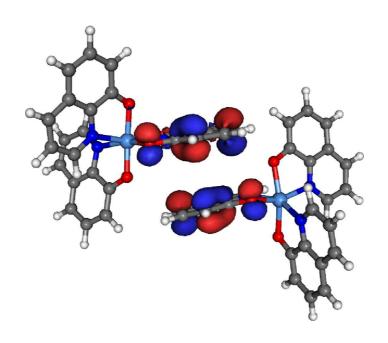# VOTCA-CTP

CHARGE TRANSPORT SIMULATIONS

USER MANUAL

compiled from: 1.5-dev (gitid: 96ef69f)

December 25, 2018

www.votca.org

ii

Disclamer

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Citations

Development of this software depends on academic research grants. If you are using the package, please cite the following papers

[1] Long-range embedding of molecular ions and excitations in a polarizable molecular environment, Carl Poelking and Denis Andrienko
J. Chem. Theory Comp. 12, 4516, 2016

[2] Modeling of spatially correlated energetic disorder in organic semiconductors, Pascal Kordt, Denis Andrienko
J. Chem. Theory Comput., 12, 36, 2016

[3] Microscopic simulations of charge transport in disordered organic semiconductors, Victor Rühle, Alexander Lukyanov, Falk May, Manuel Schrader, Thorsten Vehoff, James Kirkpatrick, Björn Baumeier and Denis Andrienko
J. Chem. Theor. Comp. 7, 3335, 2011

[4] Extracting nondispersive charge carrier mobilities of organic semiconductors from simulations of small systems, A. Lukyanov, D. Andrienko
Phys. Rev. B, 82, 193202, 2010

[5] Density-functional based determination of intermolecular charge transfer properties for large-scale morphologies, Björn Baumeier, James Kirkpatrick, and Denis Andrienko
Phys. Chem. Chem. Phys. 12, 11103, 2010

[6] Versatile Object-oriented Toolkit for Coarse-graining Applications, Victor Rühle, Christoph Junghans, Alexander Lukyanov, Kurt Kremer and Denis Andrienko
J. Chem. Theor. Comp. 5, 3211, 2009

[?] An approximate method for calculating transfer integrals based on the ZINDO Hamiltonian, James Kirkpatrick,
Int. J. Quantum Chem. 108, 51, 2007

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Chapter 1

Introduction

Charge carrier dynamics in an organic semiconductor can often be described in terms of charge hopping between localized states. The hopping rates depend on electronic coupling elements, reorganization energies, and site energies, which vary as a function of position and orientation of the molecules. The purpose of the VOTCA-CTP package [3] is to simplify the workflow for charge transport simulations, provide a uniform error-control for the methods, flexible platform for their development, and eventually allow *in silico* prescreening of organic semiconductors for specific applications.

The toolkit is implemented using modular concepts introduced earlier in the Versatile Object-oriented Toolkit for Coarse-graining Applications (VOTCA) [6]. It contains different programs, which execute specific tasks implemented in calculators representing an individual step in the workflow. Figure 1.1 summarizes a typical chain of commands to perform a charge transport simulation: First, the VOTCA code structures are adapted to reading atomistic trajectories, mapping them onto conjugated segments and rigid fragments, and substituting (if needed) rigid fragments with the optimized copies (`ctp_map`). The programs `ctp_run` and `ctp_parallel` (for heavy-duty tasks) are then used to calculate all bimolecular charge hopping rates (via precalculation of all required ingredients). Site energies (or energetic disorder) can be determined as a combination of internal (ionization potentials/electron affinities of single molecules) as well as electrostatic and polarization contributions within the molecular environment. The calculation of electronic coupling elements between conjugated segments from the corresponding molecular orbitals can be performed using a dimer-projection technique based on density-functional theory (DFT). This requires explicit calculations using quantum-chemistry software for which we provide interfaces to `Gaussian`, `Turbomole`, and `NWChem`. Alternatively, the molecular orbital overlap module calculates electronic coupling elements relying on the semi-empirical INDO Hamiltonian and molecular orbitals in the format provided by the `Gaussian` package.

The kinetic Monte Carlo module reads in the neighbor list, site coordinates, and hopping rates and performs charge dynamics simulations using either periodic boundary conditions or charge sources and sinks.

The toolkit is written as a combination of modular C++ code and scripts. The data transfer between programs is implemented via a state file (sql database), which is also used to restart simulations. Analysis functions and most of the calculation routines are encapsulated by using the observer pattern [7] which allows the implementation of new functions as individual modules.

In the following chapter 2, we summarize the theoretical background of the workflow of charge transport simulations and in particular its individual steps. Chapter 3 describes the structure and content of input and output files, while a full reference of programs and calculators is available in chapter 4. For a hands-on tutorial, the reader is referred to the VOTCA-CTP project page at http://code.google.com/p/votca-ctp/.

Input files:
`conf.gro`
 GROMACS trajectory
`topol.tpr`
 GROMACS topology
`map.xml`
 mapping and energies
`options.xml`
 options for calculators

Output files:
`state.sql`
 sqlite3 database file for
 data transfer between
 modules

Mapping, generation of the sql database
Converts and partitions atomistic GROMACS trajectory

`ctp_map  -t topol.tpr  -c traj.xtc  -s map.xml  -f state.sql`
Useful tools:
toolpdb2map

Neighbor list
Indentifies close molecular pairs between which charge transfer rates will be calculated

`ctp_run  -o options.xml  -f state.sql  -e neighborlist`

Site energies
Calculates electrostatic and polarization contribution to site energies

`ctp_run  -o options.xml  -f state.sql  -e emultipole`

Internal site and reorganization energies
Imports internal site energy (IP, EA) and reorganization energies for charging and discharging to `state.sql`

`ctp_run  -o options.xml  -f state.sql  -e einternal`

DFT   Transfer integrals   ZINDO

Monomers with DFT
Calculate the relevant transport orbitals of monomers

`ctp_parallel  -o options.xml  -f state.sql`
`-e edft  -j "write run"`

Transfer integrals with ZINDO
Calculate electronic coupling elements for all pairs in the neighbor list

`ctp_run  -o options.xml  -f state.sql  -e`
`izindo`

One can choose between quantum-chemical (computationally expensive) or semi-empirical (fast, but not always sufficiently accurate) evaluation of transfer integrals.

Transfer integrals with DFT
Calculate electronic coupling elements for all pairs in the neighbor list

`ctp_parallel  -o options.xml  -f state.sql`
`-e idft  -j "write run read"`

Outersphere reorganization energies
Contribution to reorganization of surrounding molecules due to polarization. (optional for Marcus rates)

`ctp_run  -o options.xml  -f state.sql  -e outersphere`

Charge transfer rates
Calculates rates for charge transfer among all pairs in the neighborlist

`ctp_run  -o options.xml  -f state.sql  -e rates`

Charge dynamics via kMC
Hopping of charge carriers simulated via kinetic Monte Carlo

`kmc_run  -o options.xml  -f state.sql  -e kmcmultiple`

Get list of available calculators: `ctp_run`/`ctp_parallel`/`kmc_run  -l`
Get help and list of options for a calculator: `ctp_run`/`ctp_parallel`/`kmc_run  -d neighborlist`

Figure 1.1: A practical workflow of charge transport simulations using VOTCA-CTP. The theoretical background of the individual steps is given in chapter 2. Chapter 3 describes the content of input and output files, while a full reference of programs and calculators is available in chapter 4. fig:summary

Chapter 2

Theoretical background

2.1 Workflow

A typical workflow of charge transport simulations is depicted in figure 2.1. The first step is the simulation of an atomistic morphology, which is then partitioned on hopping sites. The coordinates of the hopping sites are used to construct a list of pairs of molecules, or neighbor list.

Figure 2.1: Workflow for microscopic simulations of charge transport.

For each pair an electronic coupling element, a reorganization energy, a driving force, and eventually the hopping rate are evaluated. The neighbor list and hopping rates define a directed graph. The corresponding master equation is solved using the kinetic Monte Carlo method, which allows to explicitly monitor the charge dynamics in the system as well as to calculate time or ensemble averages of occupation probabilities, charge fluxes, correlation functions, and field-dependent mobilities.

2.2 Material morphology

There is no generic recipe on how to predict a large-scale atomistically-resolved morphology of an organic semiconductor. The required methods are system-specific: for ultra-pure crystals, for

Figure 2.2: The concept of conjugated segments and rigid fragments. Dashed lines indicate conjugated segments while colors denote rigid fragments. (a) Hexabenzocoronene: the π-conjugated system is both a rigid fragment and a conjugated segment. (b) Alq_3: the Al atom and each ligand are rigid fragments while the whole molecule is a conjugated segment. (c) Polythiophene: each repeat unit is a rigid fragment. A conjugated segment consists of one or more rigid fragments. One molecule can have several conjugated segments.

example, density-functional methods can be used provided the crystal structure is known from experiment. For partially disordered organic semiconductors, however, system sizes much larger than a unit cell are required. Classical molecular dynamics or Monte Carlo techniques are then the methods of choice.

In molecular dynamics, atoms are represented by point masses which interact via empirical potentials prescribed by a force-field. Force-fields are parametrized for a limited set of compounds and their refinement is often required for new molecules. In particular, special attention shall be paid to torsion potentials between successive repeat units of conjugated polymers or between functional groups and the π-conjugated system. First-principles methods can be used to characterize the missing terms of the potential energy function.

Self-assembling materials, such as soluble oligomers, discotic liquid crystals, block copolymers, partially crystalline polymers, etc., are the most complicated to study. The morphology of such systems often has several characteristic length scales and can be kinetically arrested in a thermodynamically non-equilibrium state. For such systems, the time- and length-scales of atomistic simulations might be insufficient to equilibrate or sample desired morphologies. In this case, systematic coarse-graining can be used to enhance sampling [6]. Note that the coarse-grained representation must reflect the structure of the atomistic system and allow for back-mapping to the atomistic resolution.

Here we assume that the morphology is already known, that is we know how the topology and the coordinates of all atoms in the systems at a given time. VOTCA-CTP can read standard `GROMACS` topology files. Custom definitions of atomistic topology via XML files are also possible. Since the description of the atomistic topology is the first step in the charge transport simulations, it is important to follow simple conventions on how the system is partitioned on molecules, residues, and how atoms are named in the topology. Required input files are described in section atomistic topology.

2.3 Conjugated segments and rigid fragments

With the morphology at hand, the next step is partitioning the system on hopping sites, or conjugated segments, and calculating charge transfer rates between them. Physically intuitive arguments can be used for the partitioning, which reflects the localization of the wave function of a charge. For most organic semiconductors, the molecular architecture includes relatively rigid, planar π-conjugated systems, which we will refer to as rigid fragments. A conjugated segment can contain one or more of such rigid fragments, which are linked by bonded degrees of freedom.

The dynamics of these degrees of freedom evolves on timescales much slower than the frequency of the internal promoting mode. In some cases, e.g. glasses, it can be 'frozen' due to non-bonded interactions with the surrounding molecules.

To illustrate the concept of conjugated segments and rigid fragments, three representative molecular architectures are shown in figure 2.2. The first one is a typical discotic liquid crystal, hexabenzocoronene. It consists of a conjugated core to which side chains are attached to aid self-assembly and solution processing. In this case the orbitals localized on side chains do not participate in charge transport and the conjugated π-system is both, a rigid fragment and a conjugated segment. In $\mathrm{Alq_3}$, a metal-coordinated compound, a charge carrier is delocalized over all three ligands. Hence, the whole molecule is one conjugated segment. Individual ligands are relatively rigid, while energies of the order of $k_\mathrm{B}T$ are sufficient to reorient them with respect to each other. Thus the Al atom and the three ligands are rigid fragments. In the case of a conjugated polymer, one molecule can consist of several conjugated segments, while each backbone repeat unit is a rigid fragment. Since the conjugation along the backbone can be broken due to large out-of-plane twists between two repeat units, an empirical criterion, based on the dihedral angle, can be used to partition the backbone on conjugated segments [8]. However, such intuitive partitioning is, to some extent, arbitrary and shall be validated by other methods [9–11].

After partitioning, an additional step is often required to remove bond length fluctuations introduced by molecular dynamics simulations, since they are already integrated out in the derivation of the rate expression. This is achieved by substituting respective molecular fragments with rigid, planar π-systems optimized using first-principles methods. Centers of mass and gyration tensors are used to align rigid fragments, though a custom definition of local axes is also possible. Such a procedure also minimizes discrepancies between the force-field and first-principles-based ground state geometries of conjugated segments, which might be important for calculations of electronic couplings, reorganization energies, and intramolecular driving forces.

To partition the system on hopping sites and substitute rigid fragments with the corresponding ground-state geometries `ctp_map` program is used:

> **ℹ Mapping the `GROMACS` trajectory**
> `ctp_map -t topol.tpr -c traj.xtc -s map.xml -f state.sql`

It reads in the `GROMACS` topology (`topol.tpr`) and trajectory (`traj.xtc`) files, definitions of conjugated segments and rigid fragments (`map.xml`) and outputs coordinates of conjugated segments (hopping sites) and rigid fragments (as provided in the MD trajectory and after rigidification) to the state file (`state.sql`). In order to do this, a mapping file `map.xml` has to be provided, which specifies the corresponding atoms in the different representations. After this step, all information (frame number, dimensions of the simulation box, etc) are stored in the state file and only this file is used for further calculations.

> **⚠ Attention**
> VOTCA-CTP requires a wrapped trajectory for mapping the segments and fragments, so all molecules should be whole in the frame.

In order to visually check the mapping one can use either the `tdump` calculator or the programm `ctp_dump` with the calculator `trajectory2pdb`.

> **ℹ Writing a mapped trajectroy with `ctp_dump`**
> `ctp_dump -f state.sql -e trajectory2pdb`

122   It reads in the state file created by `ctp_map` and outputs two trajectory files corresponding to
123   the original and rigidified atom coordinates. To check the mapping, it is useful to superimpose
124   the three outputs (original atomistic, atomistic stored in the state file, and rigidified according to
125   ground state geometries), e.g., with `VMD`.

> 🛈 **Writing a mapped trajectroy with `tdump`**
> `ctp_run -f state.sql -o options.xml -e tdump`

126   It also reads in the state file but appends the coordinates to a pdb. file. So make sure to delete old
127   QM.pdb and MD.pdb if you want to create a new imagef

2.4   Neighbor list

129   A list of neigboring conjugated segments, or neighbor list, contains all pairs of conjugated seg-
130   ments for which coupling elements, reorganization energies, site energy differences, and rates
131   are evaluated.
132   Two segments are added to this list if the distance between centers of mass of any of their rigid
133   fragments is below a certain cutoff. This allows neighbors to be selected on a criterion of min-
134   imum distance of approach rather than center of mass distance, which is useful for molecules
135   with anisotropic shapes.
136   The neighbor list can be generated from the atomistic trajectory by using the neighborlist
137   calculator. This calculator requires a cutoff, which can be specified in the options.xml file. The
138   list is saved to the state.sql file:

> 🛈 **Generating a neighbor list**
> `ctp_run -o options.xml -f state.sql -e neighborlist`

2.5   Reorganization energy

140   The reorganization energy λ_{ij} takes into account the change in nuclear (and dielectric) degrees of
141   freedom as the charge moves from donor i to acceptor j. It has two contributions: intramolecular,
142   $\lambda_{ij}^{\text{int}}$, which is due to reorganization of nuclear coordinates of the two molecules forming the
143   charge transfer complex, and intermolecular (outersphere), $\lambda_{ij}^{\text{out}}$, which is due to the relaxation of
144   the nuclear coordinates of the environment. In what follows we discuss how these contributions
145   can be calculated.

