Companion software for "Volker Ziemann, *Hands-on Accelerator physics using MATLAB, CRCPress, 2019*" (https://www.crcpress.com/9781138589940)

Pill-box cavity with beam pipe (Section 6.4.2)

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Important: this example requires the PDE toolbox!

In this example we calculate the eigenfrequencies of a pill-box cavity with a beam pipe connected to its center. This will perturb the eigenmodes and eigenfrequencies of a pure cavity. We therefore use the PDE toolbox to find the TM eigenmodes and note that the system is still rotationally symmetric such that we can use a 2D model to find the rotationally symmetric modes.

We first define the beam pipe radius rr, the cavity radius R, and the cavity length L to specify the cavity dimensions.

```
clear;
rr=0.01 % Slider for the beam pipe radius

rr = 0.0100

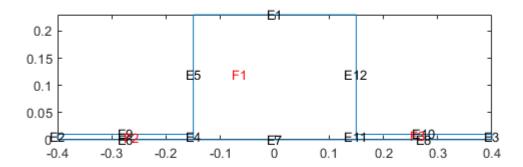
R=0.23 % Slider for the cavity radius

R = 0.2300

L=0.3; % Cavity length
```

Now we use these parameters to define the cavity and the beam pipe as three polygons Cav, Pl, and Pr, assemble them in gd, define names ns and specify the relation sf among the regions, before we construct the geometry g and create the model for one parameter $u = E_z$. Then we add the geometry g to it and inspect the model to note the EdgeLabels and the SubDomainLabels.

```
Cav = [2;4; -L/2;L/2;L/2;-L/2;
                              0.0;0.0;R;R];
                                                   % cavity
Pl=[2;4; -0.4;-L/2;-L/2;-0.4; 0.0;0.0;rr;rr];
                                                  % pipe left
Pr=[2;4; L/2;0.4;0.4;L/2; 0.0;0.0;rr;rr];
                                                   % pipe right
gd=[Cav,Pl,Pr];
                               % assemble geometry
ns=char('Cavity','Pl','Pr')'; % names of the regions
sf='Cavity+Pl+Pr';
                              % geometry relations
g=decsg(gd,sf,ns);
model=createpde(1);
geometryFromEdges(model,g);
pdegplot(model, 'EdgeLabels', 'on', 'SubDomainLabels', 'on'); axis tight
```



Since the longitudinal electric field $u=E_z$ must vanish on the boundaries where it is parallel to the boundaries E1, E9, E10 and on the exit from the beam pipe E2 and E3, we select Dirichlet boundary conditions. On the verical walls E5 and E12 the electric field is given by the derivative of the longitudinal field E_z , which implies the use of von-Neumann boundary conditions. Moreover, we also use von-Neumann conditions on the symmetry axis E6, E7, and E8.

The rotational symmetry requires to use the Laplace operator in cylindrical coordinates, which has to be brought into Matlab's generic form to specify a partial differential equation in divergence form, which is discussed further in Section 6.4.2. In the code this is accomplished by using functions cfun() and dfun() when specifying the material properties.

```
cfun=@(location,state)-location.y;
dfun=@(location,state)-location.y;
specifyCoefficients(model,'m',0,'d',dfun,'c',cfun,'a',0,'f',0);
```

At this point we can mesh the geometry...

```
generateMesh(model,'Hmax',0.01);
%figure; pdemesh(model); axis tight
```

...and solve the eigenvalue system with <code>solvepdeeig()</code>, which returns the structure <code>result</code> that contains the <code>Eigenvalues</code> and the <code>Eigenvectors</code> as members.

```
result=solvepdeeig(model,[1,800]);

Basis= 10, Time= 0.09, New conv eig= 0
Basis= 12, Time= 0.11, New conv eig= 0
Basis= 14, Time= 0.12, New conv eig= 2
Basis= 16, Time= 0.12, New conv eig= 3
Basis= 18, Time= 0.13, New conv eig= 3
Basis= 20, Time= 0.13, New conv eig= 4
Basis= 22, Time= 0.14, New conv eig= 5
End of sweep: Basis= 22, Time= 0.14, New conv eig= 5
Basis= 15, Time= 0.17, New conv eig= 0
End of sweep: Basis= 15, Time= 0.17, New conv eig= 0

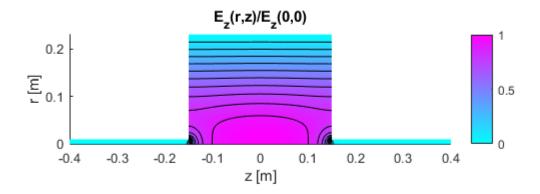
k2=result.Eigenvalues;
Ez=result.Eigenvectors;
```

