

# Computational Physics (PHYS6350)

Lecture 13: Classical mechanics problems

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_{i}}\right) - \frac{\partial L}{\partial q_{i}} = 0$$

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Instructor: Volodymyr Vovchenko (<a href="wvovchenko@uh.edu">wvovchenko@uh.edu</a>)

Course materials: <a href="https://github.com/vlvovch/PHYS6350-ComputationalPhysics/tree/spring2025">https://github.com/vlvovch/PHYS6350-ComputationalPhysics/tree/spring2025</a>

# Three-body problem

Exercise 8.10 (M. Newman, Computational Physics)

#### Three stars interacting through gravitational force

The equations of motion are

$$\frac{d^{2}\mathbf{r}_{1}}{dt} = Gm_{2} \frac{\mathbf{r}_{2} - \mathbf{r}_{1}}{|\mathbf{r}_{2} - \mathbf{r}_{1}|^{3}} + Gm_{3} \frac{\mathbf{r}_{3} - \mathbf{r}_{1}}{|\mathbf{r}_{3} - \mathbf{r}_{1}|^{3}},$$

$$\frac{d^{2}\mathbf{r}_{2}}{dt} = Gm_{1} \frac{\mathbf{r}_{1} - \mathbf{r}_{2}}{|\mathbf{r}_{1} - \mathbf{r}_{2}|^{3}} + Gm_{3} \frac{\mathbf{r}_{3} - \mathbf{r}_{2}}{|\mathbf{r}_{3} - \mathbf{r}_{2}|^{3}},$$

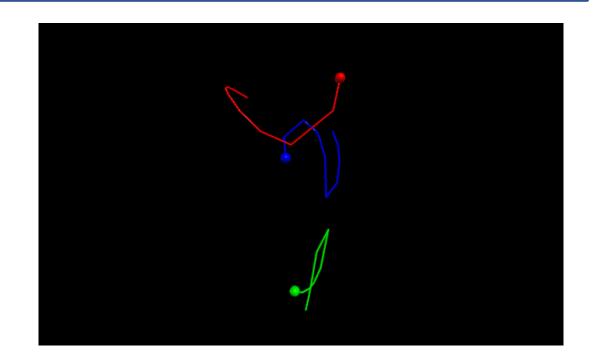
$$\frac{d^{2}\mathbf{r}_{3}}{dt} = Gm_{1} \frac{\mathbf{r}_{1} - \mathbf{r}_{3}}{|\mathbf{r}_{1} - \mathbf{r}_{3}|^{3}} + Gm_{2} \frac{\mathbf{r}_{2} - \mathbf{r}_{3}}{|\mathbf{r}_{2} - \mathbf{r}_{3}|^{3}}.$$

Name	Mass	X	у
Star 1	150.	3	1
Star 2	200.	-1	-2
Star 3	250.	-1	1

Cannot be solved analytically!

Take G = 1 (dimensionless)

Initially at rest and move in plane z=0Initial coordinates:  $\mathbf{r}_1=(3,1), \ \mathbf{r}_2=(-1,-2), \ \mathbf{r}_3=(-1,1)$ 



### Three-body problem

```
def fthreebody(xin, t):
    global f_evaluations
    f evaluations += 1
    x1 = xin[0]
   y1 = xin[1]
    x2 = xin[2]
    y2 = xin[3]
    x3 = xin[4]
   y3 = xin[5]
    r12 = np.sqrt((x1-x2)**2 + (y1-y2)**2)
    r13 = np.sqrt((x1-x3)**2 + (y1-y3)**2)
    r23 = np.sqrt((x2-x3)**2 + (y2-y3)**2)
    return np.array([xin[6],xin[7],xin[8],xin[9],xin[10],xin[11],
                    G * m2 * (x2 - x1) / r12**3 + G * m3 * (x3 - x1) / r13**3,
                    G * m2 * (y2 - y1) / r12**3 + G * m3 * (y3 - y1) / r13**3,
                    G * m1 * (x1 - x2) / r12**3 + G * m3 * (x3 - x2) / r23**3,
                    G * m1 * (y1 - y2) / r12**3 + G * m3 * (y3 - y2) / r23**3,
                    G * m1 * (x1 - x3) / r13**3 + G * m2 * (x2 - x3) / r23**3,
                    G * m1 * (y1 - y3) / r13**3 + G * m2 * (y2 - y3) / r23**3
```

