Monogamy-of-entanglement games Nonlocal Games Seminar

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Outline

Nonlocal games

Extended nonlocal games

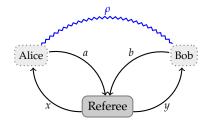
Monogamy-of-entanglement games

Supplementary material

Nonlocal games

Nonlocal games

A nonlocal game is a cooperative game played between Alice and Bob against a referee.



- 1. Question and answer sets: (Σ_A, Σ_B) and (Γ_A, Γ_B)
- 2. Distributions on question pairs: $\pi: \Sigma_A \times \Sigma_B \to [0,1]$
- 3. A predicate $V: \Gamma_A \times \Gamma_B \times \Sigma_A \times \Sigma_B \rightarrow \{0,1\}$, where

$$V(a, b|x, y) = \begin{cases} 1 & \text{if Alice and Bob win} \\ 0 & \text{if Alice and Bob lose} \end{cases}$$

Strategies and values for nonlocal games

Alice and Bob could use different types of strategies:

- ▶ Classical strategies: Alice and Bob answer deterministically, determined by functions of $x \in \Sigma_A$ and $y \in \Sigma_B$.
- ▶ Quantum strategies: Alice and Bob share a joint quantum system $\rho \in D(\mathcal{A} \otimes \mathcal{B})$ and allow their answers to be outcomes of measurements on this shared system.

Strategies and values for nonlocal games

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The *value* of a nonlocal game is the maximal winning probability for the players to win over all strategies of a specified type.

For a nonlocal game, G, we denote the classical and quantum values as

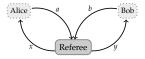
- ▶ Classical value: $\omega(G)$,
- Quantum value: $\omega^*(G)$.

Example: The CHSH game

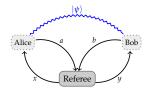
The CHSH game (G_{CHSH}). Winning condition iff $a \oplus b = x \wedge y$.

$$\omega(G_{\mathsf{CHSH}}) < \omega^*(G_{\mathsf{CHSH}})$$

• $\omega(G_{\text{CHSH}}) = \frac{3}{4} = 0.75$:



• $\omega^*(G_{CHSH}) = \cos^2(\frac{\pi}{8}) \approx 0.8536$:

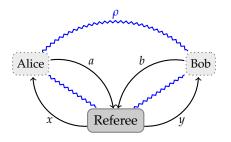


Demo Time: CHSH Game in QETLAB CHSH_GAME.M

Extended nonlocal games

Extended nonlocal games

An extended nonlocal game is a nonlocal game where the referee also holds a quantum system the he measures provided by Alice and Bob.



- 1. Question and answer sets (Σ_A, Σ_B) and (Γ_A, Γ_B) .
- 2. Distribution on question pairs: $\pi: \Sigma_A \times \Sigma_B \to [0,1]$.
- 3. A measurement operator $V : \Gamma_A \times \Gamma_B \times \Sigma_A \times \Sigma_B \to Pos(\mathcal{R})$.

Extended nonlocal games: Winning and losing probabilities

At the end of the protocol, the referee has:

1. The state at the end of the protocol:

$$\rho_{a,b}^{x,y} \in D(\mathcal{R}).$$

2. A measurement the referee makes on its part of the state ρ :

$$V(a, b|x, y) \in Pos(\mathcal{R}).$$

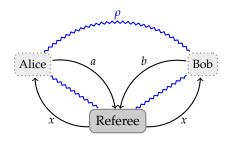
The respective winning and losing probabilities are given by

$$\left\langle V(a,b|x,y),
ho_{a,b}^{x,y} \right
angle \quad \text{and} \quad \left\langle \mathbb{1} - V(a,b|x,y),
ho_{a,b}^{x,y} \right
angle.$$

Monogamy-of-entanglement games

Monogamy-of-entanglement games

Monogamy-of-entanglement games \P , are a special type of extended nonlocal game.



- 1. Same question and answer sets: $\Sigma = \Sigma_A = \Sigma_B$ and $\Gamma = \Gamma_\Delta = \Gamma_B$.
- 2. Alice and Bob get the same question: $\pi(x,y) = 0$ for $x \neq y$.
- 3. Referee's measurement operator: $R: \Sigma \times \Gamma \to \operatorname{Pos}(\mathcal{R})$.
- 4. Winning condition: Iff Alice's output, Bob's output, and the referee's measurement output are the *same*.

¶[Tomamichel, Fehr, Kaniewski, Wehner, (2013)]

Standard quantum strategies for monogamy-of-entanglement games

► The expected *pay-off* for a monogamy-of-entanglement game, *G* using a standard quantum strategy is:

$$\sum_{x \in \Sigma} \pi(x) \sum_{a \in \Gamma} \left\langle R(a|x) \otimes A_a^x \otimes B_a^x, \rho \right\rangle.$$

Since ρ just needs to be a valid density matrix, we use convexity to assume that ρ is pure (rank-one):

$$\omega^*(G) = \left\| \sum_{\mathbf{x} \in \Sigma} \pi(\mathbf{x}) \sum_{\mathbf{a} \in \Gamma} R(\mathbf{a}|\mathbf{x}) \otimes A_{\mathbf{a}}^{\mathsf{x}} \otimes B_{\mathbf{a}}^{\mathsf{x}} \right\|$$

Unentangled strategies for monogamy-of-entanglement games

Alice and Bob only win when their outputs agree, and we assume that the measurements of the referee are positive semidefinite (from the definition for monogamy-of-entanglement games).

