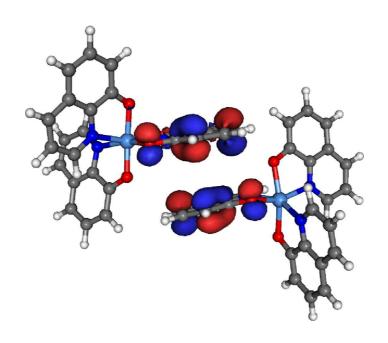
# VOTCA-CTP

## CHARGE TRANSPORT SIMULATIONS

USER MANUAL



compiled from: 1.5-dev (gitid: 6ca5dd5)

December 21, 2018 www.votca.org

#### Disclamer

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#### **Citations**

Development of this software depends on academic research grants. If you are using the package, please cite the following papers

[1] Long-range embedding of molecular ions and excitations in a polarizable molecular environment, Carl Poelking and Denis Andrienko

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J. Chem. Theory Comp. 12, 4516, 2016
```

[2] Modeling of spatially correlated energetic disorder in organic semiconductors, Pascal Kordt, Denis Andrienko

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J. Chem. Theory Comput., 12, 36, 2016
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[3] Microscopic simulations of charge transport in disordered organic semiconductors, Victor Rühle, Alexander Lukyanov, Falk May, Manuel Schrader, Thorsten Vehoff, James Kirkpatrick, Björn Baumeier and Denis Andrienko

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J. Chem. Theor. Comp. 7, 3335, 2011
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[4] Extracting nondispersive charge carrier mobilities of organic semiconductors from simulations of small systems, A. Lukyanov, D. Andrienko

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Phys. Rev. B, 82, 193202, 2010
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[5] Density-functional based determination of intermolecular charge transfer properties for large-scale morphologies, Björn Baumeier, James Kirkpatrick, and Denis Andrienko

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Phys. Chem. Chem. Phys. 12, 11103, 2010
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[6] Versatile Object-oriented Toolkit for Coarse-graining Applications, Victor Rühle, Christoph Junghans, Alexander Lukyanov, Kurt Kremer and Denis Andrienko

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J. Chem. Theor. Comp. 5, 3211, 2009
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[?] An approximate method for calculating transfer integrals based on the ZINDO Hamiltonian, James Kirkpatrick,

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Int. J. Quantum Chem. 108, 51, 2007
```

## Copyright

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## Chapter 1

## Introduction

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Charge carrier dynamics in an organic semiconductor can often be described in terms of charge hopping between localized states. The hopping rates depend on electronic coupling elements, reorganization energies, and site energies, which vary as a function of position and orientation of the molecules. The purpose of the VOTCA-CTP package [3] is to simplify the workflow for charge transport simulations, provide a uniform error-control for the methods, flexible platform for their development, and eventually allow *in silico* prescreening of organic semiconductors for specific applications.

The toolkit is implemented using modular concepts introduced earlier in the Versatile Objectoriented Toolkit for Coarse-graining Applications (VOTCA) [6]. It contains different programs, which execute specific tasks implemented in calculators representing an individual step in the workflow. Figure 1.1 summarizes a typical chain of commands to perform a charge transport simulation: First, the VOTCA code structures are adapted to reading atomistic trajectories, mapping them onto conjugated segments and rigid fragments, and substituting (if needed) rigid fragments with the optimized copies (ctp\_map). The programs ctp\_run and ctp\_parallel (for heavy-duty tasks) are then used to calculate all bimolecular charge hopping rates (via precalculation of all required ingredients). Site energies (or energetic disorder) can be determined as a combination of internal (ionization potentials/electron affinities of single molecules) as well as electrostatic and polarization contributions within the molecular environment. The calculation of electronic coupling elements between conjugated segments from the corresponding molecular orbitals can be performed using a dimer-projection technique based on density-functional theory (DFT). This requires explicit calculations using quantum-chemistry software for which we provide interfaces to Gaussian, Turbomole, and NWChem. Alternatively, the molecular orbital overlap module calculates electronic coupling elements relying on the semi-empirical INDO Hamiltonian and molecular orbitals in the format provided by the Gaussian package.

The kinetic Monte Carlo module reads in the neighbor list, site coordinates, and hopping rates and performs charge dynamics simulations using either periodic boundary conditions or charge sources and sinks.

The toolkit is written as a combination of modular C++ code and scripts. The data transfer between programs is implemented via a state file (sql database), which is also used to restart simulations. Analysis functions and most of the calculation routines are encapsulated by using the
observer pattern [7] which allows the implementation of new functions as individual modules.

In the following chapter 2, we summarize the theoretical background of the workflow of charge
transport simulations and in particular its individual steps. Chapter 3 describes the structure and
content of input and output files, while a full reference of programs and calculators is available
in chapter 4. For a hands-on tutorial, the reader is referred to the VOTCA-CTP project page at
http://code.google.com/p/votca-ctp/.

Input files:

#### conf.grc GROMACS trajectory topol.tpr GROMACS topology Mapping, generation of the sql database map.xml Converts and partitions atomistic GROMACS trajectory mapping and energies ctp\_map -t topol.tpr -c traj.xtc -s map.xml -f state.sql Useful tools: options for calculators toolpdb2map **Output files:** state.sql sqlite3 database file for Neighbor list data transfer between Indentifies close molecular pairs between which charge transfer rates will be calculated ctp\_run -o options.xml -f state.sql -e neighborlist Calculates electrostatic and polarization contribution to site energies ctp\_run -o options.xml -f state.sql -e emultipole Internal site and reorganization energies Imports internal site energy (IP, EA) and reorganization energies for charging and discharging to state.sql ctp\_run -o options.xml -f state.sql -e einternal ZINDO Transfer integrals Monomers with DFT Transfer integrals with ZINDO Calculate the relevant transport orbitals of monomers Calculate electronic coupling elements for all pairs in ctp\_parallel -o options.xml -f state.sql -e edft -j "write run" the neighbor list ctp\_run -o options.xml -f state.sql -e Transfer integrals with DFT One can choose between quantum-chemical Calculate electronic coupling elements for all pairs in (computationally expensive) or semi-empirical the neighbor list (fast, but not always sufficiently accurate) evaluation of transfer integrals. ctp\_parallel -o options.xml -f state.sql -e idft -j "write run read" Outersphere reorganization energies Contribution to reorganization of surrounding molecules due to polarization. (optional for Marcus rates) ctp\_run -o options.xml -f state.sql -e outersphere Charge transfer rates Calculates rates for charge transfer among all pairs in the neighborlist ctp\_run -o options.xml -f state.sql -e rates Charge dynamics via kMC Hopping of charge carriers simulated via kinetic Monte Carlo kmc\_run -o options.xml -f state.sql -e kmcmultiple

Figure 1.1: A practical workflow of charge transport simulations using VOTCA-CTP. The theoretical background of the individual steps is given in chapter 2. Chapter 3 describes the content of input and output files, while a full reference of programs and calculators is available in chapter 4. figsummary

Get help and list of options for a calculator: ctp\_run/ctp\_parallel/kmc\_run -d neighborlist

Get list of available calculators: ctp\_run/ctp\_parallel/kmc\_run -1

## Chapter 2

# . Theoretical background

#### sec:theor

#### 2.1 Workflow

#### sec:wokf

A typical workflow of charge transport simulations is depicted in figure 2.1. The first step is the simulation of an atomistic morphology, which is then partitioned on hopping sites. The coordinates of the hopping sites are used to construct a list of pairs of molecules, or neighbor list.

fig:workflow



Figure 2.1: Workflow for microscopic simulations of charge transport.

- 45 For each pair an electronic coupling element, a reorganization energy, a driving force, and even-
- tually the hopping rate are evaluated. The neighbor list and hopping rates define a directed
- 47 graph. The corresponding master equation is solved using the kinetic Monte Carlo method,
- which allows to explicitly monitor the charge dynamics in the system as well as to calculate time
- 49 or ensemble averages of occupation probabilities, charge fluxes, correlation functions, and field-
- 50 dependent mobilities.

## 2.2 Material morphology

#### sec:morpholo

- There is no generic recipe on how to predict a large-scale atomistically-resolved morphology of
  - 5

an organic semiconductor. The required methods are system-specific: for ultra-pure crystals, for



Figure 2.2: The concept of conjugated segments and rigid fragments. Dashed lines indicate conjugated segments while colors denote rigid fragments. (a) Hexabenzocoronene: the  $\pi$ -conjugated system is both a rigid fragment and a conjugated segment. (b)  $\mathrm{Alq}_3$ : the Al atom and each ligand are rigid fragments while the whole molecule is a conjugated segment. (c) Polythiophene: each repeat unit is a rigid fragment. A conjugated segment consists of one or more rigid fragments. One molecule can have several conjugated segments.

fig:segment

example, density-functional methods can be used provided the crystal structure is known from experiment. For partially disordered organic semiconductors, however, system sizes much larger than a unit cell are required. Classical molecular dynamics or Monte Carlo techniques are then the methods of choice.

In molecular dynamics, atoms are represented by point masses which interact via empirical potentials prescribed by a force-field. Force-fields are parametrized for a limited set of compounds and their refinement is often required for new molecules. In particular, special attention shall be paid to torsion potentials between successive repeat units of conjugated polymers or between functional groups and the  $\pi$ -conjugated system. First-principles methods can be used to characterize the missing terms of the potential energy function.

Self-assembling materials, such as soluble oligomers, discotic liquid crystals, block copolymers, partially crystalline polymers, etc., are the most complicated to study. The morphology of such systems often has several characteristic length scales and can be kinetically arrested in a thermodynamically non-equilibrium state. For such systems, the time- and length-scales of atomistic simulations might be insufficient to equilibrate or sample desired morphologies. In this case, systematic coarse-graining can be used to enhance sampling [6]. Note that the coarse-grained representation must reflect the structure of the atomistic system and allow for back-mapping to the atomistic resolution.

Here we assume that the morphology is already known, that is we know how the topology and the coordinates of all atoms in the systems at a given time. VOTCA-CTP can read standard GROMACS topology files. Custom definitions of atomistic topology via XML files are also possible. Since the description of the atomistic topology is the first step in the charge transport simulations, it is important to follow simple conventions on how the system is partitioned on molecules, residues, and how atoms are named in the topology. Required input files are described in section atomistic topology.

## 2.3 Conjugated segments and rigid fragments

With the morphology at hand, the next step is partitioning the system on hopping sites, or conjugated segments, and calculating charge transfer rates between them. Physically intuitive arguments can be used for the partitioning, which reflects the localization of the wave function of a charge. For most organic semiconductors, the molecular architecture includes relatively rigid, planar  $\pi$ -conjugated systems, which we will refer to as rigid fragments. A conjugated segment can contain one or more of such rigid fragments, which are linked by bonded degrees of freedom.

The dynamics of these degrees of freedom evolves on timescales much slower than the frequency of the internal promoting mode. In some cases, e.g. glasses, it can be 'frozen' due to non-bonded interactions with the surrounding molecules.

To illustrate the concept of conjugated segments and rigid fragments, three representative molec-89 ular architectures are shown in figure 2.2. The first one is a typical discotic liquid crystal, hex-90 abenzocoronene. It consists of a conjugated core to which side chains are attached to aid selfassembly and solution processing. In this case the orbitals localized on side chains do not partic-92 ipate in charge transport and the conjugated  $\pi$ -system is both, a rigid fragment and a conjugated 93 segment. In Alq<sub>3</sub>, a metal-coordinated compound, a charge carrier is delocalized over all three ligands. Hence, the whole molecule is one conjugated segment. Individual ligands are relatively rigid, while energies of the order of  $k_BT$  are sufficient to reorient them with respect to each other. 96 Thus the Al atom and the three ligands are rigid fragments. In the case of a conjugated polymer, 97 one molecule can consist of several conjugated segments, while each backbone repeat unit is a rigid fragment. Since the conjugation along the backbone can be broken due to large out-of-plane twists between two repeat units, an empirical criterion, based on the dihedral angle, can be used 100 to partition the backbone on conjugated segments [8]. However, such intuitive partitioning is, to 101 some extent, arbitrary and shall be validated by other methods [9–11].

After partitioning, an additional step is often required to remove bond length fluctuations introduced by molecular dynamics simulations, since they are already integrated out in the derivation of the rate expression. This is achieved by substituting respective molecular fragments with rigid, planar  $\pi$ -systems optimized using first-principles methods. Centers of mass and gyration tensors are used to align rigid fragments, though a custom definition of local axes is also possible. Such a procedure also minimizes discrepancies between the force-field and first-principles-based ground state geometries of conjugated segments, which might be important for calculations of electronic couplings, reorganization energies, and intramolecular driving forces.

To partition the system on hopping sites and substitute rigid fragments with the corresponding ground-state geometries ctp\_map program is used:

```
Mapping the GROMACS trajectory
```

ctp\_map -ttopol.tpr -ctraj.xtc -smap.xml -fstate.sql

It reads in the GROMACS topology (topol.tpr) and trajectory (traj.xtc) files, definitions of conjugated segments and rigid fragments (map.xml) and outputs coordinates of conjugated segments (hopping sites) and rigid fragments (as provided in the MD trajectory and after rigidification) to the state file (state.sql). In order to do this, a mapping file map.xml has to be provided, which specifies the corresponding atoms in the different representations. After this step, all information (frame number, dimensions of the simulation box, etc) are stored in the state file and only this file is used for further calculations.



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#### Attention

VOTCA-CTP requires a wrapped trajectory for mapping the segments and fragments, so all molecules should be whole in the frame.

In order to visually check the mapping one can use either the tdump calculator or the programm ctp\_dump with the calculator trajectory2pdb.



#### Writing a mapped trajectroy with ctp\_dump

ctp\_dump -fstate.sql -etrajectory2pdb

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It reads in the state file created by ctp\_map and outputs two trajectory files corresponding to the original and rigidified atom coordinates. To check the mapping, it is useful to superimpose the three outputs (original atomistic, atomistic stored in the state file, and rigidified according to ground state geometries), e.g., with VMD.

number Writing a mapped trajectroy with tdump

ctp\_run -f state.sql -o options.xml -e tdump

It also reads in the state file but appends the coordinates to a pdb. file. So make sure to delete old QM.pdb and MD.pdb if you want to create a new imagef

### 2.4 Neighbor list

A list of neigboring conjugated segments, or neighbor list, contains all pairs of conjugated segments for which coupling elements, reorganization energies, site energy differences, and rates are evaluated.

Two segments are added to this list if the distance between centers of mass of any of their rigid fragments is below a certain cutoff. This allows neighbors to be selected on a criterion of minimum distance of approach rather than center of mass distance, which is useful for molecules with anisotropic shapes.

The neighbor list can be generated from the atomistic trajectory by using the neighborlist calculator. This calculator requires a cutoff, which can be specified in the options.xml file. The list is saved to the state.sql file:

**Generating a neighbor list** 

ctp\_run -ooptions.xml -fstate.sql -eneighborlist

## 2.5 Reorganization energy

The reorganization energy  $\lambda_{ij}$  takes into account the change in nuclear (and dielectric) degrees of freedom as the charge moves from donor i to acceptor j. It has two contributions: intramolecular,  $\lambda_{ij}^{\text{int}}$ , which is due to reorganization of nuclear coordinates of the two molecules forming the charge transfer complex, and intermolecular (outersphere),  $\lambda_{ij}^{\text{out}}$ , which is due to the relaxation of the nuclear coordinates of the environment. In what follows we discuss how these contributions can be calculated.

