## Jacobi's Principle

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We start with principle of least action, viz.

$$\delta \int_{t_1}^{t_2} L(q_i, \dot{q}_i) dt = 0 \tag{1}$$

for the path that the particle takes. In our case, we are going to introduce an arbitrary parameter  $\tau$  and treat the time as one of the position coördinates. Primes denoting differentiation<sup>1</sup> with respect to  $\tau$ , we have

$$\delta \int_{\tau_1}^{\tau_2} \left( L\left(q_i, \frac{q_i'}{t'}\right) t' \right) d\tau = 0 \tag{2}$$

Now we can apply the Euler-Lagrange equations for t to this integral to get

$$\frac{\partial(Lt')}{\partial t} - \frac{d}{d\tau} \left( \frac{\partial(Lt')}{\partial t'} \right) = 0. \tag{3}$$

Since t does not appear in the integrand,  $\frac{\partial(Lt')}{\partial t} = 0$ , and thus  $\frac{\partial(Lt')}{\partial t'}$  is constant with respect to  $\tau$ .

The generalized momentum  $p_i$  for a variable  $q_i$  will be defined as  $\frac{\partial L}{\partial \dot{q}_i}$ ; in our case we also consider the momentum associated with time, viz.  $p_t = \frac{\partial (Lt')}{\partial t'}$ . Now,

$$p_{t} = \frac{\partial(Lt')}{\partial t'}$$

$$= L - \left(\sum_{i=1}^{n} \frac{\partial L}{\partial \dot{q}_{i}} \frac{q'_{i}}{t'^{2}}\right) t'$$

$$= \left(L - \sum_{i=1}^{n} \frac{\partial L}{\partial \dot{q}_{i}} \dot{q}_{i}\right)$$

$$= \left(L - \sum_{i=1}^{n} p_{i} \dot{q}_{i}\right)$$

$$= -E.$$

 $<sup>^1</sup>$  Actually, the  $q_i'$  are independent of the  $q_i$  but we add the auxiliary conditions  $\frac{dq_i}{d\tau}=q_i'$ 

where E is the energy of the system.

Forming  $\bar{L} = Lt' - p_t t'$ , we will show that

$$\delta \int_{\tau_1}^{\tau_2} \bar{L} d\tau = 0 \tag{4}$$

Now,

$$\delta \int_{\tau_1}^{\tau_2} p_t t' d\tau. \tag{5}$$

will be true if and only if the Euler-Lagrange equations hold for all the variables. This is clearly true, since

$$\frac{\partial(p_t t')}{\partial q_i} - \frac{d}{d\tau} \left( \frac{\partial(p_t t')}{\partial q_i'} \right) = 0 \tag{6}$$

 $p_t$  being constant and the  $q_i, q'_i$  being independent of t'. Also,

$$\frac{\partial(p_t t')}{\partial t} - \frac{d}{d\tau} \left( \frac{\partial(p_t t')}{\partial t'} \right) = 0 \tag{7}$$

since  $\frac{\partial (p_t t')}{\partial t'} = p_t$ , which is constant.

We already know  $\delta \int_{\tau_1}^{\tau_2} Lt' d\tau = 0$ , so, adding the two, we get  $\delta \int_{\tau_1}^{\tau_2} \bar{L} d\tau = 0$ . From above,

$$\bar{L} = Lt' - p_t t' = (L - p_t)t' = \left(\sum_{i=1}^n p_i \dot{q}_i\right)t',$$
 (8)

so

$$\int_{\tau_1}^{\tau_2} \bar{L} d\tau = 2 \int_{\tau_1}^{\tau_2} T t' d\tau. \tag{9}$$

It was pointed out by Jacobi that this cannot be simplified to  $\int_{\tau_1}^{\tau_2} \bar{L} d\tau = 2 \int_{t_1}^{t_2} T dt$ , because t cannot be treated as an independent variable in the variational problem. Instead, we take advantage of the fact that

$$T = \frac{1}{2} \left( \frac{ds}{dt} \right)^2 \tag{10}$$

or

$$T = \frac{1}{2} \frac{\left(\frac{ds}{d\tau}\right)^2}{t'^2} \tag{11}$$

Since T = E - V,

$$t' = \frac{1}{\sqrt{2(E-V)}} \frac{ds}{d\tau},\tag{12}$$

giving finally

$$2\int_{\tau_1}^{\tau_2} Tt' d\tau = \int_{\tau_1}^{\tau_2} \sqrt{2(E-V)} \frac{ds}{d\tau} d\tau = \int_{\tau_1}^{\tau_2} \sqrt{2(E-V)} ds, \tag{13}$$

as desired. Recalling that  $\delta \int_{\tau_1}^{\tau_2} \bar{L} d\tau = 0$ , we see that

$$\delta \int_{\tau_1}^{\tau_2} \sqrt{2(E-V)} ds = 0 \tag{14}$$

as well. This condition determines the particle's path, and is known as Jacobi's principle. Note that this determines the path a particle takes in space, but says nothing about time.

Fermat's principle of least time states that light takes the path that takes the least time, that is,

$$\delta \int dt = 0 \tag{15}$$

Rewriting dt as nds, n being the index of refraction, we get

$$\delta \int nds = 0. \tag{16}$$

As can be seen, this bears a striking resemblance to Jacobi's principle. If a material has an index of refraction n(x,y,z), then light will travel through it the same way a particle would be affected by a potential field V(x,y,x) provided that  $n=\sqrt{2(E-V)}$ . This is not the entirety of the optico-mechanical analogy, but it is part of it. The optico-mechanical analogy was used in the development of the old quantum theory, including de Broglie's Nobel Prizewinning dissertation on matter waves.