

Fiber Nonlinearities

By

Dr.A.Brinth Theres

Fiber Nonlinearities

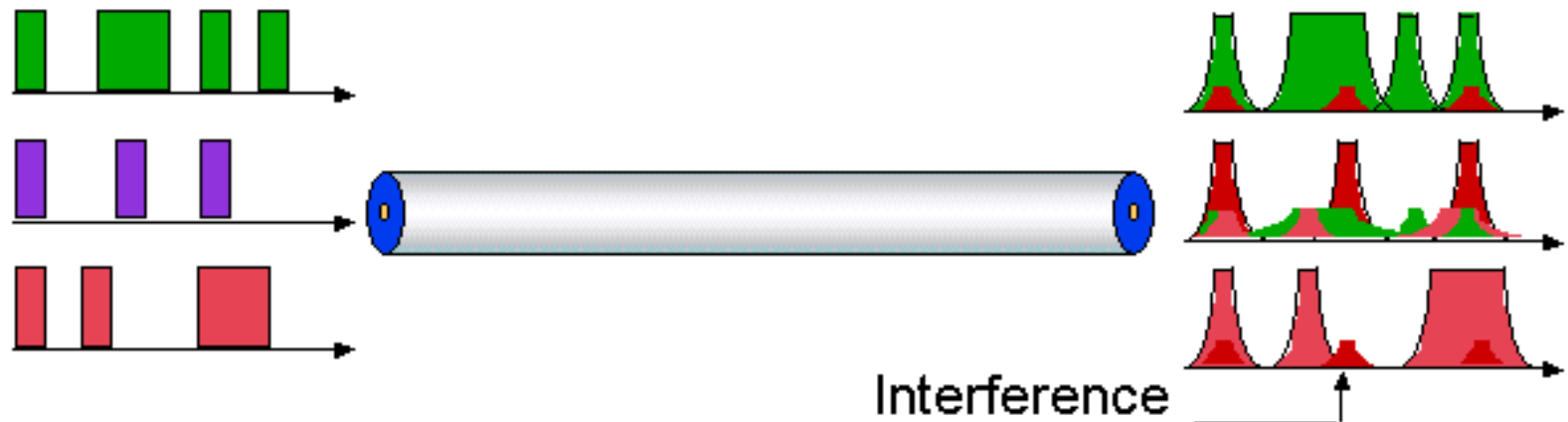
- As long as **optical power** within an optical fiber is **small**, the fiber can be treated as a **linear medium**; that is the loss and refractive index are independent of the signal power
- When **optical power** level gets fairly **high**, the fiber becomes a **nonlinear medium**; that is the loss and refractive index depend on the optical power

Effects of Nonlinearities

- A single channel's pulses interact as they travel (Self-Phase)



- Multiple channels interact as they travel (Cross Phase, FWM)



- Degradation scales as $(\text{channel power})^2$

Nonlinear Effects

- Stimulated Raman scattering
- Stimulated Brillouin scattering
- Four-wave Mixing
- Self-phase Modulation
- Cross-phase Modulation

FIBER EFFECTIVE LENGTH

- Nonlinear interaction depends on transmission length and cross-sectional area of the fiber
- The longer the length, the more the interaction and the worse the effect of the nonlinearity.
- BUT, signal propagates along link and experiences loss (from fiber attenuation) ...
...complicated to model.

Simple model: Assume power is constant over a certain effective length

P denotes power transmitted into fiber. L denotes actual fiber length

$$P(z) = P e^{-\alpha z} \quad \text{power at distance } z \text{ along link.}$$

$$PL_e = \int_{z=0}^L P(z) dz$$

$$L_e = \frac{1 - e^{-\alpha L}}{\alpha}$$

Typical:

$\alpha = 0.22 \text{ dB/km at } 1.55\mu\text{m}$

if $L \gg 1/\alpha$, then L_e approx 20 km

EFFECTIVE CROSS SECTIONAL AREA

Effect of nonlinearity grows with intensity in the fiber. This is inversely proportional to the area of the core (for a given power).

Power not evenly distributed in the cross section.
Use effective cross sectional area (for convenience).