#### 2.5.1 Intramolecular reorganization energy

If intramolecular vibrational modes of the two molecules are treated classically, the rearrangement of their nuclear coordinates after charge transfer results in the dissipation of the internal reorganization energy,  $\lambda_{ij}^{int}$ . It can be computed from four points on the potential energy surfaces (PES) of both molecules in neutral and charged states, as indicated in figure 2.3.

Adding the contributions due to discharging of molecule i and charging of molecule j yields [12]

$$\lambda_{ij}^{\rm int} = \lambda_i^{cn} + \lambda_j^{nc} = U_i^{nC} - U_i^{nN} + U_j^{cN} - U_j^{cC} \,. \tag{2.1} \quad \text{equ:lambdas}$$

Here  $U_i^{nC}$  is the internal energy of the neutral molecule i in the geometry of its charged state (small n denotes the state and capital C the geometry). Similarly,  $U_j^{cN}$  is the energy of the charged molecule j in the geometry of its neutral state. Note that the PES of the donor and acceptor are not identical for chemically different compounds or for conformers of the same molecule. In this



Figure 2.3: Potential energy surfaces of (a) donor and (b) acceptor in charged and neutral states. After the change of the charge state both molecules relax their nuclear coordinates. If all vibrational modes are treated classically, the total internal reorganization energy and the internal energy difference of the electron transfer reaction are  $\lambda_{ij}^{\text{int}} = \lambda_i^{cn} + \lambda_j^{nc}$  and  $\Delta E_{ij}^{\text{int}} = \Delta U_i - \Delta U_j$ , respectively.

fig:parabolas

case  $\lambda_i^{cn} \neq \lambda_j^{cn}$  and  $\lambda_i^{nc} \neq \lambda_j^{nc}$ . Thus  $\lambda_{ij}^{int}$  is a property of the charge transfer complex, and not of a single molecule.

Intramolecular reorganization energies for discharging ( $\lambda^{cn}$ ) and charging ( $\lambda^{nc}$ ) of a molecule 158 need to be determined using quantum-chemistry and given in map.xml. The values are written 159 to the state.sql using the calculator einternal (see also internal energy): 160

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Intramolecularl reorganization energies

| ctp\_run -ooptions.xml -fstate.sql -eeinternal

### Outersphere reorganization energy

During the charge transfer reaction, also the molecules outside the charge transfer complex reorient and polarize in order to adjust for changes in electric potential, resulting in the outersphere contribution to the reorganization energy.  $\lambda_{ij}^{\text{out}}$  is particularly important if charge transfer occurs in a polarizable environment. Assuming that charge transfer is much slower than electronic polarization but much faster than nuclear rearrangement of the environment,  $\lambda_{ii}^{out}$  can be calculated from the electric displacement fields created by the charge transfer complex [?]

$$\lambda_{ij}^{\text{out}} = \frac{c_p}{2\epsilon_0} \int_{V^{\text{out}}} dV \left[ \vec{D}_I(\vec{r}) - \vec{D}_F(\vec{r}) \right]^2 , \qquad (2.2) \text{ equilambda\_outer1}$$

where  $\epsilon_0$  is the the permittivity of free space,  $\vec{D}_{I,F}(\vec{r})$  are the electric displacement fields created by the charge transfer complex in the initial (charge on molecule i) and final (charge transferred to molecule j) states,  $V^{\text{out}}$  is the volume outside the complex, and  $c_p = \frac{1}{\epsilon_{\text{opt}}} - \frac{1}{\epsilon_{\text{s}}}$  is the Pekar factor, which is determined by the low ( $\epsilon_s$ ) and high ( $\epsilon_{opt}$ ) frequency dielectric permittivities. Eq. (2.2) can be simplified by assuming spherically symmetric charge distributions on molecules

172 i and j with total charge e. Integration over the volume  $V^{\text{out}}$  outside of the two spheres of radii  $R_i$ 173 and  $R_j$  centered on molecules i and j leads to the classical Marcus expression for the outersphere 174 reorganization energy 175

$$\lambda_{ij}^{\rm out} = \frac{c_p e^2}{4\pi\epsilon_0} \left( \frac{1}{2R_i} + \frac{1}{2R_j} - \frac{1}{r_{ij}} \right) \,, \tag{2.3} \quad \text{equ:lambda\_outer2}$$

where  $r_{ij}$  is the molecular separation. While eq. (2.3) captures the main physics, e.g. predicts smaller outer-sphere reorganization energies (higher rates) for molecules at smaller separations,

it often cannot provide quantitative estimates, since charge distributions are rarely spherically symmetric.

Alternatively, the displacement fields can be constructed using the atomic partial charges. The difference of the displacement fields at the position of an atom  $b_k$  outside the charge transfer complex (molecule  $k \neq i, j$ ) can be expressed as

$$\vec{D}_{I}(\vec{r}_{b_{k}}) - \vec{D}_{F}(\vec{r}_{b_{k}}) = \sum_{a_{i}} \frac{q_{a_{i}}^{c} - q_{a_{i}}^{n}}{4\pi} \frac{(\vec{r}_{b_{k}} - \vec{r}_{a_{i}})}{|\vec{r}_{b_{k}} - \vec{r}_{a_{i}}|^{3}} + \sum_{a_{j}} \frac{q_{a_{j}}^{n} - q_{a_{j}}^{c}}{4\pi} \frac{(\vec{r}_{b_{k}} - \vec{r}_{a_{j}})}{|\vec{r}_{b_{k}} - \vec{r}_{a_{j}}|^{3}},$$
(2.4)

where  $q_{a_i}^n$  ( $q_{a_i}^c$ ) is the partial charge of atom a of the neutral (charged) molecule i in vacuum. The partial charges of neutral and charged molecules are obtained by fitting their values to reproduce the electrostatic potential of a single molecule (charged or neutral) in vacuum. Assuming a uniform density of atoms, the integration in eq. (2.2) can be rewritten as a density-weighted sum over all atoms excluding those of the charge transfer complex.

The remaining unknown needed to calculate  $\lambda_{ij}^{\mathrm{out}}$  is the Pekar factor,  $c_p$ . In polar solvents  $\epsilon_{\mathrm{s}}\gg \epsilon_{\mathrm{opt}}\sim 1$  and  $c_p$  is of the order of 1. In most organic semiconductors, however, molecular orientations are fixed and therefore the low frequency dielectric permittivity is of the same order of magnitude as  $\epsilon_{\mathrm{opt}}$ . Hence,  $c_p$  is small and its value is very sensitive to differences in the permittivities.

Outersphere reorganization energies for all pairs of molecules in the neighbor list can be computed from the atomistic trajectory by using the eoutersphere calculator.

Two methods can be used to compute  $\lambda_{ij}^{\text{out}}$ . The first method uses the atomistic partial charges of neutral and charged molecules from files specified in map.xml and eq. (2.2). The Pekar factor  $c_p$  and a cutoff radius based on molecular centers of mass have to be specified in the options.xml file.

If this method is computationally prohibitive,  $\lambda_{ij}^{\text{out}}$  can be computed using eq. (2.3), which assumes spherical charge distributions on the molecules. In this case the radii of these spheres are specified in segments.xml, while the Pekar factor  $c_p$  is given in the options.xml file and no cutoff radius is needed.

The outer sphere reorganization energies are saved to the state.sql file:

Outersphere reorganization energy

ctp\_run -o options.xml -f state.sql -e outersphere

## 2.6 Site energies

A charge transfer reaction between molecules i and j is driven by the site energy difference,  $\Delta E_{ij} = E_i - E_j$ . Since the transfer rate,  $\omega_{ij}$ , depends exponentially on  $\Delta E_{ij}$  (see eq. (2.31)) it is important to compute its distribution as accurately as possible. The total site energy difference has contributions due to externally applied electric field, electrostatic interactions, polarization effects, and internal energy differences. In what follows we discuss how to estimate these contributions by making use of first-principles calculations and polarizable force-fields.

#### 2.6.1 Externally applied electric field

#### sec:ext\_fiel

The contribution to the total site energy difference due to an external electric field  $\vec{F}$  is given by  $\Delta E_{ij}^{\rm ext} = q \vec{F} \cdot \vec{r}_{ij}$ , where  $q = \pm e$  is the charge and  $\vec{r}_{ij} = \vec{r}_i - \vec{r}_j$  is a vector connecting molecules i and j. For typical distances between small molecules, which are of the order of 1 nm, and moderate fields of  $F < 10^8 \, {\rm V/m}$  this term is always smaller than  $0.1 \, {\rm eV}$ .

2.6. SITE ENERGIES 11

#### Internal energy 2.6.2

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The contribution to the site energy difference due to different internal energies (see figure 2.3) can be written as

$$\Delta E_{ij}^{\rm int} = \Delta U_i - \Delta U_j = \left( U_i^{cC} - U_i^{nN} \right) - \left( U_j^{cC} - U_j^{nN} \right) , \tag{2.5} \quad \text{equiconformational}$$

where  $U_i^{cC(nN)}$  is the total energy of molecule i in the charged (neutral) state and geometry.  $\Delta U_i$ corresponds to the adiabatic ionization potential (or electron affinity) of molecule i, as shown in figure 2.3. For one-component systems and negligible conformational changes  $\Delta E_{ii}^{\rm int}=0$ , while it is significant for donor-acceptor systems.

Internal energies determined using quantum-chemistry need to be specified in map.xml. The values are written to the state.sql using the calculator einternal (see also intramolecular reorganization energy):



Internal energies

| ctp\_run -ooptions.xml -fstate.sql -eeinternal

#### 2.6.3 **Electrostatic interaction energy**

sec:distributed multipole

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We represent the molecular charge density by choosing multiple expansion sites ("polar sites") per molecule in such a way as to accurately reproduce the molecular electrostatic potential (ESP), with a set of suitably chosen multipole moments  $\{Q_{lk}^a\}$  (in spherical-tensor notation) allocated to each site. The expression for the electrostatic interaction energy between two molecules A and Bin the multi-point expansion includes an implicit sum over expansion sites  $a\epsilon A$  and  $b\epsilon B$ ,

$$U_{AB} = \sum_{a \in A} \sum_{b \in B} \hat{Q}_{l_1 k_1}^a T_{l_1 k_1 l_2 k_2}^{a,b} \hat{Q}_{l_2 k_2}^b \equiv \hat{Q}_{l_1 k_1}^a T_{l_1 k_1 l_2 k_2}^{a,b} \hat{Q}_{l_2 k_2}^b, \tag{2.6}$$

where we have used the Einstein sum convention for the site indices a and b on the right-hand side of the equation, in addition to the sum convention that is in place for the multipole-moment components  $t \equiv l_1 k_1$  and  $u \equiv l_2 k_2$ . The  $T_{l_1 k_1 l_2 k_2}^{a,b}$  are tensors that mediate the interaction between a multipole component  $l_1k_1$  on site a with the moment  $l_2k_2$  on site b. If we include the molecular environment into a perturbative term W to enter in the single-molecule Hamiltonian, the above expression is exactly the first-order correction to the energy where the quantum-mechanical detail has been absorbed in classical multipole moments.

The are a number of strategies how to arrive at such a collection of distributed multipoles. They can be classified according to whether the multipoles are derived (a) from the electrostatic potential generated by the SCF charge density or (b) from a decomposition of the wavefunction itself. Here, we will only draft two of those approaches, CHELPG [13] from category (a) and DMA [14] from category (b).

The CHELPG (CHarges from ELectrostatic Potentials, Grid-based) method relies on performing a least-squares fit of atom-placed charges to reproduce the electrostatic potential as evaluated from the SCF density on a regularly spaced grid [13]. The fitted charges result from minimizing the Lagrangian function [15]

$$z(\{q_i\}) = \sum_{k=1}^{M} \left( \phi(\vec{r}_k) - \sum_{i=1}^{N} \frac{1}{4\pi\varepsilon_0} \frac{q_i}{|\vec{r}_i - \vec{r}_k|} \right) + \lambda \left( q_{\text{mol}} - \sum_{i=1}^{N} q_i \right), \tag{2.7}$$

with M grid points, N atomic sites, the set of atomic partial charges  $\{q_i\}$  and the SCF potential  $\phi$ . The Lagrange multiplier  $\lambda$  constrains the sum of the fitted charges to the molecular charge  $q_{\text{mol}}$ . The main difference from other fitting schemes [16] is the algorithm that selects the positions at which the potential is evaluated (we note that the choice of grid points can have substantial effects especially for bulky molecules). Clearly, the CHELPG method can be (and has been) extended to include higher atomic multipoles. It should be noted, however, how already the inclusion of atomic dipoles hardly improves the parametrization, and can in fact be harmful to its conformational stability.

The Distributed-Multipole-Analysis (DMA) approach [14, 17], developed by A. Stone, operates directly on the quantum-mechanical density matrix, expanded in terms of atom- and bondcentered Gaussian functions  $\chi_{\alpha} = R_{LK}(\vec{x} - \vec{s}_{\alpha}) \exp[-\zeta(\vec{x} - \vec{s}_{\alpha})^2]$ ,

$$\rho(\vec{x}) = \sum_{\alpha,\beta} \rho_{\alpha\beta} \chi_{\alpha}(\vec{x} - \vec{s}_{\alpha}) \chi_{\beta}(\vec{x} - \vec{s}_{\beta}). \tag{2.8}$$

The aim is to compute multipole moments according in a distributed fashion: If we use that the overlap product  $\chi_{\alpha}\chi_{\beta}$  of two Gaussian basis functions yields itself a Gaussian centered at  $P = (\zeta_{\alpha}\vec{s}_{\alpha} + \zeta_{\beta}\vec{s}_{\beta})/(\zeta_{\alpha} + \zeta_{\beta})$ , it is possible to proceed in two steps: First, we compute the multipole moments associated with a specific summand in the density matrix, referred to the overlap center  $\vec{P}$ :

$$Q_{LK}[\vec{P}] = -\int R_{LK}(\vec{x} - \vec{P})\rho_{\alpha\beta}\chi_{\alpha}\chi_{\beta}d^3x.$$
 (2.9)

Second, we transfer the resulting  $Q_{lk}[\vec{P}]$  to the position  $\vec{S}$  of a polar site according to the rule [14]

$$Q_{nm}[\vec{S}] = \sum_{l=0}^{L} \sum_{k=-l}^{l} \left[ \binom{n+m}{l+k} \binom{n-m}{l-k} \right]^{1/2} R_{n-l,m-k}(\vec{S} - \vec{P}) \cdot Q_{lk}[\vec{P}].$$
 (2.10)

Note how this requires a rule for the choice of the expansion site to which the multipole moment should be transferred. In the near past [17], the nearest-site algorithm, which allocates the multipole moments to the site closest to the overlap center, was replaced for diffuse functions by an algorithm based on a smooth weighting function in conjunction with grid-based integration methods in order to decrease the basis-set dependence of the resulting set of distributed multipoles.

One important advantage of the DMA approach over fitting algorithms such as CHELPG or 253 Merz-Kollman (MK) is that higher-order moments can also be derived without too large an am-254 255

The 'mps' file format used by VOTCA for the definition of distributed multipoles (as well as point polarizabilities, see subsequent section) is based on the GDMA punch format of A. Stone's GDMA program [17] (the punch output file can be immediately plugged into VOTCA without any conversions to be applied). Furthermore the log-file of different QM packages (currently Gaussian, Turbomole and NWChem) may be fed into the log2mps tool, which will subsequently generate the appropriate mps-file.