2.5.1   Intramolecular reorganization energy

147   If intramolecular vibrational modes of the two molecules are treated classically, the rearrange-
148   ment of their nuclear coordinates after charge transfer results in the dissipation of the internal
149   reorganization energy, $\lambda_{ij}^{\text{int}}$. It can be computed from four points on the potential energy surfaces
150   (PES) of both molecules in neutral and charged states, as indicated in figure 2.3.
151   Adding the contributions due to discharging of molecule i and charging of molecule j yields [12]

$$\lambda_{ij}^{\text{int}} = \lambda_i^{cn} + \lambda_j^{nc} = U_i^{nC} - U_i^{nN} + U_j^{cN} - U_j^{cC} . \tag{2.1}$$

152   Here U_i^{nC} is the internal energy of the neutral molecule i in the geometry of its charged state
153   (small n denotes the state and capital C the geometry). Similarly, U_j^{cN} is the energy of the charged
154   molecule j in the geometry of its neutral state. Note that the PES of the donor and acceptor are
155   not identical for chemically different compounds or for conformers of the same molecule. In this

Figure 2.3: Potential energy surfaces of (a) donor and (b) acceptor in charged and neutral states. After the change of the charge state both molecules relax their nuclear coordinates. If all vibrational modes are treated classically, the total internal reorganization energy and the internal energy difference of the electron transfer reaction are $\lambda_{ij}^{\text{int}} = \lambda_i^{cn} + \lambda_j^{nc}$ and $\Delta E_{ij}^{\text{int}} = \Delta U_i - \Delta U_j$, respectively.

<div style="text-align: right">fig:parabolas</div>

case $\lambda_i^{cn} \neq \lambda_j^{cn}$ and $\lambda_i^{nc} \neq \lambda_j^{nc}$. Thus $\lambda_{ij}^{\text{int}}$ is a property of the charge transfer complex, and not of a single molecule.

Intramolecular reorganization energies for discharging (λ^{cn}) and charging (λ^{nc}) of a molecule need to be determined using quantum-chemistry and given in `map.xml`. The values are written to the `state.sql` using the calculator `einternal` (see also internal energy):

> 🛈 **Intramolecularl reorganization energies**
>
> `ctp_run -o options.xml -f state.sql -e einternal`

2.5.2 Outersphere reorganization energy

During the charge transfer reaction, also the molecules outside the charge transfer complex reorient and polarize in order to adjust for changes in electric potential, resulting in the outersphere contribution to the reorganization energy. $\lambda_{ij}^{\text{out}}$ is particularly important if charge transfer occurs in a polarizable environment. Assuming that charge transfer is much slower than electronic polarization but much faster than nuclear rearrangement of the environment, $\lambda_{ij}^{\text{out}}$ can be calculated from the electric displacement fields created by the charge transfer complex [?]

$$\lambda_{ij}^{\text{out}} = \frac{c_p}{2\epsilon_0} \int_{V^{\text{out}}} dV \left[\vec{D}_I(\vec{r}) - \vec{D}_F(\vec{r}) \right]^2 , \tag{2.2}$$

<div style="text-align: right">equ:lambda_outer1</div>

where ϵ_0 is the the permittivity of free space, $\vec{D}_{I,F}(\vec{r})$ are the electric displacement fields created by the charge transfer complex in the initial (charge on molecule i) and final (charge transferred to molecule j) states, V^{out} is the volume outside the complex, and $c_p = \frac{1}{\epsilon_{\text{opt}}} - \frac{1}{\epsilon_{\text{s}}}$ is the Pekar factor, which is determined by the low (ϵ_{s}) and high (ϵ_{opt}) frequency dielectric permittivities.

Eq. (2.2) can be simplified by assuming spherically symmetric charge distributions on molecules i and j with total charge e. Integration over the volume V^{out} outside of the two spheres of radii R_i and R_j centered on molecules i and j leads to the classical Marcus expression for the outersphere reorganization energy

$$\lambda_{ij}^{\text{out}} = \frac{c_p e^2}{4\pi\epsilon_0} \left(\frac{1}{2R_i} + \frac{1}{2R_j} - \frac{1}{r_{ij}} \right) , \tag{2.3}$$

<div style="text-align: right">equ:lambda_outer2</div>

where r_{ij} is the molecular separation. While eq. (2.3) captures the main physics, e.g. predicts smaller outer-sphere reorganization energies (higher rates) for molecules at smaller separations,

178 it often cannot provide quantitative estimates, since charge distributions are rarely spherically
179 symmetric.
180 Alternatively, the displacement fields can be constructed using the atomic partial charges. The
181 difference of the displacement fields at the position of an atom b_k outside the charge transfer
182 complex (molecule $k \neq i, j$) can be expressed as

$$\vec{D}_I(\vec{r}_{b_k}) - \vec{D}_F(\vec{r}_{b_k}) = \sum_{a_i} \frac{q_{a_i}^c - q_{a_i}^n}{4\pi} \frac{(\vec{r}_{b_k} - \vec{r}_{a_i})}{|\vec{r}_{b_k} - \vec{r}_{a_i}|^3} + \sum_{a_j} \frac{q_{a_j}^n - q_{a_j}^c}{4\pi} \frac{(\vec{r}_{b_k} - \vec{r}_{a_j})}{|\vec{r}_{b_k} - \vec{r}_{a_j}|^3} \,, \qquad (2.4)$$

183 where $q_{a_i}^n$ ($q_{a_i}^c$) is the partial charge of atom a of the neutral (charged) molecule i in vacuum. The
184 partial charges of neutral and charged molecules are obtained by fitting their values to reproduce
185 the electrostatic potential of a single molecule (charged or neutral) in vacuum. Assuming a uni-
186 form density of atoms, the integration in eq. (2.2) can be rewritten as a density-weighted sum
187 over all atoms excluding those of the charge transfer complex.
188 The remaining unknown needed to calculate $\lambda_{ij}^{\text{out}}$ is the Pekar factor, c_p. In polar solvents $\epsilon_{\text{s}} \gg$
189 $\epsilon_{\text{opt}} \sim 1$ and c_p is of the order of 1. In most organic semiconductors, however, molecular orien-
190 tations are fixed and therefore the low frequency dielectric permittivity is of the same order of
191 magnitude as ϵ_{opt}. Hence, c_p is small and its value is very sensitive to differences in the permit-
192 tivities.
193 Outersphere reorganization energies for all pairs of molecules in the neighbor list can be com-
194 puted from the atomistic trajectory by using the `eoutersphere` calculator.
195 Two methods can be used to compute $\lambda_{ij}^{\text{out}}$. The first method uses the atomistic partial charges of
196 neutral and charged molecules from files specified in `map.xml` and eq. (2.2). The Pekar factor c_p
197 and a cutoff radius based on molecular centers of mass have to be specified in the `options.xml`
198 file.
199 If this method is computationally prohibitive, $\lambda_{ij}^{\text{out}}$ can be computed using eq. (2.3), which as-
200 sumes spherical charge distributions on the molecules. In this case the radii of these spheres are
201 specified in `segments.xml`, while the Pekar factor c_p is given in the `options.xml` file and no
202 cutoff radius is needed.
203 The outer sphere reorganization energies are saved to the `state.sql` file:

> 🛈 **Outersphere reorganization energy**
>
> ```
> ctp_run  -o options.xml  -f state.sql  -e outersphere
> ```

204 2.6   Site energies

205 A charge transfer reaction between molecules i and j is driven by the site energy difference,
206 $\Delta E_{ij} = E_i - E_j$. Since the transfer rate, ω_{ij}, depends exponentially on ΔE_{ij} (see eq. (2.31)) it is
207 important to compute its distribution as accurately as possible. The total site energy difference
208 has contributions due to externally applied electric field, electrostatic interactions, polarization
209 effects, and internal energy differences. In what follows we discuss how to estimate these contri-
210 butions by making use of first-principles calculations and polarizable force-fields.

211 2.6.1   Externally applied electric field

212 The contribution to the total site energy difference due to an external electric field $\vec{F}$ is given by
213 $\Delta E_{ij}^{\text{ext}} = q\vec{F} \cdot \vec{r}_{ij}$, where $q = \pm e$ is the charge and $\vec{r}_{ij} = \vec{r}_i - \vec{r}_j$ is a vector connecting molecules
214 i and j. For typical distances between small molecules, which are of the order of $1\,\text{nm}$, and
215 moderate fields of $F < 10^8\,\text{V/m}$ this term is always smaller than $0.1\,\text{eV}$.

²¹⁶ **2.6.2   Internal energy**

sec:internal_energy

²¹⁷ The contribution to the site energy difference due to different internal energies (see figure 2.3)
²¹⁸ can be written as

$$\Delta E_{ij}^{\text{int}} = \Delta U_i - \Delta U_j = \left(U_i^{cC} - U_i^{nN}\right) - \left(U_j^{cC} - U_j^{nN}\right) , \tag{2.5}$$

equ:conformational

²¹⁹ where $U_i^{cC(nN)}$ is the total energy of molecule i in the charged (neutral) state and geometry. ΔU_i
²²⁰ corresponds to the adiabatic ionization potential (or electron affinity) of molecule i, as shown
²²¹ in figure 2.3. For one-component systems and negligible conformational changes $\Delta E_{ij}^{\text{int}} = 0$,
²²² while it is significant for donor-acceptor systems.
²²³ Internal energies determined using quantum-chemistry need to be specified in `map.xml`. The
²²⁴ values are written to the `state.sql` using the calculator `einternal` (see also intramolecular
²²⁵ reorganization energy):

> 🛈 **Internal energies**
> `ctp_run -o options.xml -f state.sql -e einternal`

²²⁶ **2.6.3   Electrostatic interaction energy**

sec:distributed_multipoles

We represent the molecular charge density by choosing multiple expansion sites ("polar sites")
per molecule in such a way as to accurately reproduce the molecular electrostatic potential (ESP),
with a set of suitably chosen multipole moments $\{Q_{lk}^a\}$ (in spherical-tensor notation) allocated to
each site. The expression for the electrostatic interaction energy between two molecules A and B
in the multi-point expansion includes an implicit sum over expansion sites $a\epsilon A$ and $b\epsilon B$,

$$U_{AB} = \sum_{a\epsilon A} \sum_{b\epsilon B} \hat{Q}_{l_1 k_1}^a T_{l_1 k_1 l_2 k_2}^{a,b} \hat{Q}_{l_2 k_2}^b \equiv \hat{Q}_{l_1 k_1}^a T_{l_1 k_1 l_2 k_2}^{a,b} \hat{Q}_{l_2 k_2}^b , \tag{2.6}$$

equ:mol_distributed_U

²²⁷ where we have used the Einstein sum convention for the site indices a and b on the right-hand
²²⁸ side of the equation, in addition to the sum convention that is in place for the multipole-moment
²²⁹ components $t \equiv l_1 k_1$ and $u \equiv l_2 k_2$. The $T_{l_1 k_1 l_2 k_2}^{a,b}$ are tensors that mediate the interaction between
²³⁰ a multipole component $l_1 k_1$ on site a with the moment $l_2 k_2$ on site b. If we include the molecular
²³¹ environment into a perturbative term W to enter in the single-molecule Hamiltonian, the above
²³² expression is exactly the first-order correction to the energy where the quantum-mechanical detail
²³³ has been absorbed in classical multipole moments.
²³⁴ The are a number of strategies how to arrive at such a collection of *distributed multipoles*. They can
²³⁵ be classified according to whether the multipoles are derived (a) from the electrostatic potential
²³⁶ generated by the SCF charge density or (b) from a decomposition of the wavefunction itself. Here,
²³⁷ we will only draft two of those approaches, CHELPG [13] from category (a) and DMA [14] from
²³⁸ category (b).
The CHELPG (CHarges from ELectrostatic Potentials, Grid-based) method relies on performing
a least-squares fit of atom-placed charges to reproduce the electrostatic potential as evaluated
from the SCF density on a regularly spaced grid [13]. The fitted charges result from minimizing
the Lagrangian function [15]

$$z(\{q_i\}) = \sum_{k=1}^{M} \left(\phi(\vec{r}_k) - \sum_{i=1}^{N} \frac{1}{4\pi\varepsilon_0} \frac{q_i}{|\vec{r}_i - \vec{r}_k|}\right) + \lambda \left(q_{\text{mol}} - \sum_{i=1}^{N} q_i\right) , \tag{2.7}$$

²³⁹ with M grid points, N atomic sites, the set of atomic partial charges $\{q_i\}$ and the SCF potential ϕ.
²⁴⁰ The Lagrange multiplier λ constrains the sum of the fitted charges to the molecular charge q_{mol}.
²⁴¹ The main difference from other fitting schemes [16] is the algorithm that selects the positions

at which the potential is evaluated (we note that the choice of grid points can have substantial effects especially for bulky molecules). Clearly, the CHELPG method can be (and has been) extended to include higher atomic multipoles. It should be noted, however, how already the inclusion of atomic dipoles hardly improves the parametrization, and can in fact be harmful to its conformational stability.

The Distributed-Multipole-Analysis (DMA) approach [14, 17], developed by A. Stone, operates directly on the quantum-mechanical density matrix, expanded in terms of atom- and bond-centered Gaussian functions $\chi_\alpha = R_{LK}(\vec{x} - \vec{s}_\alpha) \exp[-\zeta(\vec{x} - \vec{s}_\alpha)^2]$,

$$\rho(\vec{x}) = \sum_{\alpha,\beta} \rho_{\alpha\beta} \chi_\alpha(\vec{x} - \vec{s}_\alpha) \chi_\beta(\vec{x} - \vec{s}_\beta). \tag{2.8}$$

The aim is to compute multipole moments according in a distributed fashion: If we use that the overlap product $\chi_\alpha \chi_\beta$ of two Gaussian basis functions yields itself a Gaussian centered at $\vec{P} = (\zeta_\alpha \vec{s}_\alpha + \zeta_\beta \vec{s}_\beta)/(\zeta_\alpha + \zeta_\beta)$, it is possible to proceed in two steps: First, we compute the multipole moments associated with a specific summand in the density matrix, referred to the overlap center $\vec{P}$:

$$Q_{LK}[\vec{P}] = -\int R_{LK}(\vec{x} - \vec{P}) \rho_{\alpha\beta} \chi_\alpha \chi_\beta d^3x. \tag{2.9}$$

Second, we transfer the resulting $Q_{lk}[\vec{P}]$ to the position $\vec{S}$ of a polar site according to the rule [14]

$$Q_{nm}[\vec{S}] = \sum_{l=0}^{L} \sum_{k=-l}^{l} \left[\begin{pmatrix} n+m \\ l+k \end{pmatrix} \begin{pmatrix} n-m \\ l-k \end{pmatrix} \right]^{1/2} R_{n-l,m-k}(\vec{S} - \vec{P}) \cdot Q_{lk}[\vec{P}]. \tag{2.10}$$

Note how this requires a rule for the choice of the expansion site to which the multipole moment should be transferred. In the near past [17], the nearest-site algorithm, which allocates the multipole moments to the site closest to the overlap center, was replaced for diffuse functions by an algorithm based on a smooth weighting function in conjunction with grid-based integration methods in order to decrease the basis-set dependence of the resulting set of distributed multipoles.

One important advantage of the DMA approach over fitting algorithms such as CHELPG or Merz-Kollman (MK) is that higher-order moments can also be derived without too large an ambiguity.

The 'mps' file format used by VOTCA for the definition of distributed multipoles (as well as point polarizabilities, see subsequent section) is based on the GDMA punch format of A. Stone's GDMA program [17] (the punch output file can be immediately plugged into VOTCA without any conversions to be applied). Furthermore the log-file of different QM packages (currently `Gaussian`, `Turbomole` and `NWChem`) may be fed into the `log2mps` tool, which will subsequently generate the appropriate mps-file.

