

$$\frac{dx}{1} = \frac{dy}{-2q} = \frac{dz}{p+2q^2} = \frac{dp}{f(zx)} = \frac{dq}{-f'(z)} = \frac{dq}{-1}$$

$$x = q + c$$

$$q = \underline{\underline{-x+a}}$$

$$p = 3x^2 - y + (x-a)^2$$

$$z = pdx + q dy$$

Solve.

# Linear Eqns with constant coeff-

$$u_{xx} - 2u_{xy} + u_{yy} = x^2 + y$$

$$D_x = \frac{\partial}{\partial x} \quad D_y = \frac{\partial}{\partial y}$$

$$D = \frac{\partial}{\partial x} \quad D' = \frac{\partial}{\partial y}$$

Theorem:-

If  $(a_i D + b_i D' + c_i)$  is a factor of  $F(D, D')$ , then a sol' of  $F(D, D') u_{(xy)} = 0$

is of form

$$u_i = e^{-c_i/a_i x} f(b_i x - a_i y)$$

Note:-

$$F(D, D') u = (D + D')(D - D') u = 0$$

then any  $'u'$  such that  $(D - D') u = 0$   
or  $(D + D') u = 0$

will be a sol' of  $\underline{\underline{F(D, D') u = 0}}$

aim

to find sol<sup>n</sup> of

$$(a_i D + b_i D' + c_i) u = 0.$$

$$a_i u_x + b_i u_y + c_i u = 0.$$

$$\rightarrow \frac{dx}{a_i} = \frac{dy}{b_i} = -\frac{du}{c_i}$$



$$b_i x - a_i y = C$$

$$\frac{du}{u} = -\frac{a_i}{b_i} dy$$

$$\frac{du}{u} = -\frac{a_i}{b_i} dx$$

$$u = \underline{\underline{e^{-C_i/b_i y}}}, C_2$$

$$u = \underline{\underline{e^{C_i/a_i x}}}, C_1$$

}

$$u_i = \underline{\underline{e^{-C_i/a_i x} f(b_i x - a_i y)}}$$

$$(2D - 5D' + 7) u = 0 \quad a_i = 2 \quad b_i = -5 \quad c_i = 7$$

$$u_1 = \underline{\underline{e^{-7/2 x} f(-5x - 2y)}} = e^{-7/2 x} g(5x + 2y)$$

$$u_2 = \underline{\underline{e^{-5x} f(-5x - 2y)}}.$$

eg:

$$(2D^2 D' - 3DD'^2 + D'^3) u = 0.$$

$$F(D, D') u = 0.$$

factorising

$$D'(2D^2 - 3DD' + D'^2)$$

$$D'(2D^2 - 2DD' - DD' + D'^2).$$

$$D'(D - D')(2D - D') u = 0.$$

$$D'u = 0; \quad Du - D'u = 0; \quad (2D - D')u = 0$$

$$(a_i = 0).$$

$$u_1 = f(x); \quad u_2 = g(x+y); \quad u_3 = h(x+2y).$$

$$(D^2 - D'^2 + 2D + 1)u = 0.$$

$$(D + D' + 1)(D - D' + 1)u = 0.$$

complementary function.

$$= e^{-x} f_1(x-y) + e^{-x} f_2(x+y).$$

Theorem 2: If  $(a_i D + b_i D' + c_i)^2$ , ( $a_i \neq 0$ ) is a factor of  $F(D, D')$  then the soln of

$F(D, D')u = 0$  is of the form  $u_i = e^{-c_i/a_i} x \left[ g(b_i x - a_i y) + \frac{x}{a_i} f(b_i x - a_i y) \right]$

proof consider  $(a_i D + b_i D' + c_i)^2 u = 0 \rightarrow$  let  $(a_i D + b_i D' + c_i) u = z$

$$\star 1 (a_i D + b_i D' + c_i) z = 0$$

$$\therefore z = e^{-c_i/a_i} x f(b_i x - a_i y).$$

$$\oplus 1 (a_i D + b_i D' + c_i) u = e^{-c_i/a_i} x f(b_i x - a_i y)$$

$$+ a_i u_x + b_i u_y = -c_i u + e^{-c_i/a_i} x f(b_i x - a_i y),$$

$$\frac{dx}{a_i} = \frac{dy}{b_i} = \frac{du}{-c_i u + e^{-c_i/a_i} x f(b_i x - a_i y)}$$