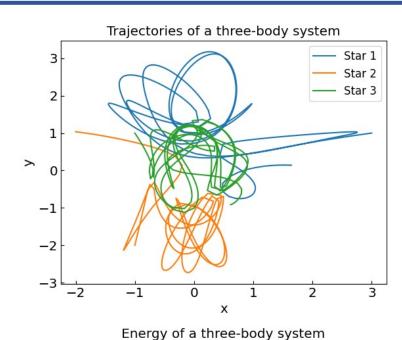
Run from t = 0 to 5

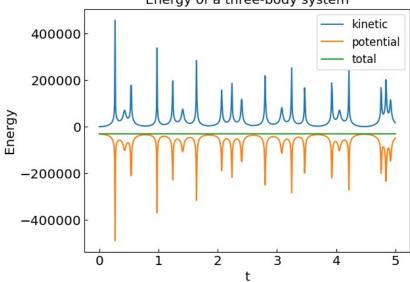
#### Three-body problem

```
def fthreebody(xin, t):
    global f_evaluations
    f evaluations += 1
    x1 = xin[0]
   y1 = xin[1]
   x2 = xin[2]
   y2 = xin[3]
   x3 = xin[4]
   y3 = xin[5]
   r12 = np.sqrt((x1-x2)**2 + (y1-y2)**2)
   r13 = np.sqrt((x1-x3)**2 + (y1-y3)**2)
    r23 = np.sqrt((x2-x3)**2 + (y2-y3)**2)
    return np.array([xin[6],xin[7],xin[8],xin[9],xin[10],xin[11],
                     G * m2 * (x2 - x1) / r12**3 + G * m3 * (x3 - x1) / r13**3,
                     G * m2 * (y2 - y1) / r12**3 + G * m3 * (y3 - y1) / r13**3,
                     G * m1 * (x1 - x2) / r12**3 + G * m3 * (x3 - x2) / r23**3,
                     G * m1 * (y1 - y2) / r12**3 + G * m3 * (y3 - y2) / r23**3,
                     G * m1 * (x1 - x3) / r13**3 + G * m2 * (x2 - x3) / r23**3,
                     G * m1 * (y1 - y3) / r13**3 + G * m2 * (y2 - y3) / r23**3
```

Run from t = 0 to 5

final project idea(?): simulate the solar system





# Lagrangian mechanics

In Lagrangian mechanics, a classical system with N degrees of freedom is characterized by generalized coordinates  $q_i$  that obey Euler-Lagrange equations of motion:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_j} = 0, \qquad j = 1 \dots N$$

L is the non-relativistic Lagrangian which is typically defined as the difference of kinetic and potential energies, L = T - V.

One can rewrite these equations using the chain rule:

$$\sum_{i=1}^{N} \frac{\partial^{2} L}{\partial \dot{q}_{j} \partial \dot{q}_{i}} \ddot{q}_{i} = -\sum_{i=1}^{N} \frac{\partial^{2} L}{\partial \dot{q}_{j} \partial q_{i}} \dot{q}_{i} - \frac{\partial^{2} L}{\partial \dot{q}_{j} \partial t} + \frac{\partial L}{\partial q_{j}}, \qquad j = 1 \dots N.$$

This is a system of N linear equations for  $\ddot{q}_i$ . When solved we have

$$\ddot{q}_i = f_i(\{q_j\}, \{\dot{q}_j\}, t).$$

or

$$\begin{aligned} \frac{dq_i}{dt} &= \dot{q}_i, \\ \frac{d\dot{q}_i}{dt} &= f_i(\{q_j\}, \{\dot{q}_j\}, t), \end{aligned}$$

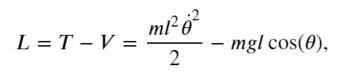
which can be solved using standard methods.