# Read in ESP charges from a QM log file | ctp\_tools -ooptions.xml -elog2mps

#### Induction energy - the Thole model

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If we in addition to the permanent set of multipole moments  $\{Q_a^a\}$  allow for induced moments  $\{\Delta Q_t^a\}$  and penalize their generation with a bilinear form (giving rise to a strictly positive contribution to the energy),

$$U_{\text{int}} = \frac{1}{2} \sum_{A} \Delta Q_t^a \eta_{tt'}^{aa'} \Delta Q_{t'}^{a'}, \tag{2.11}$$

2.6. SITE ENERGIES

it can be shown that the induction contribution to the site energy evaluates to an expression where all interactions between induced moments have cancelled out, and interactions between permanent and induced moments are scaled down by 1/2 [18]:

$$U_{pu} = \frac{1}{2} \sum_{A} \sum_{B>A} \left[ \Delta Q_t^a T_{tu}^{ab} Q_u^b + \Delta Q_t^b T_{tu}^{ab} Q_u^a \right]. \tag{2.12}$$

This term can be viewed as the second-order (induction) correction to the molecular interaction energy. The sets of  $\{Q_t^a\}$  are solved for self-consistently via

$$\Delta Q_t^a = -\sum_{B \neq A} \alpha_{tt'}^{aa'} T_{t'u}^{a'b} (Q_u^b + \Delta Q_u^b), \tag{2.13}$$
 equiself\_consistent\_dQ

where the polarizability tensors  $\alpha_{tt'}^{aa'}$  are given by the inverse of  $\eta_{tt'}^{aa'}$ .

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With eqs. 2.13 and 2.12 we have at hand expressions that allow us to compute the induction energy contribution to site energies in an iterative manner based on a set of molecular distributed multipoles  $\{Q_t^a\}$  and polarizabilities  $\{\alpha_{tt'}^{aa'}\}$ . We have drafted in the previous section how to obtain the former from a wavefunction decomposition or fitting scheme (GDMA, CHELPG). The  $\{\alpha_{tt'}^{aa'}\}$  can be derived formally (or rather: read off) from a perturbative expansion of the molecular interaction. In this work we make use of the Thole model [19? ] as a semi-empirical approach to obtain the sought-after point polarizabilities in the local dipole approximation, that is,  $[\alpha_{tt'}^{aa'}] = \alpha_{tt'}^{aa'} \delta_{t\beta} \delta_{t'\beta} \delta_{aa'}$ , where  $\beta \epsilon \{x,y,z\}$  references the dipole-moment component.

The Thole model is based on a modified dipole-dipole interaction, which can be reformulated in terms of the interaction of smeared charge densities. This has been shown to be necessary due to the divergent head-to-tail dipole-dipole interaction that otherwise results at small interseparations on the Å scale [19??]. Smearing out the charge distribution mimics the nature of the QM wavefunction, which effectively guards against this unphysical polarization catastrophe. Since the point dipoles however only react individually to the external field, any correlation effects as were still accounted for in the  $\{\alpha_{tt'}^{aa'}\}$  are lost, except perhaps those correlations that are due to the mere classical field interaction.

The smearing of the nuclei-centered multipole moments is obtained via a fractional charge density  $\rho_f(\vec{u})$  which should be normalized to unity and fall off rapidly as of a certain radius  $\vec{u} = \vec{u}(\vec{R})$ . The latter is related to the physical distance vector  $\vec{R}$  connecting two interacting sites via a linear scaling factor that takes into account the magnitude of the isotropic site polarizabilities  $\alpha^a$ . This isotropic fractional charge density gives rise to a modified potential

$$\phi(u) = -\frac{1}{4\pi\varepsilon_0} \int\limits_0^u 4\pi u' \rho(u') du' \tag{2.14}$$

We can relate the multipole interaction tensor  $T_{ij...}$  (this time in Cartesian coordinates) to the fractional charge density in two steps: First, we rewrite the tensor in terms of the scaled distance vector  $\vec{u}$ ,

$$T_{ii...}(\vec{R}) = f(\alpha^a \alpha^b) t_{ii...}(\vec{u}(\vec{R}, \alpha^a \alpha^b)), \tag{2.15}$$

where the specific form of  $f(\alpha^a \alpha^b)$  results from the choice of  $u(\vec{R}, \alpha^a \alpha^b)$ . Second, we demand that the smeared interaction tensor  $t_{ij...}$  is given as usual by the appropriate derivative of the potential in eq. 2.14,

$$t_{ij\dots}(\vec{u}) = -\partial_{u_i}\partial_{u_i}\dots\phi(\vec{u}). \tag{2.16}$$

It turns out that for a suitable choice of  $\rho_f(\vec{u})$ , the modified interaction tensors can be rewritten in such a way that powers n of the distance  $R = |\vec{R}|$  are damped with a damping function  $\lambda_n(\vec{u}(\vec{R}))$  [20].

There is a large number of fractional charge densities  $\rho_f(\vec{u})$  that have been tested for the purpose of giving best results for the molecular polarizability as well as interaction energies. Note how a great advantage of the Thole model is the exceptional transferability of the atomic polarizabilities to compounds not used for the fitting procedure [?]. In fact, for most organic molecules, a fixed set of atomic polarizabilities ( $\alpha_C=1.334$ ,  $\alpha_H=0.496$ ,  $\alpha_N=1.073$ ,  $\alpha_O=0.873$ ,  $\alpha_S=2.926$  ų) based on atomic elements yields satisfactory results.

VOTCA implements the Thole model with an exponentially-decaying fractional charge density

$$\rho(u) = \frac{3a}{4\pi} \exp(-au^3), \tag{2.17}$$

where  $\vec{u}(\vec{R},\alpha^a\alpha^b)=\vec{R}/(\alpha^a\alpha^b)^{1/6}$  and the smearing exponent a=0.39 (which can however be changed from the program options), as used in the AMOEBA force field [20]. Even though the Thole model performs very well for many organic compounds with only the above small set of element-based polarizabilities, conjugated molecules may require a more intricate parametrization. The simplest approach is to resort to scaled polarizabilities to match the effective molecular polarizable volume  $V\sim\alpha_x\alpha_y\alpha_z$  as predicted by QM calculations (here  $\alpha_x,\alpha_y,\alpha_z$  are the eigenvalues of the molecular polarizability tensor). The molpol tool assists with this task, it self-consistently calculates the Thole polarizability for an input mps-file and optimizes (if desired) the atomic polarizabilities in the above simple manner.

# Generate Thole-type polarizabilites for a segment | ctp\_tools -ooptions.xml -e molpol

The electrostatic and induction contribution to the site energy is evaluated by the emultipole calculator. Atomistic partial charges for charged and neutral molecules are taken from mps-files (extended GDMA format) specified in map.xml. Note that, in order to speed up calculations for both methods, a cut-off radius (for the molecular centers of mass) can be given in options.xml.

Threaded execution is advised.

Electrostatic and induction corrections

ctp\_run -ooptions.xml -fstate.sql -eemultipole

Furthermore available are zmultipole, which extends emultipole to allow for an electrostatic buffer layer (loosely related to the z-buffer in OpenGL, hence the name) and anisotropic point polarizabilities. For the interaction energy of charged clusters of any user-defined composition (Frenkel states, CT states, ...), xqmultipole can be used.

Interaction energy of charged molecular clusters embedded in a molecular environment

ctp\_parallel -ooptions.xml -fstate.sql -exqmultipole

## 2.7 Transfer integrals

sec:transfer integrals

The electronic transfer integral element  $J_{ij}$  entering the Marcus rates in eq. (2.31) is defined as

$$J_{ij} = \left\langle \phi_i \left| \hat{H} \right| \phi_j \right\rangle,$$
 (2.18) equ:TI

where  $\phi_i$  and  $\phi_j$  are diabatic wavefunctions, localized on molecule i and j respectively, participating in the charge transfer, and  $\hat{H}$  is the Hamiltonian of the formed dimer. Within the frozencore approximation, the usual choice for the diabatic wavefunctions  $\phi_i$  is the highest occupied

molecular orbital (HOMO) in case of hole transport, and the lowest unoccupied molecular orbital (LUMO) in the case of electron transfer, while  $\hat{H}$  is an effective single particle Hamiltonian, e.g. Fock or Kohn-Sham operator of the dimer. As such,  $J_{ij}$  is a measure of the strength of the electronic coupling of the frontier orbitals of monomers mediated by the dimer interactions. Intrinsically, the transfer integral is very sensitive to the molecular arrangement, i.e. the distance and the mutual orientation of the molecules participating in charge transport. Since this arrangement can also be significantly influenced by static and/or dynamic disorder [21? ? -23], it is essential to calculate  $J_{ij}$  explicitly for each hopping pair within a realistic morphology. Considering that the number of dimers for which eq. (2.18) has to be evaluated is proportional to the number of molecules times their coordination number, computationally efficient and at the same time quantitatively reliable schemes are required.

#### 2.7.1 Projection of monomer orbitals on dimer orbitals (DIPRO)

An approach for the determination of the transfer integral that can be used for any single-particle electronic structure method (Hartree-Fock, DFT, or semiempirical methods) is based on the projection of monomer orbitals on a manifold of explicitly calculated dimer orbitals. This dimer projection (DIPRO) technique including an assessment of computational parameters such as the basis set, exchange-correlation functionals, and convergence criteria is presented in detail in ref. [5]. A brief summary of the concept is given below.

We start from an effective Hamiltonian <sup>1</sup>

$$\hat{H}^{\text{eff}} = \sum_{i} \epsilon_{i} \hat{a}_{i}^{\dagger} \hat{a}_{i} + \sum_{j \neq i} J_{ij} \hat{a}_{i}^{\dagger} \hat{a}_{j} + c.c. \tag{2.19}$$

where  $\hat{a}_i^\dagger$  and  $\hat{a}_i$  are the creation and annihilation operators for a charge carrier located at the molecular site i. The electron site energy is given by  $\epsilon_i$ , while  $J_{ij}$  is the transfer integral between two sites i and j. We label their frontier orbitals (HOMO for hole transfer, LUMO for electron transfer)  $\phi_i$  and  $\phi_j$ , respectively. Assuming that the frontier orbitals of a dimer (adiabatic energy surfaces) result exclusively from the interaction of the frontier orbitals of monomers, and consequently expand them in terms of  $\phi_i$  and  $\phi_j$ . The expansion coefficients,  $\bar{\mathbf{C}}$ , can be determined by solving the secular equation

$$(\mathbf{H} - E\mathbf{S})\bar{\mathbf{C}} = 0 \tag{2.20} \quad \text{equidipro_eq2}$$

where  $\mathbf{H}$  and  $\mathbf{S}$  are the Hamiltonian and overlap matrices of the system, respectively. These matrices can be written explicitly as

$$\mathbf{H} = \begin{pmatrix} e_i & H_{ij} \\ H_{ij}^* & e_j \end{pmatrix} \qquad \mathbf{S} = \begin{pmatrix} 1 & S_{ij} \\ S_{ij}^* & 1 \end{pmatrix} \tag{2.21}$$

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$$\begin{aligned} e_{i} &= \left\langle \phi_{i} \right| \hat{H} \left| \phi_{i} \right\rangle & H_{ij} &= \left\langle \phi_{i} \right| \hat{H} \left| \phi_{j} \right\rangle \\ e_{j} &= \left\langle \phi_{j} \right| \hat{H} \left| \phi_{j} \right\rangle & S_{ij} &= \left\langle \phi_{j} \right| \phi_{j} \right\rangle \end{aligned} \tag{2.22} \quad \text{equidipro_eq4}$$

The matrix elements  $e_{i(j)}$ ,  $H_{ij}$ , and  $S_{ij}$  entering eq. (2.21) can be calculated via projections on the dimer orbitals (eigenfunctions of  $\hat{H}$ )  $\left\{\left|\phi_{n}^{\mathrm{D}}\right\rangle\right\}$  by inserting  $\hat{1}=\sum_{n}\left|\phi_{n}^{\mathrm{D}}\right\rangle\left\langle\phi_{n}^{\mathrm{D}}\right|$  twice. We exemplify this explicitly for  $H_{ij}$  in the following

$$H_{ij} = \sum_{nm} \left\langle \phi_i \mid \phi_n^{\rm D} \right\rangle \left\langle \phi_n^{\rm D} \mid \hat{H} \mid \phi_m^{\rm D} \right\rangle \left\langle \phi_m^{\rm D} \mid \phi_j \right\rangle. \tag{2.23}$$

The Hamiltonian is diagonal in its eigenfunctions,  $\langle \phi_n^{\rm D} | \hat{H} | \phi_m^{\rm D} \rangle = E_n \delta_{nm}$ . Collecting the projections of the frontier orbitals  $|\phi_{i(j)}\rangle$  on the n-th dimer state  $(\bar{\mathbf{V}}_{(i)})_n = \langle \phi_i | \phi_n^{\rm D} \rangle$  and  $(\bar{\mathbf{V}}_{(j)})_n = \langle \phi_i | \phi_n^{\rm D} \rangle$  respectively, into vectors we obtain

$$H_{ij} = \bar{\mathbf{V}}_{(i)} \mathbf{E} \bar{\mathbf{V}}_{(j)}^{\dagger}.$$
 (2.24) eq:dipro\_eq17

<sup>&</sup>lt;sup>1</sup>we use following notations: a - number,  $\bar{\mathbf{a}}$  - vector,  $\mathbf{A}$  - matrix,  $\hat{A}$  - operator

What is left to do is determine these projections  $\bar{\mathbf{V}}_{(k)}$ . In all practical calculations the molecular orbitals are expanded in basis sets of either plane waves or of localized atomic orbitals  $|\varphi_{\alpha}\rangle$ . We will first consider the case that the calculations for the monomers are performed using a counterpoise basis set that is commonly used to deal with the basis set superposition error (BSSE). The basis set of atom-centered orbitals of a monomer is extended to the one of the dimer by adding the respective atomic orbitals at virtual coordinates of the second monomer. We can then write the respective expansions as

$$|\phi_k\rangle = \sum_{\alpha} \lambda_{\alpha}^{(k)} |\varphi_{\alpha}\rangle$$
 and  $|\phi_n^{\rm D}\rangle = \sum_{\alpha} D_{\alpha}^{(n)} |\varphi_{\alpha}\rangle$  (2.25) eq:dipro\_eq18

where k=i,j. The projections can then be determined within this common basis set as

where S is the overlap matrix of the atomic basis functions. This allows us to finally write the elements of the Hamiltonian and overlap matrices in eq. (2.21) as:

$$H_{ij} = \bar{\lambda}_{(i)}^{\dagger} \mathcal{S} \mathbf{D} \mathbf{E} \mathbf{D}^{\dagger} \mathcal{S}^{\dagger} \bar{\lambda}_{(j)}$$

$$S_{ij} = \bar{\lambda}_{(i)}^{\dagger} \mathcal{S} \mathbf{D} \mathbf{D}^{\dagger} \mathcal{S}^{\dagger} \bar{\lambda}_{(j)}$$
(2.27) eq:dipro\_eq20

Since the two monomer frontier orbitals that form the basis of this expansion are not orthogonal in general ( $S \neq 1$ ), it is necessary to transform eq. (2.20) into a standard eigenvalue problem of the form

$$\mathbf{H}^{\mathrm{eff}}\bar{\mathbf{C}}^{\mathrm{eff}} = E\bar{\mathbf{C}}^{\mathrm{eff}}$$
 (2.28) eqxdipro\_eq7

to make it correspond to eq. (2.19). According to Löwdin such a transformation can be achieved by

$$\mathbf{H}^{\text{eff}} = \mathbf{S}^{-1/2} \mathbf{H} \mathbf{S}^{-1/2}$$
. (2.29) eq:dipro\_eq9

This then yields an effective Hamiltonian matrix in an orthogonal basis, and its entries can directly be identified with the site energies  $\epsilon_i$  and transfer integrals  $J_{ij}$ :

$$\mathbf{H}^{\text{eff}} = \begin{pmatrix} e_i^{\text{eff}} & H_{ij}^{\text{eff}} \\ H_{ii}^{*,\text{eff}} & e_i^{\text{eff}} \end{pmatrix} = \begin{pmatrix} \epsilon_i & J_{ij} \\ J_{ij}^{*} & \epsilon_j \end{pmatrix}$$
(2.30) eq.dipro\_eq11

#### 2.7.2 DFT-based transfer integrals using DIPRO

The calculation of one electronic coupling element based on DFT using the DIPRO method requires the overlap matrix of atomic orbitals  $\mathcal{S}$ , the expansion coefficients for monomer  $\bar{\lambda}_{(k)} = \{\lambda_{\alpha}^{(k)}\}$  and dimer orbitals  $\bar{\mathbf{D}}_{(n)} = \{D_{\alpha}^{(n)}\}$ , as well as the orbital energies  $E_n$  of the dimer are required as input. In practical situations, performing self-consistent quantum-chemical calculations for each individual monomer and one for the dimer to obtain this input data is extremely demanding. Several simplifications can be made to reduce the computational effort, such as using non-Counterpoise basis sets for the monomers (thereby decoupling the monomer calculations from the dimer run) and performing only a single SCF step in a dimer calculation starting from an initial guess formed from a superposition of monomer orbitals. This "noCP+noSCF" variant of DIPRO is shown in figure 2.4(a) and recommended for production runs. A detailed comparative study of the different variants can be found in [5]. The code currently contains supports evaluation of transfer integrals from quantum-chemical calculations performed with the Gaussian, Turbomole, and NWChem packages. The interfacing procedure consists of three main steps: generation of input files for monomers and dimers, performing the actual quantum-chemical calculations, and calculating the transfer integrals.