A = actual cross sectional area

$I(r, \theta)$ = actual cross sectional distribution of the intensity.

$$A_e = \frac{\left[\int_r \int_\theta r dr d\theta I(r, \theta) \right]^2}{\int_r \int_\theta r dr d\theta I(r, \theta)}$$

Most cases of interest:

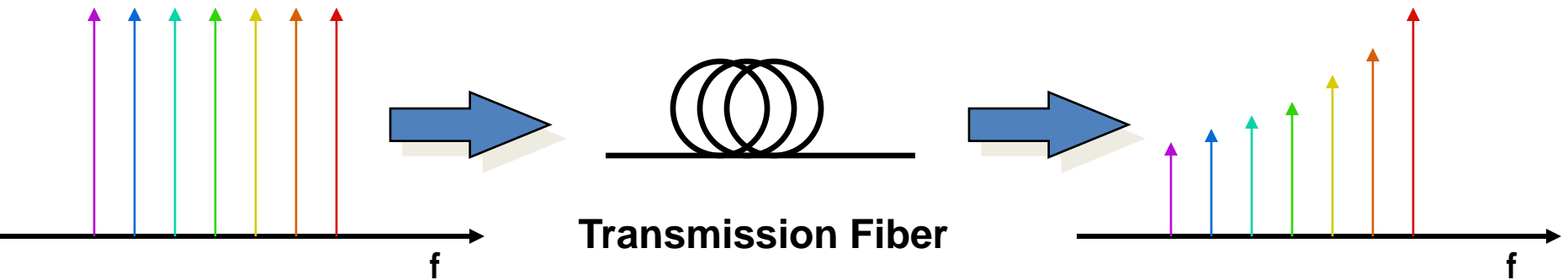
$A_e \approx$ area of single mode fiber

Stimulated Raman Scattering

- Scattering of light from vibrating silicon molecule
- If two or more signals at different wavelengths are injected into a fiber, SRS causes power to be transferred from the lower wavelength channels to the higher-wavelength channels.
- Has a broadband effect (unlike SBS)
- Gain coefficient g_R as function of channel separation.
- Both forward and reverse traveling Stokes wave.
- Coupling between channels occurs only if both channels sending a “1”.
- SRS penalty is therefore reduced by dispersion.

SRS Cont...

- N Channels equally spaced
- Channel 0 lowest wavelength is affected worst



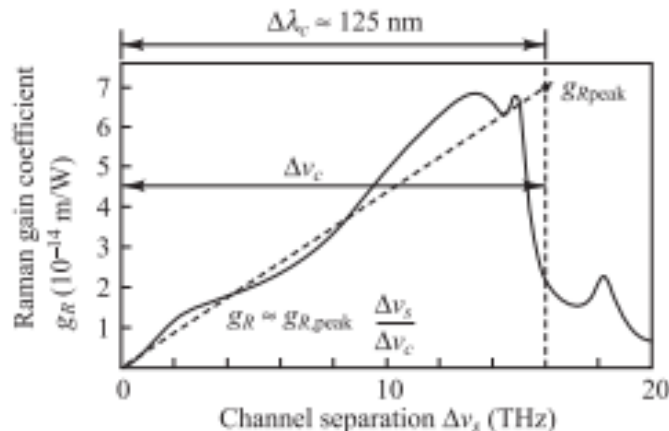
SRS Cont...

Consider a WDM system that has N channels equally spaced in a 30-nm band centered at 1545 nm. Assume that the transmitted power P is the same on all channels, that the

If $F_{\text{out}}(j)$ is the fraction of power coupled from channel 0 to channel j , then the total fraction of power coupled out of channel 0 to all the other channels is

$$F_{\text{out}} = \sum_{j=1}^{N-1} F_{\text{out}}(j) = \sum_{j=1}^{N-1} g_{R,\text{peak}} \frac{j \Delta \nu_s}{\Delta \nu_c} \frac{PL_{\text{eff}}}{2A_{\text{eff}}} = \frac{g_{R,\text{peak}} \Delta \nu_s PL_{\text{eff}}}{2 \Delta \nu_c A_{\text{eff}}} \frac{N(N-1)}{2}$$

Raman gain increases linearly as shown by the dashed line in Figure



Stimulated Raman Scattering (SRS)

1) Effect and consequences

- SRS causes a signal wavelength to behave as a “pump” for longer wavelengths, either other signal channels or spontaneously scattered Raman-shifted light. The shorter wavelengths is attenuated by this process, which amplifies longer wavelengths
- SRS takes place in the transmission fiber

2) SRS could be exploited as an advantage

- By using suitable Raman Pumps it is possible to implement a Distributed Raman Amplifier into the transmission fiber. This helps the amplification of the signal (in co-operation with the localized EDFA). The pumps are depleted and the power is transferred to the signal

Stimulated Brillouin Scattering

- Arises when a strong optical signal generate acoustic waves(Sound waves represent alternating regions of compressed material and expanded material)
- Produce refractive index variation
- Scattering is induced by index discontinuities
- Cause backwards scattering in fiber