► For any monogamy-of-entanglement game, *G*, the unentangled value is:

$$\omega(G) = \max_{f:\Sigma \to \Gamma} \left\| \sum_{x \in \Sigma} \pi(x) R(f(x)|x) \right\|.$$

for choice of measurements $\{A_a^x\}$ for Alice and $\{B_b^y\}$ for Bob.

The BB84 monogamy-of-entanglement game

The BB84 game $(G_{BB84} \text{ for short})^{\P}$ is defined by:

1. Question and answer sets:

$$\Sigma = \Gamma = \{0, 1\},\,$$

2. Uniform probability for questions:

$$\pi(0) = \pi(1) = \frac{1}{2}$$

3. Measurements defined by the BB84 bases:

For
$$x = 0$$
: $R(0|0) = |0\rangle\langle 0|$, $R(1|0) = |1\rangle\langle 1|$
For $x = 1$: $R(0|1) = |+\rangle\langle +|$, $R(1|1) = |-\rangle\langle -|$

The unentangled and standard quantum values for G_{RB84} coincide:

$$\omega(G_{\text{BB84}}) = \omega^*(G_{\text{BB84}}) = \cos^2(\pi/8) \approx 0.8536$$

Demo Time: BB84 Game BB84_GAME.M

Exhaustive search over unentangled strategies

Consider the following SDP:

Primal Problem:
$$(\gamma)$$

max: $\sum_{x \in \Sigma} \pi(x) \left\langle R(f(x)|x), \rho \right\rangle$

s.t.: $\text{Tr}(\rho) = 1$, $(\rho \text{ is pure})$. $\rho \geq 0$, $(\rho \text{ is PSD})$.

We cycle over all possible choices of $f(x) \to a$ and run the above SDP. The best we can do is represented by $\max(\gamma)$ over all such choices.

Calculating max(γ) is now not an SDP, but for small values of |Σ| and |Γ|, we can brute force over every possible combination to obtain the maximum.

Demo Time: Calculating the unentangled value UNENTANGLED_MOE_2IN_2OUT.M

A natural question for monogamy-of-entanglement games

▶ Question: For any monogamy-of-entanglement game, G, is it true that the *unentangled* and *standard quantum* values always coincide? In other words is it true that:

$$\omega(G) = \omega^*(G)$$

for all monogamy-of-entanglement games G?

Demo Time: Random Monogamy Games RANDOM_MOE_GAMES.M

A natural question for monogamy-of-entanglement games

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$$\omega(G)=\omega^*(G)$$

for all monogamy-of-entanglement games *G*?

- Answer:
 - For certain cases: Yes.
 - ► In general: No.

$$\omega(G) = \omega^*(G)$$

In general No

Monogamy-of-entanglement games where $\omega(\textit{G}) \neq \omega^*(\textit{G})$

There exists a monogamy-of-entanglement game, G, with $|\Sigma|=4$ and $|\Gamma|=3$ such that

$$\omega(G) < \omega^*(G)$$
.

1. Question and answer sets:

$$\Sigma = \{0, 1, 2, 3\}, \quad \Gamma = \{0, 1, 2\}.$$

2. Uniform probability for questions:

$$\pi(0) = \pi(1) = \pi(2) = \pi(3) = \frac{1}{4}$$

3. Measurements defined by a mutually unbiased basis \P :

$$\{R(0|x), R(1|x), R(2|x)\}.$$

 $\frac{1}{\|u_x(a)^*u_{x'}(a)\|^2} = 1/|\Gamma| \text{ for } R(a|x) = u_x(a)u_x(a)^*, R(a|x') = u_{x'}(a)u_{x'}(a)^*$

Monogamy-of-entanglement games where $\omega(G) \neq \omega^*(G)$

An exhaustive search over all unentangled strategies reveals an optimal unentangled value:

$$\omega(G)=\frac{3+\sqrt{5}}{8}\approx 0.6545.$$

▶ Alternatively, a computer search over standard quantum strategies and a heuristic approximation for the upper bound of $\omega^*(G)$ reveals that

$$2/3 \ge \omega^*(G) \ge 0.6609$$

This ability to compute upper bounds for extended nonlocal games is obtained from an adaptation of a technique known as the *NPA hierarchy*.

Demo Time: MUB game MUB_4_3_GAME.M

$$\omega(G) = \omega^*(G)$$

For certain classes, Yes.

