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## Notes on the calculation of Transition Density Matrices (TDM)

In the molecular orbital (MO) basis, the TDM between excited states **I** and **J** can be written as:

$$\rho_{\text{MO}}^{IJ} = \langle 4_I | a_p^\dagger a_q | 4_J \rangle, \quad (1)$$

where  $a_p^\dagger$  and  $a_q$  are creation and annihilation operators acting on the virtual MO  $p$  and the occupied MO  $q$ , and the excitation states wavefunctions are given by:

$$|4_J\rangle = \sum_{ia} X_{ia}^{(J)} |\Xi_i^a\rangle, \quad (2)$$

where  $|\Xi_i^a\rangle$  is a single Slater determinant representing the transition from the occupied MO  $i$  to the virtual orbital  $a$ , which can be cast in terms of the ground state (GS) solution  $|\Xi_i\rangle$  (also given by a single Slater determinant) as:

$$|\Xi_i^a\rangle = a_a^\dagger a_i |\Xi_i\rangle, \quad (3)$$

and  $X_{ia}^{(J)}$  are the corresponding configuration interaction singles (CIS) coefficients for excited state **J**. Analogously we have:

$$|4_I\rangle = \sum_{jb} X_{jb}^{(I)} |\Xi_j^b\rangle, \quad (4)$$

and:

$$|\Xi_j^b\rangle = a_b^\dagger a_j |\Xi_j\rangle$$

Combining eqs 1-5 we get:

$$\langle 4_I | a_p^\dagger a_q | 4_J \rangle = \sum_{ia} \sum_{jb} (X_{jb}^{(I)})^* \langle \Xi_b | a_b^\dagger a_p^\dagger a_q a_i^\dagger | \Xi_i \rangle X_{ia}^{(J)}, \quad (5)$$

The string of creation and annihilation operators can be calculated by means of the Wick's theorem, which allows to write it as a summation of the possible three pairs of concatenated creator and annihilator operators. Contractions including  $a_j^\dagger a_b$ ,  $a_p^\dagger a_q$ , or  $a_i^\dagger a_i$  render transitions to the same MO which have zero weight in the excited state.

expansion. The only non-trivial contributions are:

$$\overbrace{a_j^+ a_b^+ a_p^+ a_q^+} \overbrace{a_a^+ a_i^+} \rightarrow (-1)^3 \delta_{jq} \delta_{ab} \delta_{ip}, \quad (7)$$

and:

$$\overbrace{a_j^+ a_b^+ a_p^+ a_q^+} \overbrace{a_a^+ a_i^+} \rightarrow (-1)^4 \delta_{ij} \delta_{bp} \delta_{qa}, \quad (8)$$

From eq. 7 we get for the occupied block of the TDM:

$$\rho_{MO_{ij}}^{IJ} = - \sum_a (x_{ja}^{(I)})^* x_{ia}^{(J)}, \quad (9)$$

From eq. 8 we get for the virtual block of the TDM:

$$\rho_{MO_{ba}}^{IJ} = \sum_i (x_{ib}^{(I)})^* x_{ia}^{(J)}, \quad (10)$$

Taking into account that the TDM for transitions from the GS to any given state can be cast as:

$$\rho_{MO_{ia}}^{(OJ)} = \begin{pmatrix} 0_{oo} & X_{or}^{(J)} \\ 0_{vo} & 0_{vr} \end{pmatrix}, \quad (11)$$

where the blocks  $0_{oo}$ ,  $0_{vo}$ , and  $0_{vr}$  are empty block representing occupied-occupied, virtual-occupied, and virtual-virtual contributions, respectively. Then we can write eqs. 9-10 in a simplified matrix notation:

$$\rho_{MO}^{IJ} = ((\rho_{MO}^{OI})^T \rho_{MO}^{OJ} - \rho_{MO}^{OV} (\rho_{MO}^{OJ})^T), \quad (12)$$

which takes the same form in the atomic orbital (AO) basis:

$$\rho_{AO}^{IJ} = ((\rho_{AO}^{OI})^T \rho_{AO}^{OJ} - \rho_{AO}^{OV} (\rho_{AO}^{OJ})^T), \quad (13)$$