the values of $\sim 120 \text{ mV/dec}$ usually found for the 2H-polytype^{45,88} of MoS₂ and indicate the rate limiting step may be different in GaS nanosheets. However, the exchange current densities were $< 2.5 \times 10^{-6} \text{ mA/cm}^2$ for all three GaS samples, much smaller than is usually found for 2H-MoS₂. In addition, the onset potential was at least 480 mV for the samples studied here, considerably larger than the values of $\sim 300 \text{ mV}$ usually found for 2H MoS₂.^{45,88}

Thus, at first glance, GaS might appear to be a poor catalyst for hydrogen evolution in contrast to MoS₂, which is commonly accepted as a promising material. This is due to its relatively high onset potential and low exchange current densities. However, the lower Tafel slope may offer some compensation for these deficiencies. To test this, we produced 2H-MoS₂ nanosheets by liquid exfoliation and size selection using previously described procedures.³⁷ We produced a film with near identical mass ($\sim 0.7 \text{ mg/cm}^2$) to the GaS samples and slightly smaller nanosheet size ($L = 120 \text{ nm}$). Polarization curves for LPE GaS and 2H-MoS₂ are compared in Figure 5C. It is apparent that, although GaS requires large overpotentials to initiate the hydrogen evolution, its smaller Tafel slope causes

a more rapid increase of the current with potential. In fact, both materials reach the same current density at 0.6 V , above which GaS dominates. This suggests that GaS can indeed be regarded as potentially attractive hydrogen evolution catalyst, especially if problems associated with the low conductivity could be overcome, for example, by addition of nanotubes.⁸⁹

CONCLUSION

In conclusion, we have demonstrated that layered III–VI semiconductors such as GaS can be exfoliated in appropriate solvents by sonication. Dispersibility can be well described in the framework of solution thermodynamics. Although a number of solvents can be used to exfoliate GaS, we have focused on 2-propanol due to its low toxicity and boiling point. We have fully optimized sonication and centrifugation conditions to yield stable dispersions with typical concentrations of 0.2 g/L . Raman, XPS, and high resolution STEM imaging show the exfoliated nanosheets to be widely structurally perfect and free of defects except for a higher oxygen content in the vicinity of edges, as shown by EELS.

We utilized one advantage of liquid exfoliation techniques and performed size-selection by controlled centrifugation. This allowed us to produce liquid-exfoliated GaS nanosheet dispersions with mean lateral dimensions ranging from $\sim 100 \text{ nm}$ to $>400 \text{ nm}$ and mean number of layers from 10 to 40 as quantified by statistical TEM and AFM analysis. Importantly, we found optical extinction, absorbance, and scattering spectra to vary strongly as a function of size and thickness. This enabled us to establish quantitative spectroscopic metrics to accurately determine mean length, thickness, and concentration of the dispersion.

We subsequently used these metrics to produce nanosheet dispersions with known concentration and lateral dimensions that were used as electrodes and tested for hydrogen evolution catalysis. We find a clear size effect with smaller nanosheets performing much better than larger nanosheets. This suggests that catalytic sites are located at the edges of the nanosheets. Even though onset potentials are typically larger than for MoS₂ of similar lateral dimensions, the lower Tafel slopes make GaS an attractive material for hydrogen evolution, particularly in combination with other catalysts that exhibit lower onset potential.

We believe these results are general and that it should be possible to transfer these procedures almost exactly for the exfoliation of other III–VI layered semiconductors such as InS and GaSe. This will open up this whole family of 2D materials for exploitation. In addition, we think it is important to note that this paper describes using LPE to exfoliate a member (GaS) of yet another family of layered compounds (III–VI layered semiconductors). This underlines the generality and versatility of this method. We believe that many as yet untested layered materials will be exfoliated using such procedures.

MATERIALS AND METHODS

Materials. Gallium sulfide powder (99.999% GaS-05-P) was purchased from American Elements. All solvents were purchased from Sigma-Aldrich at the highest available purity, with all being $>99\%$.