> ⓘ **Read in ESP charges from a QM log file**
> `ctp_tools -o options.xml -e log2mps`

2.6.4 Induction energy - the Thole model

If we in addition to the permanent set of multipole moments $\{Q_t^a\}$ allow for induced moments $\{\Delta Q_t^a\}$ and penalize their generation with a bilinear form (giving rise to a strictly positive contribution to the energy),

$$U_{\text{int}} = \frac{1}{2} \sum_A \Delta Q_t^a \eta_{tt'}^{aa'} \Delta Q_{t'}^{a'}, \tag{2.11}$$

it can be shown that the induction contribution to the site energy evaluates to an expression where all interactions between induced moments have cancelled out, and interactions between permanent and induced moments are scaled down by $1/2$ [18]:

$$U_{pu} = \frac{1}{2} \sum_A \sum_{B>A} \left[\Delta Q_t^a T_{tu}^{ab} Q_u^b + \Delta Q_t^b T_{tu}^{ab} Q_u^a \right]. \tag{2.12}$$

equ:u_pu

This term can be viewed as the second-order (induction) correction to the molecular interaction energy. The sets of $\{Q_t^a\}$ are solved for self-consistently via

$$\Delta Q_t^a = - \sum_{B \neq A} \alpha_{tt'}^{aa'} T_{t'u}^{a'b} (Q_u^b + \Delta Q_u^b), \tag{2.13}$$

equ:self_consistent_dQ

263 where the polarizability tensors $\alpha_{tt'}^{aa'}$ are given by the inverse of $\eta_{tt'}^{aa'}$.
264 With eqs. 2.13 and 2.12 we have at hand expressions that allow us to compute the induction en-
265 ergy contribution to site energies in an iterative manner based on a set of molecular distributed
266 multipoles $\{Q_t^a\}$ and polarizabilities $\{\alpha_{tt'}^{aa'}\}$. We have drafted in the previous section how to ob-
267 tain the former from a wavefunction decomposition or fitting scheme (GDMA, CHELPG). The
268 $\{\alpha_{tt'}^{aa'}\}$ can be derived formally (or rather: read off) from a perturbative expansion of the molec-
269 ular interaction. In this work we make use of the Thole model [19?] as a semi-empirical ap-
270 proach to obtain the sought-after point polarizabilities in the local dipole approximation, that is,
271 $[\alpha_{tt'}^{aa'}] = \alpha_{tt'}^{aa'} \delta_{t\beta} \delta_{t'\beta} \delta_{aa'}$, where $\beta \epsilon \{x, y, z\}$ references the dipole-moment component.
272 The Thole model is based on a modified dipole-dipole interaction, which can be reformulated in
273 terms of the interaction of smeared charge densities. This has been shown to be necessary due
274 to the divergent head-to-tail dipole-dipole interaction that otherwise results at small intersepara-
275 tions on the Å scale [19? ?]. Smearing out the charge distribution mimics the nature of the QM
276 wavefunction, which effectively guards against this unphysical polarization catastrophe. Since
277 the point dipoles however only react individually to the external field, any correlation effects as
278 were still accounted for in the $\{\alpha_{tt'}^{aa'}\}$ are lost, except perhaps those correlations that are due to
279 the mere classical field interaction.
The smearing of the nuclei-centered multipole moments is obtained via a fractional charge den-
sity $\rho_f(\vec{u})$ which should be normalized to unity and fall off rapidly as of a certain radius $\vec{u} = \vec{u}(\vec{R})$.
The latter is related to the physical distance vector $\vec{R}$ connecting two interacting sites via a linear
scaling factor that takes into account the magnitude of the isotropic site polarizabilities α^a. This
isotropic fractional charge density gives rise to a modified potential

$$\phi(u) = -\frac{1}{4\pi\varepsilon_0} \int_0^u 4\pi u' \rho(u') du' \tag{2.14}$$

equ:mod_potential

We can relate the multipole interaction tensor $T_{ij\ldots}$ (this time in Cartesian coordinates) to the fractional charge density in two steps: First, we rewrite the tensor in terms of the scaled distance vector $\vec{u}$,

$$T_{ij\ldots}(\vec{R}) = f(\alpha^a \alpha^b) \, t_{ij\ldots}(\vec{u}(\vec{R}, \alpha^a \alpha^b)), \tag{2.15}$$

where the specific form of $f(\alpha^a \alpha^b)$ results from the choice of $u(\vec{R}, \alpha^a \alpha^b)$. Second, we demand that the smeared interaction tensor $t_{ij\ldots}$ is given as usual by the appropriate derivative of the potential in eq. 2.14,

$$t_{ij\ldots}(\vec{u}) = -\partial_{u_i} \partial_{u_j} \ldots \phi(\vec{u}). \tag{2.16}$$

280 It turns out that for a suitable choice of $\rho_f(\vec{u})$, the modified interaction tensors can be rewritten
281 in such a way that powers n of the distance $R = |\vec{R}|$ are damped with a damping function
282 $\lambda_n(\vec{u}(\vec{R}))$ [20].

There is a large number of fractional charge densities $\rho_f(\vec{u})$ that have been tested for the purpose of giving best results for the molecular polarizability as well as interaction energies. Note how a great advantage of the Thole model is the exceptional transferability of the atomic polarizabilities to compounds not used for the fitting procedure [?]. In fact, for most organic molecules, a fixed set of atomic polarizabilities ($\alpha_C = 1.334$, $\alpha_H = 0.496$, $\alpha_N = 1.073$, $\alpha_O = 0.873$, $\alpha_S = 2.926$ Å^3) based on atomic elements yields satisfactory results.

VOTCA implements the Thole model with an exponentially-decaying fractional charge density

$$\rho(u) = \frac{3a}{4\pi} \exp(-au^3), \qquad (2.17)$$

where $\vec{u}(\vec{R}, \alpha^a\alpha^b) = \vec{R}/(\alpha^a\alpha^b)^{1/6}$ and the smearing exponent $a = 0.39$ (which can however be changed from the program options), as used in the AMOEBA force field [20].

Even though the Thole model performs very well for many organic compounds with only the above small set of element-based polarizabilities, conjugated molecules may require a more intricate parametrization. The simplest approach is to resort to scaled polarizabilities to match the effective molecular polarizable volume $V \sim \alpha_x\alpha_y\alpha_z$ as predicted by QM calculations (here $\alpha_x, \alpha_y, \alpha_z$ are the eigenvalues of the molecular polarizability tensor). The molpol tool assists with this task, it self-consistently calculates the Thole polarizability for an input mps-file and optimizes (if desired) the atomic polarizabilities in the above simple manner.

Generate Thole-type polarizabilites for a segment
```
ctp_tools -o options.xml -e molpol
```

The electrostatic and induction contribution to the site energy is evaluated by the emultipole calculator. Atomistic partial charges for charged and neutral molecules are taken from mps-files (extended GDMA format) specified in map.xml. Note that, in order to speed up calculations for both methods, a cut-off radius (for the molecular centers of mass) can be given in options.xml. Threaded execution is advised.

Electrostatic and induction corrections
```
ctp_run -o options.xml -f state.sql -e emultipole
```

Furthermore available are zmultipole, which extends emultipole to allow for an electrostatic buffer layer (loosely related to the z-buffer in OpenGL, hence the name) and anisotropic point polarizabilities. For the interaction energy of charged clusters of any user-defined composition (Frenkel states, CT states, ...), xqmultipole can be used.

Interaction energy of charged molecular clusters embedded in a molecular environment
```
ctp_parallel -o options.xml -f state.sql -e xqmultipole
```

2.7 Transfer integrals

The electronic transfer integral element J_{ij} entering the Marcus rates in eq. (2.31) is defined as

$$J_{ij} = \left\langle \phi_i \left| \hat{H} \right| \phi_j \right\rangle, \qquad (2.18)$$

where ϕ_i and ϕ_j are diabatic wavefunctions, localized on molecule i and j respectively, participating in the charge transfer, and $\hat{H}$ is the Hamiltonian of the formed dimer. Within the frozen-core approximation, the usual choice for the diabatic wavefunctions ϕ_i is the highest occupied

molecular orbital (HOMO) in case of hole transport, and the lowest unoccupied molecular orbital (LUMO) in the case of electron transfer, while $\hat{H}$ is an effective single particle Hamiltonian, e.g. Fock or Kohn-Sham operator of the dimer. As such, J_{ij} is a measure of the strength of the electronic coupling of the frontier orbitals of monomers mediated by the dimer interactions.

Intrinsically, the transfer integral is very sensitive to the molecular arrangement, i.e. the distance and the mutual orientation of the molecules participating in charge transport. Since this arrangement can also be significantly influenced by static and/or dynamic disorder [21? ? –23], it is essential to calculate J_{ij} explicitly for each hopping pair within a realistic morphology. Considering that the number of dimers for which eq. (2.18) has to be evaluated is proportional to the number of molecules times their coordination number, computationally efficient and at the same time quantitatively reliable schemes are required.

2.7.1 Projection of monomer orbitals on dimer orbitals (DIPRO)

sec:dipro

An approach for the determination of the transfer integral that can be used for any single-particle electronic structure method (Hartree-Fock, DFT, or semiempirical methods) is based on the projection of monomer orbitals on a manifold of explicitly calculated dimer orbitals. This dimer projection (DIPRO) technique including an assessment of computational parameters such as the basis set, exchange-correlation functionals, and convergence criteria is presented in detail in ref. [5]. A brief summary of the concept is given below.

We start from an effective Hamiltonian [1]

$$\hat{H}^{\text{eff}} = \sum_i \epsilon_i \hat{a}_i^\dagger \hat{a}_i + \sum_{j \neq i} J_{ij} \hat{a}_i^\dagger \hat{a}_j + c.c. \tag{2.19}$$

equ:dipro_eq1

where $\hat{a}_i^\dagger$ and $\hat{a}_i$ are the creation and annihilation operators for a charge carrier located at the molecular site i. The electron site energy is given by ϵ_i, while J_{ij} is the transfer integral between two sites i and j. We label their frontier orbitals (HOMO for hole transfer, LUMO for electron transfer) ϕ_i and ϕ_j, respectively. Assuming that the frontier orbitals of a dimer (adiabatic energy surfaces) result exclusively from the interaction of the frontier orbitals of monomers, and consequently expand them in terms of ϕ_i and ϕ_j. The expansion coefficients, $\bar{\mathbf{C}}$, can be determined by solving the secular equation

$$(\mathbf{H} - E\mathbf{S})\bar{\mathbf{C}} = 0 \tag{2.20}$$

equ:dipro_eq2

where $\mathbf{H}$ and $\mathbf{S}$ are the Hamiltonian and overlap matrices of the system, respectively. These matrices can be written explicitly as

$$\mathbf{H} = \begin{pmatrix} e_i & H_{ij} \\ H_{ij}^* & e_j \end{pmatrix} \qquad\qquad \mathbf{S} = \begin{pmatrix} 1 & S_{ij} \\ S_{ij}^* & 1 \end{pmatrix} \tag{2.21}$$

equ:dipro_eq3

with

$$\begin{aligned} e_i &= \langle \phi_i | \hat{H} | \phi_i \rangle & H_{ij} &= \langle \phi_i | \hat{H} | \phi_j \rangle \\ e_j &= \langle \phi_j | \hat{H} | \phi_j \rangle & S_{ij} &= \langle \phi_j | \phi_j \rangle \end{aligned} \tag{2.22}$$

equ:dipro_eq4

The matrix elements $e_{i(j)}$, H_{ij}, and S_{ij} entering eq. (2.21) can be calculated via projections on the dimer orbitals (eigenfunctions of $\hat{H}$) $\{|\phi_n^{\text{D}}\rangle\}$ by inserting $\hat{1} = \sum_n |\phi_n^{\text{D}}\rangle \langle \phi_n^{\text{D}}|$ twice. We exemplify this explicitly for H_{ij} in the following

$$H_{ij} = \sum_{nm} \langle \phi_i | \phi_n^{D} \rangle \langle \phi_n^{D} | \hat{H} | \phi_m^{D} \rangle \langle \phi_m^{D} | \phi_j \rangle. \tag{2.23}$$

eq:dipro_eq16

The Hamiltonian is diagonal in its eigenfunctions, $\langle \phi_n^{\text{D}} | \hat{H} | \phi_m^{\text{D}} \rangle = E_n \delta_{nm}$. Collecting the projections of the frontier orbitals $|\phi_{i(j)}\rangle$ on the n-th dimer state $\left(\bar{\mathbf{V}}_{(i)} \right)_n = \langle \phi_i | \phi_n^{\text{D}} \rangle$ and $\left(\bar{\mathbf{V}}_{(j)} \right)_n = \langle \phi_j | \phi_n^{\text{D}} \rangle$ respectively, into vectors we obtain

$$H_{ij} = \bar{\mathbf{V}}_{(i)} \mathbf{E} \bar{\mathbf{V}}_{(j)}^\dagger. \tag{2.24}$$

eq:dipro_eq17

---

[1] we use following notations: a - number, $\bar{\mathbf{a}}$ - vector, $\mathbf{A}$ - matrix, $\hat{A}$ - operator

What is left to do is determine these projections $\bar{\mathbf{V}}_{(k)}$. In all practical calculations the molecular orbitals are expanded in basis sets of either plane waves or of localized atomic orbitals $|\varphi_\alpha\rangle$. We will first consider the case that the calculations for the monomers are performed using a counterpoise basis set that is commonly used to deal with the basis set superposition error (BSSE). The basis set of atom-centered orbitals of a monomer is extended to the one of the dimer by adding the respective atomic orbitals at virtual coordinates of the second monomer. We can then write the respective expansions as

$$|\phi_k\rangle = \sum_\alpha \lambda_\alpha^{(k)} |\varphi_\alpha\rangle \qquad \text{and} \qquad |\phi_n^{\mathrm{D}}\rangle = \sum_\alpha D_\alpha^{(n)} |\varphi_\alpha\rangle \qquad (2.25)$$

where $k = i, j$. The projections can then be determined within this common basis set as

$$\left(\bar{\mathbf{V}}_k\right)_n = \langle \phi_k \, | \, \phi_n^{\mathrm{D}} \rangle = \sum_\alpha \lambda_\alpha^{(k)} \langle \alpha | \sum_\beta D_\beta^{(n)} |\beta\rangle = \bar{\boldsymbol{\lambda}}_{(k)}^\dagger \mathcal{S} \bar{\mathbf{D}}_{(n)} \qquad (2.26)$$

where $\mathcal{S}$ is the overlap matrix of the atomic basis functions. This allows us to finally write the elements of the Hamiltonian and overlap matrices in eq. (2.21) as:

$$H_{ij} = \bar{\boldsymbol{\lambda}}_{(i)}^\dagger \mathcal{S} \mathbf{D} \mathbf{E} \mathbf{D}^\dagger \mathcal{S}^\dagger \bar{\boldsymbol{\lambda}}_{(j)}$$
$$S_{ij} = \bar{\boldsymbol{\lambda}}_{(i)}^\dagger \mathcal{S} \mathbf{D} \mathbf{D}^\dagger \mathcal{S}^\dagger \bar{\boldsymbol{\lambda}}_{(j)} \qquad (2.27)$$

Since the two monomer frontier orbitals that form the basis of this expansion are not orthogonal in general ($\mathbf{S} \neq \mathbf{1}$), it is necessary to transform eq. (2.20) into a standard eigenvalue problem of the form

$$\mathbf{H}^{\mathrm{eff}} \bar{\mathbf{C}}^{\mathrm{eff}} = E \bar{\mathbf{C}}^{\mathrm{eff}} \qquad (2.28)$$

to make it correspond to eq. (2.19). According to Löwdin such a transformation can be achieved by

$$\mathbf{H}^{\mathrm{eff}} = \mathbf{S}^{-1/2} \mathbf{H} \mathbf{S}^{-1/2}. \qquad (2.29)$$

This then yields an effective Hamiltonian matrix in an orthogonal basis, and its entries can directly be identified with the site energies ϵ_i and transfer integrals J_{ij}:

$$\mathbf{H}^{\mathrm{eff}} = \begin{pmatrix} e_i^{\mathrm{eff}} & H_{ij}^{\mathrm{eff}} \\ H_{ij}^{*,\mathrm{eff}} & e_j^{\mathrm{eff}} \end{pmatrix} = \begin{pmatrix} \epsilon_i & J_{ij} \\ J_{ij}^* & \epsilon_j \end{pmatrix} \qquad (2.30)$$

2.7.2 DFT-based transfer integrals using DIPRO

The calculation of one electronic coupling element based on DFT using the `DIPRO` method requires the overlap matrix of atomic orbitals $\mathcal{S}$, the expansion coefficients for monomer $\bar{\boldsymbol{\lambda}}_{(k)} = \{\lambda_\alpha^{(k)}\}$ and dimer orbitals $\bar{\mathbf{D}}_{(n)} = \{D_\alpha^{(n)}\}$, as well as the orbital energies E_n of the dimer are required as input. In practical situations, performing self-consistent quantum-chemical calculations for each individual monomer and one for the dimer to obtain this input data is extremely demanding. Several simplifications can be made to reduce the computational effort, such as using non-Counterpoise basis sets for the monomers (thereby decoupling the monomer calculations from the dimer run) and performing only a single SCF step in a dimer calculation starting from an initial guess formed from a superposition of monomer orbitals. This "noCP+noSCF" variant of `DIPRO` is shown in figure 2.4(a) and recommended for production runs. A detailed comparative study of the different variants can be found in [5].