Figure 2.4: Schematics of the DIPRO method. (a) General workflow of the projection technique. (b) Strategy of the efficient noCP+noSCF implementation, in which the monomer calculations are performed independently form the dimer configurations (noCP), using the edft calculator. The dimer Hamiltonian is subsequently constructed based on an initial guess formed from monomer orbitals and only diagonalized once (noSCF) before the transfer integral is calculated by projection. This second step is performed by the idft calculator.

fig:dipro\_scheme

#### Monomer calculations

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First, hopping sites and a neighbor list need to be generated from the atomistic topology and trajectory and written to the state.sql file. Then the parallel edft calculator manages the calculation of the monomer properties required for the determination of electronic coupling elements. Specifically, the individual steps it performs are:

1. Creation of a job file containing the list of molecules to be calculated with DFT

```
Writing job file for edft

| ctp_parallel -o options.xml -f state.sql -e edft -j write
```

2. Running of all jobs in job file

```
Running all edft jobs

I ctp_parallel -o options.xml -f state.sql -e edft -j run
```

which includes

• creating the input files for the DFT calculation (using the package specified in options.xml) in the directory

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OR\_FILES/package/frame\_F/mol\_M

where F is the index of the frame in the trajectory, M is the index of a molecule in this frame,

- executing the DFT run, and
- after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients), and saving it in compressed form to

OR\_FILES/molecules/frame\_F/molecule\_M.orb

#### Calculating the transfer integrals

After the momomer calculations have been completed successfully, the respective runs for dimers from the neighborlist can be performed using the parallel idft calculator, which manages the DFT runs for the hopping pairs and determines the coupling element using DIPRO. Again, several steps are required:

1. Creation of a job file containing the list of pairs to be calculated with DFT

```
Writing job file for idft

| ctp_parallel -o options.xml -f state.sql -e idft -j write
```

2. Running of all jobs in job file

```
Running all idft jobs

| ctp_parallel -o options.xml -f state.sql -e idft -j run
```

#### which includes

• creating the input files (including the merged guess for a noSCF calculation, if requested) for the DFT calculation (using the package specified in options.xml) in the directory

```
OR_FILES/package/frame_F/pair_M_N
```

where M and N are the indices of the molecules in this pair,

- executing the DFT run, and
- after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients and energies, atomic orbital overlap matrix), and saving the pair information in compressed form to

```
OR_FILES/pairs/frame_F/pair_M_N.orb
```

- loading the monomer orbitals from the previously saved \*.orb files.
- calculating the coupling elements and write them to the job file
- 3. Reading the coupling elements from the job file and saving them to the state.sql file

```
Saving idft results from job file to state.sql

| ctp_parallel -o options.xml -f state.sql -e idft -j read
```

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#### 2.7.3 ZINDO-based transfer integrals using MOO

An approximate method based on Zerner's Intermediate Neglect of Differential Overlap (ZINDO) has been described in Ref. [24]. This semiempirical method is substantially faster than first-principles approaches, since it avoids the self-consistent calculations on each individual monomer and dimer. This allows to construct the matrix elements of the ZINDO Hamiltonian of the dimer from the weighted overlap of molecular orbitals of the two monomers. Together with the introduction of rigid segments, only a single self-consistent calculation on one isolated conjugated segment is required. All relevant molecular overlaps can then be constructed from the obtained molecular orbitals.

The main advantage of the molecular orbital overlap (MOO) library is *fast* evaluation of electronic coupling elements. Note that MOO is based on the ZINDO Hamiltonian which has limited applicability. The general advice is to first compare the accuracy of the MOO method to the DFT-based calculations.

MOO can be used both in a standalone mode and as an izindo calculator of VOTCA-CTP.

Since MOO constructs the Fock operator of a dimer from the molecular orbitals of monomers by translating and rotating the orbitals of rigid fragments, the optimized geometry of all conjugated segments and the coefficients of the molecular orbitals are required as its input in addition to the state file (state.sql) with the neighbor list. Coordinates are stored in geometry.xyz files with four columns, first being the atom type and the next three atom coordinates. This is a standard xyz format without a header. Note that the atom order in the geometry.xyz files can be different from that of the mapping files. The correspondence between the two is established in the map.xml file.



Izindo requires the specification of orbitals for hole and electron transport in map.xml. They are the HOMO and LUMO respectively and can be retrieved from the log file from which the zindo.orb file is generated. The number of alpha electrons is the HOMO, the LUMO is HOMO+1

The calculated transfer integrals are immediately saved to the state.sql file.



## 2.8 Charge transfer rate

Charge transfer rates can be postulated based on intuitive physical considerations, as it is done in the Gaussian disorder models [25–28]. Alternatively, charge transfer theories can be used to evaluate rates from quantum chemical calculations [3, 5, 12, 29–31]. In spite of being significantly more computationally demanding, the latter approach allows to link the chemical and electronic structure, as well as the morphology, to charge dynamics.

#### 2.8.1 Classical charge transfer rate

sec:rate\_classic

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The high temperature limit of classical charge transfer theory [32, 33] is often used as a tradeoff between theoretical rigor and computational complexity. It captures key parameters which influence charge transport while at the same time providing an analytical expression for the rate. Within this limit, the transfer rate for a charge to hop from a site i to a site j reads

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{J_{ij}^2}{\sqrt{4\pi \lambda_{ij} k_{\rm B} T}} \exp\left[-\frac{(\Delta E_{ij} - \lambda_{ij})^2}{4\lambda_{ij} k_{\rm B} T}\right],\tag{2.31}$$

where T is the temperature,  $\lambda_{ij} = \lambda_{ij}^{\text{int}} + \lambda_{ij}^{\text{out}}$  is the reorganization energy, which is a sum of intraand inter-molecular (outersphere) contributions,  $\Delta E_{ij}$  is the site-energy difference, or driving force, and  $J_{ij}$  is the electronic coupling element, or transfer integral.

#### 2.8.2 Semi-classical bimolecular rate

The main assumptions in eq. (2.31) are non-adiabaticity (small electronic coupling and charge transfer between two diabatic, non-interacting states), and harmonic promoting modes, which are treated classically. At ambient conditions, however, the intramolecular promoting mode, which roughly corresponds to C-C bond stretching, has a vibrational energy of  $\hbar\omega\approx0.2\,\mathrm{eV}\gg1.00$  kb and should be treated quantum-mechanically. The outer-sphere (slow) mode has much lower vibrational energy than the intramolecular promoting mode, and therefore can be treated classically. The weak interaction between molecules also implies that each molecule has its own, practically independent, set of quantum mechanical degrees of freedom.

A more general, quantum-classical expression for a bimolecular multi-channel rate is derived in the Supporting Information of ref. [3] and has the following form

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{|J_{ij}|^2}{\sqrt{4\pi\lambda_{ij}^{\text{out}}k_{\text{B}}T}} \sum_{l',m'=0}^{\infty} |\langle \chi_{i0}^c | \chi_{il'}^n \rangle|^2 |\langle \chi_{j0}^n | \chi_{jm'}^c \rangle|^2 \exp\left\{-\frac{\left[\Delta E_{ij} - \hbar(l'\omega_i^n + m'\omega_j^c) - \lambda_{ij}^{\text{out}}\right]^2}{4\lambda_{ij}^{\text{out}}k_{\text{B}}T}\right\}. \tag{2.32}$$

If the curvatures of intramolecular PES of charged and neutral states of a molecule are different, that is  $\omega_i^c \neq \omega_i^n$ , the corresponding reorganization energies,  $\lambda_i^{cn} = \frac{1}{2}[\omega_i^n(q_i^n-q_i^c)]^2$  and  $\lambda_i^{nc} = \frac{1}{2}[\omega_i^c(q_i^n-q_i^c)]^2$ , will also differ. In this case the Franck-Condon (FC) factors for discharging of molecule i read [34]

$$|\langle \chi_{i0}^{c} | \chi_{il'}^{n} \rangle|^{2} = \frac{2}{2^{l'} l'!} \frac{\sqrt{\omega_{i}^{c} \omega_{i}^{n}}}{(\omega_{i}^{c} + \omega_{i}^{n})} \exp\left(-|s_{i}|\right) \left[ \sum_{k=0}^{l'} \binom{l'}{k} \left(\frac{2\omega_{i}^{c}}{\omega_{i}^{c} + \omega_{i}^{n}}\right)^{k/2} \frac{k!}{(k/2)!} H_{l'-k} \left(\frac{s_{i}}{\sqrt{2S_{i}^{cn}}}\right) \right]^{2},$$
(2.33)

where  $H_n(x)$  is a Hermite polynomial,  $s_i = 2\sqrt{\lambda_i^{nc}\lambda_i^{cn}}/\hbar(\omega_i^c + \omega_i^n)$ , and  $S_i^{cn} = \lambda_i^{cn}/\hbar\omega_i^c$ . The FC factors for charging of molecule j can be obtained by substituting  $(s_i, S_i^{cn}, \omega_i^c)$  with  $(-s_j, S_j^{nc}, \omega_j^n)$ . In order to evaluate the FC factors, the internal reorganization energy  $\lambda_i^{cn}$  can be computed from the intramolecular PES.

#### 2.8.3 Semi-classical rate

sec:rate semiclassica

One can also use the quantum-classical rate with a common set of vibrational coordinates [?]

$$\omega_{ij} = \frac{2\pi}{\hbar} \frac{|J_{ij}|^2}{\sqrt{4\pi\lambda_{ij}^{\text{out}} k_{\text{B}} T}} \sum_{N=0}^{\infty} \frac{1}{N!} \left(\frac{\lambda_{ij}^{\text{int}}}{\hbar\omega^{\text{int}}}\right)^N \exp\left(-\frac{\lambda_{ij}^{\text{int}}}{\hbar\omega^{\text{int}}}\right) \exp\left\{-\frac{\left[\Delta E_{ij} - \hbar N\omega^{\text{int}} - \lambda_{ij}^{\text{out}}\right]^2}{4\lambda_{ij}^{\text{out}} k_{\text{B}} T}\right\}. \tag{2.34}$$

Numerical estimates show that if  $\lambda_{ij}^{\rm int} \approx \lambda_{ij}^{\rm out}$  and  $|\Delta E_{ij}| \ll \lambda_{ij}^{\rm out}$  the rates are similar to those of eq. (2.31). In general, there is no robust method to compute  $\lambda_{ij}^{\rm out}$  [35] and both reorganization energies are often assumed to be of the same order of magnitude. In this case the second condition also holds, unless there are large differences in electron affinities or ionization potentials of neighboring molecules, e.g. in donor-acceptor blends.

To calculate rates of the type specified in options.xml for all pairs in the neighbor list and to save them into the state.sql file, run the rates calculator. Note that all required ingredients (reorganization energies, transfer integrals, and site energies have to be calculated before).

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#### Calculation of transfer rates

ctp\_run -ooptions.xml -fstate.sql -erates

### 2.9 Master equation

Having determined the list of conjugated segments (hopping sites) and charge transfer rates between them, the next task is to solve the master equation which describes the time evolution of the system

where  $P_{\alpha}$  is the probability of the system to be in a state  $\alpha$  at time t and  $\Omega_{\alpha\beta}$  is the transition rate

$$\frac{\partial P_{\alpha}}{\partial t} = \sum_{\beta} P_{\beta} \Omega_{\beta\alpha} - \sum_{\beta} P_{\alpha} \Omega_{\alpha\beta}, \tag{2.35}$$

from state  $\alpha$  to state  $\beta$ . A state  $\alpha$  is specified by a set of site occupations,  $\{\alpha_i\}$ , where  $\alpha_i = 1(0)$ for an occupied (unoccupied) site i, and the matrix  $\hat{\Omega}$  can be constructed from rates  $\omega_{ij}$ . 487 The solution of eq. (2.35) is be obtained by using kinetic Monte Carlo (KMC) methods. KMC 488 explicitly simulates the dynamics of charge carriers by constructing a Markov chain in state space 489 and can find both stationary and transient solutions of the master equation. The main advantage 490 of KMC is that only states with a direct link to the current state need to be considered at each step. 49 Since these can be constructed solely from current site occupations, extensions to multiple charge carriers (without the mean-field approximation), site-occupation dependent rates (needed for 493 the explicit treatment of Coulomb interactions), and different types of interacting particles and 494 processes, are straightforward. To optimize memory usage and efficiency, a combination of the 495 variable step size method [36] and the first reaction method is implemented. To obtain the dynamics of charges using KMC, the program kmc\_run executes a specific calcula-

tor after reading its options (charge carrier type, runtime, numer of carriers etc.) from options.xml.

KMC for a single carrier in periodic boundary conditions

| kmc\_run -o options.xml -f state.sql -e kmcsingle

KMC for multiple carriers of the same type in periodic boundary conditions

| kmc\_run -o options.xml -f state.sql -e kmcmultiple

#### 2.9.1 Extrapolation to nondispersive mobilities

Predictions of charge-carrier mobilities in partially disordered semiconductors rely on charge transport simulations in systems which are only several nanometers thick. As a result, simulated charge transport might be dispersive for materials with large energetic disorder [37, 38] and simulated mobilities are system-size dependent. In time-of-flight (TOF) experiments, however, a typical sample thickness is in the micrometer range and transport is often nondispersive. In order to link simulation and experiment, one needs to extract the nondispersive mobility from simulations of small systems, where charge transport is dispersive at room temperature. Such extrapolation is possible if the temperature dependence of the nondispersive mobility is known in a wide temperature range. For example, one can use analytical results derived for one-dimensional models [39–41]. The mobility-temperature dependence can then be parametrized by

simulating charge transport at elevated temperatures, for which transport is nondispersive even for small system sizes. This dependence can then be used to extrapolate to the nondispersive mobility at room temperature [4].