Production of GaS Nanosheets. Gallium sulfide powder was sonicated in a solvent using an ultrasonic bath (P30 H Ultrasonic from Fischer scientific). The sonication was performed with an amplitude of 100% and a frequency of 37

kHz in 50 mL plastic centrifuge tubes. The water in the sonic bath was cooled by a water cooling system to 20–30 °C (depending on length of sonication) enabled by cold water being pumped through piping that was wrapped around the interior of the bath. Once sonicated, the dispersion was centrifuged in a Hettich Mikro 220R centrifuge with a fixed-angle rotor 1060 (NB: For this centrifuge, rpm are related to *g*-force via $RCF = 106.4f^2$, where *f* is the rotation rate in krpm). The top 60% was then taken from the centrifuged dispersion (supernatant) for analysis and the sediment was discarded. For the solvent screening, 1 g/L of GaS was sonicated in 20 mL of each solvent for 6 h and centrifuged for 180 min at 2.5 krpm. The supernatant was decanted and the absorption and extinction were measured in a 4 mm path length cuvette using a PerkinElmer Lambda 650 spectrometer (details see below). The following solvents were used: isopropanol ($\delta_s = 23.6 \text{ MPa}^{1/2}$), *N*-methyl-2-pyrrolidone ($\delta_s = 23 \text{ MPa}^{1/2}$), methanol ($\delta_s = 29.6 \text{ MPa}^{1/2}$), chloroform ($\delta_s = 19 \text{ MPa}^{1/2}$), *N*-cyclohexyl-2-pyrrolidone ($\delta_s = 20.5 \text{ MPa}^{1/2}$), dimethylformamide ($\delta_s = 24.9 \text{ MPa}^{1/2}$), cyclopentane ($\delta_s = 16.5 \text{ MPa}^{1/2}$), heptane ($\delta_s = 15.3 \text{ MPa}^{1/2}$), hexane ($\delta_s = 14.9 \text{ MPa}^{1/2}$), pentane ($\delta_s = 14.4 \text{ MPa}^{1/2}$), acetone ($\delta_s = 19.9 \text{ MPa}^{1/2}$), acetonitrile ($\delta_s = 24.3 \text{ MPa}^{1/2}$), 1,3-dioxolane ($\delta_s = 21.4 \text{ MPa}^{1/2}$), benzonitrile ($\delta_s = 22.5 \text{ MPa}^{1/2}$), isopropoxyethanol ($\delta_s = 21.4 \text{ MPa}^{1/2}$). The final, optimized exfoliation was performed as follows: 45 g/L of GaS were sonicated for 6 h in isopropanol in 20 mL vials and then centrifuged for 180 min at 2.5 krpm.

Size Selection. We used controlled centrifugation with subsequently increasing rotation speeds as previously reported.^{6,40} Then, 10 g/L GaS in 90 mL of isopropanol was sonicated for 6 h in three centrifuge tubes. The 90 mL of sonicated dispersion was centrifuged at 0.5 krpm for 60 min. The sediment was discarded and the supernatant was centrifuged again at 1 krpm for 60 min. The sediment after this centrifugation step was redispersed in fresh IPA (5 min bath sonication) producing the largest size. The supernatant after the 1 krpm centrifugation step was centrifuged at 1.5 krpm for 60 min, producing the second largest size in the redispersed sediment. These steps were repeated in further increments of 2 krpm, 2.5 krpm, and 3 krpm, thus producing five sizes.

Characterization and Equipment. Optical extinction and absorbance was measured on a PerkinElmer 650 spectrometer in quartz cuvettes with a path length of 0.4 cm. To distinguish between contributions from scattering and absorbance to the extinction spectra, dispersions were measured in an integrating sphere using a home-built sample holder to place the cuvette in the center of the sphere (NB cuvettes need to be transparent to all sides and correct positioning is important). The absorbance spectrum is obtained from the measurement inside the sphere. A second measurement on each dispersion was performed outside the sphere in the standard configuration to obtain the extinction spectrum. This allows calculation of the scattering spectrum (extinction minus absorbance).

Low-resolution bright field transmission electron microscopy imaging was performed using a JEOL 2100, operated at 200 kV. Holey carbon grids (400 mesh) were purchased from Agar Scientific and prepared by diluting a dispersion to a low concentration and drop casting onto a grid placed on a filter membrane to wick away excess solvent. Statistical analysis was performed of the flake dimensions by measuring the longest axis of the nanosheet and assigning it as “length”, *L*. STEM imaging and STEM EELS was carried out on a Nion

UltraSTEM100 aberration-corrected dedicated STEM, equipped with a Gatan Enfina electron energy loss (EEL) spectrometer, at the SuperSTEM Laboratory in Daresbury, U.K. The microscope was operated at an acceleration voltage of 60 kV.

