



Robotic Systems I

Lecture 3: Orientation Representations and Lie Groups Introduction

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How can we represent orientations?

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- Rotation Matrices

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- Rotation Matrices
- Euler Angles
- Axis-Angle
- Quaternions

Axis-Angle:

$$\boldsymbol{\theta} = \mathbf{r}\theta$$

where $\mathbf{r} \in \mathbb{R}^3$ is the “axis of rotation” unit vector, and $\theta \in \mathbb{R}$ is the rotation angle.

- **Singularity** at $\theta = \pm\pi$
- $\boldsymbol{\theta}_{wb} = (\log \mathbf{R}_{wb})^\vee$
- $\mathbf{R}_{wb} = \exp(\hat{\boldsymbol{\theta}}_{wb})$

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- $\mathbf{R}_{wb} = \exp(\hat{\boldsymbol{\theta}}_{wb})$
- We can compute the above more efficiently than regular exp/log matrix operations

Euler Angles:

$$\begin{aligned} R &= R_1 R_2 R_3 \\ &= \underbrace{R_{\alpha_1}(\theta_1)}_{R_1} \underbrace{R_{\alpha_2}(\theta_2)}_{R_2} \underbrace{R_{\alpha_3}(\theta_3)}_{R_3} \end{aligned}$$

where α_i are the axes of rotation and θ_i the corresponding angles of rotation.

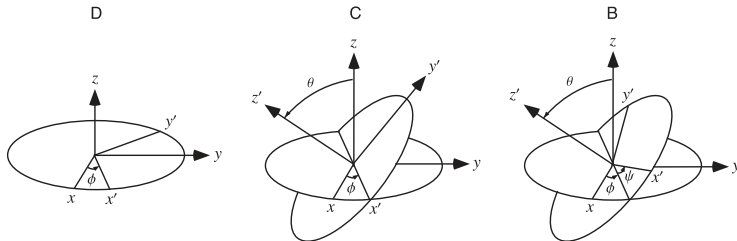


Figure: Taken from WolframAlpha

Euler Angles (2)

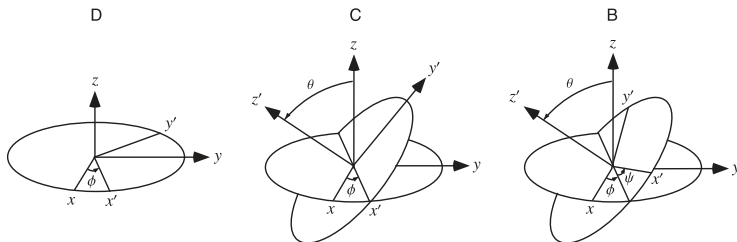


Figure: Taken from WolframAlpha

- There are many conventions: ZYX (yaw-pitch-roll) and XYZ (roll-pitch-yaw) the most popular ones
 - $\theta_1 \in [-\pi, \pi], \theta_2 \in [-\frac{\pi}{2}, \frac{\pi}{2}], \theta_3 \in [-\pi, \pi]$
 - **Singularity** at $\theta = \pm \frac{\pi}{2}$
- We can convert from/to rotation matrices

Axis-Angle:

$$\boldsymbol{\theta} = \boldsymbol{r}\theta$$

$$\boldsymbol{R} = \exp(\boldsymbol{\theta}_{wb}^\wedge)$$

$$= \boldsymbol{I} + \sin \theta \hat{\boldsymbol{r}} + (1 - \cos \theta) \hat{\boldsymbol{r}}^2, \text{ (Rodrigues formula)} \Rightarrow$$

$$\dot{\boldsymbol{R}} = \dot{\theta} \cos \theta \hat{\boldsymbol{r}} + \sin \theta \dot{\hat{\boldsymbol{r}}} + \dot{\theta} \sin \theta \hat{\boldsymbol{r}}^2 + (1 - \cos \theta)(\dot{\hat{\boldsymbol{r}}} \hat{\boldsymbol{r}} + \hat{\boldsymbol{r}} \dot{\hat{\boldsymbol{r}}}) \Rightarrow$$

...