Monogamy games that obey $\omega(G) = \omega^*(G)$

Theorem

For any monogamy-of-entanglement game, G, for which $|\Sigma|=2$:

$$\omega(G) = \omega^*(G).$$

Proof: Monogamy games that obey $\omega(G) = \omega^*(G)$

Recall that for any monogamy-of-entanglement, G, the standard quantum value may be written as

$$\omega^*(G) = \left\| \lambda \sum_{a \in \Gamma} R(a|0) \otimes A_a^0 \otimes B_a^0 + (1 - \lambda) \sum_{b \in \Gamma} R(b|1) \otimes A_b^1 \otimes B_b^1 \right\|$$

Proof: Monogamy games that obey $\omega(G) = \omega^*(G)$

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Since $||P|| \le ||Q||$ if $P \le Q$ for any $P, Q \ge 0$:

$$\omega^*(G) \leq \left\| \lambda \sum_{a \in \Gamma} R(a|0) \otimes A_a^0 \otimes \mathbb{1}_{\mathcal{B}} + (1-\lambda) \sum_{b \in \Gamma} R(b|1) \otimes \mathbb{1}_{\mathcal{A}} \otimes B_b^1 \right\|$$

Proof: Monogamy games that obey $\omega(G) = \omega^*(G)$

Recall that for any monogamy-of-entanglement, G, the standard quantum value may be written as

$$\omega^*(G) = \left\| \lambda \sum_{a \in \Gamma} R(a|0) \otimes A_a^0 \otimes B_a^0 + (1 - \lambda) \sum_{b \in \Gamma} R(b|1) \otimes A_b^1 \otimes B_b^1 \right\|$$

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Since $\sum_{a \in \Gamma} A_a^x = \sum_{b \in \Gamma} B_b^y = 1$ the above quantity is equal to:

$$\left\|\lambda \sum_{(a,b)\in\Gamma} R(a|0) \otimes A_a^0 \otimes B_b^1 + (1-\lambda) \sum_{(a,b)\in\Gamma} R(b|1) \otimes A_a^0 \otimes B_b^1 \right\|.$$

Monogamy games that obey $\omega({\sf G})=\omega^*({\sf G})$

(Previous slide):

$$\bigg\|\lambda \sum_{(a,b)\in\Gamma} R(a|0)\otimes A_a^0\otimes B_b^1 + (1-\lambda) \sum_{(a,b)\in\Gamma} R(b|1)\otimes A_a^0\otimes B_b^1\bigg\|.$$

Monogamy games that obey $\omega(G) = \omega^*(G)$

(Previous slide):

$$\left\|\lambda \sum_{(a,b)\in\Gamma} R(a|0)\otimes A_a^0\otimes B_b^1 + (1-\lambda) \sum_{(a,b)\in\Gamma} R(b|1)\otimes A_a^0\otimes B_b^1\right\|.$$

Since $\{A_a^0\otimes B_b^1: a,b\in\Gamma\}$ are pairwise orthogonal projections ($\langle A_a^0\otimes B_b^1,A_{a'}\otimes B_{b'}\rangle=0$ for $a\neq a'$ and $b\neq b'$) we have

$$\bigg\| \sum_{(a,b)\in\Gamma} (\lambda R(a|0) + (1-\lambda)R(b|1)) \otimes A_a^0 \otimes B_b^1 \bigg\|.$$

Monogamy games that obey $\omega(G) = \omega^*(G)$

(Previous slide):

$$\Bigg\|\lambda \sum_{(a,b)\in\Gamma} R(a|0)\otimes A_a^0\otimes B_b^1 + (1-\lambda) \sum_{(a,b)\in\Gamma} R(b|1)\otimes A_a^0\otimes B_b^1\Bigg\|.$$

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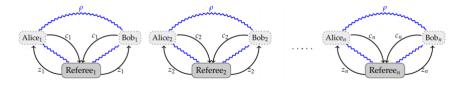
Every entangled strategy is equivalent to one where Alice and Bob use projective measurements so

$$\max_{a,b\in\Gamma} \left\| \lambda R(a|0) + (1-\lambda)R(b|1) \right\| = \omega(G).$$

Parallel Repetition of Monogamy-of-entanglement Games

Parallel repetition of monogamy-of-entanglement games

- ▶ Parallel repetition: Run a monogamy-of-entanglement game, G, for n times in parallel, denoted as G^n .
- Strong parallel repetition: $\omega(G^n) = \omega(G)^n$



Question: Do all monogamy-of-entanglement games obey strong parallel repetition?

Parallel repetition of monogamy-of-entanglement games

► Recall:

$$\omega(G_{\text{BB84}}) = \omega^*(G_{\text{BB84}}) = \cos^2(\pi/8) \approx 0.8536.$$

▶ G_{BB84} obeys strong parallel repetition¶:

$$\omega(G_{\text{BB84}}^n) = \omega^*(G_{\text{BB84}}^n) = (\cos^2(\pi/8))^n$$
.