The code currently contains supports evaluation of transfer integrals from quantum-chemical calculations performed with the `Gaussian`, `Turbomole`, and `NWChem` packages. The interfacing procedure consists of three main steps: generation of input files for monomers and dimers, performing the actual quantum-chemical calculations, and calculating the transfer integrals.

Figure 2.4: Schematics of the DIPRO method. (a) General workflow of the projection technique. (b) Strategy of the efficient noCP+noSCF implementation, in which the monomer calculations are performed independently form the dimer configurations (noCP), using the `edft` calculator. The dimer Hamiltonian is subsequently constructed based on an initial guess formed from monomer orbitals and only diagonalized once (noSCF) before the transfer integral is calculated by projection. This second step is performed by the `idft` calculator.

fig:dipro_scheme

Monomer calculations

sec:edft

First, hopping sites and a neighbor list need to be generated from the atomistic topology and trajectory and written to the `state.sql` file. Then the parallel `edft` calculator manages the calculation of the monomer properties required for the determination of electronic coupling elements. Specifically, the individual steps it performs are:

1. Creation of a job file containing the list of molecules to be calculated with DFT

   > ⓘ **Writing job file for `edft`**
   > ```
   > ctp_parallel -o options.xml -f state.sql -e edft -j write
   > ```

2. Running of all jobs in job file

   > ⓘ **Running all `edft` jobs**
   > ```
   > ctp_parallel -o options.xml -f state.sql -e edft -j run
   > ```

   which includes

   - creating the input files for the DFT calculation (using the package specified in `options.xml`) in the directory

```
OR_FILES/package/frame_F/mol_M
```

where `F` is the index of the frame in the trajectory, `M` is the index of a molecule in this frame,

- executing the DFT run, and

- after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients), and saving it in compressed form to

```
OR_FILES/molecules/frame_F/molecule_M.orb
```

Calculating the transfer integrals

After the momomer calculations have been completed successfully, the respective runs for dimers from the neighborlist can be performed using the parallel `idft` calculator, which manages the DFT runs for the hopping pairs and determines the coupling element using `DIPRO`. Again, several steps are required:

1. Creation of a job file containing the list of pairs to be calculated with DFT

   > ℹ️ **Writing job file for `idft`**
   > `ctp_parallel -o options.xml -f state.sql -e idft -j write`

2. Running of all jobs in job file

   > ℹ️ **Running all `idft` jobs**
   > `ctp_parallel -o options.xml -f state.sql -e idft -j run`

   which includes

   - creating the input files (including the merged guess for a noSCF calculation, if requested) for the DFT calculation (using the package specified in `options.xml`) in the directory

   ```
   OR_FILES/package/frame_F/pair_M_N
   ```

   where `M` and `N` are the indices of the molecules in this pair,

   - executing the DFT run, and

   - after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients and energies, atomic orbital overlap matrix), and saving the pair information in compressed form to

   ```
   OR_FILES/pairs/frame_F/pair_M_N.orb
   ```

   - loading the monomer orbitals from the previously saved `*.orb files`.
   - calculating the coupling elements and write them to the job file

3. Reading the coupling elements from the job file and saving them to the `state.sql` file

   > ℹ️ **Saving `idft` results from job file to `state.sql`**
   > `ctp_parallel -o options.xml -f state.sql -e idft -j read`

2.7.3 ZINDO-based transfer integrals using MOO

sec:izindo

An approximate method based on Zerner's Intermediate Neglect of Differential Overlap (ZINDO) has been described in Ref. [24]. This semiempirical method is substantially faster than first-principles approaches, since it avoids the self-consistent calculations on each individual monomer and dimer. This allows to construct the matrix elements of the ZINDO Hamiltonian of the dimer from the weighted overlap of molecular orbitals of the two monomers. Together with the introduction of rigid segments, only a single self-consistent calculation on one isolated conjugated segment is required. All relevant molecular overlaps can then be constructed from the obtained molecular orbitals.

The main advantage of the molecular orbital overlap (MOO) library is *fast* evaluation of electronic coupling elements. Note that MOO is based on the ZINDO Hamiltonian which has limited applicability. The general advice is to first compare the accuracy of the MOO method to the DFT-based calculations.

MOO can be used both in a standalone mode and as an `izindo` calculator of VOTCA-CTP.

Since MOO constructs the Fock operator of a dimer from the molecular orbitals of monomers by translating and rotating the orbitals of rigid fragments, the optimized geometry of all conjugated segments and the coefficients of the molecular orbitals are required as its input in addition to the state file (`state.sql`) with the neighbor list. Coordinates are stored in `geometry.xyz` files with four columns, first being the atom type and the next three atom coordinates. This is a standard `xyz` format without a header. Note that the atom order in the `geometry.xyz` files can be different from that of the mapping files. The correspondence between the two is established in the `map.xml` file.

> ⚠️ **Attention**
>
> Izindo requires the specification of orbitals for hole and electron transport in `map.xml`. They are the HOMO and LUMO respectively and can be retrieved from the `log` file from which the `zindo.orb` file is generated. The number of `alpha electrons` is the HOMO, the LUMO is HOMO+1

The calculated transfer integrals are immediately saved to the `state.sql` file.

> ℹ️ **Transfer integrals from `izindo`**
>
> `ctp_run -o options.xml -f state.sql -e izindo`

2.8 Charge transfer rate

sec:rates

Charge transfer rates can be postulated based on intuitive physical considerations, as it is done in the Gaussian disorder models [25–28]. Alternatively, charge transfer theories can be used to evaluate rates from quantum chemical calculations [3, 5, 12, 29–31]. In spite of being significantly more computationally demanding, the latter approach allows to link the chemical and electronic structure, as well as the morphology, to charge dynamics.

2.8.1 Classical charge transfer rate

sec:rate_classical

The high temperature limit of classical charge transfer theory [32, 33] is often used as a trade-off between theoretical rigor and computational complexity. It captures key parameters which influence charge transport while at the same time providing an analytical expression for the rate. Within this limit, the transfer rate for a charge to hop from a site i to a site j reads

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{J_{ij}^2}{\sqrt{4\pi\lambda_{ij}k_{\mathrm{B}}T}} \exp\left[-\frac{(\Delta E_{ij} - \lambda_{ij})^2}{4\lambda_{ij}k_{\mathrm{B}}T}\right], \qquad (2.31)$$

equ:marcus

where T is the temperature, $\lambda_{ij} = \lambda_{ij}^{\mathrm{int}} + \lambda_{ij}^{\mathrm{out}}$ is the reorganization energy, which is a sum of intra- and inter-molecular (outersphere) contributions, ΔE_{ij} is the site-energy difference, or driving force, and J_{ij} is the electronic coupling element, or transfer integral.

2.8.2   Semi-classical bimolecular rate

The main assumptions in eq. (2.31) are non-adiabaticity (small electronic coupling and charge transfer between two diabatic, non-interacting states), and harmonic promoting modes, which are treated classically. At ambient conditions, however, the intramolecular promoting mode, which roughly corresponds to C-C bond stretching, has a vibrational energy of $\hbar\omega \approx 0.2\,\mathrm{eV} \gg k_{\mathrm{B}}T$ and should be treated quantum-mechanically. The outer-sphere (slow) mode has much lower vibrational energy than the intramolecular promoting mode, and therefore can be treated classically. The weak interaction between molecules also implies that each molecule has its own, practically independent, set of quantum mechanical degrees of freedom.

A more general, quantum-classical expression for a bimolecular multi-channel rate is derived in the Supporting Information of ref. [3] and has the following form

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{|J_{ij}|^2}{\sqrt{4\pi\lambda_{ij}^{\mathrm{out}}k_{\mathrm{B}}T}} \sum_{l',m'=0}^{\infty} |\langle\chi_{i0}^c|\chi_{il'}^n\rangle|^2 |\langle\chi_{j0}^n|\chi_{jm'}^c\rangle|^2 \exp\left\{-\frac{\left[\Delta E_{ij} - \hbar(l'\omega_i^n + m'\omega_j^c) - \lambda_{ij}^{\mathrm{out}}\right]^2}{4\lambda_{ij}^{\mathrm{out}}k_{\mathrm{B}}T}\right\}.$$

$$(2.32) \quad \text{equ:jjortner}$$

If the curvatures of intramolecular PES of charged and neutral states of a molecule are different, that is $\omega_i^c \neq \omega_i^n$, the corresponding reorganization energies, $\lambda_i^{cn} = \frac{1}{2}[\omega_i^n(q_i^n - q_i^c)]^2$ and $\lambda_i^{nc} = \frac{1}{2}[\omega_i^c(q_i^n - q_i^c)]^2$, will also differ. In this case the Franck-Condon (FC) factors for discharging of molecule i read [34]

$$|\langle\chi_{i0}^c|\chi_{il'}^n\rangle|^2 = \frac{2}{2^{l'}l'!} \frac{\sqrt{\omega_i^c\omega_i^n}}{(\omega_i^c + \omega_i^n)} \exp\left(-|s_i|\right) \left[\sum_{\substack{k=0\\k\,\mathrm{even}}}^{l'} \binom{l'}{k}\left(\frac{2\omega_i^c}{\omega_i^c + \omega_i^n}\right)^{k/2} \frac{k!}{(k/2)!} H_{l'-k}\left(\frac{s_i}{\sqrt{2S_i^{cn}}}\right)\right]^2,$$

$$(2.33)$$

where $H_n(x)$ is a Hermite polynomial, $s_i = 2\sqrt{\lambda_i^{nc}\lambda_i^{cn}}/\hbar(\omega_i^c + \omega_i^n)$, and $S_i^{cn} = \lambda_i^{cn}/\hbar\omega_i^c$. The FC factors for charging of molecule j can be obtained by substituting $(s_i, S_i^{cn}, \omega_i^c)$ with $(-s_j, S_j^{nc}, \omega_j^n)$. In order to evaluate the FC factors, the internal reorganization energy λ_i^{cn} can be computed from the intramolecular PES.

2.8.3   Semi-classical rate

One can also use the quantum-classical rate with a common set of vibrational coordinates [?]

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{|J_{ij}|^2}{\sqrt{4\pi\lambda_{ij}^{\mathrm{out}}k_{\mathrm{B}}T}} \sum_{N=0}^{\infty} \frac{1}{N!} \left(\frac{\lambda_{ij}^{\mathrm{int}}}{\hbar\omega^{\mathrm{int}}}\right)^N \exp\left(-\frac{\lambda_{ij}^{\mathrm{int}}}{\hbar\omega^{\mathrm{int}}}\right) \exp\left\{-\frac{\left[\Delta E_{ij} - \hbar N\omega^{\mathrm{int}} - \lambda_{ij}^{\mathrm{out}}\right]^2}{4\lambda_{ij}^{\mathrm{out}}k_{\mathrm{B}}T}\right\}.$$

$$(2.34) \quad \text{equ:jortner}$$

Numerical estimates show that if $\lambda_{ij}^{\mathrm{int}} \approx \lambda_{ij}^{\mathrm{out}}$ and $|\Delta E_{ij}| \ll \lambda_{ij}^{\mathrm{out}}$ the rates are similar to those of eq. (2.31). In general, there is no robust method to compute $\lambda_{ij}^{\mathrm{out}}$ [35] and both reorganization energies are often assumed to be of the same order of magnitude. In this case the second condition also holds, unless there are large differences in electron affinities or ionization potentials of neighboring molecules, e.g. in donor-acceptor blends.

To calculate rates of the type specified in `options.xml` for all pairs in the neighbor list and to
save them into the `state.sql` file, run the `rates` calculator. Note that all required ingredients (reorganization energies, transfer integrals, and site energies have to be calculated before).

> **ⓘ Calculation of transfer rates**
> `ctp_run -o options.xml -f state.sql -e rates`

2.9 Master equation

Having determined the list of conjugated segments (hopping sites) and charge transfer rates between them, the next task is to solve the master equation which describes the time evolution of the system

$$\frac{\partial P_\alpha}{\partial t} = \sum_\beta P_\beta \Omega_{\beta\alpha} - \sum_\beta P_\alpha \Omega_{\alpha\beta}, \tag{2.35}$$

where P_α is the probability of the system to be in a state α at time t and $\Omega_{\alpha\beta}$ is the transition rate
from state α to state β. A state α is specified by a set of site occupations, $\{\alpha_i\}$, where $\alpha_i = 1(0)$
for an occupied (unoccupied) site i, and the matrix $\hat{\Omega}$ can be constructed from rates ω_{ij}.
The solution of eq. (2.35) is be obtained by using kinetic Monte Carlo (KMC) methods. KMC
explicitly simulates the dynamics of charge carriers by constructing a Markov chain in state space
and can find both stationary and transient solutions of the master equation. The main advantage
of KMC is that only states with a direct link to the current state need to be considered at each step.
Since these can be constructed solely from current site occupations, extensions to multiple charge
carriers (without the mean-field approximation), site-occupation dependent rates (needed for
the explicit treatment of Coulomb interactions), and different types of interacting particles and
processes, are straightforward. To optimize memory usage and efficiency, a combination of the
variable step size method [36] and the first reaction method is implemented.
To obtain the dynamics of charges using KMC, the program `kmc_run` executes a specific calculator after reading its options (charge carrier type, runtime, numer of carriers etc.) from `options.xml`.