SBS, continued

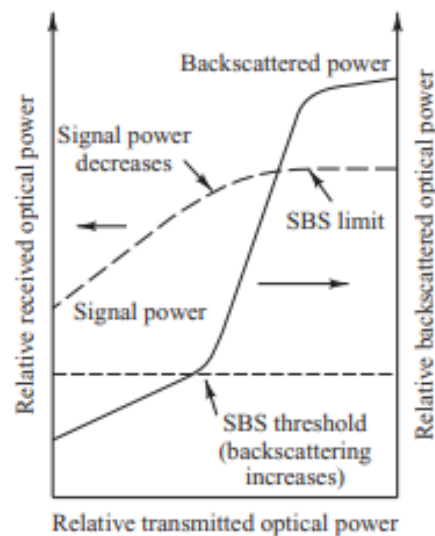
- Brillouin frequency shift equal to $2nV_s/\lambda$, where n is the mode index and v is the speed of sound in the material
- For fiber, scattered light is 11 GHz lower in frequency than signal wavelength (speed of sound is 5.96 km/s)
- System impairment starts when amplitude of the scattered wave is comparable to the signal power

SBS

- SBS threshold Power

$$P_{th} = 21 \frac{A_{eff} b}{g_B L_{eff}} \left(1 + \frac{\Delta \nu_{source}}{\Delta \nu_B} \right)$$

- The effect of SBS on signal power in an optical fiber



Schemes for reducing the SBS

- Optical isolators
- Keeping optical power per WDM below threshold
- Increasing the linewidth

Self Phase Modulation

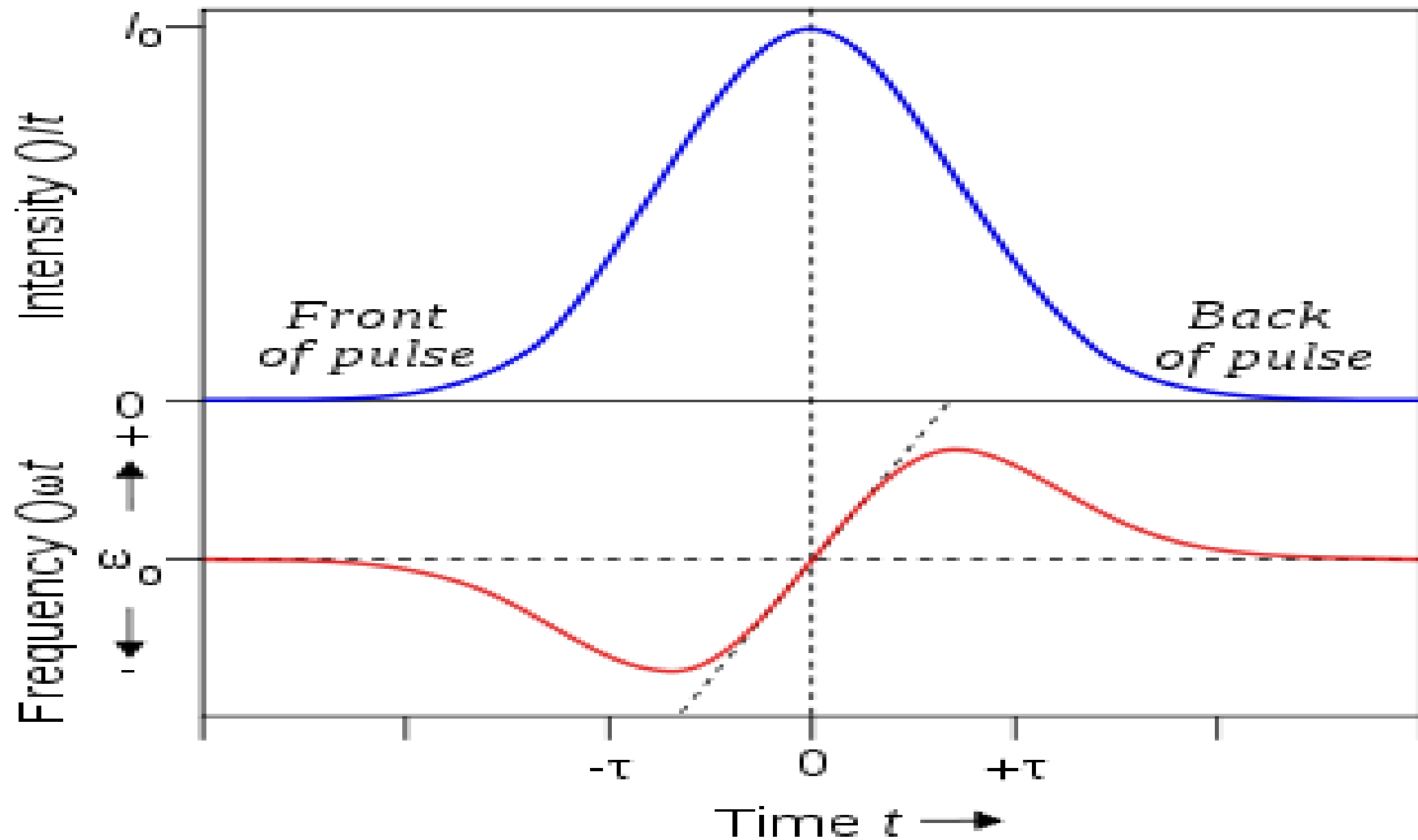
- The refractive index of many materials depends on optical intensity I
- $n = n_o + n_2 I$
 $= n_o + n_2 P/A_{\text{eff}}$