For  $Alq_3$ , the charge carrier mobility of a periodic system of 512 molecules was shown to be more than three orders of magnitude higher than the nondispersive mobility of an infinitely large system [4]. Furthermore, it was shown that the transition between the dispersive and nondispersive transport has a logarithmic dependence on the number of hopping sites N. Hence, a brute-force increase of the system size cannot resolve the problem for compounds with large energetic disorder  $\sigma$ , since N increases exponentially with  $\sigma^2$ .

### 2.10 Macroscopic observables

Spatial distributions of charge and current densities can provide a better insight in the microscopic mechanisms of charge transport. If O is an observable which has a value  $O_{\alpha}$  in a state  $\alpha$ , its ensemble average at time t is a sum over all states weighted by the probability  $P_{\alpha}$  to be in a state  $\alpha$  at time t

$$\langle O \rangle = \sum_{\alpha} O_{\alpha} P_{\alpha}.$$
 (2.36) equiensemble

If O does not explicitly depend on time, the time evolution of  $\langle O \rangle$  can be calculated as

$$\frac{d\langle O\rangle}{dt} = \sum_{\alpha,\beta} \left[ P_{\beta} \Omega_{\beta\alpha} - P_{\alpha} \Omega_{\alpha\beta} \right] O_{\alpha} = \sum_{\alpha,\beta} P_{\beta} \Omega_{\beta\alpha} \left[ O_{\alpha} - O_{\beta} \right]. \tag{2.37}$$

If averages are obtained from KMC trajectories,  $P_{\alpha} = s_{\alpha}/s$ , where  $s_{\alpha}$  is the number of Markov chains ending in the state  $\alpha$  after time t, and s is the total number of chains.

Alternatively, one can calculate time averages by analyzing a single Markov chain. If the total occupation time of the state  $\alpha$  is  $\tau_{\alpha}$  then

$$\overline{O} = \frac{1}{\tau} \sum_{\alpha} O_{\alpha} \tau_{\alpha} \,, \tag{2.38}$$

where  $au = \sum_{lpha} au_{lpha}$  is the total time used for time averaging.

For ergodic systems and sufficient sampling times, ensemble and time averages should give identical results. In many cases, the averaging procedure reflects a specific experimental technique. For example, an ensemble average over several KMC trajectories with different starting conditions corresponds to averaging over injected charge carriers in a time-of-flight experiment. In what follows, we focus on the single charge carrier (low concentration of charges) case.

#### 2.10.1 Charge density

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For a specific type of particles, the microscopic charge density of a site i is proportional to the occupation probability of the site,  $p_i$ 

$$\rho_i = e p_i / V_i \,, \tag{2.39}$$

where, for an irregular lattice, the effective volume  $V_i$  can be obtained from a Voronoi tessellation of space. For reasonably uniform lattices (uniform site densities) this volume is almost independent of the site and a constant volume per cite,  $V_i = V/N$ , can be assumed. In the macroscopic limit, the charge density can be calculated using a smoothing kernel function, i.e. a distance-weighted average over multiple sites. Site occupations  $p_i$  can be obtained from eq. (2.36) or eq. (2.38) by using the occupation of site i in state  $\alpha$  as an observable.

If the system is in thermodynamic equilibrium, that is without sources or sinks and without circular currents (and therefore no net flux) a condition, known as detailed balance, holds

$$p_j \omega_{ji} = p_i \omega_{ij},$$
 (2.40) equivalent equivalence

It can be used to test whether the system is ergodic or not by correlating  $\log p_i$  and the site energy  $E_i$ . Indeed, if  $\lambda_{ij} = \lambda_{ji}$  the ratios of forward and backward rates are determined solely by the energetic disorder,  $\omega_{ji}/\omega_{ij} = \exp(-\Delta E_{ij}/k_BT)$  (see eq. (2.31)).

#### **2.10.2** Current

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If the position of the charge,  $\vec{r}$ , is an observable, the time evolution of its average  $\langle \vec{r} \rangle$  is the total current in the system

$$\vec{J} = e \langle \vec{v} \rangle = e \frac{d \langle \vec{r} \rangle}{dt} = e \sum_{i,j} p_j \omega_{ji} (\vec{r}_i - \vec{r}_j). \tag{2.41}$$

Symmetrizing this expression we obtain

$$\vec{J} = \frac{1}{2}e\sum_{i,j} (p_j\omega_{ji} - p_i\omega_{ij})\vec{r}_{ij}, \tag{2.42}$$

where  $\vec{r}_{ij} = \vec{r}_i - \vec{r}_j$ . Symmetrization ensures equal flux splitting between neighboring sites and absence of local average fluxes in equilibrium. It allows to define a local current through site i as

$$\vec{J}_i = \frac{1}{2}e\sum_j \left(p_j\omega_{ji} - p_i\omega_{ij}\right)\vec{r}_{ij}. \tag{2.43}$$

A large value of the local current indicates that the site contributes considerably to the total current. A collection of such sites thus represents most favorable charge pathways [42].

#### 2.10.3 Mobility and diffusion constant

For a single particle, e.g. a charge or an exciton, a zero-field mobility can be determined by studying particle diffusion in the absence of external fields. Using the particle displacement squared,  $\Delta r_i^2$ , as an observable we obtain

$$2dD_{\gamma\delta} = \frac{d\left\langle \Delta r_{i,\gamma} \Delta r_{i,\delta} \right\rangle}{dt} = \sum_{\substack{i,j\\i \neq j}} p_j \omega_{ji} \left( \Delta r_{i,\gamma} \Delta r_{i,\delta} - \Delta r_{j,\gamma} \Delta r_{j,\delta} \right) = \sum_{\substack{i,j\\i \neq j}} p_j \omega_{ji} \left( r_{i,\gamma} r_{i,\delta} - r_{j,\gamma} r_{j,\delta} \right) \,. \tag{2.44}$$

Here  $\vec{r_i}$  is the coordinate of the site i,  $D_{\gamma\delta}$  is the diffusion tensor,  $\gamma, \delta = x, y, z$ , and d = 3 is the system dimension. Using the Einstein relation,

$$D_{\gamma\delta} = k_{\rm B}T\mu_{\gamma\delta}\,,\tag{2.45}$$

one can, in principle, obtain the zero-field mobility tensor  $\mu_{\gamma\delta}$ . Eq. (2.44), however, does not take into account the use of periodic boundary conditions when simulating charge dynamics. In this case, the simulated occupation probabilities can be compared to the solution of the Smoluchowski equation with periodic boundary conditions (see the supporting information for details).

Alternatively, one can directly analyze time-evolution of the KMC trajectory and obtain the diffusion tensor from a linear fit to the mean square displacement,  $\overline{\Delta r_{i,\gamma}\Delta r_{i,\delta}}=2dD_{\gamma\delta}t$ .

The charge carrier mobility tensor,  $\hat{\mu}$ , for any value of the external field can be determined either from the average charge velocity defined in eq. (2.41)

$$\langle \vec{v} \rangle = \sum_{i,j} p_j \omega_{ji} (\vec{r}_i - \vec{r}_j) = \hat{\mu} \vec{F} , \qquad (2.46)$$

or directly from the KMC trajectory. In the latter case the velocity is calculated from the unwrapped (if periodic boundary conditions are used) charge displacement vector divided by the total simulation time. Projecting this velocity on the direction of the field  $\vec{F}$  yields the charge carrier mobility in this particular direction. In order to improve statistics, mobilities can be averaged over several KMC trajectories and MD snapshots.

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#### 2.10.4 Spatial correlations of energetic disorder

Long-range, e.g. electrostatic and polarization, interactions often result in spatially correlated disorder [43], which affects the onset of the mobility-field (Poole-Frenkel) dependence [39, 44, 45]. To quantify the degree of correlation, one can calculate the spatial correlation function of  $E_i$  and  $E_j$  at a distance  $r_{ij}$ 

$$C(r_{ij}) = \frac{\langle (E_i - \langle E \rangle) (E_j - \langle E \rangle) \rangle}{\langle (E_i - \langle E \rangle)^2 \rangle}, \tag{2.47}$$

where  $\langle E \rangle$  is the average site energy.  $C(r_{ij})$  is zero if  $E_i$  and  $E_j$  are uncorrelated and 1 if they are fully correlated. For a system of randomly oriented point dipoles, the correlation function decays as 1/r at large distances [46].

For systems with spatial correlations, variations in site energy differences,  $\Delta E_{ij}$ , of pairs of molecules from the neighbor list are smaller than variations in site energies,  $E_i$ , of all individual molecules. Since only neighbor list pairs affect transport, the distribution of  $\Delta E_{ij}$  rather than that of individual site energies,  $E_i$ , should be used to characterize energetic disorder.

Note that the eanalyze calculator takes into account all contributions to the site energies



#### 2.10.5 DFT-based transfer integrals using DIPRO

The calculation of one electronic coupling element based on DFT using the DIPRO method requires the overlap matrix of atomic orbitals  $\mathcal{S}$ , the expansion coefficients for monomer  $\bar{\lambda}_{(k)} = \{\lambda_{\alpha}^{(k)}\}$  and dimer orbitals  $\bar{\mathbf{D}}_{(n)} = \{D_{\alpha}^{(n)}\}$ , as well as the orbital energies  $E_n$  of the dimer are required as input. In practical situations, performing self-consistent quantum-chemical calculations for each individual monomer and one for the dimer to obtain this input data is extremely demanding. Several simplifications can be made to reduce the computational effort, such as using non-Counterpoise basis sets for the monomers (thereby decoupling the monomer calculations from the dimer run) and performing only a single SCF step in a dimer calculation starting from an initial guess formed from a superposition of monomer orbitals. This "noCP+noSCF" variant of DIPRO is shown in figure 2.4(a) and recommended for production runs. A detailed comparative study of the different variants can be found in [5]. The code currently contains supports evaluation of transfer integrals from quantum-chemical calculations performed with the Gaussian, Turbomole, and NWChem packages. The interfacing procedure consists of three main steps: generation of input files for monomers and dimers, performing the actual quantum-chemical calculations, and calculating the transfer integrals.

#### Monomer calculations

First, hopping sites and a neighbor list need to be generated from the atomistic topology and trajectory and written to the <code>state.sql</code> file. Then the parallel <code>edft</code> calculator manages the calculation of the monomer properties required for the determination of electronic coupling elements. Specifically, the individual steps it performs are:

1. Creation of a job file containing the list of molecules to be calculated with DFT

```
Writing job file for edft

| ctp_parallel -o options.xml -f state.sql -e edft -j write
```

#### 2. Running of all jobs in job file

## Running all edft jobs

```
ctp_parallel -o options.xml -f state.sql -e edft -j run
```

#### which includes

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• creating the input files for the DFT calculation (using the package specified in options.xml) in the directory

```
OR_FILES/package/frame_F/mol_M
```

where F is the index of the frame in the trajectory, M is the index of a molecule in this frame,

- executing the DFT run, and
- after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients), and saving it in compressed form to

```
OR_FILES/molecules/frame_F/molecule_M.orb
```

#### Calculating the transfer integrals

After the momomer calculations have been completed successfully, the respective runs for dimers from the neighborlist can be performed using the parallel idft calculator, which manages the DFT runs for the hopping pairs and determines the coupling element using DIPRO. Again, several steps are required:

1. Creation of a job file containing the list of pairs to be calculated with DFT

```
Writing job file for idft

| ctp_parallel -o options.xml -f state.sql -e idft -j write
```

2. Running of all jobs in job file

```
Running all idft jobs

| ctp_parallel -o options.xml -f state.sql -e idft -j run
```

#### which includes

creating the input files (including the merged guess for a noSCF calculation, if requested) for the DFT calculation (using the package specified in options.xml) in the directory

```
OR_FILES/package/frame_F/pair_M_N
```

where M and N are the indices of the molecules in this pair,

- executing the DFT run, and
- after completion of this run, parsing the output (number of electrons, basis set, molecular orbital expansion coefficients and energies, atomic orbital overlap matrix), and saving the pair information in compressed form to

```
OR_FILES/pairs/frame_F/pair_M_N.orb
```

- loading the monomer orbitals from the previously saved \*.orb files.
- calculating the coupling elements and write them to the job file
- 3. Reading the coupling elements from the job file and saving them to the state.sql file

Saving idft results from job file to state.sql

| ctp\_parallel -ooptions.xml -f state.sql -e idft -j read

## **Chapter 3**

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# Input and output files

## 3.1 Atomistic topology

If you are using GROMACS for generating atomistic configurations, it is possible to directly use the topology file provided by GROMACS (topology.tpr). In this case the GROMACS residue and atom names should be used to specify the coarse-grained topology and conjugated segments. A custom topology can also be defined using an XML file. Moreover, it s possible to partially overwrite the information provided in, for example, GROMACS topology file. We will illustrate how to create a custom topology file using DCV2T. The structure of DCV2T, together with atom type definitions, is shown in fig. 3.1. DCV2T has two thiophene (THI) and two dicyanovinyl (NIT) residues. The pdb file which contains residue types, residue numbering, atom names, atom types, and atom coordinates is shown in listing 3.1.



Figure 3.1: (a) DCV2T with atoms labelled according to residue\_number:residue\_name:atom\_name. There are four residues and two residue types: thiophene (THI) and dicyanovinyl (NIT). The corresponding pdb file is shown in listing 3.1. Atom numbering is used to split conjugated segments on rigid fragments and to link atomistic ((b) from GROMACS topology) and quantum descriptions (c).

fig:dcv2t

Listing 3.1: pdb file of DCV2T.

054							1				
654 655	HETATM	1	N1	NIT	1	2.388	8.533	11.066	1.00	4.14	N
656	HETATM	2	CN1	NIT	1	1.984	9.553	10.718	1.00	2.54	C
657	HETATM	3	N2	NIT	1	-1.138	10.872	10.087	1.00	3.24	N
658	HETATM	4	CN2	NIT	1	0.003	10.871	10.213	1.00	2.37	С
659	HETATM	5	CC1	NIT	1	1.441	10.824	10.327	1.00	1.91	С
660	HETATM	6	C1	NIT	1	2.193	11.939	10.071	1.00	1.61	С
661	HETATM	7	HN1	NIT	1	1.715	12.710	9.872	1.00	1.97	H
662	HETATM	8	S1	THI	2	4.758	10.743	10.130	1.00	1.52	S
663	HETATM	9	CA1	THI	2	3.613	12.024	9.948	1.00	1.22	C
664	HETATM	10	CA2	THI	2	6.099	11.836	9.997	1.00	1.30	С
665	HETATM	11	CB1	THI	2	4.251	13.243	9.782	1.00	1.39	C
666	HETATM	12	CB2	THI	2	5.658	13.131	9.818	1.00	1.45	C
667	HETATM	13	HC1	THI	2	3.800	14.047	9.660	1.00	1.66	H
668	HETATM	14	HC2	THI	2	6.230	13.860	9.731	1.00	1.74	H
669	HETATM	15	S1	THI	3	8.803	12.414	9.882	1.00	1.38	S
670	HETATM	16	CA1	THI	3	7.456	11.347	10.094	1.00	1.37	C
671	HETATM	17	CA2	THI	3	9.940	11.122	10.152	1.00	1.42	C
672	HETATM	18	CB1	THI	3	7.873	10.048	10.355	1.00	1.73	C
673	HETATM	19		THI	3	9.267	9.926	10.399	1.00	1.82	C
674	HETATM	20		THI	3	7.288	9.335	10.487	1.00	2.05	Н
675	HETATM	21	HC2	THI	3	9.704	9.123	10.576	1.00	2.21	Н
676	HETATM	22	N1	NIT	4	11.235	14.572	9.094	1.00	3.08	N
677	HETATM	23		NIT	4	11.665	13.566	9.441	1.00	2.04	C
678	HETATM	24	N2	NIT	4	14.733	12.005	10.009	1.00	2.17	N
679	HETATM	25		NIT	4	13.590	12.149	9.933	1.00	1.77	C
680	HETATM	26		NIT	4	12.156	12.282	9.861	1.00	1.71	C
681	HETATM	27	C1	NIT	4	11.363	11.220	10.154	1.00	1.59	С
88 <del>3</del>	HETATM	28	HN1	NIT	4	11.813	10.440	10.389	1.00	1.89	H

3.2. MAPPING FILE 29

### 3.2 Mapping file

secxmimap

The mapping file (referred here as map.xml) is used by the program ctp\_map to convert an atomistic trajectory to a trajectory with conjugated segments and rigid fragments. This trajectory is stored in a state file and contains positions, names, types of atoms belonging to rigid fragments. The description of the mapping options is given in table 3.1. An example of map.xml for a DCV2T molecule is shown in listing 3.2.