$$\Rightarrow \boldsymbol{\omega}_w = \dot{\theta} \boldsymbol{r} + \sin \theta \dot{\boldsymbol{r}} + \hat{\boldsymbol{r}} \dot{\boldsymbol{r}} (1 - \cos \theta)$$

$$\Rightarrow \boldsymbol{\omega}_b = \dot{\theta} \boldsymbol{r} + \sin \theta \dot{\boldsymbol{r}} - \hat{\boldsymbol{r}} \dot{\boldsymbol{r}} (1 - \cos \theta)$$

Euler Angles and Angular Velocity

ZYX Euler Angles:

$$\omega_b = \begin{bmatrix} -\sin \theta_2 & 0 & 1 \\ \cos \theta_2 \sin \theta_3 & \cos \theta_3 & 0 \\ \cos \theta_2 \cos \theta_3 & -\sin \theta_3 & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \dot{\theta}_3 \end{bmatrix}$$

$$\omega_w = \begin{bmatrix} 0 & -\sin \theta_1 & \cos \theta_2 \cos \theta_1 \\ 0 & \cos \theta_1 & \cos \theta_2 \sin \theta_1 \\ 1 & 0 & -\sin \theta_2 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \dot{\theta}_3 \end{bmatrix}$$

XYZ Euler Angles:

⋮

⋮

Let's start with the **Axis-Angle**:

$$\boldsymbol{\theta} = \mathbf{r}\theta$$

A **unit quaternion** describing the same orientation/rotation is given by:

$$\mathbf{q} = \begin{bmatrix} \cos \frac{\theta}{2} \\ \mathbf{r} \sin \frac{\theta}{2} \end{bmatrix} \in \mathbb{R}^4$$

- $\cos \frac{\theta}{2}$ is the **scalar** part, and $\mathbf{r} \sin \frac{\theta}{2}$ is the **vector** part
- For \mathbf{q} to be a valid orientation/rotation, we also need $\|\mathbf{q}\| = 1$
- \mathbf{q} and $-\mathbf{q}$ describe the **same orientation**, but *different rotations!* (double cover)

Quaternion Math

- \mathbf{q}_{wb} is a rotation from “body frame” to “world frame”
- $\mathbf{q}_{wa} = \mathbf{q}_{wb} \odot \mathbf{q}_{ba}$

$$\mathbf{q}_{wa} = \begin{bmatrix} s_{wb} \\ \mathbf{v}_{wb} \end{bmatrix} \odot \begin{bmatrix} s_{ba} \\ \mathbf{v}_{ba} \end{bmatrix} = \begin{bmatrix} s_{wb}s_{ba} - \mathbf{v}_{wb}^T \mathbf{v}_{ba} \\ s_{wb}\mathbf{v}_{ba} + s_{ba}\mathbf{v}_{wb} + \mathbf{v}_{wb} \times \mathbf{v}_{ba} \end{bmatrix}$$

- **Quaternion Conjugate (Inverse):**

$$\mathbf{q}_{wb}^\dagger = \begin{bmatrix} s_{wb} \\ -\mathbf{v}_{wb} \end{bmatrix}$$

- $\mathbf{q}_{wb}^\dagger \odot \mathbf{q}_{wb} = \mathbf{q}_0, \mathbf{q}_0 = \begin{bmatrix} 1 \\ \mathbf{0} \end{bmatrix}$
- $\begin{bmatrix} 0 \\ \mathbf{x}_w \end{bmatrix} = \mathbf{q}_{wb} \odot \begin{bmatrix} 0 \\ \mathbf{x}_b \end{bmatrix} \odot \mathbf{q}_{wb}^\dagger, \text{ (rotating a vector/point)}$

Time derivative of Quaternion:

$$\begin{aligned}\dot{\mathbf{q}} &= \frac{1}{2} \mathbf{q} \odot \begin{bmatrix} 0 \\ \boldsymbol{\omega}_b \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 0 \\ \boldsymbol{\omega}_w \end{bmatrix} \odot \mathbf{q}\end{aligned}$$

- We can now integrate quaternions

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- **Beware!** Integrating naively with $\dot{\mathbf{q}}$ (e.g. with Euler) results in a non unit-norm quaternion!