Demo Time: Strong parallel repetition of BB84 BB84_PARALLEL_REP.M

Strong parallel repetition for certain monogamy-of-entanglement games

Theorem (Johnston, Mittal, R, Watrous)

Let $G=(\pi,R)$ be a monogamy-of-entanglement game for which $\Sigma_A=\{0,1\}$, π is uniform over Σ_A , and R(a|x) is a projection operator. It holds that

$$\omega^*(G^n) = \omega^*(G)^n = \left(\frac{1}{2} + \frac{1}{2}\sqrt{c(G)}\right)^n.$$

where c(G) is the "maximal overlap of measurements" of the referee

$$c(G) = \max_{\substack{x,y \in \Sigma_{A} \\ x \neq y}} \max_{a,b \in \Gamma_{A}} \left\| \sqrt{R(a|x)} \sqrt{R(b|y)} \right\|^{2}$$

A key proposition and lemma

Proposition

Let $G=(\pi,R)$ be a monogamy-of-entanglement game for which $\Sigma=\{0,1\}$, π is uniform over Σ , and R(a|x) is a projection operator for each $x\in\Sigma$ and $a\in\Gamma$. It holds that

$$\omega(G) = \frac{1}{2} + \frac{1}{2} \max_{a,b \in \Gamma} \left\| R(a|0)R(b|1) \right\|.$$

Lemma

Let Π_0 and Π_1 be nonzero projection operators on \mathbb{C}^n . It holds that

$$\|\Pi_0 + \Pi_1\| = 1 + \|\Pi_0\Pi_1\|.$$

Proof of proposition

Assuming the lemma stating $\|\Pi_0+\Pi_1\|=1+\|\Pi_0\Pi_1\|,$ we have that

$$\omega(G) = \max_{a,b \in \Gamma} \left\| \frac{R(a|0) + R(b|1)}{2} \right\| = \frac{1}{2} + \frac{1}{2} \max_{a,b \in \Gamma} \left\| R(a|0)R(b|1) \right\|.$$

The lower bound on $\omega(G)^*$ is obtained from [TFWK13]¶. The upper bound follows from the fact that Alice and Bob can just play the optimal strategy for every n:

$$\omega^*(G^n) \ge \omega(G^n) \ge \left(\frac{1}{2} + \frac{1}{2} \max_{a,b \in \Gamma} \left\| R(a|0)R(b|1) \right\| \right)^n = \left(\frac{1}{2} + \frac{1}{2}\sqrt{c(G)}\right)^n.$$

¶[Tomamichel, Fehr, Kaniweski, Wehner (2013)]

Open questions

Unentangled vs. standard quantum strategies for monogamy-of-entanglement games

Inputs (Σ)	Outputs (Γ)	$\omega^*(G) = \omega(G)$	$\omega^*(G^n) = \omega^*(G)^n$	$\omega_{ns}(G^n) = \omega_{ns}(G)^n$
2	$ \Gamma \geq 1$	yes	yes¶	no
3	$ \Gamma \geq 1$?	?	no
4	3	no	?	no

Question: What about $|\Sigma| = 3$?

- ▶ Proof technique fails for $|\Sigma| > 2$.
- Computational search:
 - ▶ Generate random monogamy-of-entanglement games where $|\Sigma|=3$ and $|\Gamma|\geq 2$.
 - ▶ 10⁸ random games generates, no counterexamples found.

[¶]So long as the measurements used by the referee are projective and the probability distribution, π , from which the questions are asked is uniform.

Supplementary material

Supplementary material: Extended nonlocal games

Winning probability for standard quantum strategies

The winning probability is given by the following equation:

$$\sum_{x,y} \pi(x,y) \sum_{a,b} \frac{\langle V(a,b|x,y), \operatorname{Tr}_{\mathcal{A} \otimes \mathcal{B}} (\mathbb{1}_{\mathcal{R}} \otimes \mathcal{A}_{a}^{\mathsf{x}} \otimes \mathcal{B}_{b}^{\mathsf{y}}) \rho) \rangle}{\operatorname{Tr} (\mathbb{1}_{\mathcal{R}} \otimes \mathcal{A}_{a}^{\mathsf{x}} \otimes \mathcal{B}_{b}^{\mathsf{y}}) \rho)} \operatorname{Tr} (\mathbb{1}_{\mathcal{R}} \otimes \mathcal{A}_{a}^{\mathsf{x}} \otimes \mathcal{B}_{b}^{\mathsf{y}}) \rho)$$

The probabilities cancel giving

$$\sum_{x,y} \pi(x,y) \sum_{a,b} \mathsf{Tr} \left(V(a,b|x,y) \, \mathsf{Tr}_{\mathcal{A} \otimes \mathcal{B}} \left(\mathbb{1}_{\mathcal{R}} \otimes A_a^x \otimes B_b^y \right) \rho \right)$$

The trace operator slips past the $\mathbb{1}_{\mathcal{R}}$ giving

$$\sum_{\mathsf{x},\mathsf{y}} \pi(\mathsf{x},\mathsf{y}) \sum_{\mathsf{a},\mathsf{b}} \mathsf{Tr} \left(V(\mathsf{a},\mathsf{b}|\mathsf{x},\mathsf{y}) \otimes (A^\mathsf{x}_\mathsf{a} \otimes B^\mathsf{y}_\mathsf{b}) \rho \right)$$