# Non-linear pendulum

The generalized coordinate is the displacement angle heta

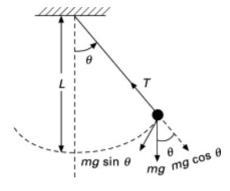
$$x = l\sin(\theta),$$
  
$$y = -l\cos(\theta),$$

Lagrangian

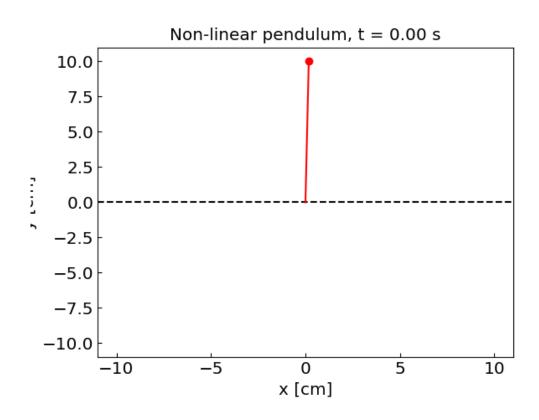


Equations of motion

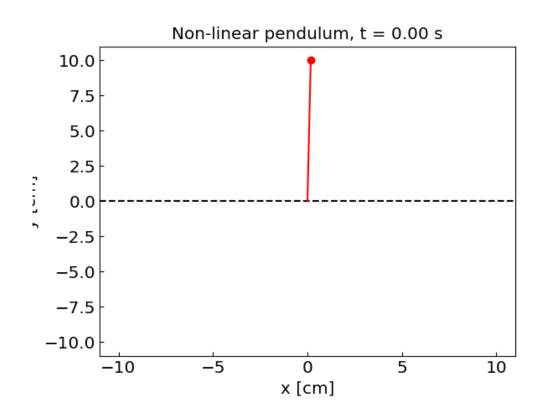
$$ml\ddot{\theta} = -mgl\sin\theta.$$

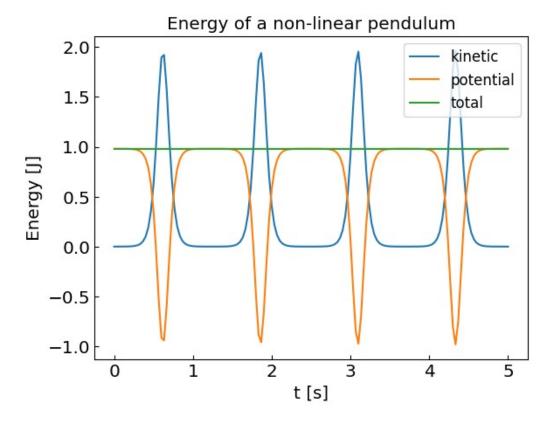


# Non-linear pendulum



# Non-linear pendulum





Double pendulum is the simplest system exhibiting chaotic motion – deterministic chaos

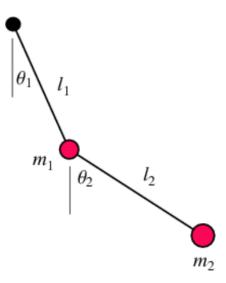
Two degrees of freedom: the displacement angles  $\theta_1$  and  $\theta_2$ 

```
x_1 = l_1 \sin(\theta_1),

y_1 = -l_1 \cos(\theta_1),

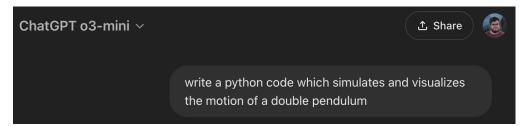
x_2 = l_1 \sin(\theta_1) + l_2 \sin(\theta_2),

y_2 = -l_1 \cos(\theta_1) - l_2 \cos(\theta_2).
```