$$b_i x - a_i y = \text{constant}$$

$$\frac{dx}{a_i} = \frac{du}{-c_i u + e^{-c_i/a_i} x f(b_i x - a_i y)}$$

$$+ \frac{du}{dx} + \frac{G}{a_i} u = \frac{1}{a_i} e^{-c_i/a_i} x f(d_i).$$

$$+ \frac{du}{dx} e^{c_i/a_i x} + \frac{c_i}{a_i} u e^{c_i/a_i x} = \frac{1}{a_i} f(d_i),$$

$$+ \frac{du}{dx} (u e^{c_i/a_i x}) = \frac{1}{a_i} f(d_i)$$

$$+ u e^{c_i/a_i x} = \frac{x}{a_i} f(d_i) + d_2$$

$$+ \text{general soln.}$$

$$u = e^{-\frac{c_i}{a_i} x} \left[ g(b_i x - a_i y) + \frac{x}{a_i} f(b_i x - a_i y) \right]$$

Corollary

If  $(a_i D + b_i D' + c_i)^n$ ,  $a_i \neq 0$  is a factor, then

$u = e^{-c_i/a_i} x \left( \sum_{j=1}^n x^{j-1} f_{ij}(b_i x - a_i y) \right)$  will be sol<sup>n</sup> of  $F(D, D') u = 0$

Corollary 2) If  $a_i = 0$ ,  $b_i \neq 0$  and  $(b_i D' + c_i)^n$  is a factor then  $e^{-c_i/b_i} y \sum_{j=1}^n (x)^{j-1} f_{ij}(b_i y)$  will be the sol<sup>n</sup> of  $F(D, D') u = 0$

Eg) ①  $(D^3 + 3D^2 D' - 4D'^3) u = 0$ .

②  $(D^3 - 3D^2 D' + 3DD'^2 - D'^3) u = 0$

③  $(2D + D' + 1)^2 (D - 2D' + 2)^3 u = 0$

Sol<sup>n</sup>

①  $(D^3 + 3D^2 D' - 4D'^3) u$

$$(D^3 - D'^3 + 3D^2 D' - 3D'^3) u = 0$$

$$(D - D')(D^2 + D'^2 + DD' + 3D'(D + D')) u = 0.$$

~~So  $(D - D')(D^2 + D'^2 + DD' + 3D'D + 3D'^2) u = 0$~~

$$(D - D')(D^2 + 4DD' + 4D'^2) u = 0$$

$$(D - D')(D + 2D')^2 u = 0$$

$$u_1 = \underline{f(x+y)}$$

$$u_2 = \underline{\alpha f_1(2x-y) + g_1(2x-y)}$$

$$\text{Sol}^n \leftarrow a = \underline{\alpha} u_1 + \underline{\beta} u_2$$

②  $(D^3 - 3D^2 D' + 3DD'^2 - D'^3) u = 0$

$$(D - D')^3 u = 0$$

$$u = \underline{\alpha^2 f(x+y) + \alpha g(x+y) + h(x+y)}$$

$$(D^2 + 2D - 1)^2 (D - 2D^2 + 2)^3 u = 0.$$

$$u_1 = \pi e^{+2x} f_{11}(2x-y) + e^{-2x} f_{12}(2x-y)$$

$$u_2 = x^2 e^{-2x} f_{21}(2x+y) + x e^{-2x} f_{22}(2x+y) + e^{-2x} f_{23}(2x+y)$$

$$\underline{u = A u_1 + B u_2}$$

(no need to write A, B. it will be absorbed in  $u_1, u_2$ ).