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- We can now integrate quaternions
- **Beware!** Integrating naively with $\dot{\mathbf{q}}$ (e.g. with Euler) results in a non unit-norm quaternion!
- **Lots of “custom” functions to use the quaternions!**

$$\mathbf{q}_{wa} = \begin{bmatrix} s_{wb} \\ \mathbf{v}_{wb} \end{bmatrix} \odot \begin{bmatrix} s_{ba} \\ \mathbf{v}_{ba} \end{bmatrix} = \begin{bmatrix} s_{wb}s_{ba} - \mathbf{v}_{wb}^T \mathbf{v}_{ba} \\ s_{wb}\mathbf{v}_{ba} + s_{ba}\mathbf{v}_{wb} + \mathbf{v}_{wb} \times \mathbf{v}_{ba} \end{bmatrix}$$

How can we write the above such that we do it via matrix multiplication?

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$$\begin{aligned} \mathbf{q}_{wa} &= \begin{bmatrix} s_{wb} \\ \mathbf{v}_{wb} \end{bmatrix} \odot \begin{bmatrix} s_{ba} \\ \mathbf{v}_{ba} \end{bmatrix} = \underbrace{\begin{bmatrix} s_{wb} & -\mathbf{v}_{wb}^T \\ \mathbf{v}_{wb} & s_{wb}\mathbf{I} + \hat{\mathbf{v}}_{wb} \end{bmatrix}}_{L(\mathbf{q}_{wb})} \begin{bmatrix} s_{ba} \\ \mathbf{v}_{ba} \end{bmatrix} \\ &= \underbrace{\begin{bmatrix} s_{ba} & -\mathbf{v}_{ba}^T \\ \mathbf{v}_{ba} & s_{ba}\mathbf{I} - \hat{\mathbf{v}}_{ba} \end{bmatrix}}_{R(\mathbf{q}_{ba})} \begin{bmatrix} s_{wb} \\ \mathbf{v}_{wb} \end{bmatrix} \end{aligned}$$

A few interesting relations:

$$\begin{aligned}\mathbf{q}_{wb} \odot \mathbf{q}_{ba} &= L(\mathbf{q}_{wb})\mathbf{q}_{ba} \\ &= R(\mathbf{q}_{ba})\mathbf{q}_{wb}\end{aligned}$$

$$\begin{aligned}L(\mathbf{q}^\dagger) &= L(\mathbf{q})^T = L(\mathbf{q})^{-1} \\ R(\mathbf{q}^\dagger) &= R(\mathbf{q})^T = R(\mathbf{q})^{-1}\end{aligned}$$

Quaternion Magic (3)

We can use the same logic in all quaternion operations:

$$\blacksquare \mathbf{q}^\dagger = \begin{bmatrix} s \\ -\mathbf{v} \end{bmatrix} = \underbrace{\mathbf{T} \mathbf{q}}_{\mathbf{T}} = \begin{bmatrix} 1 & 0 \\ 0 & -\mathbf{I} \end{bmatrix} \begin{bmatrix} s \\ \mathbf{v} \end{bmatrix}$$

■ Rotate a vector/point:

$$\begin{aligned} \begin{bmatrix} 0 \\ \mathbf{x}_w \end{bmatrix} &= \mathbf{q}_{wb} \odot \begin{bmatrix} 0 \\ \mathbf{x}_b \end{bmatrix} \odot \mathbf{q}_{wb}^\dagger \\ \underbrace{\begin{bmatrix} 0 \\ \mathbf{I} \end{bmatrix}}_{\mathbf{H} \in \mathbb{R}^{4 \times 3}} \mathbf{x}_w &= \mathbf{q}_{wb} \odot \mathbf{H} \mathbf{x}_b \odot \mathbf{q}_{wb}^\dagger \\ \mathbf{H} \mathbf{x}_w &= L(\mathbf{q}_{wb}) (R(\mathbf{q}_{wb})^T \mathbf{H} \mathbf{x}_b) \\ \mathbf{x}_w &= \underbrace{\mathbf{H}^T L(\mathbf{q}_{wb}) R(\mathbf{q}_{wb})^T \mathbf{H}}_{\mathbf{R}_{wb}} \mathbf{x}_b \end{aligned}$$

