Free particle in 1 D box:

Here free particle refers to electron and 1D box can be an atom or a piece of metal wire. As the electron is completely free to move within the box, its potential energy is zero. And the electron moves only along x-axis, the following partial differential equation (from Schrodinger's time independent equation) will be written as compete differential equation.

$$\nabla^2 \Psi + \frac{2m(E-V)}{\hbar^2} \Psi$$

As V=0 between the walls, the equation becomes

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{2m}{\hbar^2} E \Psi = 0$$

 $\frac{d^2\Psi}{dx^2} + k^2 = 0$ is the wave equation for electron's travel along one dimension.

x=0 and x=L are the boundaries of the one dimensional potential box. In the following figure the potential energy inside the box is zero; and the electron may face up to infinite potential by the walls of the box in such a way that the electron never gets outside the box. Thus, the potential function is defined in the following way:

$$V(x)=0$$
 for L>x>0
 $V(x)=\infty$ for x≥L, x≤0

$$: k^2 = \frac{2mE}{\hbar^2}$$

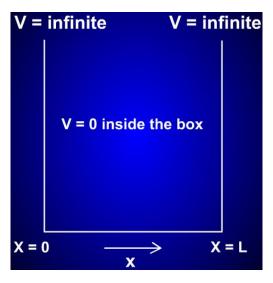
The solution of the second order differential equation is as follows,

$$\Psi(x) = A \sin kx + B \cos kx$$

The wave function is zero at the boundaries so,

$$\Psi = 0$$
 at $x = 0$

After applying the first boundary condition,



$$0 = A \times 0 + B \times 1$$

$$= 0 + B$$

$$\Rightarrow B = 0$$
Thus, $\Psi(x) = A \sin kx$

 $\Psi=0$ at x=a (the second boundary condition)

After applying the second boundary condition

$$0 = A \sin ka$$

$$ka = n\pi, n = 1, 2, 3, \dots$$

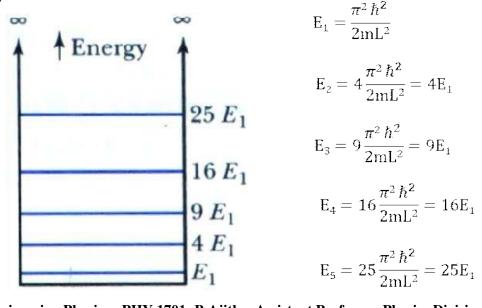
$$k = \frac{n\pi}{a}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}}, \text{ and also } k = \frac{n\pi}{a}$$

$$k^2a^2 = \frac{2mE_na^2}{\hbar^2} = n^2\pi^2$$

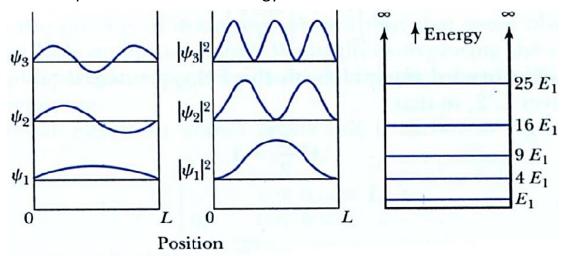
$$E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}, n = 1, 2, 3 \dots$$

Thus, for each value of n, the possible energy of the particle is given by the above equation



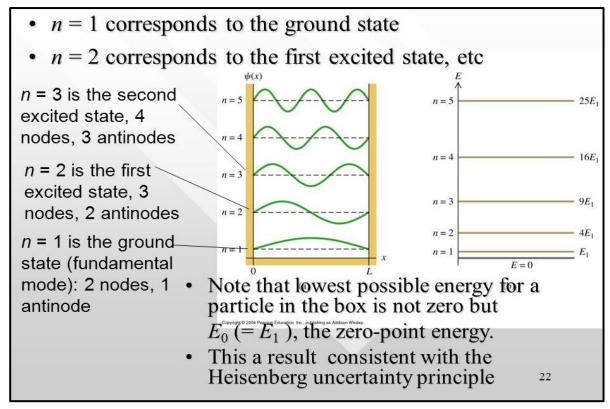
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The following figure explains the wave functions, probability densities and energy values of the particle at different energy levels.



At a particular point, the probability of the particle being present is different for different quantum numbers. For example, $|\Psi_1|^2$ has its maximum value at L/2 in the middle. Thus, a particle in the lowest energy level of n=1 is most likely to be in the middle while a particle in the next higher state of n=2 is never there.

Nodes and Antinodes:



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Eigen functions:

Eigen means special or particular. After including particular energy value function, the wave function is called as Eigen function.

$$\Psi_n=A\sin kx=A\sin\sqrt{\frac{2mE_n}{\hbar^2}}x$$
 , where, $n=1,2,3$
$$\Psi_n=A\sin\frac{n\pi x}{a}$$

 $\int_{-\infty}^{+\infty} |\Psi_n|^2 dx = 1$ (within the box, the hundred percentage of the particle is existing, since boundaries prevent them strongly from escaping outside)

$$\begin{split} \int\limits_0^a |\Psi_n|^2 dx &= \int\limits_0^a A^2 sin^2 \left(\frac{n\pi x}{a}\right) dx = 1 \\ \int\limits_0^a sin^2 \left(\frac{n\pi x}{a}\right) dx &= \frac{a}{2} \text{(standard integral)} \\ \int\limits_0^a |\Psi_n|^2 dx &= A^2 \frac{a}{2} = 1 \\ A^2 &= \frac{2}{a} \end{split}$$

$$A = \sqrt{\frac{2}{a}} \text{ is the normalised constant}$$

Thus by applying the normal condition (100% existence of the particle within the box) the constant A has been found out. So the constant A is called as normalised constant.

Normalised Eigen functions:

Eigen function with the normalised constant is called as normalized Eigen function.

$$\Psi_n(x) = A \sin kx = \sqrt{\frac{2}{a}} \sin \left(\frac{n\pi}{a}x\right)$$
, where, $n = 1,2,3...$