The file map.xml contains the whole electrostatic information about the molecules as well as the structural information. The toolpdb2map creates a map.xml from a pdb file and is a good starting point for further refinement.

tab:ma

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Table 3.1: Description of the XML mapping file (map.xml).

list:ma

Listing 3.2: Examle of map.xml for DCV2T. Each rigid fragment (coarse-grained bead) is defined by a list of atoms. Atom numbers, names, and residue names should correspond to those used in GROMACS topology (see the corresponing listing 3.1 of the pdb file).

```
this file is used to conver an atomistic trajectory to conjugated segments -->
     <topology> <!--
695
     <molecules>
     <molecule>
696
         <name>DCV2T-MOL</name> <!-- name of the conjugated molecule</pre>
697
         <mdname>Protein</mdname> <!-- corresponding name of this molecule in the MD trajectory , should be</pre>
699
               the same as the name given at the end of topol.top-->
700
         <segments>
701
         <segment>
               <name>DCV</name> <!-- name of the conjugated segment within the molecule --->
<qmcoords>QC_FILES/DCV2T.xyz</qmcoords> <!-- QM coordinates of the conjugated segment --->
702
703
704
                     <!-- IZINDO INPUT -->
706
               <basisset>INDO</basisset>
707
                $$ \ensuremath{$<$}$ orbitals>QC_FILES/DCV2T.orb</r> $$ \ensuremath{$<$}$ orbital_h>50</torbital_h><!-- Number of the HOMO Orbital (e.g. alpha electrons, can be alpha electrons). $$
708
709
                     found in the log-file belonging to DCV2T.orb) -->
710
711
                    <!-- EMULTIPOLE INPUT -->
712
               <multipoles_n>MP_FILES/DCV2T_h.mps</multipoles_n><!-- Multipole file for neutral state -->
<multipoles_h>MP_FILES/DCV2T_h.mps</multipoles_h><!-- Multipole file for hole state -->
713
714
715
               716
718
719
720
                          <!-- EINTERNAL INPUT --->
721

    <U_cC_nN_h>0.0</U_cC_nN_h> <!-- Site energy -->
    <U_nC_nN_h>0.1</U_nC_nN_h> <!-- Reorg. discharge -->

722
               <U_cN_cC_h>0.1</U_cN_cC_h> <!-- Reorg. charge
725
                    <!-- MD QM MP Mapping -->
726
               <fragments>
727
                <fragment>
728
                 <name>NI1</name> <!-- name of the rigid fragment within the segment --->
729
                 <!-- list of atoms in the fragment resnum:resname:atomname -
730
                 <mdatoms>1:NIT:N1 1:NIT:CN1 1:NIT:N2 1:NIT:CN2 1:NIT:CC1 1:NIT:C1 1:NIT:HN1</mdatoms>
731
                 <!-- corresponding ground state geometry atomnumber:atomtype read from .xyz file--> <qmatoms> 20:N 19:C 14:N 13:C 12:C 11:C 23:H </qmatoms> <!-- corresponding group state geometry multipoles read from .mps files --->
732
733
734
                <mpoles> 20:N 19:C 14:N 13:C 12:C 11:C
<!-- weights to determine the fragment center (here CoM is used) -</pre>
735
                                                                                                         23:H </mpoles>
               737
738
739
                 <localframe> 20 19 14 </localframe>
740
                <!-- Optional parameters (if not set <localframe> is used): used when atom labels in the .mps
                       and .xyz file differ or more sites in the .mps file are used, so refers to <mpoles>
742
743
                 <localframe_mps> 20 19 14 </localframe_mps>
                <!-- Optional parameters (if not set <localframe> is used): weights to determine the fragment center (here CoM is used), used when atom labels in the .mps and .xyz file differ or additional sites in the .mps file are used -->
744
745
746
                <weights_mps>
                                     14
                                                 12
                                                              14
                    weights mps>
```

```
<\!\!!-- Optional\ flag:\ says \quad if\ a\ site\ is\ virtual\ or\ not\ ,\ (virtual=1,\ real=0)--\!\!>
749
             <virtual_mps>
                 virtual_mps>
751
752
            </fragment>
753
754
           <fragment>
              <name>TH1</name>
755
              <mdatoms>2:THI:S1 2:THI:CA1 2:THI:CA2 2:THI:CB1 2:THI:B2 2:THI:HC1 2:THI:HC2</mdatoms>
756
                                    8:C 6:C 9:C 10:C 24:H
8:C 6:C 9:C 10:C 24:H
                                                                                24:H 25:H </qmatoms>
24:H 25:H </mpoles>
1 1 </weights>
              <qmatoms>
                          7:S
              <mpoles>
758
                                   8:C
                          32
                                     12
759
              <weights>
                                              12
                                                                               1
              <localframe> 7 8 6 </localframe>
760
           </fragment>
761
762
764
              <name>TH2</name>
              <mdatoms>3:THI:S1 3:THI:CA1 3:THI:CA2 3:THI:CB2 3:THI:BC3 3:THI:HC1 3:THI:HC2/mdatoms>
765
              <qmatoms> 3:S 4:C 2:C
<weights> 32 12 12
                                                               1:C 26:H 27:H 
                                                      5:C
766
                                                           12
767
              <localframe> 3 4 2 </localframe>
768
           </fragment>
770
771
           <fragment>
              <name>NT2</name>
772
              <mdatoms>4:NIT:N1 4:NIT:CN1 4:NIT:N2 4:NIT:CN2 4:NIT:CC1 4:NIT:C1 4:NIT:HN1//mdatoms>
773
                                   21:C
                                             18:N 17:C 16:C
18:N 17:C 16:C
                                                                             15:C
15:C
              <qmatoms> 22:N
774
              <mpoles> 22:N
                                             18:N
                                                                                       28:H </mpoles>
775
              <weights>
                                                                                           1 </weights>
              <localframe> 22 21 18 </localframe>
777
           </fragment>
778
          </fragments>
779
        </segment>
780
781
        </segments>
    </molecule>
783
    </molecules>
    </topology>
```

## 3.3 Conjugated segments

The file describing hopping sites, or conjugated segments, is used by practically all programs and calculators. It links the coarse-grained trajectory (positions and orientations of rigid fragments) and quantum-mechanical descriptions of all conjugated segments. The description of this XML file (segments.xml) is given in table 3.2. An example for DCV2T is shown in listing 3.3.

Table 3.2: Description of conjugated segments (segments.xml).

tab:segments

790 list:segments

787

Listing 3.3: XML file describing conjugated segments. Note that the mapping and weights for each segment are separated by a colon.

```
<segments>
793
                 <segment> <!-- DCV2T here is one conjugated segment -
                           <coordinates>DCV2T.xyz</coordinates> <!-</pre>
                            <orbitals>DCV2T.fort.7</orbitals> <!-- ZINDO orbitals [GAUSSIAN] -->
795
796
                            <basisset>INDO</basisset>
                            <torbital>50</torbital> <!-- HOMO for hole conduction --->
797
                           <edischarging>0.084</edischarging><echarging>0.086</pr
798
799
                           <qneutral>DCV2T_neutr esp</qneutral> <!-- partial charges of a neutral molecule -->
<qcharged>DCV2T_catio.esp</qcharged> <!-- partial charges of a cation -->
800
                           <energy>0.0</energy> <!-- cite energy -->
<name>DCV2T</name> <!-- name of the conjugated segment -->
<map> <!-- rigid fragments separated by a colon -->
802
803
804
                                      22 21 18 17 16 15 28:

3 2 4 1 5 27 26:

7 6 8 10 9 25 24:
805
806
807
                                      20 19 14 13 12 11 23
809
                            <weights> <!--
                                                for centers of rigid fragments -->
810
                                      14 12 14 12 12 12 1:
811
812
                                      32 12 12 12 12 1 1:
```

```
14 12 14 12 12 12 1
814
                       </weights>
815
              </segment>
816
     </segments>
818
```

#### 3.4 Molecular orbitals

If the semi-empirical method is used to calculate electronic coupling elements, molecular orbitals of all molecules must be supplied. They can be generated using Gaussian program. The Gaussian input file for DCV2T is shown in listing 3.4. Provided with this input, Gaussian will generate fort. 7 file which contains the molecular orbitals of a DCV2T. This file can be renamed to DCV2T.orb. Note that the order of the atoms in the input file and the order of coefficients should always match. Therefore, the coordinate part of the input file must be supplied together with the orbitals. We will assume the coordinates, in the format atom type: x y z, is saved to the DCV2T.xyz file.



#### Attention

Izindo requires the specification of orbitals for hole and electron transport in map.xml. They are the HOMO and LUMO respectively and can be retrieved from the log file from which the DCV2T.orb file is generated. The number of alpha electrons is the HOMO, the LUMO is HOMO+1

list:zindo\_orbitals

820

822

823

826

827

Listing 3.4: Gaussian input file get\_orbitals.com used for generating molecular orbitals. The first line contains the name of the check file, the second the requested RAM. int=zindos requests the method ZINDO, punch=mo states that the molecular orbitals ought to be written to the fort. 7 file, nosymm forbids use of symmetry and is necessary to ensure correct position of orbitals with respect to the provided coordinates. The two integer numbers correspond to the charge and multiplicity of the system: 01 corresponds to a neutral system with a multiplicity of one. They are followed by the types and coordinates of all atoms in the molecule.

```
828
    %chk=DCV2T.chk
    %mem=100Mb
830
831
    #p int=zindos punch=mo nosymm
832
    DCV2T molecular orbitals
833
834
835
    0 1
              -1.44650
                              2.12185
                                              0.00135
    S
836
837
    C
              -2.43098
                               0.58936
                                              -0.00048
              -1.59065
                              -0.51859
                                              -0.00146
838
              -0.21222
                              -0.22233
                                              -0.00095
839
    С
               0.07761
                               1.13376
                                               0.00040
840
    S
               2.87651
                               0.79316
                                               0.00148
841
842
               3.86099
                               2.32565
                                               0.00235
               3.02066
                               3.43359
                                               0.00231
843
               1.64223
                               3.13733
                                               0.00162
    C
844
845
               1.35240
                               1.78125
                                               0.00114
              -3.85350
                               0.52245
                                              -0.00081
846
                                              -0.00008
              -4.79569
                               1.52479
847
              -6.18500
                               1.18622
                                              -0.00117
848
    С
              -4.47544
                               2.91565
                                               0.00081
849
850
               5.28350
                               2.39256
                                               0.00296
               6.22569
                               1.39020
                                               0.00327
851
               7.61500
                               1.72876
                                               0.00432
852
853
    C
               5.90542
                              -0.00064
                                               0.00333
854
              -7.32389
                               0.89743
                                               -0.00195
              -4.21872
                               4.06274
    Ν
                                               0.00142
855
856
    Ν
               8.75389
                               2.01754
                                               0.00510
857
               5.64864
                              -1.14772
              -1.98064
                              -1.52966
                                              -0.00256
858
    H
             0.55785 -0.98374 -0.00169
```

```
4.44466 0.00272
              3.41065
860
              0.87216
                            3.89874
                                          0.00147
861
   H
862
   Н
             -4.24640
                            -0.49192
                                          -0.00188
   Н
              5.67641
                            3.40692
                                           0.00337
863
```

### 3.5 Monomer calculations for DFT transfer integrals

list:edft\_gaussian\_xm

Listing 3.5: Example package.xml file for the Gaussian package required in the options of the edft calculator for the monomer calculations as preparation for the determination of transfer integrals using DIPRO.

```
<package>
867
      <name>gaussian</name>
      <executable>g09</executable>
869
      <checkpoint></checkpoint>
870
      <scratch></scratch>
871
872
873
      <charge>0</charge>
      <spin>1</spin>
874
      <options># pop=minimal pbepbe/6-311g** scf=tight punch=mo nosymm test/options>
875
876
      <memory>1Gb</memory>
      <threads>2</threads>
877
878
879
      <cleanup></cleanup>
    </package>
889
```

list:edft\_turbomole\_xml

Listing 3.6: Example package.xml file for the Turbomole package required in the options of the edft calculator for the monomer calculations as preparation for the determination of transfer integrals using DIPRO.

```
882
     <package>
883
       <name>turbomole</name>
884
885
       <executable>ridft</executable>
       <scratch>/tmp</scratch>
886
887
888
       <options>
    TITLE
889
890
    a coord
891
    no
892
893
    b all def-TZVP
894
    eht
895
    0
897
898
    dft
899
900
    on
901
    func
    pbe
902
903
    grid
904
    m3
905
906
    ri
907
    on
    m 300
908
909
910
    scf
    conv
911
912
    iter
913
    200
914
915
```

```
916 marij
917
918 q
919 </options>
920
921 <cleanup></cleanup>
922 </package>
```

list:edft nwchem x

Listing 3.7: Example package.xml file for the NWChem package required in the options of the edft calculator for the monomer calculations as preparation for the determination of transfer integrals using DIPRO.

```
924
    <package>
925
926
      <name>nwchem</name>
927
      <executable>nwchem</executable>
      <checkpoint></checkpoint>
928
      <scratch>/tmp/nwchem</scratch>
      <charge>0</charge>
930
931
      <spin>1</spin>
      <threads>1</threads>
932
      <memory></memory>
933
934
      <options>
    start
935
936
    basis
     * library 6-311gss
938
    end
    memory 1500 mb
939
940
    dft.
941
942
     xc xpbe96 cpbe96
     direct
943
     iterations 100
944
945
     noprint "final vectors analysis"
    end
946
947
    task dft
948
    </options>
      <cleanup></cleanup>
949
    </package>
850
```