We can use the same logic in all quaternion operations:

$$\begin{aligned} & \vdots \\ \dot{\mathbf{q}} &= \frac{1}{2} \mathbf{q} \odot \begin{bmatrix} 0 \\ \boldsymbol{\omega}_b \end{bmatrix} \\ &= \frac{1}{2} L(\mathbf{q}) \mathbf{H} \boldsymbol{\omega}_b \end{aligned}$$

Rodrigues Parameters:¹

$$\boldsymbol{\phi} = \boldsymbol{r} \tan \frac{\theta}{2} \in \mathbb{R}^3$$


where $\boldsymbol{r} \in \mathbb{R}^3$ is the “axis of rotation” unit vector, and $\theta \in \mathbb{R}$ is the rotation angle. $\boldsymbol{\phi}$ is referred to as “Rodrigues” or “Gibbs” vector.

Quaternions and Rodrigues Parameters:


$$\boldsymbol{q} = \varphi(\boldsymbol{\phi}) = \frac{1}{\sqrt{1 + \boldsymbol{\phi}^T \boldsymbol{\phi}}} \begin{pmatrix} 1 \\ \boldsymbol{\phi} \end{pmatrix}$$
$$\boldsymbol{\phi} = \varphi^{-1}(\boldsymbol{q}) = \frac{\boldsymbol{v}}{s}$$

¹https://en.wikipedia.org/wiki/Rotation_formalisms_in_three_dimensions

Orientation Representations - Code Example

 jupyter 2-orientations Last Checkpoint: 4 minutes ago

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 Code ▾

```
theta = 1.  
axis = np.array([[0., 0., 1.]]).T  
  
# quaternion  
q = np.block([[np.cos(theta/2.)], [axis * np.sin(theta/2.)]])  
  
# rotation matrix  
R = quat_to_rotation_matrix(q)  
  
omega_b = np.array([[1., 0., 0.]]).T
```

- The **Special Orthogonal Group**, representing orientations and rotations, is simply the set of valid rotation matrices:

$$SO(3) = \{\mathbf{R} \in \mathbb{R}^{3 \times 3} \mid \mathbf{R}\mathbf{R}^T = \mathbf{I}, \det(\mathbf{R}) = 1\}$$

- The set of all 3×3 matrices is a vectorspace. What does it mean?

- The **Special Orthogonal Group**, representing orientations and rotations, is simply the set of valid rotation matrices:

$$SO(3) = \{\mathbf{R} \in \mathbb{R}^{3 \times 3} \mid \mathbf{R}\mathbf{R}^T = \mathbf{I}, \det(\mathbf{R}) = 1\}$$

- The set of all 3×3 matrices is a vectorspace. What does it mean?
- $SO(3)$ is not:
 - 1 The zero matrix $\mathbf{0}$ is not part of the space!
 - 2 $\mathbf{R}_1, \mathbf{R}_2 \in SO(3) \not\Rightarrow \mathbf{R}_1 + \mathbf{R}_2 \in SO(3)$

What is a Group? A set non-empty \mathcal{G} along with a binary operator \circ that satisfies four properties:

- 1 Closure:** $\forall \mathcal{X}, \mathcal{Y} \in \mathcal{G}, \mathcal{X} \circ \mathcal{Y} \in \mathcal{G}$
- 2 Associativity:** $\forall \mathcal{X}, \mathcal{Y}, \mathcal{Z} \in \mathcal{G}, (\mathcal{X} \circ \mathcal{Y}) \circ \mathcal{Z} = \mathcal{X} \circ (\mathcal{Y} \circ \mathcal{Z})$
- 3 Identity:** $\exists \varepsilon \in \mathcal{G}, \forall \mathcal{X} \in \mathcal{G}, \varepsilon \circ \mathcal{X} = \mathcal{X} \circ \varepsilon = \mathcal{X}$
- 4 Inverse:** $\forall \mathcal{X} \in \mathcal{G}, \exists \mathcal{X}^{-1} \in \mathcal{G}, \mathcal{X} \circ \mathcal{X}^{-1} = \mathcal{X}^{-1} \circ \mathcal{X} = \varepsilon$

The operator \circ is not necessarily commutative!