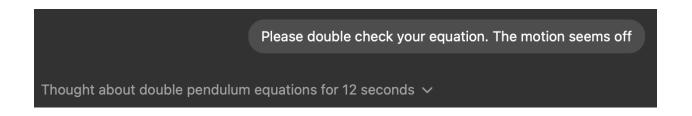


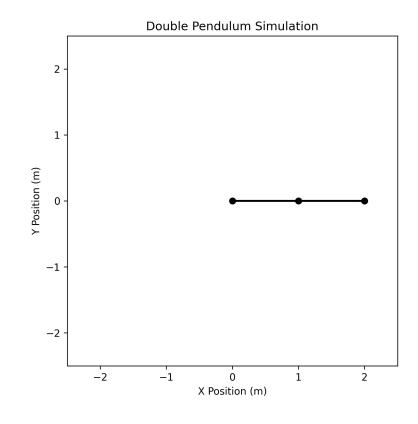
# Double pendulum vs ChatGPT

#### First approach nowadays: try LLM



**Double Pendulum Simulation** 2 Y Position (m) -1-2 -2 -1X Position (m)





Double pendulum is the simplest system exhibiting chaotic motion – deterministic chaos

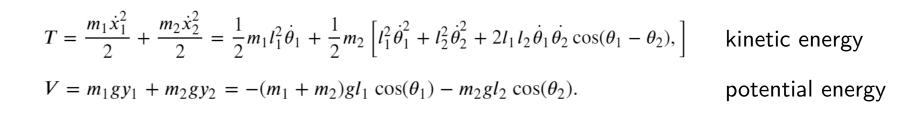
Two degrees of freedom: the displacement angles  $\theta_1$  and  $\theta_2$ 

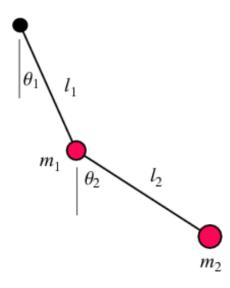
$$x_1 = l_1 \sin(\theta_1),$$
  

$$y_1 = -l_1 \cos(\theta_1),$$
  

$$x_2 = l_1 \sin(\theta_1) + l_2 \sin(\theta_2),$$
  

$$y_2 = -l_1 \cos(\theta_1) - l_2 \cos(\theta_2).$$





Double pendulum is the simplest system exhibiting chaotic motion – deterministic chaos

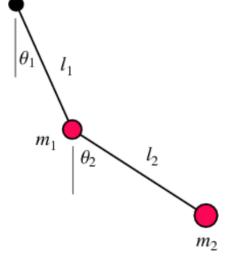
Two degrees of freedom: the displacement angles  $\theta_1$  and  $\theta_2$ 

$$x_1 = l_1 \sin(\theta_1),$$
  

$$y_1 = -l_1 \cos(\theta_1),$$
  

$$x_2 = l_1 \sin(\theta_1) + l_2 \sin(\theta_2),$$
  

$$y_2 = -l_1 \cos(\theta_1) - l_2 \cos(\theta_2).$$



$$T = \frac{m_1 \dot{x}_1^2}{2} + \frac{m_2 \dot{x}_2^2}{2} = \frac{1}{2} m_1 l_1^2 \dot{\theta}_1 + \frac{1}{2} m_2 \left[ l_1^2 \dot{\theta}_1^2 + l_2^2 \dot{\theta}_2^2 + 2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \cos(\theta_1 - \theta_2), \right]$$
 kinetic energy 
$$V = m_1 g y_1 + m_2 g y_2 = -(m_1 + m_2) g l_1 \cos(\theta_1) - m_2 g l_2 \cos(\theta_2).$$
 potential energy