n_o = ordinary refractive index

n_2 = nonlinear index coefficient

Kerr nonlinearity produce self phase modulation
called kerr effect

SPM Frequency Chirping



SPM

- Magnitude of nonlinear effect in SPM

$$\gamma = \frac{2\pi}{\lambda} \frac{n_2}{A_{eff}}$$

- The frequency shift

$$\Delta\omega = \frac{d\omega}{dt} = \gamma L_{eff} \frac{dP}{dt}$$

Cross Phase modulation

- Similar to SPM appear in WDM
- Refractive index nonlinearity converts optical intensity fluctuation in one wavelength channel to phase fluctuation in another channel

$$\Delta\varphi = \frac{d\varphi}{dt} = 2\gamma L_{eff} \frac{dP}{dt}$$

SPM

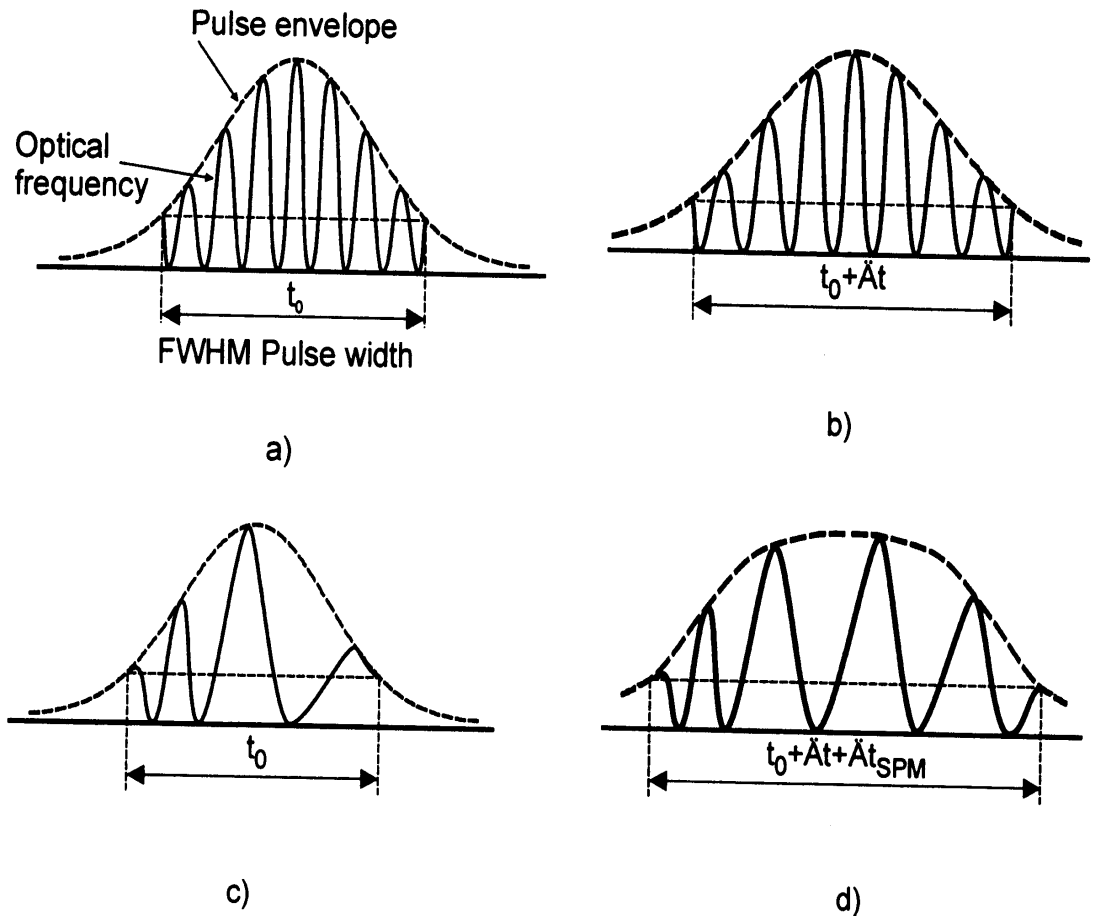
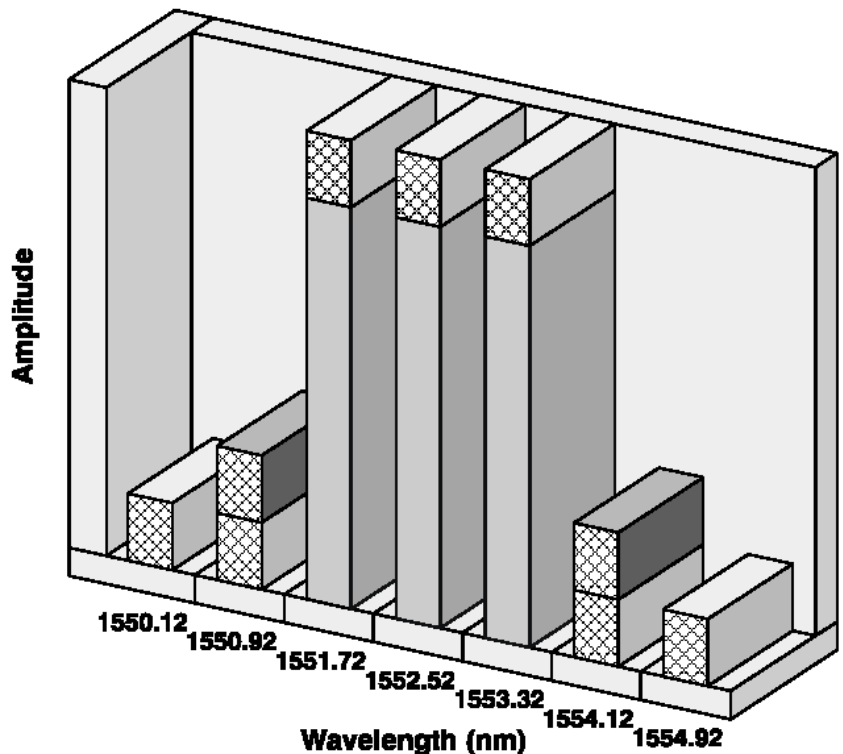


Figure 6.22 Self-phase modulation effect: spreading of chirped pulse: (a) Regular unchirped pulse entering the link; (b) the same pulse distorted after traveling distance L along the fiber; (c) chirped pulse entering the link; (d) chirped pulse broadens after traveling distance L .