What is a Manifold? A Manifold, \mathcal{M} , is a topological space that locally resembles Euclidean space, \mathbb{R}^n , near each point \mathcal{X} :

- The neighborhood of \mathcal{X} lies on the tangent space $T_{\mathcal{X}}\mathcal{M}$ of the manifold \mathcal{M} at point \mathcal{X}
- The neighborhood of \mathcal{X} and the corresponding tangent space are homeomorphic to an open subset of \mathbb{R}^n : $T_{\mathcal{X}}\mathcal{M} \cong \mathbb{R}^n$
- A smooth manifold is infinitely differentiable at each point

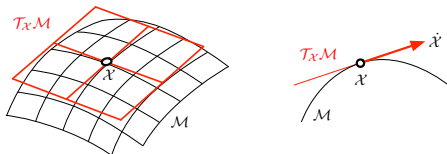


Figure: Sola, Joan, Jeremie Deray, and Dinesh Atchuthan. "A micro lie theory for state estimation in robotics." *arXiv:1812.01537* (2018).

What is a Lie Group? A Lie Group is a group that is also a smooth manifold:

- The group operations are smooth: $\mathcal{G} \times \mathcal{G} \rightarrow \mathcal{G}$,
 $(\mathcal{X}, \mathcal{Y}) \mapsto \mathcal{X} \circ \mathcal{Y}$
- The inverse operation is smooth: $\mathcal{G} \rightarrow \mathcal{G}$, $\mathcal{X} \mapsto \mathcal{X}^{-1}$
- The identity element is smooth: $\mathcal{G} \rightarrow \mathcal{G}$, $\mathcal{X} \mapsto \varepsilon$

Why are Lie Groups important? They are the natural setting for the study of continuous symmetries of differential equations.

$SO(3)$ is a Matrix Lie Group

$SO(3)$ forms a *matrix Lie group* (\circ is the regular matrix multiplication, ε is the identity matrix):

- $R_1, R_2 \in SO(3) \Rightarrow R_1 R_2 \in SO(3)$
- $(R_1 R_2) R_3 = R_1 (R_2 R_3) = R_1 R_2 R_3$
- $R, I \in SO(3)$ and $RI = IR = R$
- $R^{-1} = R^T \in SO(3)$

Given a lie group \mathcal{M} and a set \mathcal{V} , a *Lie group action*, \cdot , is a map $\cdot : \mathcal{M} \times \mathcal{V} \rightarrow \mathcal{V}$ that satisfies:

1 Identity: $\varepsilon \cdot \mathbf{v} = \mathbf{v}$

2 Compatibility: $\mathcal{X} \cdot (\mathcal{Y} \cdot \mathbf{v}) = (\mathcal{X} \circ \mathcal{Y}) \cdot \mathbf{v}$

where $\mathcal{X}, \mathcal{Y} \in \mathcal{M}$ and $\mathbf{v} \in \mathcal{V}$

Example: The rotation of a vector \mathbf{v} by a rotation matrix \mathbf{R} :

$$\mathbf{R} \cdot \mathbf{v} = \mathbf{R}\mathbf{v}$$

Tangent Spaces of a Lie Group

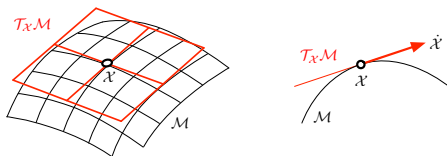


Figure: Sola, Joan, Jeremie Deray, and Dinesh Atchuthan. "A micro lie theory for state estimation in robotics." *arXiv:1812.01537* (2018).

Given a lie group \mathcal{M} and a point $\mathcal{X} \in \mathcal{M}$, the *tangent space* at \mathcal{X} is written as $T_{\mathcal{X}}\mathcal{M}$.