From the eigenvectors, which are $k^2 = \omega^2/c^2$ we derive the frequencies of the eigenmodes in GHz.

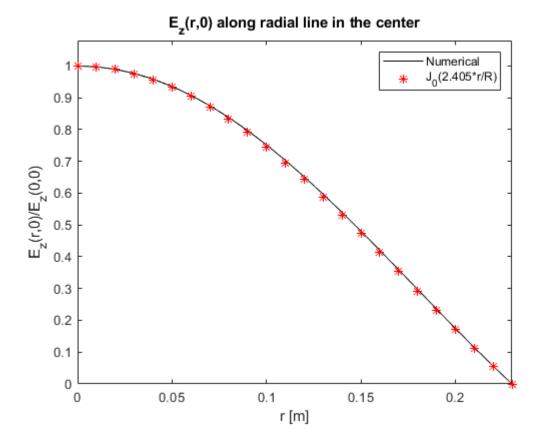
```
freq_GHz=sqrt(k2')*3e8/(2*pi*1e9)

freq_GHz = 1x5
    0.5138    0.7271    1.1253    1.1655    1.2768
```

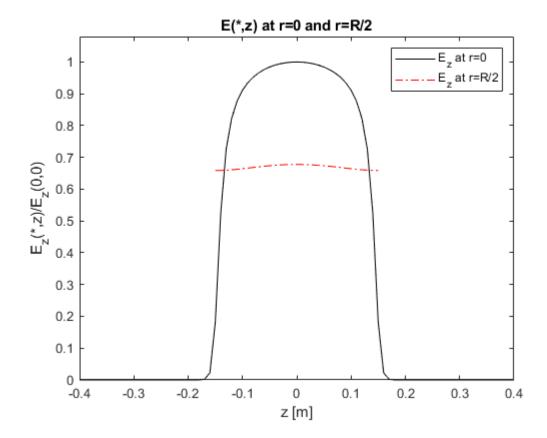
Since we are interested in the accelerating mode around 500 MHz, we choose to inspect the lowest-frequency mode=1 and first use meshToPet() to obtain information about the points p, edges e, and triagnles t of the mesh, such that we can use pdeInterpolant() to prepare a interpolant structure F that allows us to find the field at points that are not on grid points. We use it to find the field Ez0 in the center of the cavity and then use that to normalize $u = E_z$ to its maximum value at the center and display it with equipotential lines superimposed. We observe that the equipotential lines at large radius are almost parallel to the z-axis, which is what we would also observe in the analytical solution from Section 5.1. At smaller radii, however, we find the equipotential lines successively more distorted at smaller radii.



In the next figure we use the interpolant ${\mathbb F}$ to evaluate E_z at z=0 along the radial direction and compare this to the analytical result from Equation 5.21 $E_z \propto J_0(2.405r/R_{cav})$, where 2.405 is the first zero of $J_0(x)$. The figure shows good agreement between the two curves, because the center of the cavity is far away from the beam pipe entry points.



Finally, we calculate the accelerating field E_z that the beam 'sees', which is along the center line at r=0. We also calculate E_z at half the cavity radius R/2. For comparison we note that the analytic solution from Equation 5.21 for $p=k_zL=0$ is constant along the z-axis. Here, on the other hand, we see a significant distortion. Especially at r=0 the field decreases towards the cavity ends at $z=\pm L/2$ and even 'leaks' somewhat into the beam pipe. The fields E_z at R/2, shown as the red line, is not constant either, but much less distorted.



And now it's time to play around with this model. First you may increase the beam pipe radius and second, you might want to change the radius *R* of the cavity to tune it to 500 MHz. In the present configuration it is a bit off-frequency.