> **ⓘ KMC for a single carrier in periodic boundary conditions**
> `kmc_run -o options.xml -f state.sql -e kmcsingle`

> **ⓘ KMC for multiple carriers of the same type in periodic boundary conditions**
> `kmc_run -o options.xml -f state.sql -e kmcmultiple`

2.9.1 Extrapolation to nondispersive mobilities

Predictions of charge-carrier mobilities in partially disordered semiconductors rely on charge
transport simulations in systems which are only several nanometers thick. As a result, simulated charge transport might be dispersive for materials with large energetic disorder [37, 38]
and simulated mobilities are system-size dependent. In time-of-flight (TOF) experiments, however, a typical sample thickness is in the micrometer range and transport is often nondispersive.
In order to link simulation and experiment, one needs to extract the nondispersive mobility from
simulations of small systems, where charge transport is dispersive at room temperature.
Such extrapolation is possible if the temperature dependence of the nondispersive mobility is
known in a wide temperature range. For example, one can use analytical results derived for one-dimensional models [39–41]. The mobility-temperature dependence can then be parametrized by

510   simulating charge transport at elevated temperatures, for which transport is nondispersive even
511   for small system sizes. This dependence can then be used to extrapolate to the nondispersive
512   mobility at room temperature [4].
513   For Alq_3, the charge carrier mobility of a periodic system of 512 molecules was shown to be more
514   than three orders of magnitude higher than the nondispersive mobility of an infinitely large sys-
515   tem [4]. Furthermore, it was shown that the transition between the dispersive and nondispersive
516   transport has a logarithmic dependence on the number of hopping sites N. Hence, a brute-force
517   increase of the system size cannot resolve the problem for compounds with large energetic disor-
518   der σ, since N increases exponentially with σ^2.

2.10   Macroscopic observables

520   Spatial distributions of charge and current densities can provide a better insight in the micro-
521   scopic mechanisms of charge transport. If O is an observable which has a value O_α in a state α,
522   its ensemble average at time t is a sum over all states weighted by the probability P_α to be in a
523   state α at time t

$$\langle O \rangle = \sum_\alpha O_\alpha P_\alpha. \tag{2.36}$$

524   If O does not explicitly depend on time, the time evolution of $\langle O \rangle$ can be calculated as

$$\frac{d\langle O \rangle}{dt} = \sum_{\alpha,\beta} \left[P_\beta \Omega_{\beta\alpha} - P_\alpha \Omega_{\alpha\beta} \right] O_\alpha = \sum_{\alpha,\beta} P_\beta \Omega_{\beta\alpha} \left[O_\alpha - O_\beta \right]. \tag{2.37}$$

525   If averages are obtained from KMC trajectories, $P_\alpha = s_\alpha/s$, where s_α is the number of Markov
526   chains ending in the state α after time t, and s is the total number of chains.
      Alternatively, one can calculate time averages by analyzing a single Markov chain. If the total
      occupation time of the state α is τ_α then

$$\overline{O} = \frac{1}{\tau} \sum_\alpha O_\alpha \tau_\alpha , \tag{2.38}$$

527   where $\tau = \sum_\alpha \tau_\alpha$ is the total time used for time averaging.
528   For ergodic systems and sufficient sampling times, ensemble and time averages should give iden-
529   tical results. In many cases, the averaging procedure reflects a specific experimental technique.
530   For example, an ensemble average over several KMC trajectories with different starting condi-
531   tions corresponds to averaging over injected charge carriers in a time-of-flight experiment. In
532   what follows, we focus on the single charge carrier (low concentration of charges) case.

2.10.1   Charge density

534   For a specific type of particles, the microscopic charge density of a site i is proportional to the
535   occupation probability of the site, p_i

$$\rho_i = ep_i/V_i , \tag{2.39}$$

536   where, for an irregular lattice, the effective volume V_i can be obtained from a Voronoi tessellation
537   of space. For reasonably uniform lattices (uniform site densities) this volume is almost indepen-
538   dent of the site and a constant volume per cite, $V_i = V/N$, can be assumed. In the macroscopic
539   limit, the charge density can be calculated using a smoothing kernel function, i.e. a distance-
540   weighted average over multiple sites. Site occupations p_i can be obtained from eq. (2.36) or
541   eq. (2.38) by using the occupation of site i in state α as an observable.
542   If the system is in thermodynamic equilibrium, that is without sources or sinks and without
543   circular currents (and therefore no net flux) a condition, known as detailed balance, holds

$$p_j \omega_{ji} = p_i \omega_{ij}, \tag{2.40}$$

544 It can be used to test whether the system is ergodic or not by correlating $\log p_i$ and the site energy
545 E_i. Indeed, if $\lambda_{ij} = \lambda_{ji}$ the ratios of forward and backward rates are determined solely by the
546 energetic disorder, $\omega_{ji}/\omega_{ij} = \exp(-\Delta E_{ij}/k_\mathrm{B}T)$ (see eq. (2.31)).

2.10.2 Current

548 If the position of the charge, $\vec{r}$, is an observable, the time evolution of its average $\langle \vec{r} \rangle$ is the total
549 current in the system

$$\vec{J} = e \langle \vec{v} \rangle = e \frac{d \langle \vec{r} \rangle}{dt} = e \sum_{i,j} p_j \omega_{ji} (\vec{r}_i - \vec{r}_j). \tag{2.41}$$

equ:current_def

550 Symmetrizing this expression we obtain

$$\vec{J} = \frac{1}{2} e \sum_{i,j} (p_j \omega_{ji} - p_i \omega_{ij}) \, \vec{r}_{ij}, \tag{2.42}$$

equ:current

551 where $\vec{r}_{ij} = \vec{r}_i - \vec{r}_j$. Symmetrization ensures equal flux splitting between neighboring sites and
552 absence of local average fluxes in equilibrium. It allows to define a local current through site i as

$$\vec{J}_i = \frac{1}{2} e \sum_{j} (p_j \omega_{ji} - p_i \omega_{ij}) \, \vec{r}_{ij}. \tag{2.43}$$

equ:site_current

553 A large value of the local current indicates that the site contributes considerably to the total cur-
554 rent. A collection of such sites thus represents most favorable charge pathways [42].

2.10.3 Mobility and diffusion constant

556 For a single particle, e.g. a charge or an exciton, a zero-field mobility can be determined by
557 studying particle diffusion in the absence of external fields. Using the particle displacement
558 squared, $\Delta \boldsymbol{r}_i^2$, as an observable we obtain

$$2dD_{\gamma\delta} = \frac{d \langle \Delta r_{i,\gamma} \Delta r_{i,\delta} \rangle}{dt} = \sum_{\substack{i,j \\ i \neq j}} p_j \omega_{ji} (\Delta r_{i,\gamma} \Delta r_{i,\delta} - \Delta r_{j,\gamma} \Delta r_{j,\delta}) = \sum_{\substack{i,j \\ i \neq j}} p_j \omega_{ji} (r_{i,\gamma} r_{i,\delta} - r_{j,\gamma} r_{j,\delta}) \, .$$

$$\tag{2.44}$$

equ:diffusion

559 Here $\vec{r}_i$ is the coordinate of the site i, $D_{\gamma\delta}$ is the diffusion tensor, $\gamma, \delta = x, y, z$, and $d = 3$ is the
560 system dimension. Using the Einstein relation,

$$D_{\gamma\delta} = k_\mathrm{B} T \mu_{\gamma\delta} \, , \tag{2.45}$$

561 one can, in principle, obtain the zero-field mobility tensor $\mu_{\gamma\delta}$. Eq. (2.44), however, does not take
562 into account the use of periodic boundary conditions when simulating charge dynamics. In this
563 case, the simulated occupation probabilities can be compared to the solution of the Smoluchowski
564 equation with periodic boundary conditions (see the supporting information for details).
565 Alternatively, one can directly analyze time-evolution of the KMC trajectory and obtain the dif-
566 fusion tensor from a linear fit to the mean square displacement, $\overline{\Delta r_{i,\gamma} \Delta r_{i,\delta}} = 2dD_{\gamma\delta} t$.
567 The charge carrier mobility tensor, $\hat{\mu}$, for any value of the external field can be determined either
568 from the average charge velocity defined in eq. (2.41)

$$\langle \vec{v} \rangle = \sum_{i,j} p_j \omega_{ji} (\vec{r}_i - \vec{r}_j) = \hat{\mu} \vec{F} \, , \tag{2.46}$$

569 or directly from the KMC trajectory. In the latter case the velocity is calculated from the un-
570 wrapped (if periodic boundary conditions are used) charge displacement vector divided by the
571 total simulation time. Projecting this velocity on the direction of the field $\vec{F}$ yields the charge car-
572 rier mobility in this particular direction. In order to improve statistics, mobilities can be averaged
573 over several KMC trajectories and MD snapshots.

2.10.4   Spatial correlations of energetic disorder

Long-range, e.g. electrostatic and polarization, interactions often result in spatially correlated disorder [43], which affects the onset of the mobility-field (Poole-Frenkel) dependence [39, 44, 45]. To quantify the degree of correlation, one can calculate the spatial correlation function of E_i and E_j at a distance r_{ij}

$$C(r_{ij}) = \frac{\langle (E_i - \langle E \rangle)(E_j - \langle E \rangle) \rangle}{\langle (E_i - \langle E \rangle)^2 \rangle},$$ (2.47)

where $\langle E \rangle$ is the average site energy. $C(r_{ij})$ is zero if E_i and E_j are uncorrelated and 1 if they are fully correlated. For a system of randomly oriented point dipoles, the correlation function decays as $1/r$ at large distances [46].

For systems with spatial correlations, variations in site energy differences, ΔE_{ij}, of pairs of molecules from the neighbor list are smaller than variations in site energies, E_i, of all individual molecules. Since only neighbor list pairs affect transport, the distribution of ΔE_{ij} rather than that of individual site energies, E_i, should be used to characterize energetic disorder.

Note that the `eanalyze` calculator takes into account *all* contributions to the site energies

> 🛈 **Analyze distribution and correlations of site energeies**
> `ctp_run -o options.xml -f state.sql -e eanalyze`

2.10.5   DFT-based transfer integrals using DIPRO

The calculation of one electronic coupling element based on DFT using the `DIPRO` method requires the overlap matrix of atomic orbitals $\mathcal{S}$, the expansion coefficients for monomer $\bar{\boldsymbol{\lambda}}_{(k)} = \{\lambda_\alpha^{(k)}\}$ and dimer orbitals $\bar{\mathbf{D}}_{(n)} = \{D_\alpha^{(n)}\}$, as well as the orbital energies E_n of the dimer are required as input. In practical situations, performing self-consistent quantum-chemical calculations for each individual monomer and one for the dimer to obtain this input data is extremely demanding. Several simplifications can be made to reduce the computational effort, such as using non-Counterpoise basis sets for the monomers (thereby decoupling the monomer calculations from the dimer run) and performing only a single SCF step in a dimer calculation starting from an initial guess formed from a superposition of monomer orbitals. This "noCP+noSCF" variant of `DIPRO` is shown in figure 2.4(a) and recommended for production runs. A detailed comparative study of the different variants can be found in [5].

The code currently contains supports evaluation of transfer integrals from quantum-chemical calculations performed with the `Gaussian`, `Turbomole`, and `NWChem` packages. The interfacing procedure consists of three main steps: generation of input files for monomers and dimers, performing the actual quantum-chemical calculations, and calculating the transfer integrals.

Monomer calculations

First, hopping sites and a neighbor list need to be generated from the atomistic topology and trajectory and written to the `state.sql` file. Then the parallel `edft` calculator manages the calculation of the monomer properties required for the determination of electronic coupling elements. Specifically, the individual steps it performs are:

1. Creation of a job file containing the list of molecules to be calculated with DFT

> 🛈 **Writing job file for `edft`**
> `ctp_parallel -o options.xml -f state.sql -e edft -j write`

609   2. Running of all jobs in job file

> **Running all `edft` jobs**
> | `ctp_parallel -o options.xml -f state.sql -e edft -j run`

610   which includes

611   • creating the input files for the DFT calculation (using the package specified in `options.xml`)
612   in the directory

613   `OR_FILES/package/frame_F/mol_M`

614   where `F` is the index of the frame in the trajectory, `M` is the index of a molecule in this
615   frame,

616   • executing the DFT run, and

617   • after completion of this run, parsing the output (number of electrons, basis set, molec-
618   ular orbital expansion coefficients), and saving it in compressed form to

619   `OR_FILES/molecules/frame_F/molecule_M.orb`

620   **Calculating the transfer integrals**

sec:idft

621   After the momomer calculations have been completed successfully, the respective runs for dimers
622   from the neighborlist can be performed using the parallel `idft calculator`, which manages the
623   DFT runs for the hopping pairs and determines the coupling element using `DIPRO`. Again, sev-
624   eral steps are required:

625   1. Creation of a job file containing the list of pairs to be calculated with DFT

> **Writing job file for `idft`**
> | `ctp_parallel -o options.xml -f state.sql -e idft -j write`

626   2. Running of all jobs in job file

> **Running all `idft` jobs**
> | `ctp_parallel -o options.xml -f state.sql -e idft -j run`

627   which includes

628   • creating the input files (including the merged guess for a noSCF calculation, if re-
629   quested) for the DFT calculation (using the package specified in `options.xml`) in the
630   directory

631   `OR_FILES/package/frame_F/pair_M_N`

632   where `M` and `N` are the indices of the molecules in this pair,

633   • executing the DFT run, and

634   • after completion of this run, parsing the output (number of electrons, basis set, molec-
635   ular orbital expansion coefficients and energies, atomic orbital overlap matrix), and
636   saving the pair information in compressed form to

637   `OR_FILES/pairs/frame_F/pair_M_N.orb`

638   • loading the monomer orbitals from the previously saved `*.orb files`.

639   • calculating the coupling elements and write them to the job file

640   3. Reading the coupling elements from the job file and saving them to the `state.sql` file

641

ⓘ **Saving `idft` results from job file to `state.sql`**
```
ctp_parallel -o options.xml -f state.sql -e idft -j read
```

Chapter 3

Input and output files

3.1 Atomistic topology

If you are using `GROMACS` for generating atomistic configurations, it is possible to directly use the topology file provided by `GROMACS` (`topology.tpr`). In this case the `GROMACS` residue and atom names should be used to specify the coarse-grained topology and conjugated segments.

A custom topology can also be defined using an XML file. Moreover, it s possible to partially overwrite the information provided in, for example, `GROMACS` topology file. We will illustrate how to create a custom topology file using DCV2T. The structure of DCV2T, together with atom type definitions, is shown in fig. 3.1. DCV2T has two thiophene (THI) and two dicyanovinyl (NIT) residues. The pdb file which contains residue types, residue numbering, atom names, atom types, and atom coordinates is shown in listing 3.1.

Figure 3.1: (a) DCV2T with atoms labelled according to `residue_number:residue_name:atom_name`. There are four residues and two residue types: thiophene (THI) and dicyanovinyl (NIT). The corresponding pdb file is shown in listing 3.1. Atom numbering is used to split conjugated segments on rigid fragments and to link atomistic ((b) from GROMACS topology) and quantum descriptions (c).