Writing the trace in terms of the inner product, we have that

$$\sum_{x,y} \pi(x,y) \sum_{a,b} \left\langle V(a,b|x,y) \otimes A_a^x \otimes B_b^y, \rho \right\rangle.$$

Values of nonlocal games

The value of a game is the maximal winning probability of the players, where the winning probability for a choice of a, b, x, y is

Classical:

$$\omega(G) = \max_{(a,b)\in\Gamma} \sum_{(x,y)\in\Sigma} V(f(x),g(y)|x,y),$$

where the max is over all functions $f: \Sigma_A \to \Gamma_A$ and $g: \Sigma_B \to \Gamma_B$.

Quantum:

$$\omega^*(G) = \sup \sum_{(x,y) \in \Sigma} \pi(x,y) \sum_{(a,b) \in \Gamma} V(a,b|x,y)$$

Measurements may be assumed to be projective

WLOG, we may assume that Alice and Bob's measurements are projective since:

- Alice and Bob may extend the sizes of their Hilbert spaces,
- Naimark's theorem states that any strategy using non-projective measurements can be simulated by a strategy with projective measurements.

Supplementary material: Monogamy-of-entanglement games

Pure strategies are sufficient

Recall that the pay-off for a standard quantum strategy is given by

$$\sum_{(x,y)\in \Sigma_{\mathsf{A}}\times \Sigma_{\mathsf{B}}} \pi(x,y) \sum_{(a,b)\in \Gamma_{\mathsf{A}}\times \Gamma_{\mathsf{B}}} \bigg\langle V(a,b|x,y)\otimes A_{\mathsf{a}}^{\mathsf{x}}\otimes B_{\mathsf{b}}^{\mathsf{y}},\rho \bigg\rangle.$$

WLOG, ρ can be assumed to be pure and the measurement operators projective since one could increase the dimension of the Hilbert space, i.e.:

$$\sum_{(x,y)\in \Sigma_{\mathsf{A}}\times \Sigma_{\mathsf{B}}} \pi(x,y) \sum_{(a,b)\in \Gamma_{\mathsf{A}}\times \Gamma_{\mathsf{B}}} u^* \left(V(a,b|x,y)\otimes A_a^x\otimes B_b^y\right) u,$$

where $\rho = uu^*$ with $u \in \mathcal{R} \otimes \mathcal{A} \otimes \mathcal{B}$. This follows from Naimark's theorem (next slide).

Pure strategies are sufficient: Naimark's theorem

Let $u \in \mathcal{R} \otimes \mathcal{A} \otimes \mathcal{B}$, $w_0 \in \mathbb{C}^{\Gamma_A}$. Define

$$v = u \otimes w_0 \in \mathcal{R} \otimes \mathcal{A} \otimes \mathcal{B} \otimes \mathbb{C}^{\Gamma_A}$$
 and $\widetilde{A}_a^x = U^* (\mathbb{1}_{\mathcal{A}} \otimes E_{a,a}) U$,

where $U \in U(A \otimes \mathbb{C}^{\Gamma_A})$ such that

$$U = w \otimes w_0 \to \sum_{a \in \Gamma_A} \sqrt{A_a^x} w \otimes e_a \qquad (\forall w \in \mathcal{A}).$$

It holds that $\widetilde{A}_a^x \in \operatorname{Proj}(A)$ and hence

$$\widetilde{A}_{a}^{x}v=U^{*}\left(\mathbb{1}_{\mathcal{A}}\otimes \mathsf{E}_{\mathsf{a},\mathsf{a}}\right)U\left(u\otimes w_{0}\right)=U^{*}\sqrt{A_{\mathsf{a}}^{x}}u\otimes \mathsf{e}_{\mathsf{a}},$$

and
$$\widetilde{A}_a^x = \left(\widetilde{A}_a^x\right)^* \left(\widetilde{A}_a^x\right)$$
. Note¶.

[¶]A similar argument holds for Bob's measurements.

The BB84 game: Unentangled value

For any monogamy-of-entanglement game, G, the unentangled value is:

$$\omega(G) = \max_{f:\Sigma \to \Gamma} \left\| \sum_{x \in \Sigma} \pi(x) R(f(x)|x) \right\|$$

Recall, GBB84 has

$$\pi(0) = \pi(1) = 1/2$$
 and $\Sigma = \Gamma = \{0, 1\}.$

So we have that:

$$\omega(G_{\text{BB84}}) = \max_{f:\Sigma\to\Gamma} \left\| \frac{1}{2} R(f(x)|0) + \frac{1}{2} R(f(x)|1) \right\|$$

$$= \frac{1}{2} \||0\rangle\langle 0| + |+\rangle\langle +|\| + \frac{1}{2} \||1\rangle\langle 1| + |-\rangle\langle -|\|$$

$$= \frac{1}{2} + \frac{1}{2\sqrt{2}} = \cos^2(\pi/8).$$

The BB84 game: Standard quantum value

For any monogamy-of-entanglement game, G, the standard quantum value is:

$$\omega^*(G) = \left\| \sum_{x \in \Sigma} \pi(x) \sum_{a \in \Gamma} R(a|x) \otimes A_a^{\mathsf{x}} \otimes B_a^{\mathsf{x}} \right\|.$$