# 3.6 Pair calculations for DFT transfer integrals

list:idft\_gaussian\_xm

Listing 3.8: Example package.xml file for the Gaussian package required in the options of the idft calculator for the pair calculations and the determination of transfer integrals using DIPRO.

```
953
    <package>
      <name>gaussian</name>
955
956
      <executable>g09</executable>
      <checkpoint></checkpoint>
957
      <scratch></scratch>
958
959
      <charge>0</charge>
960
      <spin>1</spin>
961
      <options># pop=minimal pbepbe/6-311g** nosymm IOp(3/33=1,3/36=-1) punch=mo guess=cards scf=(maxcycle=1,
      <memory>1Gb</memory>
963
      <threads>1</threads>
964
965
      <cleanup></cleanup>
966
    </package>
967
```

list:idft\_turbomole\_xml

Listing 3.9: Example package.xml file for the Turbomole package required in the options of the idft calculator for the pair calculations and the determination of transfer integrals using DIPRO.

```
<name>turbomole</name>
971
        <executable>ridft</executable>
972
973
        <scratch>/tmp</scratch>
974
        <options>
975
976
     $intsdebug cao
     a coord
977
978
979
     no
     b all def-TZVP
980
981
     eht
982
983
     0
984
985
     dft
986
987
     on
     func
988
989
     pbe
     grid
990
991
     m3
992
     ri
993
994
     on
995
     m 300
996
997
     scf
998
     conv
999
     iter
1000
1001
     diis
1002
1003
1004
     damp
1005
     0.00
1006
1007
1008
     marij
1009
1010
1011
        </options>
1012
1013
1014
        <cleanup></cleanup>
     </package>
1815
```

list:idft\_nwchem\_xml

Listing 3.10: Example package.xml file for the NWChem package required in the options of the idft calculator for the pair calculations and the determination of transfer integrals using DIPRO.

```
1017
     <package>
1019
       <name>nwchem</name>
       <executable>nwchem</executable>
1020
       <checkpoint></checkpoint>
1021
       <scratch>/tmp/nwchem</scratch>
1022
       <charge>0</charge>
1023
1024
       <spin>1</spin>
       <memory></memory>
1025
       <threads>1</threads>
1026
1027
       <options>
    start
1028
1029
    basis
     * library 6-311gss
1030
1031 end
    memory 1500 mb
1032
1033
1034
    print "ao overlap"
```

```
xc xpbe96 cpbe96
1036
1037
      direct
1038
      iterations 1
      convergence nodamping nodiis
1039
      noprint "final vectors analysis"
1040
1041
      vectors input system.movecs
    end
1042
    task dft
1043
     </options>
1044
      <cleanup></cleanup>
1045
     </package>
1849
```

# 3.7 DFT transfer integrals

list:TI\_xml

Listing 3.11: Example TI.xml file created as the output of a DIPRO calculation. Due to slightly different implementations, the orbitals indices refer to monomer indices in a Gaussian run but to indices in the merged dimer guess in a Turbomole run.

```
<pair name="pair_100_155">
1050
         <parameters>
1051
1052
            <HOMO_A>162</HOMO_A>
            <NoccA>1</NoccA>
1053
1054
            <LUMO_A>164</LUMO_A>
1055
            <NvirtA>1</NvirtA>
            <hOMO_B>161</hOMO_B>
1056
            <NoccB>1</NoccB>
1057
            <LUMO_B>163</LUMO_B>
1058
            <NvirtB>1</NvirtB>
1059
         </parameters>
1060
          <transport name="hole">
1061
              <channel name="single">
1062
                   <J>1.546400416750696E-003</J>
1063
                   <e_A>-6.30726450715697</e_A>
1064
                   <e_B>-6.36775613794166</e_B>
1065
               </channel>
1066
               <channel name="multi">
1067
                  <molecule name="A">
1068
                      <e_HOMOm0>-6.30726450715697</e_HOMOm0>
1069
                  </molecule>
1070
                  <molecule name="B">
1071
                      <e_HOMOm0>-6.36775613794166</e_HOMOm0>
1072
1073
                  </molecule>
                      <dimer name="integrals">
1074
                            <T_00>1.546400416750696E-003</T_00>
1075
1076
                            <J_sq_degen>2.391354248926727E-006</J_sq_degen>
                            <J_sq_boltz>2.391354248926727E-006</J_sq_boltz>
1077
                      </dimer>
1078
               </channel>
1079
          </transport>
1080
          <transport name="electron">
1081
               <channel name="single">
1082
                   <J>-2.797473760331286E-003</J>
1083
1084
                   <e_A>-4.50318366770689</e_A>
                   <e_B>-4.53143397059021</e_B>
1085
1086
               </channel>
               <channel name="multi">
1087
                      <molecule name="A">
1088
                            <e_LUMOp0>-4.50318366770689</e_LUMOp0>
1089
1090
                       </molecule>
                      <molecule name="B">
1091
1092
                            <e_LUMOp0>-4.53143397059021</e_LUMOp0>
                       </molecule>
1093
                      <dimer name="integrals">
1094
                           <T_00>-2.797473760331286E-003</T_00>
```

### 3.8 State file

```
1103
sec:statefile
```

1106

All data structures are saved to the state.sql file in sqlite3 format, see http://www.sqlite.org/. They are available in form of tables in the state.sql file as can be seen by the command sqlite3 state.sql " .tables "

An example of such a table are molecules. The full table can be displayed using the command (similar for the other tables)

1109 sqlite3 state.sql " SELECT \* FROM molecules "

The meaning of all the entries in the table can be displayed by a command like

1111 sqlite3 state.sql " .SCHEMA molecules "

The first and second entry are integers for internal and regular id of the molecule and the third entry is the name. A single field from the table like the name of the molecule can be displayed by a command like

1115 sqlite3 state.sql " SELECT name FROM molecules "

Besides molecules, the following tables are stored in the state.sql:

1117 conjseq\_properties:

Conjugated segments are stored with id, name and x,y,z coordinates of the center of mass in nm.

1119 conjsegs:

Reorganization energies for charging or discharging a conjugated segment are stored together with the coulomb energy and any other user defined energy contribution (in eV) and occupation probabilities.

1123 pairs:

The pairs from the neighborlist are stored with the pair id, the id of the first and second segment, the rate from the first to the second , the rate from the second to the first (both in  $s^{-1}$ ) and the x,y,z coordinates in nm of the distance between the first and the second segment.

1127 pairintegrals:

Transfer integrals for all pairs are stored in the following way: The pair id , the number for counting possible different electronic overlaps (e.g if only the frontier orbitals are taken into account
this is always zero, while an effective value is stored in addition to the different overlaps of e.g.
HOMO-1 and HOMO-1 if more frontier orbitals are taken into account) and the integral in eV.

1132 pairproperties:

The outer sphere reorganization energy of all pairs is stored by an id, the pair id, a string lambda\_outer and the energy in eV.

1135 conjsegs:

Conjugated segments are saved in the following way: The id, the name, the type, the molecule id, the time frame, the x,y,z coordinates in nm and the occupation probability.

1138 conjseg\_properties:

Properties of the conjugated segments like reorganization energies for charging or discharging a charge unit or the coulomb contribution to the site energy are stored by: id, conjugated segment id, a string like lambda\_intra\_charging, lambda\_intra\_discharging or energy\_coulomb and a corresponding value in eV.

The tables rigidfrag\_properties, rigidfrags and frames offer information about rigid fragments and time frames including periodic boundary conditions.

The data in the state.sql file can also be modified by the user. Here is an example how to modify the transfer integral between the conjugated segments number one and two assuming that they are in the neighborlist. Their pair id can be found by the command

3.8. STATE FILE

```
pair_ID='sqlite3state.sql"SELECT _id FROM pairs WHERE conjseq1=1 AND conjseq2=2"'
    The old value of the transfer integral can be deleted using
1149
    sqlite3 state.sql "DELETE FROM pair_integrals WHERE pair=$pair_ID"
    Finally the new transfer integral J can be written to the state.sql file by the command
    sqlite3 state.sql "INSERT INTO pair_integrals (pair,num, J) VALUES ($pair_ID,0,$J)"
1152
    Here the num=0 indicates that only the effective transfer integrals is written to the file, while other
1153
    values of num would correspond to overlap between other orbitals than the frontier orbitals.
1154
    In a similar way the coulomb contribution to the site energy of the first conjugated segment can
    be overwritten by first getting its id
1156
    c_ID='sqlite3 state.sql "SELECT _id from conjseg_properties where conjseg=1 AND
1157
    key =\"energy_coulomb\""
1158
    Then deleting the old value
1159
    sqlite3 state.sql "DELETE FROM from conjseq_properties WHERE _id=$c_ID"
1160
    Then the new coulomb energy E can be written to this id
1161
    sqlite3 state.sql "INSERT INTO conjseg_properties (_id,conjseg,key,value)
    VALUES ($c_ID,1,\"energy_coulomb\",$E)"
    Finally the resulting coulomb contribution to all conjugated segments can be displayed by
1164
    sqlite3 state.sql "SELECT * from conjseg_properties WHERE key=\"energy_coulomb\""
1165
```

# Chapter 4

# Reference

sec:reference

```
Programs
       4.1
   1169
       Programs execute specific tasks (calculators).
       4.1.1 ctp_map
   1171
prog:ctp_map
       Generates QM | MD topology
   1172
             -h [ --help ] display this help and exit
   1173
             -v [ --verbose ] be loud and noisy
   1174
             -t [ --topology ] arg topology
   1175
                [ --coordinates ] arg coordinates or trajectory
             -s [ --segments ] arg definition of segments and fragments
   1177
             -f [ --file ] arg state file
   1178
             --man output man-formatted manual pages
   1179
             --tex output tex-formatted manual pages
       4.1.2 ctp_dump
       Extracts information from the state file
             -h [ --help ] display this help and exit
   1183
             -v [ --verbose ] be loud and noisy
   1184
             -o [ --options ] arg calculator options
   1186
                [ --file ] arg sqlight state file, *.sql
             -i [ --first-frame ] arg (=1) start from this frame
   1187
             -n [ --nframes ] arg (=1) number of frames to process
   1188
             -t [ --nthreads ] arg (=1) number of threads to create
             -s [ --save ] arg (=1) whether or not to save changes to state file
   1190
             -e [ --extract ] arg List of extractors separated by ',' or ''
   1191
             -l [ --list ] Lists all available extractors
   1192
             -d [ --description ] arg Short description of an extractor
             --man output man-formatted manual pages
   1194
             --tex output tex-formatted manual pages
   1195
       4.1.3 ctp_tools
   1196
prog:ctp_tools
       Runs charge transport tools
   1197
             -h [ --help ] display this help and exit
   1198
             -v [ --verbose ] be loud and noisy
```

```
-t [ --nthreads ] arg (=1) number of threads to create
     1200
               -o [ --options ] arg calculator options
     120
               --man output man-formatted manual pages
               --tex output tex-formatted manual pages
               -e [ --execute ] arg List of tools separated by ',' or ''
     1204
               -l [ --list ] Lists all available tools
     1205
               -d [ --description ] arg Short description of a tool
     1206
          4.1.4 ctp_run
     1207
          Runs charge transport calculators
               -h [ --help ] display this help and exit
                     --verbose ] be loud and noisy
     1210
               -o [ --options ] arg calculator options
     1211
               -f [ --file ] arg sqlight state file, *.sql
     1212
               -i [ --first-frame ] arg (=1) start from this frame
     1213
               -n [ --nframes ] arg (=1) number of frames to process
     1214
                  [ --nthreads ] arg (=1) number of threads to create
     1215
                   [ --save ] arg (=1) whether or not to save changes to state file
                   [ --execute ] arg List of calculators separated by ',' or ''
     1217
               -1 [ --list ] Lists all available calculators
     1218
               -d [ --description ] arg Short description of a calculator
     1219
               --man output man-formatted manual pages
               --tex output tex-formatted manual pages
     122
          4.1.5 ctp_parallel
     1222
progecto paralle
          Runs job-based heavy-duty calculators
     1223
               -h [ --help ] display this help and exit
     1224
     1225
               -v [ --verbose ] be loud and noisy
               -o [ --options ] arg calculator options
     1226
               -f [ --file ] arg sqlite state file, *.sql
     1227
               -i [ --first-frame ] arg (=1) start from this frame
     1228
               -n [ --nframes ] arg (=1) number of frames to process
                   [ --nthreads ] arg (=1) number of threads to create
               -s [ --save ] arg (=1) whether or not to save changes to state file
     123
               -r [ --restart ] arg restart pattern: 'host(pc1:234) stat(FAILED)'
     1232
               -c [ --cache ] arg (=8) assigns jobs in blocks of this size
     1233
               -j [ --jobs ] arg (=run) task(s) to perform: input, run, import
     1234
               -m [ --maxjobs ] arg (=-1) maximum number of jobs to process (-1 = inf)
     1235
               -e [ --execute ] arg List of calculators separated by ',' or '
     1236
               -l [ --list ] Lists all available calculators
               -d [ --description ] arg Short description of a calculator
               --man output man-formatted manual pages
     1239
               --tex output tex-formatted manual pages
     1240
          4.1.6 moo_overlap
prog:moo_overlap
               -h [ --help ] display this help and exit
     1242
               -v [ --verbose ] be loud and noisy
               --man output man-formatted manual pages
     1244
               --tex output tex-formatted manual pages
     1245
               --conjseg arg xml file describing two conjugated segments
     1246
```

--pos1 arg position and orientation of molecule 1

```
    1248 --pos2 arg position and orientation of molecule 2
    1249 --pdb arg (=geometry.pdb) pdb file of two molecules
```

# 4.2 Calculators

#### sæf:calculato

1250

list:calc

1275

Calculator is a piece of code which computes specific system properties, such as site energies, transfer integrals, etc. ctp\_run, kmc\_run are wrapper programs which executes such calculators. The generic syntax is

```
ctp_run -e "calc1, calc2, ..." -o options.xml
```

File options.xml lists all options needed to run a specific calculator. The format of this file is explained in listing 4.1. A complete list of calculators is given in the calculators reference section.

#### Listing 4.1: A part of the options.xml file with options for the calculator\_name {1, 2} calculators.

```
1257
     <calculator_name1>
1258
                 <option1>value1</option1>
1259
                <option2>value2</option2>
1260
1261
     </calculator_name1>
1262
1263
1264
     <calculator_name2>
                <option1>value1</option1>
1265
                <option2>value2</option2>
1266
1267
     </calculator_name2>
1268
1298
```

A list of all calculators and their short descriptions can be obtain using

```
1272 ctp_run --list
```

A detailed description of all options of a specific calculator(s) is available via

```
ctp_run --desc calc1, calc2,...
```

# 4.2.1 coupling

Electronic couplings from log and orbital files (GAUSSAIN, TURBOMOLE, NWChem)

option	default	unit	description
package			First-principles package
output	coupling.out.xr	ď	Output file
degeneracy	0	eV	Criterium for the degeneracy of two levels
moleculeA			
log	A.log		Log file of molecule A
orbitals	A.orb		Orbitals file
levels	3		Output HOMO,, HOMO-levels; LUMO,, LUMO+levels
trim	2		
moleculeB			
log	B.log		Log file of molecule B
orbitals	B.orb		Orbitals file
levels	3		Output HOMO,, HOMO-levels; LUMO,, LUMO+levels
trim	2		
dimerAB			
log	AB.log		Log file of dimer AB
orbitals	A.orb		Orbitals file

Return to the description of coupling.

# 4.2.2

# log2mps

Generates an mps-file (with polar-site definitions) from a QM log-file

option	default	unit	description
package			QM package
logfile			Log-file generated by QM package, with population/esp-fit data

Return to the description of log2mps.

# 4.2.3 molpol

1281

Molecular polarizability calculator (and optimizer)

option	default	unit	description
mpsfiles			
input			mps input file
output			mps output file
polar			xml file with infos on polarizability tensor
induction			
expdamp			Thole sharpness parameter
wSOR			mixing factor for convergence
maxiter			maximum number of iterations
tolerance			rel. tolerance for induced moments
target			
optimize			if 'true', refine atomic polarizabilities to match molecular polarizable volume specified in target.molpol
molpol			target polarizability tensor in format $xx xy xz yy yz zz$ (this should be in the eigen-frame, hence $xy = xz = yz = 0$ ), if optimize=true the associated polarizable volume will be matched iteratively and the resulting set of polar sites written to mpsfiles.output
tolerance			relative tolerance when optimizing the polarizable volume

Return to the description of molpol.

# 4.2.4 pdb2map



Converts MD + QM files to VOTCA mapping. Combinations: pdb+xyz,gro+xyz,pdb

option	default	unit	description
pdb	conf.pdb		Input pdb file
gro	conf.gro		Input gro file
xyz	conf.xyz		Input xyz file
xml	conf.xml		Resulting xml file

Return to the description of pdb2map.