EELS were acquired with a convergence semiangle of 32 mrad, a collection semiangle of approximately 37 mrad, a dispersion of 0.5 eV/channel, and an energy resolution over vacuum of 0.3 eV. STEM EELS maps were acquired with 0.1 s acquisition time per spectra per pixel. Subsequently, the spectra were aligned to the oxygen K-edge. The spectra shown here are representative of a set of 10 STEM EELS maps that were acquired over similar regions.

Atomic force microscopy (AFM) was carried out on a Veeco Nanoscope-IIIa (Digital Instruments) system equipped with a E-head (13 μm scanner) in tapping mode after depositing a drop of the dispersion (10 μL) on a preheated (120 °C) Si/SiO₂ wafer with an oxide layer of 300 nm. Typical image sizes were 3–10 μm at scan rates of 0.4–0.6 Hz.

Raman spectroscopy was performed using a WITec alpha 300 with 532 nm excitation laser in air under ambient conditions. The Raman emission was collected by an Olympus 100 \times objective (N.A. = 0.8) and dispersed by 600 lines mm^{-1} gratings. The laser energy was kept below 0.2 mW. The mean of 20 spectra is displayed.

X-ray Photoelectron Spectroscopy was performed under ultrahigh vacuum conditions ($<5 \times 10^{-10}$ mbar) using monochromated Al K α X-rays from an Omicron XM1000 MkII X-ray source and an Omicron EA125 energy analyzer. The analyzer pass energy was set to 100 eV for survey and 20 eV for core-level spectra, yielding a maximum energy resolution of ~ 0.65 eV. An electron flood gun was used for charge compensation and the binding energy scale was referenced to the adventitious carbon 1s core-level at 284.8 eV. After subtraction of a Shirley background, the core-level spectra were fitted with Gaussian–Lorentzian line shapes. Samples were prepared by vacuum-filtering the dispersions using porous cellulose filter membranes (MF-Millipore membrane, mixed cellulose esters, hydrophilic, 0.025 μm , 47 mm) to give thin films.

Scanning electron microscopy was performed with a Carl Zeiss Ultra SEM operating at 2 kV. Images were acquired using the secondary electron detector. Powder XRD was performed using a Siemens D500 X-ray Diffractometer equipped with a Cu K α emission source ($\lambda = 1.54056 \text{ \AA}$) filtered through a graphite monochromator at ambient temperature.

Hydrogen Evolution Catalysis. Dispersions of GaS in IPA were vacuum-filtered using porous cellulose filter membranes (MF-Millipore membrane, mixed cellulose esters, hydrophilic, 0.025 μm) to give uniform thin films. The deposited films were then cut into pieces (0.65 cm^2) and transferred on PyC electrodes using heat and pressure as described previously.^{24,90} The cellulose filter membrane was then removed by treatment with acetone vapor and subsequent acetone liquid baths followed by an isopropanol rinse to remove the acetone residue. The mass per area for all samples was approximately 0.65 mg/cm^2 .

Pyrolytic carbon was grown by chemical vapor deposition (CVD) from an acetylene feedstock at 950 °C for 30 min to a thickness of 300–400 nm on 300 nm thermal SiO₂ on Si substrates in a hot wall quartz tube furnace as previously reported.⁹¹

Electrochemical measurements were carried out to evaluate the performance of GaS on PyC films as electrodes for HER.

The measurements were performed in 0.5 M H₂SO₄ solution using a three-electrode electrochemical cell, with a reversible hydrogen electrode (RHE) reference electrode and a graphite rod counter electrode. Electrochemical tests consisted of linear sweep voltammetry and electrochemical impedance spectroscopy using a Gamry Reference 3000 potentiostat. Before each test, samples were conditioned at a given voltage for 3 min. The electrocatalysis was measured using linear sweeping from 0 V to −0.8 V (vs RHE) with a scan rate of 5 mV/S. The AC impedance was measured within the frequency range of 0.1 to 10 mHz with perturbation voltage amplitude of 10 mV. The equivalent series resistance of the system was measured by impedance spectroscopy from the high frequency intercept with the real impedance axis and all the data were corrected by *iR* compensation.

■ ASSOCIATED CONTENT

■ Supporting Information

This includes more detailed descriptions of the basic characterization, size selection, optical characterization, and hydrogen evolution. This material is available free of charge via the Internet at <http://pubs.acs.org/>.

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Notes

The authors declare no competing financial interest.

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