1. Alice and Bob send the following state to the referee:

$$v_{\pm} = \cos(\pi/8)|0\rangle \pm \sin(\pi/8)|1\rangle.$$

2. Alice and Bob always output a=0 and have measurements:

$$A_0^0 = v_+ v_+^*, \quad A_0^1 = v_- v_-^*, \quad B_0^0 = B_0^1 = 1$$

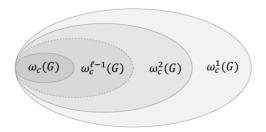
Plugging this into the standard quantum value formula:

$$\omega^*(G) = \left\| \frac{1}{2} \left(|0\rangle\langle 0| \otimes A_0^0 \otimes B_0^0 + |1\rangle\langle 1| \otimes A_0^0 \otimes B_0^0 \right) + \frac{1}{2} \left(|+\rangle\langle +| \otimes A_0^1 \otimes B_0^1 + |-\rangle\langle -| \otimes A_0^1 \otimes B_0^1 \right) \right\| = \cos^2(\pi/8)$$

Supplementary material: Upper bounds for extended nonlocal games

Upper bounds for nonlocal games

- ► The NPA hierarchy sis a method of placing *upper bounds* on the *quantum value* of nonlocal games.
- ▶ Hierarchy of semidefinite programs is *guaranteed* to converge to the commuting measurement value for some finite level, ℓ of the hierarchy.
- ▶ The commuting measurement value is an upper bound on the quantum value, $\omega^*(G) \leq \omega_c(G)$, for all nonlocal games, G.



 $^{^\}P[\mathsf{Navascu\acute{e}s},\,\mathsf{Pironio},\,\mathsf{and}\,\,\mathsf{Ac\'{i}n},\,(2008)]$

NPA hierarchy: Level 1

For a nonlocal game, the pay-off for a commuting measurement strategy is \P

$$\sum_{(x,y)\in\Sigma_{\mathsf{A}}\times\Sigma_{\mathsf{B}}}\pi(x,y)\sum_{(a,b)\in\Gamma_{\mathsf{A}}\times\Gamma_{\mathsf{B}}}V(a,b|x,y)\langle A_{\mathsf{a}}^{\mathsf{x}}B_{\mathsf{b}}^{\mathsf{y}},\rho\rangle.$$

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Define a Gram matrix with entries:

$$C((x, a), (y, b)) = \langle A_a^x B_b^y, \rho \rangle.$$

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Define a Gram matrix with entries:

$$C((x,a),(y,b)) = \langle A_a^x B_b^y, \rho \rangle.$$

The full block-matrix is

$$C = \left(\frac{\langle A_a^{\mathsf{x}} A_{a'}^{\mathsf{x'}}, \rho \rangle \mid \langle A_a^{\mathsf{x}} B_b^{\mathsf{y}}, \rho \rangle}{\langle B_b^{\mathsf{y}} A_a^{\mathsf{x}}, \rho \rangle \mid \langle B_b^{\mathsf{y}} B_{b'}^{\mathsf{y'}}, \rho \rangle} \right)$$

- ▶ If the entries in *C* come from a commuting measurement strategy, *C* will satisfy *certain properties*.
- ► These properties are verifiable via an SDP.
- ¶ for $[A_a^x, B_b^y] = 0$, $\{A_a^x\}, \{B_b^y\} \subset Pos(\mathcal{H})$, and $\rho \in D(\mathcal{H})$.

Matrix constraints

As mentioned, the matrix C satisfies a number of constraints:

- It is positive semidefinite (by definition from the fact that it is a Gram matrix)
- Normalization: C(1,1)=1.
- Commutation:

$$C((x, a), (y, b)) = C((y, b), (x, a))$$

Measurements sum to identity:

$$\sum_{(x,a)} C((x,a),(y,b)) = C(1,(y,b))$$
$$\sum_{(y,b)} C((x,a),(y,b)) = C((x,a),1).$$

Measurements are projective:

$$C(1,(y,b)) = C((y,b),(y,b))$$

 $C((x,a),1) = C((x,a),(x,a))$

Pseudo commuting measurement assemblages

If these properties are satisfied, we say that C is a 1-st order pseudo commuting measurement assemblage.

- ▶ By imposing more structure on this matrix, we get closer to the set of commuting measurement operators.
- ▶ Indexing correspondence between strings and operators:

$$A_a^x \leftrightarrow (x, a)$$
 and $A_{a_1}^{x_1} \cdots A_{a_n}^{x_n} \leftrightarrow (x_1, a_1) \cdots (x_n, a_n)$

Analogous for B_b^y operators.

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Analogous for B_b^y operators.