Non Linear Effects: Four Wave Mixing (FWM) cont...

- Consider a simple three wavelength (λ_1 , λ_2 & λ_3)
- Let's assume that the input wavelengths are $\lambda_1 = 1551.72$ nm, $\lambda_2 = 1552.52$ nm & $\lambda_3 = 1553.32$ nm. The interfering wavelengths that are of most concern in our hypothetical three wavelength system are:

- $\lambda_1 + \lambda_2 - \lambda_3 = 1550.92$ nm
- $\lambda_1 - \lambda_2 + \lambda_3 = 1552.52$ nm
- $\lambda_2 + \lambda_3 - \lambda_1 = 1554.12$ nm
- $2\lambda_1 - \lambda_2 = 1550.92$ nm
- $2\lambda_1 - \lambda_3 = 1550.12$ nm
- $2\lambda_2 - \lambda_1 = 1553.32$ nm
- $2\lambda_2 - \lambda_3 = 1551.72$ nm
- $2\lambda_3 - \lambda_1 = 1554.92$ nm
- $2\lambda_3 - \lambda_2 = 1554.12$ nm



FWM

- K nonlinear interaction constant

$$P_{ijk}(L) = \eta(Dk)^2 P_i(0) P_j(0) P_k(0) \exp(-\alpha L)$$

$$K = \frac{32\pi^3 \chi_{1111}}{n_2 \lambda c} \left(\frac{Leff}{Aeff} \right)$$