# 4.2.5 pdb2top

1288 Generates fake Gromacs topology file .top

option	default	unit	description
num	1		Num of mols in the box
pdb	conf.pdb		Input pdb file
gro	conf.gro		Input gro file

Return to the description of pdb2top.

#### 4.2.6 ptopreader

calc:ptopreade 1291

1290

Reads binary .ptop-files (serialized from ewdbgpol) and processes them into something readable

option	default	unit	description
ptop_file			Binary archive .ptop-file

Return to the description of ptopreader.

# 4.2.7 eanalyze

eanalyze: 1294

Histogram and correlation function of site energies and pair energy differences

option	default	unit	description
resolution_sites		eV	Bin size for site energy histogram
resolution_pairs		eV	Bin size for pair energy histogram
resolution_space		eV	Bin size for site energy correlation
states			?

Return to the description of eanalyze.

# 4.2.8 eimport

1296 alc:eimport

Imports site energies from the output file of emultipole and writes them to the state file

Return to the description of eimport.

#### 1299 calc:einterna 1300

# 4.2.9 einternal

Reads in site and reorganosation energies and writes them to the state file

option	default	unit	description
energiesXML			XML input file with vacuum site, reorganization (charging, discharging) energies

Return to the description of einternal.

# 4.2.10 emultipole

1302 calc:emultipol

Evaluates polarization contribution based on the Thole model

option	default	unit	description
multipoles			Polar Site Definitions in GDMA punch-file format
control			Control options for induction computation
induce	1		Enter '1' / '0' to toggle induction on / off
first			First segment for which to compute site energies
last			Last segment for which to compute site energies
output			File to write site energies to. Site energies are also stored
output			in the state file
check			Check mapping of polar sites to fragment
tholeparam			Thole parameters required for charge-smearing
cutoff		nm	Cut-off beyond which all interactions are neglected

cutoff2	nm	Cut-off beyond which polarization is neglected
expdamp		Damping exponent used in exponential damping function
scaling		1-n interaction scaling, currently not in use
esp		Control options for potential calculation
calcESP		Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
cube		
grid		XYZ file specifying grid points for potential evaluation
output		File to write grid-point potential to
esf		Control options for field calculation
calcESF		Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
grid		XYZ file specifying grid points for field evaluation
output		File to write grid-point field to
alphamol		Control options for molecular-polarizability calculation
calcAlpha		Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
output		File to write polarizability tensor in global frame and in diagonal form to
convparam		Convergence parameters for self-consistent field calculation
wSOR_N		Mixing factor for successive overrelaxation of neutral system, usually between 0.3 and 0.5
wSOR_C		Mixing factor for successive overrelaxation of charged system, usually between 0.3 and 0.5
tolerance		Convergence criterion, fulfilled if relative change smaller than tolerance
maxiter		Maximum number of iterations in the convergence loop

Return to the description of emultipole.

# 4.2.11 eoutersphere

 $\begin{array}{cc} \mbox{\tiny {\bf calc:eoutersphere}} \\ \mbox{\tiny {\bf 1306}} \end{array} \ \ Evaluates \ outersphere \ reorganization \ energy \\$ 

option	default	unit	description
multipoles			XML allocation polar sites
method			Type of the method: **constant** - all pairs have value **lambda**. **spheres** - molecules are treated as spheres with radii **radius** and Pekar factor **pekar**. **dielectric** - with Pekar factor **pekar** and partial charges from resulting dielectric fields
lambdaconst		eV	The value for all pairs in the **constant** method
pekar			Pekar factor used for methods **spheres** and **dielectric**
segment			
type			
radius			
segment			
type			
radius			
cutoff		nm	Cutoff radius in between pair and the exterior molecule. Can be used in **spheres** and **dielectric**

<sup>1307</sup> Return to the description of eoutersphere.

# 4.2.12 ewdbgpol

lc:ewdbgpo

Calculates background polarisation needed for ewald calc

option	default	unit	description
multipoles			
control			
mps_table			
pdb_check			
coulombmethod			
method			
cutoff			
shape			
polarmethod			
method			
induce			
cutoff			
convergence			
energy			
kfactor			
rfactor			

Return to the description of ewdbgpol.

# 4.2.13 ianalyze

1311 alc:ianalyz

Evaluates a histogram of a logarithm of squared couplings

option	default	unit	description
resolution_logJ2			Bin size of histogram log(J2)
states			States for which to calculate the histogram. Example: 1 -1

Return to the description of ianalyze.

# 1314 4.2.14 iimport

alc:iimpoi

Imports electronic couplings from xml of ctp-dipro using folders of pairdump

option	default	unit	description
idft jobs file			idft jobs file

Return to the description of iimport.

#### 4.2.15 izindo

1317 lc:izindo

Semiempirical electronic coupling elements for all neighbor list pairs

option	default	unit	description
orbitalsXML			File with paths to .orb files

Return to the description of izindo.

# 4.2.16 jobwriter

alc:jobwrite

Writes list of jobs for a parallel execusion

option	default	unit	description
--------	---------	------	-------------

keys			job type
states	n e h		hole, electron, nuetral: mps file is required
single_id			Segment ID as argument for mps.single
kmc_cutoff		nm	Pair-interaction cut-off as argument for mps.kmc

Return to the description of jobwriter.

# 4.2.17 neighborlist

calc:neighborlist

1323

Constructs a list of neighboring conjugated segments

option	default	unit	description
constant	0.5	nm	If provided, this value is used for all segment types
segments			A pair of segment types
type			Types of two segments. For types A and B this can be A A, A B or B B
cutoff		nm	Cutoff radius for centers of mass of rigid fragments

Return to the description of neighborlist.

# 4.2.18 pairdump

alc:pairdump 1327

Coordinates of molecules and pairs from the neighbor list

option	default	unit	description
molecules			If **true** outputs single molecules, otherwise only pairs

1328 Return to the description of pairdump.

# 4.2.19 profile

1329 calc:profile

Density and site energy profiles

option	default	unit	description
axis			Axis along which to calculate density and energy profiles
direction	0 0 1		Axis direction
min		nm	Minimal projected position for manual binning
max		nm	Maximal projected position for manual binning
bin	0.1	nm	Spatial resolution of the profile
auto	1		'0' for manual binning using min and max, '1' for automated
particles			
type	segments		What centers of mass to use: 'segments' or 'atoms'
first	1		ID of the first segment
last	-1		ID of the last segment, -1 is the list end
output			
density	density.dat		Density profile file
energy	energy.dat		Energy profile file

Return to the description of profile.

#### **4.2.20** rates

1332 alc:rate

Hopping rates using classical or semi-classical expression

option def	fault unit	description
------------	------------	-------------

field			Field in x y z direction
temperature		K	Temperature for rates
method			Method chosen to compute rates. Can either be **marcus** or **jortner**. The first is the high temperature limit of Marcus theory, the second is the rate proposed by Jortner and Bixon
nmaxvib	20		If the method of choice is **jortner**, the maximal number of excited vibrations on the molecules has to be specified as an integer for the summation
omegavib	0.2	eV	If the method of choice is **jortner**, the vibration frequency of the quantum mode has to be given in units of eV. The default value is close to the CC bond-stretch at 0.2eV

Return to the description of rates.

# 1335 **4.2.21** sandbox

1338

1341

1336 Sandbox to test ctp classes

option	default	unit	description
ID			Not in use

1337 Return to the description of sandbox.

#### 4.2.22 stateserver

Export SQLite file to human readable format

option	default	unit	description
out			Output file name
pdb			PDB coordinate file name
keys			Sections to write to readable format (topology, segments, pairs, coordinates)

Return to the description of stateserver.

# 4.2.23 tdump

Coarse-grained and back-mapped (using rigid fragments) trajectories

option	default	unit	description
md	MD.pdb		Name of the coarse-grained trajectory
qm	QM.pdb		Name of the trajectory with back-substituted rigid frag- ments
frames	1		Number of frames to output

1343 Return to the description of tdump.

#### 4.2.24 vaverage

Computes site-centered velocity averages from site occupancies

option	default	unit	description
carriers			Carrier types for which to compute velocity averages
tabulate			Tabulate 'atoms' or 'segments'

Return to the description of vaverage.

# 4.2.25 zmultipole

Evaluates polarization contribution based on the Thole model

option	default	unit	description
multipoles			Polar Site Definitions in GDMA punch-file format
control			Control options for induction computation
induce	1		Enter '1' / '0' to toggle induction on / off
first			First segment for which to compute site energies
last			Last segment for which to compute site energies
output			File to write site energies to. Site energies are also stored in the state file
check			Check mapping of polar sites to fragment
tholeparam			Thole parameters required for charge-smearing
cutoff		nm	Cut-off beyond which all interactions are neglected
cutoff2		nm	Cut-off beyond which polarization is neglected
expdamp			Damping exponent used in exponential damping function
scaling			1-n interaction scaling, currently not in use
esp			Control options for potential calculation
calcESP			Enter '1' / '0' to toggle on / off. If '1', site energies will not be evaluated
cube			
grid			XYZ file specifying grid points for potential evaluation
output			File to write grid-point potential to
esf			Control options for field calculation
1 PCP			Enter '1' / '0' to toggle on / off. If '1', site energies will
calcESF			not be evaluated
grid			XYZ file specifying grid points for field evaluation
output			File to write grid-point field to
alphamol			Control options for molecular-polarizability calculation
•			Enter '1' / '0' to toggle on / off. If '1', site energies will
calcAlpha			not be evaluated
output			File to write polarizability tensor in global frame and in diagonal form to
convparam			Convergence parameters for self-consistent field calculation
wSOR_N			Mixing factor for successive overrelaxation of neutral system, usually between 0.3 and 0.5
wSOR_C			Mixing factor for successive overrelaxation of charged system, usually between 0.3 and 0.5
tolerance			Convergence criterion, fulfilled if relative change smaller than tolerance
maxiter			Maximum number of iterations in the convergence loop

Return to the description of zmultipole.

# 4.2.26 edft

1350

<sup>1351</sup> A wrapper for first principles based single site calculations

option	default	unit	description
job			Job options
tasks	input,run,parse		What to run
store	orbitals		What to store

Return to the description of edft.

# 4.2.27 idft

1353

1356

Projection method for electronic couplings. Requires edft otput

option	default	unit	description
tasks	input,run,parse,		What to do
store	orbitals,overlap		What to store
degeneracy	0	eV	Criterium for the degeneracy of two levels
levels	3		Output between HOMO,, HOMO-levels; LUMO,, LUMO+levels
trim	2		Use trim*occupied of virtual orbitals

Return to the description of idft.

# 4.2.28 pewald3d

Evaluates site energies in a periodic setting

option	default	unit	description
jobcontrol	delaar	dilit	description
job_file			
multipoles			
mapping			
mps_table			
pdb_check			
coulombmethod			
method			
cutoff			
shape			
polarmethod method			
induce			
cutoff			
tasks			
calculate_fields			
polarize_fg			
evaluate_energy			
coarsegrain			
cg_background			
cg_foreground			
cg_radius			
cg_anisotropic			
convergence			
energy			
kfactor			
rfactor			

 $^{1358}$  Return to the description of pewald3d.

# 4.2.29 qmmm

1360 (

QM/MM with the Thole MM model

option	default	unit	description
op tron	or creation to	011111	description

control			
pdb_check			PDB file of polar sites
write_chk	dipoles.xyz		XYZ file with dipoles split onto point charges
format_chk	xyz		format, gaussian or xyz
split_dpl	1		'0' do not split dipoles onto point charges, '1' do split
dpl_spacing	1e-3	nm	Spacing to be used when splitting dipole onto point charges: $d = q * a$
qmpackage			
package			QM package to use for the QM region
gwbse			Specify if GW/BSE excited state calculation ist needed
gwbse_options			GW/BSE options file
state			Number of excited state, which is to be calculated
type			Character of the excited state to be calculated
			Filter with which to find the excited state after each calcu-
filter			lation
oscilla-			Oscillator strength filter, only states with higher oscillator
tor_strength			strength are considered
_			Charge transfer filter , only states with charge transfer
charge_transfer			above threshold are consdered
qmmmconvg			convergence criteria for the QM/MM
dR	0.001	nm	RMS of coordinates
dQ	0.001	e	RMS of charges
dE_QM	0.0001	eV	Energy change of the QM region
dE_MM	0.0001	eV	Energy change of the MM region
max iter	10		Number of iterations
coulombmethod			Options for the MM embedding
method	cut-off		Method for evaluation of electrostatics
cutoff1			Cut-off for the polarizable MM1 shell
cutoff2			Cut-off for the static MM2 shell
tholemodel			Parameters for teh Thole model
induce			'1' - induce '0' - no induction
			'1' - include mutual interaction of induced dipoles in the
induce_intra_pair			QM region. '0' - do not
exp_damp	0.39		Sharpness parameter
scaling	0.05		Bond scaling factors
Seming			Convergence parameters for the MM1 (polarizable) re-
convergence			gion
			Mixing factor for the succesive overrelaxation algorithm
wSOR_N			for a neutral QM region
			Mixing factor for the succesive overrelaxation algorithm
wSOR_C			for a charged QM region
			Maximal number of iterations to converge induced
max_iter	512		dipoles
tolerance			Maximum RMS change allowed in induced dipoles
tolerance		1	waxiii ani kwis change anowed in induced dipoles

Return to the description of qmmm.

# 4.2.30 xqmultipole

Electrostatic interaction and induction energy of charged molecular clusters

option	default	unit	description
multipoles			Polar-site mapping definition
control			
job_file			Job file
emp_file			Polar-background definition, allocation of mps-files to segments

pdb_check		Whether or not to output a pdb-file of the mapped polar sites
format_chk		Format for check-file: 'xyz' or 'gaussian'
split_dpl		Split dipoles onto point charges in check-file
dpl_spacing	nm	Spacing between point charges for check-file output
coulombmethod		
method		Currently only cut-off supported
cutoff1	nm	Full-interaction radius cut-off
cutoff2	nm	Radius of electrostatic buffer
tholemodel		
induce		Induce - or not
induce_intra_pair		Induce mutually within the charged cluster
exp_damp		Thole sharpness parameter
scaling		Bond scaling parameters, currently not used
convergence		
wSOR_N		SOR mixing factor for overall neutral clusters
wSOR_C		SOR mixing factor for overall charged clusters
max_iter		Maximum number of iterations
tolerance		Relative tolerance as convergence criterion

Return to the description of xqmultipole.

# 4.2.31 energy2xml

calc:energy2xml

Write out energies from SQL file

	option	default	unit	description
--	--------	---------	------	-------------

Return to the description of energy2xml.

# 4.2.32 integrals2xml

calc:integrals2xm 1369

Write out transfer integrals from SQL file

option default unit description	
---------------------------------	--

Return to the description of integrals2xml.

# 4.2.33 occupations2xml

calc:occupations2xm

Write out site occupation probabilities from SQL file

option	default	unit	description

Return to the description of occupations2xml.

# 4.2.34 pairs2xml

calc:pairs2xn

1374

Write out neighbourlist from SQL file

option	default	unit	description
--------	---------	------	-------------

Return to the description of pairs2xml.

#### 1377 calc:rates2xm

1378

# 4.2.35 rates2xml

Write out charge transfer rates from SQL file

option | default | unit | description

1379 Return to the description of rates2xml.

# 4.2.36 segments2xml

calc:segments2xm 1381

1380

Write out segment data from SQL file

option | default | unit | description

Return to the description of segments2xml.

# 4.2.37 trajectory2pdb

1384

1383

Generate PDB files for the mapped MD/QM topology

option default unit description

Return to the description of trajectory2pdb.

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