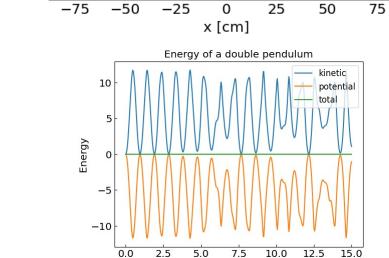
The Lagrange equations of motion read

$$(m_1 + m_2)l_1\ddot{\theta}_1 + m_2l_2\cos(\theta_1 - \theta_2)\ddot{\theta}_2 = -m_2l_1\dot{\theta}_2^2\sin(\theta_1 - \theta_2) - (m_1 + m_2)g\sin(\theta_1),$$
  

$$m_2l_1\cos(\theta_1 - \theta_2)\ddot{\theta}_1 + m_2l_2\ddot{\theta}_2 = m_2l_1\dot{\theta}_1^2\sin(\theta_1 - \theta_2) - m_2g\sin(\theta_2).$$

This is a system of two linear equations for  $\ddot{\theta}_{1,2}$  that can be solved straightforwardly.

```
g = 9.81
11 = 0.4
12 = 0.4
m1 = 1.0
m2 = 1.0
def fdoublependulum(xin, t):
    global f evaluations
    f_evaluations += 1
    theta1 = xin[0]
    theta2 = xin[1]
    omega1 = xin[2]
    omega2 = xin[3]
    a1 = (m1 + m2)*11
    b1 = m2*12*np.cos(theta1 - theta2)
    c1 = m2*12*omega2*omega2*np.sin(theta1 - theta2) + (m1 + m2)*g*np.sin(theta1)
    a2 = m2*l1*np.cos(theta1 - theta2)
    b2 = m2*12;
    c2 = -m2*l1*omega1*omega1*np.sin(theta1 - theta2) + m2*g*np.sin(theta2) # + k
    domega1 = - (c2/b2 - c1/b1) / (a2/b2 - a1/b1)
    domega2 = - (c2/a2 - c1/a1) / (b2/a2 - b1/a1)
    return np.array([omega1,
                     omega2,
                     domega1,
                     domega2
```



t [s]

Double pendulum, t = 0.00 s

80

60

40

20

-20

-40

-60

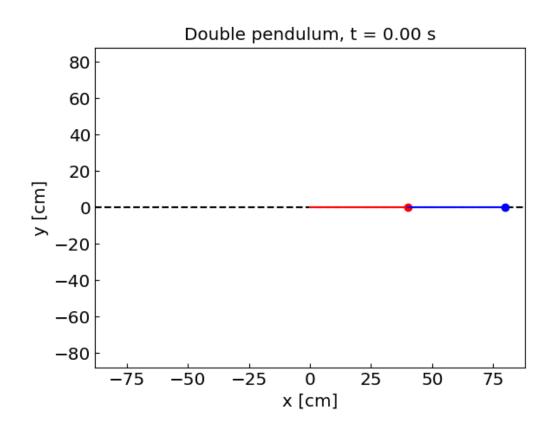
-80

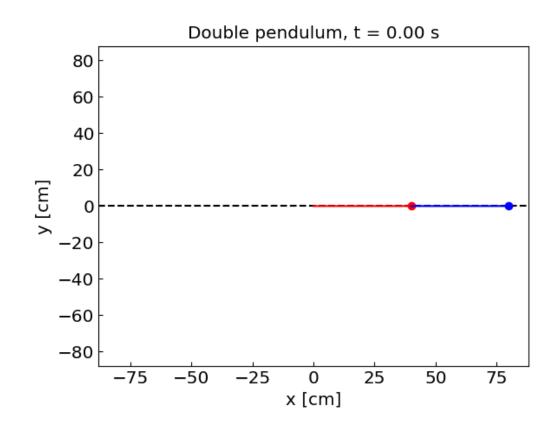
y [cm]

# Double pendulum: chaotic behavior

$$\theta_1^0 = \theta_2^0 = \pi/2$$

$$\theta_1^0 = \theta_2^0 = \pi/2 + 10^{-4}$$





# Double pendulum: chaotic behavior

$$\theta_1^0 = \theta_2^0 = \pi/2$$

$$\theta_1^0 = \theta_2^0 = \pi/2 + 10^{-4}$$