- Let's assume a trajectory $\mathcal{Y}(t)$ in \mathcal{M} that passes through the point \mathcal{X} at some t_0 (from now on, $t_0 = 0$ for simplicity)
- $\dot{\mathcal{Y}}(t_0) = \dot{\mathcal{X}}(t_0)$ must be in the tangent space $T_{\mathcal{X}}\mathcal{M}$
- The tangent space is constructed by the velocities $\dot{\mathcal{Y}}(t_0)$ of all possible trajectories $\mathcal{Y}(t)$ that pass through \mathcal{X} at $t = t_0$
- The tangent space is a vectorspace, and the elements of $T_{\mathcal{X}}\mathcal{M}$ are called *tangent vectors*

The tangent space at the identity element ε is called the *Lie algebra* of the Lie group. A *Lie algebra* consists of a vectorspace, \mathbb{V} , over some field, \mathbb{F} , together with a binary operation, $[\cdot, \cdot]$, called the *Lie bracket* (of the algebra) and that satisfies four properties:

- closure: $[\mathbf{X}, \mathbf{Y}] \in \mathbb{V}$
- bilinearity:

$$[a\mathbf{X} + b\mathbf{Y}, \mathbf{Z}] = a[\mathbf{X}, \mathbf{Y}] + b[\mathbf{Y}, \mathbf{Z}]$$

$$[\mathbf{Z}, a\mathbf{X} + b\mathbf{Y}] = a[\mathbf{Z}, \mathbf{X}] + b[\mathbf{Z}, \mathbf{Y}]$$

- alternating: $[\mathbf{X}, \mathbf{X}] = \mathbf{0}$
- Jacobi identity: $[\mathbf{X}, [\mathbf{Y}, \mathbf{Z}]] + [\mathbf{Y}, [\mathbf{Z}, \mathbf{X}]] + [\mathbf{Z}, [\mathbf{X}, \mathbf{Y}]] = \mathbf{0}$

for all $\mathbf{X}, \mathbf{Y}, \mathbf{Z} \in \mathbb{V}$ and $a, b \in \mathbb{F}$.

Lie Algebra Representations

There exist two isomorphic representations of the tangent space:

- The Lie algebra $\mathfrak{m} = T_{\mathcal{X}}\mathcal{M}$ (\mathfrak{m} has m dimensions)
- The corresponding Cartesian space \mathbb{R}^m : $\mathfrak{m} \cong \mathbb{R}^m$
- We can go from one representation to the other using two mutually inverse linear maps called “hat” and “vee”:

$$\begin{aligned}\tau^\wedge &\in \mathfrak{m} \text{ “hat”} \\ \tau &= (\tau^\wedge)^\vee \in \mathbb{R}^m \text{ “vee”}\end{aligned}$$

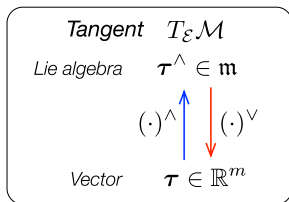


Figure: Sola, Joan, Jeremie Deray, and Dinesh Atchuthan. "A micro lie theory for state estimation in robotics." *arXiv:1812.01537* (2018).

$SO(3)$'s Lie Algebra

$SO(3)$'s Lie Algebra:

vectorspace: $\mathfrak{so}(3) = \{\Phi = \phi^\wedge \mid \phi \in \mathbb{R}^3\}$

field: \mathbb{R}

Lie bracket: $[\Phi_1, \Phi_2] = \Phi_1 \Phi_2 - \Phi_2 \Phi_1$

where ϕ^\wedge is the skew symmetric matrix operator:

$$\phi^\wedge = \begin{bmatrix} 0 & -\phi_3 & \phi_2 \\ \phi_3 & 0 & -\phi_1 \\ -\phi_2 & \phi_1 & 0 \end{bmatrix} \in \mathbb{R}^{3 \times 3}$$

We also use ϕ^\vee as the inverse operator:

$$\Phi = \phi^\wedge \Rightarrow \phi = \Phi^\vee$$

Lie Algebra structure

In Lie Groups, we can get the structure of the tangent space (or the Lie algebra) by differentiating the inverse group constraint. For groups with multiplicative binary operator, we get:

$$\tau^\wedge = \chi^{-1} \dot{\chi} = -\dot{\chi}^{-1} \chi$$

For $SO(3)$ specifically, we get:

$$\phi^\wedge = \mathbf{R}^T \dot{\mathbf{R}} = -\dot{\mathbf{R}}^T \mathbf{R}$$

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What does this reminds us of?

$$\omega_w^\wedge = \dot{R}_{wb} R_{wb}^T$$

$$\omega_b^\wedge = R_{wb}^T \dot{R}_{wb}$$

!!!!