Listing 3.1: pdb file of DCV2T.

```
654
655   HETATM    1   N1   NIT    1      2.388    8.533   11.066   1.00   4.14              N
656   HETATM    2   CN1  NIT    1      1.984    9.553   10.718   1.00   2.54              C
657   HETATM    3   N2   NIT    1     -1.138   10.872   10.087   1.00   3.24              N
658   HETATM    4   CN2  NIT    1      0.003   10.871   10.213   1.00   2.37              C
659   HETATM    5   CC1  NIT    1      1.441   10.824   10.327   1.00   1.91              C
660   HETATM    6   C1   NIT    1      2.193   11.939   10.071   1.00   1.61              C
661   HETATM    7   HN1  NIT    1      1.715   12.710    9.872   1.00   1.97              H
662   HETATM    8   S1   THI    2      4.758   10.743   10.130   1.00   1.52              S
663   HETATM    9   CA1  THI    2      3.613   12.024    9.948   1.00   1.22              C
664   HETATM   10   CA2  THI    2      6.099   11.836    9.997   1.00   1.30              C
665   HETATM   11   CB1  THI    2      4.251   13.243    9.782   1.00   1.39              C
666   HETATM   12   CB2  THI    2      5.658   13.131    9.818   1.00   1.45              C
667   HETATM   13   HC1  THI    2      3.800   14.047    9.660   1.00   1.66              H
668   HETATM   14   HC2  THI    2      6.230   13.860    9.731   1.00   1.74              H
669   HETATM   15   S1   THI    3      8.803   12.414    9.882   1.00   1.38              S
670   HETATM   16   CA1  THI    3      7.456   11.347   10.094   1.00   1.37              C
671   HETATM   17   CA2  THI    3      9.940   11.122   10.152   1.00   1.42              C
672   HETATM   18   CB1  THI    3      7.873   10.048   10.355   1.00   1.73              C
673   HETATM   19   CB2  THI    3      9.267    9.926   10.399   1.00   1.82              C
674   HETATM   20   HC1  THI    3      7.288    9.335   10.487   1.00   2.05              H
675   HETATM   21   HC2  THI    3      9.704    9.123   10.576   1.00   2.21              H
676   HETATM   22   N1   NIT    4     11.235   14.572    9.094   1.00   3.08              N
677   HETATM   23   CN1  NIT    4     11.665   13.566    9.441   1.00   2.04              C
678   HETATM   24   N2   NIT    4     14.733   12.005   10.009   1.00   2.17              N
679   HETATM   25   CN2  NIT    4     13.590   12.149    9.933   1.00   1.77              C
680   HETATM   26   CC1  NIT    4     12.156   12.282    9.861   1.00   1.71              C
681   HETATM   27   C1   NIT    4     11.363   11.220   10.154   1.00   1.59              C
682   HETATM   28   HN1  NIT    4     11.813   10.440   10.389   1.00   1.89              H
683
```

3.2 Mapping file

The mapping file (referred here as `map.xml`) is used by the program `ctp_map` to convert an atomistic trajectory to a trajectory with conjugated segments and rigid fragments. This trajectory is stored in a state file and contains positions, names, types of atoms belonging to rigid fragments. The description of the mapping options is given in table 3.1. An example of `map.xml` for a DCV2T molecule is shown in listing 3.2.

The file `map.xml` contains the whole electrostatic information about the molecules as well as the structural information. The toolpdb2map creates a `map.xml` from a pdb file and is a good starting point for further refinement.

Table 3.1: Description of the XML mapping file (`map.xml`).

Listing 3.2: Exale of `map.xml` for DCV2T. Each rigid fragment (coarse-grained bead) is defined by a list of atoms. Atom numbers, names, and residue names should correspond to those used in GROMACS topology (see the corresponding listing 3.1 of the pdb file).

```xml
<topology> <!-- this file is used to conver an atomistic trajectory to conjugated segments -->
<molecules>
<molecule>
   <name>DCV2T-MOL</name> <!-- name of the conjugated molecule -->
   <mdname>Protein</mdname> <!-- corresponding name of this molecule in the MD trajectory, should be
       the same as the name given at the end of topol.top-->
   <segments>
   <segment>
       <name>DCV</name> <!-- name of the conjugated segment within the molecule -->
       <qmcoords>QC_FILES/DCV2T.xyz</qmcoords>  <!-- QM coordinates of the conjugated segment -->


           <!-- IZINDO INPUT -->
       <basisset>INDO</basisset>
       <orbitals>QC_FILES/DCV2T.orb</orbitals>
       <torbital_h>50</torbital_h><!-- Number of the HOMO Orbital (e.g. alpha electrons, can be
           found in the log-file belonging to DCV2T.orb) -->

           <!-- EMULTIPOLE INPUT -->
       <multipoles_n>MP_FILES/DCV2T.mps</multipoles_n><!-- Multipole file for neutral state -->
       <multipoles_h>MP_FILES/DCV2T_h.mps</multipoles_h><!-- Multipole file for hole state -->

       <map2md>0</map2md><!-- specifies if planar QM coordinates (map2md=0) or MD coordinates (
           map2md=1) of atoms are used for distribution of partial charges. For MD coordinates the
           order and numbering in <mdatoms> and <mpoles> must be identical it has no impact on the
           qm e.g. DFT or GWBSE calculations-->

           <!-- EINTERNAL INPUT -->
       <U_cC_nN_h>0.0</U_cC_nN_h> <!-- Site energy      -->
       <U_nC_nN_h>0.1</U_nC_nN_h> <!-- Reorg. discharge -->
       <U_cN_cC_h>0.1</U_cN_cC_h> <!-- Reorg. charge    -->

           <!-- MD QM MP Mapping -->
       <fragments>
       <fragment>
        <name>NI1</name> <!-- name of the rigid fragment within the segment -->
        <!-- list of atoms in the fragment resnum:resname:atomname -->
        <mdatoms>1:NIT:N1 1:NIT:CN1 1:NIT:N2  1:NIT:CN2 1:NIT:CC1 1:NIT:C1  1:NIT:HN1</mdatoms>
        <!-- corresponding ground state geometry atomnumber:atomtype  read from .xyz file-->
        <qmatoms>   20:N     19:C     14:N     13:C     12:C     11:C     23:H </qmatoms>
        <!-- corresponding group state geometry multipoles read from .mps files -->
        <mpoles>   20:N     19:C     14:N     13:C     12:C     11:C     23:H </mpoles>
        <!-- weights to determine the fragment center (here CoM is used) -->
        <weights>   14     12      14      12      12      12       1 </weights>
        <!-- three atoms: define a cartesian local frame, two atoms: fragment is assumed to be
            rotationally invariant around the axis, one atom: fragment is assumed isotropic -->
        <localframe> 20 19 14 </localframe>
        <!-- Optional parameters (if not set <localframe> is used): used when atom labels in the .mps
            and .xyz file differ or more sites in the .mps file are used, so refers to <mpoles> -->
        <localframe_mps> 20 19 14 </localframe_mps>
        <!-- Optional parameters (if not set <localframe> is used): weights to determine the
            fragment center (here CoM is used), used when atom labels in the .mps and .xyz file
            differ or additional sites in the .mps file are used -->
        <weights_mps>   14     12      14      12      12      12       1 </
            weights_mps>
```

```
749          <!-- Optional flag: says  if a site is virtual or not, (virtual=1, real=0)-->
750          <virtual_mps>      0        0        0        0       0        0       0  </
751              virtual_mps>
752        </fragment>
753
754      <fragment>
755          <name>TH1</name>
756          <mdatoms>2:THI:S1 2:THI:CA1 2:THI:CA2 2:THI:CB1 2:THI:CB2 2:THI:HC1 2:THI:HC2</mdatoms>
757          <qmatoms>    7:S       8:C      6:C       9:C     10:C      24:H      25:H </qmatoms>
758          <mpoles>     7:S       8:C      6:C       9:C     10:C      24:H      25:H </mpoles>
759          <weights>     32        12       12        12       12        1        1 </weights>
760          <localframe> 7 8 6 </localframe>
761      </fragment>
762
763      <fragment>
764          <name>TH2</name>
765          <mdatoms>3:THI:S1 3:THI:CA1 3:THI:CA2 3:THI:CB1 3:THI:CB2 3:THI:HC1 3:THI:HC2</mdatoms>
766          <qmatoms>    3:S       4:C      2:C       5:C      1:C      26:H      27:H </qmatoms>
767          <weights>     32        12       12        12       12        1        1 </weights>
768          <localframe> 3 4 2 </localframe>
769      </fragment>
770
771      <fragment>
772          <name>NI2</name>
773          <mdatoms>4:NIT:N1 4:NIT:CN1 4:NIT:N2  4:NIT:CN2 4:NIT:CC1 4:NIT:C1  4:NIT:HN1</mdatoms>
774          <qmatoms>   22:N      21:C      18:N      17:C     16:C     15:C      28:H </qmatoms>
775          <mpoles>    22:N      21:C      18:N      17:C     16:C     15:C      28:H </mpoles>
776          <weights>    14        12        14        12       12       12        1 </weights>
777          <localframe> 22 21 18 </localframe>
778      </fragment>
779      </fragments>
780
781    </segments>
782 </molecule>
783 </molecules>
784 </topology>
785
```

3.3 Conjugated segments

The file describing hopping sites, or conjugated segments, is used by practically all programs and calculators. It links the coarse-grained trajectory (positions and orientations of rigid fragments) and quantum-mechanical descriptions of all conjugated segments. The description of this XML file (segments.xml) is given in table 3.2. An example for DCV2T is shown in listing 3.3.

Table 3.2: Description of conjugated segments (segments.xml).

Listing 3.3: XML file describing conjugated segments. Note that the mapping and weights for each segment are separated by a colon.

```
791
792 <segments>
793      <segment> <!-- DCV2T here is one conjugated segment -->
794          <coordinates>DCV2T.xyz</coordinates> <!-- xyz coordinates -->
795          <orbitals>DCV2T.fort.7</orbitals> <!-- ZINDO orbitals [GAUSSIAN] -->
796          <basisset>INDO</basisset>
797          <torbital>50</torbital> <!-- HOMO for hole conduction -->
798          <edischarging>0.084</edischarging> <!-- reorganization energy -->
799          <echarging>0.086</echarging> <!-- reorganization energy -->
800          <qneutral>DCV2T_neutr.esp</qneutral> <!-- partial charges of a neutral molecule -->
801          <qcharged>DCV2T_catio.esp</qcharged> <!-- partial charges of a cation -->
802          <energy>0.0</energy> <!-- cite energy -->
803          <name>DCV2T</name> <!-- name of the conjugated segment -->
804          <map> <!-- rigid fragments separated by a colon -->
805                  22 21 18 17 16 15 28:
806                   3  2  4  1  5 27 26:
807                   7  6  8 10  9 25 24:
808                  20 19 14 13 12 11 23
809          </map>
810          <weights> <!-- for centers of rigid fragments -->
811                  14 12 14 12 12 12 1:
812                  32 12 12 12 12 1 1:
813                  32 12 12 12 12 1 1:
```

```
814                         14 12 14 12 12 12 1
815                    </weights>
816
817
818      </segments>
```

3.4 Molecular orbitals

If the semi-empirical method is used to calculate electronic coupling elements, molecular orbitals of all molecules must be supplied. They can be generated using `Gaussian` program. The `Gaussian` input file for DCV2T is shown in listing 3.4. Provided with this input, `Gaussian` will generate `fort.7` file which contains the molecular orbitals of a DCV2T. This file can be renamed to `DCV2T.orb`. Note that the order of the atoms in the input file and the order of coefficients should always match. Therefore, the coordinate part of the input file must be supplied together with the orbitals. We will assume the coordinates, in the format `atom_type:  x y z`, is saved to the `DCV2T.xyz` file.

> ⚠️ **Attention**
>
> Izindo requires the specification of orbitals for hole and electron transport in `map.xml`. They are the HOMO and LUMO respectively and can be retrieved from the `log` file from which the `DCV2T.orb` file is generated. The number of `alpha electrons` is the HOMO, the LUMO is HOMO+1

Listing 3.4: `Gaussian` input file `get_orbitals.com` used for generating molecular orbitals. The first line contains the name of the check file, the second the requested RAM. `int=zindos` requests the method ZINDO, `punch=mo` states that the molecular orbitals ought to be written to the `fort.7` file, `nosymm` forbids use of symmetry and is necessary to ensure correct position of orbitals with respect to the provided coordinates. The two integer numbers correspond to the charge and multiplicity of the system: 0 1 corresponds to a neutral system with a multiplicity of one. They are followed by the types and coordinates of all atoms in the molecule.

```
828
829   %chk=DCV2T.chk
830   %mem=100Mb
831   #p int=zindos punch=mo nosymm
832
833   DCV2T molecular orbitals
834
835   0 1
836   S        -1.44650        2.12185        0.00135
837   C        -2.43098        0.58936       -0.00048
838   C        -1.59065       -0.51859       -0.00146
839   C        -0.21222       -0.22233       -0.00095
840   C         0.07761        1.13376        0.00040
841   S         2.87651        0.79316        0.00148
842   C         3.86099        2.32565        0.00235
843   C         3.02066        3.43359        0.00231
844   C         1.64223        3.13733        0.00162
845   C         1.35240        1.78125        0.00114
846   C        -3.85350        0.52245       -0.00081
847   C        -4.79569        1.52479       -0.00008
848   C        -6.18500        1.18622       -0.00117
849   C        -4.47544        2.91565        0.00081
850   C         5.28350        2.39256        0.00296
851   C         6.22569        1.39020        0.00327
852   C         7.61500        1.72876        0.00432
853   C         5.90542       -0.00064        0.00333
854   N        -7.32389        0.89743       -0.00195
855   N        -4.21872        4.06274        0.00142
856   N         8.75389        2.01754        0.00510
857   N         5.64864       -1.14772        0.00361
858   H        -1.98064       -1.52966       -0.00256
859   H         0.55785       -0.98374       -0.00169
```

```
860  H            3.41065          4.44466          0.00272
861  H            0.87216          3.89874          0.00147
862  H           -4.24640         -0.49192         -0.00188
863
864  H            5.67641          3.40692          0.00337
```

3.5   Monomer calculations for DFT transfer integrals

Listing 3.5: Example `package.xml` file for the `Gaussian` package required in the options of the `edft` calculator for the monomer calculations as preparation for the determination of transfer integrals using DIPRO.

```
866
867  <package>
868    <name>gaussian</name>
869    <executable>g09</executable>
870    <checkpoint></checkpoint>
871    <scratch></scratch>
872
873    <charge>0</charge>
874    <spin>1</spin>
875    <options># pop=minimal pbepbe/6-311g** scf=tight punch=mo nosymm test</options>
876    <memory>1Gb</memory>
877    <threads>2</threads>
878
879    <cleanup></cleanup>
880  </package>
881
```

Listing 3.6: Example `package.xml` file for the `Turbomole` package required in the options of the `edft` calculator for the monomer calculations as preparation for the determination of transfer integrals using DIPRO.

```
882
883  <package>
884    <name>turbomole</name>
885    <executable>ridft</executable>
886    <scratch>/tmp</scratch>
887
888    <options>
889  TITLE
890  a coord
891  *
892  no
893  b all def-TZVP
894  *
895  eht
896  y
897  0
898  y
899  dft
900  on
901  func
902  pbe
903  grid
904  m3
905  *
906  ri
907  on
908  m 300
909  *
910  scf
911  conv
912  7
913  iter
914  200
915
```

```
916  marij
917
918  q
919    </options>
920
921    <cleanup></cleanup>
922
923  </package>
```

Listing 3.7: Example `package.xml` file for the `NWChem` package required in the options of the `edft` calculator for the monomer calculations as preparation for the determination of transfer integrals using `DIPRO`.

```
924
925  <package>
926    <name>nwchem</name>
927    <executable>nwchem</executable>
928    <checkpoint></checkpoint>
929    <scratch>/tmp/nwchem</scratch>
930    <charge>0</charge>
931    <spin>1</spin>
932    <threads>1</threads>
933    <memory></memory>
934    <options>
935  start
936  basis
937   * library 6-311gss
938  end
939  memory 1500 mb
940
941  dft
942   xc xpbe96 cpbe96
943   direct
944   iterations 100
945   noprint "final vectors analysis"
946  end
947  task dft
948  </options>
949    <cleanup></cleanup>
950
951  </package>
```

3.6 Pair calculations for DFT transfer integrals

Listing 3.8: Example `package.xml` file for the `Gaussian` package required in the options of the `idft` calculator for the pair calculations and the determination of transfer integrals using `DIPRO`.

```
953
954  <package>
955    <name>gaussian</name>
956    <executable>g09</executable>
957    <checkpoint></checkpoint>
958    <scratch></scratch>
959
960    <charge>0</charge>
961    <spin>1</spin>
962    <options># pop=minimal pbepbe/6-311g** nosymm  IOp(3/33=1,3/36=-1) punch=mo guess=cards scf=(maxcycle=1,
963    <memory>1Gb</memory>
964    <threads>1</threads>
965
966    <cleanup></cleanup>
967
968  </package>
```