Define the alphabets

$$\Delta_A^{\leq \ell} = \{\epsilon\} \cup (\Sigma_A \times \Gamma_A)^{\leq \ell} \quad \text{and} \quad \Delta_B^{\leq \ell} = \{\epsilon\} \cup (\Sigma_B \times \Gamma_B)^{\leq \ell} \,.$$

The ℓ -th order pseudo commuting measurement assemblage is

$$C^{(\ell)} = \left(\begin{array}{c|c} \Delta_{A}^{\leq \ell} \cup \Delta_{A}^{\leq \ell} & \Delta_{A}^{\leq \ell} \cup \Delta_{B}^{\leq \ell} \\ \hline \Delta_{B}^{\leq \ell} \cup \Delta_{A}^{\leq \ell} & \Delta_{B}^{\leq \ell} \cup \Delta_{B}^{\leq \ell} \end{array} \right)$$

NPA hierarchy theorem

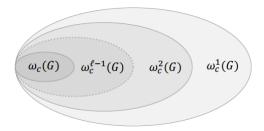
For some finite level ℓ , the pay-off for a commuting measurement strategy for a nonlocal game can be defined by

$$\sum_{(x,y)\in\Sigma_{\mathsf{A}}\times\Sigma_{\mathsf{B}}}\pi(x,y)\sum_{(a,b)\in\Gamma_{\mathsf{A}}\times\Gamma_{\mathsf{B}}}V(a,b|x,y)C^{(\ell)}.$$

Upper bounds for extended nonlocal games

Extended NPA hierarchy:

- ▶ Uses the same idea as the NPA hierarchy. (For dim(\mathcal{R}) = 1, the NPA hierarchy is a special case.)
- ► Enables one to compute *upper bounds* on the *standard* quantum value for extended nonlocal games.
- Same idea as before, only now we need to take into account the actions of the referee.



Commuting measurement strategies

A commuting measurement strategy consists of a finite-dimensional complex Euclidean space \mathcal{H} as well as the following:

- ▶ Shared state: $\rho \in \mathcal{R} \otimes \mathcal{H}$.
- ▶ Measurements: $\{A_a^x\} \subset \operatorname{Pos}(\mathcal{H}), \{B_b^y\} \subset \operatorname{Pos}(\mathcal{H}),$ where $[A_a^x, B_b^y] = 0$ for all x, y, a, b.

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The expected *pay-off* for a commuting measurement strategy is given by:

$$\sum_{(x,y)\in\Sigma_{\mathsf{A}}\times\Sigma_{\mathsf{B}}}\pi(x,y)\sum_{(a,b)\in\Gamma_{\mathsf{A}}\times\Gamma_{\mathsf{B}}}\left\langle V(a,b|x,y)\otimes A_{a}^{x}B_{b}^{y},\rho\right\rangle$$

The commuting measurement value, denoted as $\omega_c(G)$, is the supremum of the pay-off over all commuting measurement strategies.

Extended NPA hierarchy

Same idea, but now we're taking into account the referee, and therefore have a larger matrix.

For each ℓ , now consider block matrices

$$M^{(\ell)} = \begin{pmatrix} M_{1,1}^{(\ell)} & \cdots & M_{1,m}^{(\ell)} \\ \vdots & \ddots & \vdots \\ M_{m,1}^{(\ell)} & \cdots & M_{m,m}^{(\ell)} \end{pmatrix}$$

where each block takes the form $M_{i,j}^{(\ell)}: \Sigma^{\leq \ell} \times \Sigma^{\leq \ell} \to \mathbb{C}$.

- Each submatrix has similar properties to what we saw for the NPA hierarchy.
- ► The overall matrix also has some structure, which is unique to this case.

Assemblages

Another natural way to think about the commuting measurement value is in terms of assemblages:

$$K(a, b|x, y) = \operatorname{Tr}_{\mathcal{H}} \left((\mathbb{1}_{\mathcal{R}} \otimes A_a^x B_b^y) \rho \right)$$

- ▶ For a particular choice of *x*, *y*, *a*, *b*, an assemblage corresponds to the *unnormalized state* in the referee's hands at the end of the game.
- ► The function K completely determines the performance of Alice and Bob's strategy:

$$\sum_{(x,y)\in \Sigma_{\mathsf{A}}\times \Sigma_{\mathsf{B}}} \pi(x,y) \sum_{(a,b)\in \Gamma_{\mathsf{A}}\times \Gamma_{\mathsf{B}}} \bigg\langle V(a,b|x,y), K(a,b|x,y) \bigg\rangle.$$

Extended NPA hierarchy theorem

Let
$$dim(\mathcal{R}) = m$$
. Then

$$\left\langle V(a,b|x,y),K(a,b|x,y)\right\rangle = \left\langle V(a,b|x,y),M^{(\ell)}((x,a),(y,b))\right\rangle$$

for all $m, \ell \geq 1$, $(x, y) \in \Sigma_A \times \Sigma_B$ and $(a, b) \in \Gamma_A \times \Gamma_B$.