Relation between Lie Algebra and the Manifold

For groups with multiplicative group binary operator:

$$\mathbf{v}^\wedge = \mathcal{X}^{-1} \dot{\mathcal{X}}$$

$$\dot{\mathcal{X}} = \mathcal{X} \mathbf{v}^\wedge$$

This is a linear differential equation. We can solve it:

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$$\mathcal{X}(t) = \mathcal{X}(0) \exp(\mathbf{v}^\wedge t)$$

But $\mathcal{X}(0) = \varepsilon = I$! Why?

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But $\mathcal{X}(0) = \varepsilon = I$! Why? Setting $\boldsymbol{\tau}^\wedge = \mathbf{v}^\wedge t$, we get:

$$\mathcal{X} = \exp(\boldsymbol{\tau}^\wedge)$$

Now, we can go back to the manifold using the exponential map!
We can do the inverse operation using the logarithm map:

$$\boldsymbol{\tau}^\wedge = \log(\mathcal{X})$$

Relation between $SO(3)$ and $\mathfrak{so}(3)$

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- $\mathbf{R} = \exp(\phi^\wedge)$, \exp is the matrix exponential!
- This is a many to one relationship: aka we can generate the same \mathbf{R} from multiple ϕ
- And back?

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- And back?
- $\phi = \log(\mathbf{R})^\vee$, \log is the matrix logarithm!
- What does this reminds us of?
- **Angle-Axis!**
- We can compute the above more efficiently than regular \exp/\log matrix operations

Exp/Log Map for Rotations

We can compute the exp/log operations efficiently:

$$\begin{aligned}\exp(\phi^{\wedge}) &= \exp(\phi \mathbf{a}^{\wedge}) = \sum_{n=0}^{\infty} \frac{1}{n!} (\phi^{\wedge})^n \\ &= \mathbf{I} + \phi \mathbf{a}^{\wedge} + \frac{1}{2!} \phi^2 \mathbf{a}^{\wedge} \mathbf{a}^{\wedge} + \dots \\ &= \dots \\ &= \mathbf{I} \cos \phi + (1 - \cos \phi) \mathbf{a} \mathbf{a}^T + \sin \phi \mathbf{a}^{\wedge} \\ &= \mathbf{I} + \sin \phi \mathbf{a}^{\wedge} + (1 - \cos \phi) \mathbf{a}^{\wedge} \mathbf{a}^{\wedge} \quad (\text{Rodrigues formula}) \\ &= \mathbf{R}\end{aligned}$$

Similarly:¹

$$\begin{aligned}\log(\mathbf{R})^{\vee} &= \phi = \phi \mathbf{a} \\ \mathbf{a}^{\wedge} &= \frac{1}{2 \sin \phi} (\mathbf{R} - \mathbf{R}^T), \phi = \cos^{-1} \left(\frac{\text{tr}(\mathbf{R}) - 1}{2} \right)\end{aligned}$$

¹We also need to take care of some special cases!

A micro Lie theory for state estimation in robotics, *Joan Sola, Jeremie Deray, Dinesh Atchuthan*, 2018, arXiv. [pdf](#)

Sections 7.1.1-7.1.3 from **State Estimation for Robotics**, *Timothy D. Barfoot*, 2022, Cambridge University Press. [ebook](#)

Chapter 3 from **Modern Robotics: Mechanics, Planning, and Control**, *Kevin M. Lynch and Frank C. Park*, 2017, Cambridge University Press. [ebook](#)

Thank you

- Any Questions?
- Office Hours:
 - Tue-Wed (10:00-12:00)
 - 24/7 by email (costashatz@upatras.gr, subject: *ECE_RSI_AM*)
- Material and Announcements



Laboratory of Automation & Robotics