Listing 3.9: Example `package.xml` file for the `Turbomole` package required in the options of the `idft` calculator for the pair calculations and the determination of transfer integrals using `DIPRO`.

```
969
970  <package>
```

```
971      <name>turbomole</name>
972      <executable>ridft</executable>
973      <scratch>/tmp</scratch>
974
975      <options>
976    $intsdebug cao
977    a coord
978    *
979    no
980    b all def-TZVP
981    *
982    eht
983    y
984    0
985    y
986    dft
987    on
988    func
989    pbe
990    grid
991    m3
992    *
993    ri
994    on
995    m 300
996    *
997    scf
998    conv
999    7
1000   iter
1001   1
1002   diis
1003   3
1004   damp
1005   0.00
1006
1007
1008
1009   marij
1010
1011   q
1012     </options>
1013
1014     <cleanup></cleanup>
1015   </package>
1016
```

Listing 3.10: Example `package.xml` file for the `NWChem` package required in the options of the `idft` calculator for the pair calculations and the determination of transfer integrals using `DIPRO`.

```
1017
1018   <package>
1019     <name>nwchem</name>
1020     <executable>nwchem</executable>
1021     <checkpoint></checkpoint>
1022     <scratch>/tmp/nwchem</scratch>
1023     <charge>0</charge>
1024     <spin>1</spin>
1025     <memory></memory>
1026     <threads>1</threads>
1027     <options>
1028   start
1029   basis
1030    * library 6-311gss
1031   end
1032   memory 1500 mb
1033
1034   dft
1035    print "ao overlap"
```

```
1036   xc xpbe96 cpbe96
1037   direct
1038   iterations 1
1039   convergence nodamping nodiis
1040   noprint "final vectors analysis"
1041   vectors input system.movecs
1042   end
1043 task dft
1044 </options>
1045   <cleanup></cleanup>
1046 </package>
1047
```

3.7   DFT transfer integrals

list:TI_xml

Listing 3.11: Example `TI.xml` file created as the output of a `DIPRO` calculation. Due to slightly different implementations, the orbitals indices refer to monomer indices in a `Gaussian` run but to indices in the merged dimer guess in a `Turbomole` run.

```
1049
1050  <pair name="pair_100_155">
1051     <parameters>
1052        <HOMO_A>162</HOMO_A>
1053        <NoccA>1</NoccA>
1054        <LUMO_A>164</LUMO_A>
1055        <NvirtA>1</NvirtA>
1056        <HOMO_B>161</HOMO_B>
1057        <NoccB>1</NoccB>
1058        <LUMO_B>163</LUMO_B>
1059        <NvirtB>1</NvirtB>
1060     </parameters>
1061     <transport name="hole">
1062        <channel name="single">
1063           <J>1.546400416750696E-003</J>
1064           <e_A>-6.30726450715697</e_A>
1065           <e_B>-6.36775613794166</e_B>
1066        </channel>
1067        <channel name="multi">
1068           <molecule name="A">
1069              <e_HOMOm0>-6.30726450715697</e_HOMOm0>
1070           </molecule>
1071           <molecule name="B">
1072              <e_HOMOm0>-6.36775613794166</e_HOMOm0>
1073           </molecule>
1074           <dimer name="integrals">
1075              <T_00>1.546400416750696E-003</T_00>
1076              <J_sq_degen>2.391354248926727E-006</J_sq_degen>
1077              <J_sq_boltz>2.391354248926727E-006</J_sq_boltz>
1078           </dimer>
1079        </channel>
1080     </transport>
1081     <transport name="electron">
1082        <channel name="single">
1083           <J>-2.797473760331286E-003</J>
1084           <e_A>-4.50318366770689</e_A>
1085           <e_B>-4.53143397059021</e_B>
1086        </channel>
1087        <channel name="multi">
1088           <molecule name="A">
1089              <e_LUMOp0>-4.50318366770689</e_LUMOp0>
1090           </molecule>
1091           <molecule name="B">
1092              <e_LUMOp0>-4.53143397059021</e_LUMOp0>
1093           </molecule>
1094           <dimer name="integrals">
1095              <T_00>-2.797473760331286E-003</T_00>
```

```
1096                    <J_sq_degen>7.825859439742066E-006</J_sq_degen>
1097                    <J_sq_boltz>7.825859439742066E-006</J_sq_boltz>
1098              </dimer>
1099        </channel>
1100    </transport>
1101
1102  </pair>
```

3.8   State file

All data structures are saved to the `state.sql` file in sqlite3 format, see http://www.sqlite.org/. They are available in form of tables in the `state.sql` file as can be seen by the command

`sqlite3 state.sql " .tables "`

An example of such a table are `molecules`. The full table can be displayed using the command (similar for the other tables)

`sqlite3 state.sql " SELECT * FROM molecules "`

The meaning of all the entries in the table can be displayed by a command like

`sqlite3 state.sql " .SCHEMA molecules "`

The first and second entry are integers for internal and regular id of the molecule and the third entry is the name. A single field from the table like the name of the molecule can be displayed by a command like

`sqlite3 state.sql " SELECT name FROM molecules "`

Besides `molecules`, the following tables are stored in the `state.sql`:

`conjseg_properties`:
Conjugated segments are stored with id, name and x,y,z coordinates of the center of mass in nm.

`conjsegs`:
Reorganization energies for charging or discharging a conjugated segment are stored together with the coulomb energy and any other user defined energy contribution (in eV) and occupation probabilities.

`pairs`:
The pairs from the neighborlist are stored with the pair id, the id of the first and second segment, the rate from the first to the second , the rate from the second to the first (both in s^{-1}) and the x,y,z coordinates in nm of the distance between the first and the second segment.

`pairintegrals`:
Transfer integrals for all pairs are stored in the following way: The pair id , the number for counting possible different electronic overlaps (e.g if only the frontier orbitals are taken into account this is always zero, while an effective value is stored in addition to the different overlaps of e.g. HOMO-1 and HOMO-1 if more frontier orbitals are taken into account) and the integral in eV.

`pairproperties`:
The outer sphere reorganization energy of all pairs is stored by an id, the pair id, a string `lambda_outer` and the energy in eV.

`conjsegs`:
Conjugated segments are saved in the following way: The id, the name, the type, the molecule id, the time frame, the x,y,z coordinates in nm and the occupation probability.

`conjseg_properties`:
Properties of the conjugated segments like reorganization energies for charging or discharging a charge unit or the coulomb contribution to the site energy are stored by: id, conjugated segment id, a string like `lambda_intra_charging`, `lambda_intra_discharging` or `energy_coulomb` and a corresponding value in eV.

The tables `rigidfrag_properties`, `rigidfrags` and `frames` offer information about rigid fragments and time frames including periodic boundary conditions.

The data in the `state.sql` file can also be modified by the user. Here is an example how to modify the transfer integral between the conjugated segments number one and two assuming that they are in the neighborlist. Their pair id can be found by the command

```
1148  pair_ID=`sqlite3 state.sql "SELECT _id FROM pairs WHERE conjseg1=1 AND conjseg2=2"`
```
1149 The old value of the transfer integral can be deleted using
```
1150  sqlite3 state.sql "DELETE FROM pair_integrals WHERE pair=$pair_ID"
```
1151 Finally the new transfer integral J can be written to the `state.sql` file by the command
```
1152  sqlite3 state.sql "INSERT INTO pair_integrals (pair,num,J) VALUES ($pair_ID,0,$J)"
```
1153 Here the `num=0` indicates that only the effective transfer integrals is written to the file, while other
1154 values of `num` would correspond to overlap between other orbitals than the frontier orbitals.
1155 In a similar way the coulomb contribution to the site energy of the first conjugated segment can
1156 be overwritten by first getting its id
```
1157  c_ID=`sqlite3 state.sql "SELECT _id from conjseg_properties  where conjseg=1 AND
1158  key =\"energy_coulomb\""
```
1159 Then deleting the old value
```
1160  sqlite3 state.sql "DELETE FROM from conjseg_properties WHERE _id=$c_ID"
```
1161 Then the new coulomb energy E can be written to this id
```
1162  sqlite3 state.sql "INSERT INTO conjseg_properties (_id,conjseg,key,value)
1163  VALUES ($c_ID,1,\"energy_coulomb\",$E)"
```
1164 Finally the resulting coulomb contribution to all conjugated segments can be displayed by
```
1165  sqlite3 state.sql "SELECT * from conjseg_properties WHERE key=\"energy_coulomb\""
1166
```

Chapter 4

Reference

4.1 Programs

Programs execute specific tasks (calculators).

4.1.1 ctp_map

Generates QM | MD topology

```
-h [ --help ]    display this help and exit
-v [ --verbose ]    be loud and noisy
-t [ --topology ] arg topology
-c [ --coordinates ] arg coordinates or trajectory
-s [ --segments ] arg definition of segments and fragments
-f [ --file ] arg state file
--man   output man-formatted manual pages
--tex   output tex-formatted manual pages
```

4.1.2 ctp_dump

Extracts information from the state file

```
-h [ --help ]    display this help and exit
-v [ --verbose ]    be loud and noisy
-o [ --options ] arg calculator options
-f [ --file ] arg sqlight state file, *.sql
-i [ --first-frame ] arg (=1) start from this frame
-n [ --nframes ] arg (=1) number of frames to process
-t [ --nthreads ] arg (=1) number of threads to create
-s [ --save ] arg (=1) whether or not to save changes to state file
-e [ --extract ] arg List of extractors separated by ',' or ' '
-l [ --list ]    Lists all available extractors
-d [ --description ] arg Short description of an extractor
--man   output man-formatted manual pages
--tex   output tex-formatted manual pages
```

4.1.3 ctp_tools

Runs charge transport tools

```
-h [ --help ]    display this help and exit
-v [ --verbose ]    be loud and noisy
```

1200   `-t [ --nthreads ] arg (=1)` number of threads to create
1201   `-o [ --options ] arg` calculator options
1202   `--man` output man-formatted manual pages
1203   `--tex` output tex-formatted manual pages
1204   `-e [ --execute ] arg` List of tools separated by ',' or ' '
1205   `-l [ --list ]` Lists all available tools
1206   `-d [ --description ] arg` Short description of a tool

4.1.4 ctp_run
1207

prog:ctp_run

1208 Runs charge transport calculators
1209   `-h [ --help ]` display this help and exit
1210   `-v [ --verbose ]` be loud and noisy
1211   `-o [ --options ] arg` calculator options
1212   `-f [ --file ] arg` sqlight state file, *.sql
1213   `-i [ --first-frame ] arg (=1)` start from this frame
1214   `-n [ --nframes ] arg (=1)` number of frames to process
1215   `-t [ --nthreads ] arg (=1)` number of threads to create
1216   `-s [ --save ] arg (=1)` whether or not to save changes to state file
1217   `-e [ --execute ] arg` List of calculators separated by ',' or ' '
1218   `-l [ --list ]` Lists all available calculators
1219   `-d [ --description ] arg` Short description of a calculator
1220   `--man` output man-formatted manual pages
1221   `--tex` output tex-formatted manual pages

4.1.5 ctp_parallel
1222

prog:ctp_parallel

1223 Runs job-based heavy-duty calculators
1224   `-h [ --help ]` display this help and exit
1225   `-v [ --verbose ]` be loud and noisy
1226   `-o [ --options ] arg` calculator options
1227   `-f [ --file ] arg` sqlite state file, *.sql
1228   `-i [ --first-frame ] arg (=1)` start from this frame
1229   `-n [ --nframes ] arg (=1)` number of frames to process
1230   `-t [ --nthreads ] arg (=1)` number of threads to create
1231   `-s [ --save ] arg (=1)` whether or not to save changes to state file
1232   `-r [ --restart ] arg` restart pattern: 'host(pc1:234) stat(FAILED)'
1233   `-c [ --cache ] arg (=8)` assigns jobs in blocks of this size
1234   `-j [ --jobs ] arg (=run)` task(s) to perform: input, run, import
1235   `-m [ --maxjobs ] arg (=-1)` maximum number of jobs to process (-1 = inf)
1236   `-e [ --execute ] arg` List of calculators separated by ',' or ' '
1237   `-l [ --list ]` Lists all available calculators
1238   `-d [ --description ] arg` Short description of a calculator
1239   `--man` output man-formatted manual pages
1240   `--tex` output tex-formatted manual pages

4.1.6 moo_overlap
1241

prog:moo_overlap

1242   `-h [ --help ]` display this help and exit
1243   `-v [ --verbose ]` be loud and noisy
1244   `--man` output man-formatted manual pages
1245   `--tex` output tex-formatted manual pages
1246   `--conjseg arg` xml file describing two conjugated segments
1247   `--pos1 arg` position and orientation of molecule 1

1248   `--pos2 arg` position and orientation of molecule 2

1249   `--pdb arg (=geometry.pdb)` pdb file of two molecules

1250 4.2 Calculators

1251 Calculator is a piece of code which computes specific system properties, such as site energies,

1252 transfer integrals, etc. `ctp_run`, `kmc_run` are wrapper programs which executes such calcula-

1253 tors. The generic syntax is

1254 `ctp_run -e "calc1, calc2, ..." -o options.xml`

1255 File `options.xml` lists all options needed to run a specific calculator. The format of this file is

1256 explained in listing 4.1. A complete list of calculators is given in the calculators reference section.

Listing 4.1: A part of the `options.xml` file with options for the `calculator_name{1,2}` calculators.

```
1258  <calculator_name1>
1259          <option1>value1</option1>
1260          <option2>value2</option2>
1261          ...
1262  </calculator_name1>
1263
1264  <calculator_name2>
1265          <option1>value1</option1>
1266          <option2>value2</option2>
1267          ...
1268  </calculator_name2>
1269  ...
1270
```

1271 A list of all calculators and their short descriptions can be obtain using

1272 `ctp_run --list`

1273 A detailed description of all options of a specific calculator(s) is available via

1274 `ctp_run --desc calc1,calc2,...`

1275 4.2.1 coupling

1276 Electronic couplings from log and orbital files (GAUSSAIN, TURBOMOLE, NWChem)

option	default	unit	description
package			First-principles package
output	coupling.out.xml		Output file
degeneracy	0	eV	Criterium for the degeneracy of two levels
moleculeA			
log	A.log		Log file of molecule A
orbitals	A.orb		Orbitals file
levels	3		Output HOMO, ..., HOMO-levels; LUMO, ..., LUMO+levels
trim	2		
moleculeB			
log	B.log		Log file of molecule B
orbitals	B.orb		Orbitals file
levels	3		Output HOMO, ..., HOMO-levels; LUMO, ..., LUMO+levels
trim	2		
dimerAB			
log	AB.log		Log file of dimer AB
orbitals	A.orb		Orbitals file

1277 Return to the description of `coupling`.

4.2.2   log2mps

Generates an mps-file (with polar-site definitions) from a QM log-file

option	default	unit	description
package			QM package
logfile			Log-file generated by QM package, with population/esp-fit data

Return to the description of `log2mps`.

4.2.3   molpol

Molecular polarizability calculator (and optimizer)

option	default	unit	description
mpsfiles			
input			mps input file
output			mps output file
polar			xml file with infos on polarizability tensor
induction			
expdamp			Thole sharpness parameter
wSOR			mixing factor for convergence
maxiter			maximum number of iterations
tolerance			rel. tolerance for induced moments
target			
optimize			if 'true', refine atomic polarizabilities to match molecular polarizable volume specified in target.molpol
molpol			target polarizability tensor in format xx xy xz yy yz zz (this should be in the eigen-frame, hence xy = xz = yz = 0), if optimize=true the associated polarizable volume will be matched iteratively and the resulting set of polar sites written to mpsfiles.output
tolerance			relative tolerance when optimizing the polarizable volume

Return to the description of `molpol`.