Supplementary material: Lower bounds for extended nonlocal games

Lower bounds for extended nonlocal games

Key idea: Fixing measurements on one system yields the optimal measurements of the other system via an SDP^\P

Lower bounds for extended nonlocal games

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Iterative "see-saw" algorithm between two SDPs:

- ► SDP-1: Fix Bob's measurements. Optimize over Alice's measurements.
- ► SDP-2: Fix Alice's measurements (from SDP-1). Optimize over Bob's measurements.
- Repeat.

Not guaranteed to give optimal value, as the algorithm can get stuck in a local minimum.

Lower bounds for extended nonlocal games

Define $\{\rho_a^x : x \in \Sigma_A, \ a \in \Gamma_A\} \subset \operatorname{Pos}(\mathcal{R} \otimes \mathcal{B})$ as the residual states acting on the referee and Bob's systems and let

$$f = \sum_{(x,y) \in \Sigma_{A} \times \Sigma_{B}} \pi(x,y) \sum_{(a,b) \in \Gamma_{A} \times \Gamma_{B}} \left\langle V(a,b|x,y) \otimes B_{b}^{y}, \rho_{a}^{x} \right\rangle$$

$$\underline{\text{Lower bound (SDP-1)}} \qquad \underline{\text{Lower bound (SDP-2)}}$$

$$\max: \quad f \qquad \qquad \max: \quad f$$

$$\text{s.t.:} \quad \sum_{a \in \Gamma_{A}} \rho_{a}^{x} = \tau, \qquad \qquad \text{s.t.:} \quad \sum_{b \in \Gamma_{B}} B_{b}^{y} = \mathbb{1}_{\mathcal{B}},$$

$$\tau \in \mathrm{D}(\mathcal{R} \otimes \mathcal{B}). \qquad \qquad B_{b}^{y} \in \mathrm{Pos}(\mathcal{B}).$$

▶ Iterate between SDP-1 and SDP-2 until desired numerical precision is reached.

Supplementary material: Parallel repetition for monogamy-of-entanglement games

By definition

$$\|\Pi_0 + \Pi_1\| = \max\{v^*(\Pi_0 + \Pi_1)v : v \in \mathcal{S}\}.$$

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Recall the definition of the ∞ -norm for some operator $A \in \mathrm{L}(\mathcal{X},\mathcal{Y})$:

$$||A|| = \max\{||Av|| : v \in \mathcal{S}(\mathcal{X})\}.$$

We can then rewrite our expression as

$$\max\{\|\Pi_0 v\|^2 + \|\Pi_1 v\|^2 : v \in \mathcal{S}\}.$$

Define unit vectors $u_0, u_1 \in \mathcal{S}(\mathbb{C}^n)$, we can write

$$\Pi_0=\mathit{u}_0\mathit{u}_0^*\quad\text{and}\quad \Pi_1=\mathit{u}_1\mathit{u}_1^*.$$

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 and $\Pi_1 = u_1 u_1^*$.

Therefore

$$\Pi_0 v = u_0 u_0^* v = \langle u_0, v \rangle u_0$$

and that

$$\|\Pi_0 v\|^2 = |\langle u_0, v \rangle|^2 \|u_0\| = |\langle u_0, v \rangle|^2$$

(and similarly for Π_1 and u_1).

Therefore, we have that

$$\max\{|\langle u_0, v \rangle| + |\langle u_1, v \rangle|^2 : v \in \mathcal{S}, u_0 \in \mathcal{S}_0, u_1 \in \mathcal{S}_1\},\$$

where S_0 is unit sphere of im(Π_0) (similarly for S_1 and Π_1).

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$$|\langle u_0, v \rangle|^2 = ||u_0|| ||v||$$

(similarly for u_1 and v). Equality is achieved since u_0 and v are linearly dependent.

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Again using definition of norm

$$\max\{v^*(u_0u_0^*+u_1u_1^*)v:v\in\mathcal{S},u_0\in\mathcal{S}_0,u_1\in\mathcal{S}_1\}.$$

Therefore

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Note that for every unit vector $u_0, u_1 \in \mathbb{C}^n$ it holds that

$$||u_0u_0^* + u_1u_1^*|| = 1 + |\langle u_0, u_1 \rangle|,$$

since $u_0u_0^*+u_1u_1^*$ has at most two nonzero eigenvalues of $1\pm |\langle u_0,u_1\rangle|$. Therefore

$$\max\{1+|\langle \textit{u}_0,\textit{u}_1\rangle|:\textit{u}_0\in\mathcal{S}_0,\textit{u}_1\in\mathcal{S}_1\}.$$

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$$\max\{1+|\langle u_0, u_1\rangle|: u_0 \in \mathcal{S}_0, u_1 \in \mathcal{S}_1\}.$$

Finally, since

$$||A||_{p} = \max\{|\langle B, A \rangle| : B \in L(\mathcal{X}, \mathcal{Y}), ||B||_{p^*} \le 1\},$$

it holds that

$$1 + \|\Pi_0\Pi_1\|.$$