4.2.4   pdb2map

Converts MD + QM files to VOTCA mapping. Combinations: pdb+xyz,gro+xyz,pdb

option	default	unit	description
pdb	conf.pdb		Input pdb file
gro	conf.gro		Input gro file
xyz	conf.xyz		Input xyz file
xml	conf.xml		Resulting xml file

Return to the description of `pdb2map`.

4.2.5   pdb2top

Generates fake Gromacs topology file .top

option	default	unit	description
num	1		Num of mols in the box
pdb	conf.pdb		Input pdb file
gro	conf.gro		Input gro file

1289 Return to the description of `pdb2top`.

4.2.6 ptopreader

calc:ptopreader

1291 Reads binary .ptop-files (serialized from ewdbgpol) and processes them into something readable

option	default	unit	description
ptop_file			Binary archive .ptop-file

1292 Return to the description of `ptopreader`.

4.2.7 eanalyze

calc:eanalyze

1294 Histogram and correlation function of site energies and pair energy differences

option	default	unit	description
resolution_sites		eV	Bin size for site energy histogram
resolution_pairs		eV	Bin size for pair energy histogram
resolution_space		eV	Bin size for site energy correlation
states			?

1295 Return to the description of `eanalyze`.

4.2.8 eimport

calc:eimport

1297 Imports site energies from the output file of emultipole and writes them to the state file

option	default	unit	description

1298 Return to the description of `eimport`.

4.2.9 einternal

calc:einternal

1300 Reads in site and reorganosation energies and writes them to the state file

option	default	unit	description
energiesXML			XML input file with vacuum site, reorganization (charging, discharging) energies

1301 Return to the description of `einternal`.

4.2.10 emultipole

calc:emultipole

1303 Evaluates polarization contribution based on the Thole model

option	default	unit	description
multipoles			Polar Site Definitions in GDMA punch-file format
control			Control options for induction computation
induce	1		Enter '1' / '0' to toggle induction on / off
first			First segment for which to compute site energies
last			Last segment for which to compute site energies
output			File to write site energies to. Site energies are also stored in the state file
check			Check mapping of polar sites to fragment
tholeparam			Thole parameters required for charge-smearing
cutoff		nm	Cut-off beyond which all interactions are neglected

cutoff2		nm	Cut-off beyond which polarization is neglected
expdamp			Damping exponent used in exponential damping function
scaling			1-n interaction scaling, currently not in use
esp			Control options for potential calculation
calcESP			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
cube			
grid			XYZ file specifying grid points for potential evaluation
output			File to write grid-point potential to
esf			Control options for field calculation
calcESF			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
grid			XYZ file specifying grid points for field evaluation
output			File to write grid-point field to
alphamol			Control options for molecular-polarizability calculation
calcAlpha			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
output			File to write polarizability tensor in global frame and in diagonal form to
convparam			Convergence parameters for self-consistent field calculation
wSOR_N			Mixing factor for successive overrelaxation of neutral system, usually between 0.3 and 0.5
wSOR_C			Mixing factor for successive overrelaxation of charged system, usually between 0.3 and 0.5
tolerance			Convergence criterion, fulfilled if relative change smaller than tolerance
maxiter			Maximum number of iterations in the convergence loop

1304 Return to the description of `emultipole`.

4.2.11   eoutersphere
1305

calc:eoutersphere
1306 Evaluates outersphere reorganization energy

option	default	unit	description
multipoles			XML allocation polar sites
method			Type of the method: **constant** - all pairs have value **lambda**. **spheres** - molecules are treated as spheres with radii **radius** and Pekar factor **pekar**. **dielectric** - with Pekar factor **pekar** and partial charges from resulting dielectric fields
lambdaconst		eV	The value for all pairs in the **constant** method
pekar			Pekar factor used for methods **spheres** and **dielectric**
segment			
type			
radius			
segment			
type			
radius			
cutoff		nm	Cutoff radius in between pair and the exterior molecule. Can be used in **spheres** and **dielectric**

1307 Return to the description of `eoutersphere`.

1308 ## 4.2.12 ewdbgpol

calc:ewdbgpol

1309 Calculates background polarisation needed for ewald calc

option	default	unit	description
multipoles			
control			
mps_table			
pdb_check			
coulombmethod			
method			
cutoff			
shape			
polarmethod			
method			
induce			
cutoff			
convergence			
energy			
kfactor			
rfactor			

1310 Return to the description of `ewdbgpol`.

1311 ## 4.2.13 ianalyze

calc:ianalyze

1312 Evaluates a histogram of a logarithm of squared couplings

option	default	unit	description
resolution_logJ2			Bin size of histogram log(J2)
states			States for which to calculate the histogram. Example: 1 -1

1313 Return to the description of `ianalyze`.

1314 ## 4.2.14 iimport

calc:iimport

1315 Imports electronic couplings from xml of ctp-dipro using folders of pairdump

option	default	unit	description
idft_jobs_file			idft jobs file

1316 Return to the description of `iimport`.

1317 ## 4.2.15 izindo

calc:izindo

1318 Semiempirical electronic coupling elements for all neighbor list pairs

option	default	unit	description
orbitalsXML			File with paths to .orb files

1319 Return to the description of `izindo`.

1320 ## 4.2.16 jobwriter

calc:jobwriter

1321 Writes list of jobs for a parallel execusion

option	default	unit	description

keys			job type
states	n e h		hole, electron, nuetral: mps file is required
single_id			Segment ID as argument for mps.single
kmc_cutoff		nm	Pair-interaction cut-off as argument for mps.kmc

1322   Return to the description of `jobwriter`.

4.2.17 neighborlist

calc:neighborlist

1324   Constructs a list of neighboring conjugated segments

option	default	unit	description
constant	0.5	nm	If provided, this value is used for all segment types
segments			A pair of segment types
type			Types of two segments. For types A and B this can be A A, A B or B B
cutoff		nm	Cutoff radius for centers of mass of rigid fragments

1325   Return to the description of `neighborlist`.

4.2.18 pairdump

calc:pairdump

1327   Coordinates of molecules and pairs from the neighbor list

option	default	unit	description
molecules			If **true** outputs single molecules, otherwise only pairs

1328   Return to the description of `pairdump`.

4.2.19 profile

calc:profile

1330   Density and site energy profiles

option	default	unit	description
axis			Axis along which to calculate density and energy profiles
direction	0 0 1		Axis direction
min		nm	Minimal projected position for manual binning
max		nm	Maximal projected position for manual binning
bin	0.1	nm	Spatial resolution of the profile
auto	1		'0' for manual binning using min and max, '1' for automated
particles			
type	segments		What centers of mass to use: 'segments' or 'atoms'
first	1		ID of the first segment
last	-1		ID of the last segment, -1 is the list end
output			
density	density.dat		Density profile file
energy	energy.dat		Energy profile file

1331   Return to the description of `profile`.

4.2.20 rates

calc:rates

1333   Hopping rates using classical or semi-classical expression

option	default	unit	description

field			Field in x y z direction
temperature		K	Temperature for rates
method			Method chosen to compute rates. Can either be **marcus** or **jortner**. The first is the high temperature limit of Marcus theory, the second is the rate proposed by Jortner and Bixon
nmaxvib	20		If the method of choice is **jortner**, the maximal number of excited vibrations on the molecules has to be specified as an integer for the summation
omegavib	0.2	eV	If the method of choice is **jortner**, the vibration frequency of the quantum mode has to be given in units of eV. The default value is close to the CC bond-stretch at 0.2eV

1334 Return to the description of `rates`.

4.2.21 sandbox

1336 Sandbox to test ctp classes

option	default	unit	description
ID			Not in use

1337 Return to the description of `sandbox`.

4.2.22 stateserver

1339 Export SQLite file to human readable format

option	default	unit	description
out			Output file name
pdb			PDB coordinate file name
keys			Sections to write to readable format (topology, segments, pairs, coordinates)

1340 Return to the description of `stateserver`.

4.2.23 tdump

1342 Coarse-grained and back-mapped (using rigid fragments) trajectories

option	default	unit	description
md	MD.pdb		Name of the coarse-grained trajectory
qm	QM.pdb		Name of the trajectory with back-substituted rigid fragments
frames	1		Number of frames to output

1343 Return to the description of `tdump`.

4.2.24 vaverage

1345 Computes site-centered velocity averages from site occupancies

option	default	unit	description
carriers			Carrier types for which to compute velocity averages
tabulate			Tabulate 'atoms' or 'segments'

1346   Return to the description of `vaverage`.

4.2.25   zmultipole

1348   Evaluates polarization contribution based on the Thole model

option	default	unit	description
multipoles			Polar Site Definitions in GDMA punch-file format
control			Control options for induction computation
induce	1		Enter '1' / '0' to toggle induction on / off
first			First segment for which to compute site energies
last			Last segment for which to compute site energies
output			File to write site energies to. Site energies are also stored in the state file
check			Check mapping of polar sites to fragment
tholeparam			Thole parameters required for charge-smearing
cutoff		nm	Cut-off beyond which all interactions are neglected
cutoff2		nm	Cut-off beyond which polarization is neglected
expdamp			Damping exponent used in exponential damping function
scaling			1-n interaction scaling, currently not in use
esp			Control options for potential calculation
calcESP			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
cube			
grid			XYZ file specifying grid points for potential evaluation
output			File to write grid-point potential to
esf			Control options for field calculation
calcESF			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
grid			XYZ file specifying grid points for field evaluation
output			File to write grid-point field to
alphamol			Control options for molecular-polarizability calculation
calcAlpha			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
output			File to write polarizability tensor in global frame and in diagonal form to
convparam			Convergence parameters for self-consistent field calculation
wSOR_N			Mixing factor for successive overrelaxation of neutral system, usually between 0.3 and 0.5
wSOR_C			Mixing factor for successive overrelaxation of charged system, usually between 0.3 and 0.5
tolerance			Convergence criterion, fulfilled if relative change smaller than tolerance
maxiter			Maximum number of iterations in the convergence loop

1349   Return to the description of `zmultipole`.

4.2.26   edft

1351   A wrapper for first principles based single site calculations

option	default	unit	description
job			Job options
tasks	input,run,parse		What to run
store	orbitals		What to store

1352 Return to the description of `edft`.

4.2.27 idft

1353

calc:idft

1354 Projection method for electronic couplings. Requires edft otput

option	default	unit	description
tasks	input,run,parse,		What to do
store	orbitals,overlap		What to store
degeneracy	0	eV	Criterium for the degeneracy of two levels
levels	3		Output between HOMO, ..., HOMO-levels; LUMO, ..., LUMO+levels
trim	2		Use trim*occupied of virtual orbitals

1355 Return to the description of `idft`.

4.2.28 pewald3d

1356

calc:pewald3d

1357 Evaluates site energies in a periodic setting

option	default	unit	description
jobcontrol			
job_file			
multipoles			
mapping			
mps_table			
pdb_check			
coulombmethod			
method			
cutoff			
shape			
polarmethod			
method			
induce			
cutoff			
tasks			
calculate_fields			
polarize_fg			
evaluate_energy			
coarsegrain			
cg_background			
cg_foreground			
cg_radius			
cg_anisotropic			
convergence			
energy			
kfactor			
rfactor			

1358 Return to the description of `pewald3d`.

4.2.29 qmmm

1359

calc:qmmm

1360 QM/MM with the Thole MM model

option	default	unit	description

control				
pdb_check			PDB file of polar sites	
write_chk	dipoles.xyz		XYZ file with dipoles split onto point charges	
format_chk	xyz		format, gaussian or xyz	
split_dpl	1		'0' do not split dipoles onto point charges, '1' do split	
dpl_spacing	1e-3	nm	Spacing to be used when splitting dipole onto point charges: d = q * a	
qmpackage				
package			QM package to use for the QM region	
gwbse			Specify if GW/BSE excited state calculation ist needed	
gwbse_options			GW/BSE options file	
state			Number of excited state, which is to be calculated	
type			Character of the excited state to be calculated	
filter			Filter with which to find the excited state after each calculation	
oscillator_strength			Oscillator strength filter, only states with higher oscillator strength are considered	
charge_transfer			Charge transfer filter , only states with charge transfer above threshold are consdered	
qmmmconvg			convergence criteria for the QM/MM	
dR	0.001	nm	RMS of coordinates	
dQ	0.001	e	RMS of charges	
dE_QM	0.0001	eV	Energy change of the QM region	
dE_MM	0.0001	eV	Energy change of the MM region	
max_iter	10		Number of iterations	
coulombmethod			Options for the MM embedding	
method	cut-off		Method for evaluation of electrostatics	
cutoff1			Cut-off for the polarizable MM1 shell	
cutoff2			Cut-off for the static MM2 shell	
tholemodel			Parameters for teh Thole model	
induce			'1' - induce '0' - no induction	
induce_intra_pair			'1' - include mutual interaction of induced dipoles in the QM region. '0' - do not	
exp_damp	0.39		Sharpness parameter	
scaling			Bond scaling factors	
convergence			Convergence parameters for the MM1 (polarizable) region	
wSOR_N			Mixing factor for the succesive overrelaxation algorithm for a neutral QM region	
wSOR_C			Mixing factor for the succesive overrelaxation algorithm for a charged QM region	
max_iter	512		Maximal number of iterations to converge induced dipoles	
tolerance			Maximum RMS change allowed in induced dipoles	

1361 Return to the description of qmmm.

4.2.30   xqmultipole

1363 Electrostatic interaction and induction energy of charged molecular clusters

option	default	unit	description
multipoles			Polar-site mapping definition
control			
job_file			Job file
emp_file			Polar-background definition, allocation of mps-files to segments

pdb_check			Whether or not to output a pdb-file of the mapped polar sites
format_chk			Format for check-file: 'xyz' or 'gaussian'
split_dpl			Split dipoles onto point charges in check-file
dpl_spacing		nm	Spacing between point charges for check-file output
coulombmethod			
method			Currently only cut-off supported
cutoff1		nm	Full-interaction radius cut-off
cutoff2		nm	Radius of electrostatic buffer
tholemodel			
induce			Induce - or not
induce_intra_pair			Induce mutually within the charged cluster
exp_damp			Thole sharpness parameter
scaling			Bond scaling parameters, currently not used
convergence			
wSOR_N			SOR mixing factor for overall neutral clusters
wSOR_C			SOR mixing factor for overall charged clusters
max_iter			Maximum number of iterations
tolerance			Relative tolerance as convergence criterion

1364   Return to the description of `xqmultipole`.

4.2.31   energy2xml

1365

calc:energy2xml

1366   Write out energies from SQL file

option	default	unit	description

1367   Return to the description of `energy2xml`.

4.2.32   integrals2xml

1368

calc:integrals2xml

1369   Write out transfer integrals from SQL file

option	default	unit	description

1370   Return to the description of `integrals2xml`.

4.2.33   occupations2xml

1371

calc:occupations2xml

1372   Write out site occupation probabilities from SQL file

option	default	unit	description

1373   Return to the description of `occupations2xml`.

4.2.34   pairs2xml

1374

calc:pairs2xml

1375   Write out neighbourlist from SQL file

option	default	unit	description

1376   Return to the description of `pairs2xml`.

1377 **4.2.35 rates2xml**

calc:rates2xml

1378 Write out charge transfer rates from SQL file

option	default	unit	description

1379 Return to the description of rates2xml.

1380 **4.2.36 segments2xml**

calc:segments2xml

1381 Write out segment data from SQL file

option	default	unit	description

1382 Return to the description of segments2xml.

1383 **4.2.37 trajectory2pdb**

calc:trajectory2pdb

1384 Generate PDB files for the mapped MD/QM topology

option	default	unit	description

1385 Return to the description of trajectory2pdb.

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