## # Irreducible Equations.

$$① (2D^2 - D'^2 + P)u = 0$$

Consider  $F(D, D')u = 0$  where

$$F(D, D') = \sum_{\alpha} \sum_{\beta} C_{\alpha \beta} D^{\alpha} D'^{\beta}$$

$$\Rightarrow u = \sum_{j=1}^{\infty} d_j e^{(q_j x + b_j y)} \text{ is a soln.}$$

$$\text{provided } F(q_j, b_j) = 0.$$

$$u = \sum_{j=1}^{\infty} d_j e^{q_j x + b_j y}$$

$$\text{where } q_j, b_j \text{ are } 2q_j^2 - b_j^2 + q_j = 0$$

$$b_j = \pm \sqrt{2q_j^2 + q_j}$$

$$\therefore u = \sum_{j=1}^{\infty} d_j e^{q_j x \pm \sqrt{2q_j^2 + q_j} y}$$

$$\text{Eq. } (D - 2D' + 5)(D^2 + D' + 3)u = 0.$$

$$\text{CF of } (D - 2D' + 5)v = 0 \text{ is } v = e^{-5x} f_1(2x+y)$$

$$\text{CF of } (D^2 + D' + 3)u_1 = 0 \text{ is } u_1 = \sum_{j=1}^{\infty} d_j e^{q_j x + b_j y}$$

$$q_j^2 + b_j^2 + 3 = 0. \quad \therefore b_j = -3 - q_j^2$$

$$\therefore u_1 = \sum_{j=1}^{\infty} d_j e^{q_j x - (3 + q_j^2)y}$$

$$\therefore u = e^{-5x} f_1(2x+y) + \sum_{j=1}^{\infty} d_j e^{(q_j x - (3 + q_j^2)y)}$$

# Finding particular integral of non-homogeneous eqn.

(1)

①  $(D^2 + 2D' + D'^2) u = 12xy$ .

$$u_C = \frac{f(x-y) + g(x-y)}{(D+D')^2}$$

$$u_p = \frac{1}{D^2 + 2D' + D'^2} 12xy.$$

$$= \frac{1}{(D+D')^2} 12xy = \frac{12}{D^2} \left(1 + \frac{D'}{D}\right)^{-2} xy.$$

$$= \frac{12}{D^2} \left(1 - \frac{2D'}{D} + \frac{3D'^2}{D^2} - \dots\right) (xy)$$

(both differentiation and integration  
in one equation)

$$\underline{\underline{\frac{12}{D^2} (xy - x^2)}}.$$

(further terms are 0).

$$= 12 \left( \frac{x^3}{6} y - \frac{x^4}{12} \right) = \underline{\underline{2x^3 y - x^4}}.$$

②  $(D+D'-1)(D+2D'-3)u = 4+3x+6y.$

$$u_C = e^{x(D+D')} f_1(x-y) + e^{3x(D+D')} f_2(2x-y)$$

$$u_p = \frac{1}{(D+D'-1)(D+2D'-3)} (4+3x+6y).$$

$$= [1 - (D+D')^{-1}]^{-1} \frac{1}{3} [1 - (\frac{D+2D'}{3})]^{-1} (4+3x+6y).$$

$$= \frac{1}{3} [1 + (D+D') + (D+D')^2 + \dots] [1 + \frac{D+2D'}{3} + \dots] (4+3x+6y)$$

$$= \frac{1}{3} [1 + (D+D') + (D+D')^2 + \dots] \underbrace{(4+3x+6y + \frac{3+12}{3})}_g$$

$$= \frac{1}{3} (9+3x+6y + 3+6)$$

$$= \underline{\underline{x+2y+6}}.$$

Theorem

If  $f(x,y)$  is of the form  $e^{ax+by}$ , then P.I. of  $F(D, D') = f(x,y)$

is  $\frac{1}{F(D, D')} e^{ax+by} = \frac{1}{F(a, b)} e^{ax+by}$ .

provided  $F(a, b) \neq 0$ .

Ex  $(D^2 - 3DD' + 2D'^2) u = e^{2x-y}$   $a=2$   $b=-1$

$$u_p = \frac{1}{D^2 - 3DD' + 2D'^2} e^{2x-y}, \quad a=2$$

$$F(a, b) = (2^2 - 3 \cdot 2 \cdot (-1) + 2 \cdot (-1)^2) \neq 0.$$
$$= 12$$

$$\therefore u_p = \frac{1}{12} e^{2x-y}$$

Ex  $(D^2 - D') u = e^{x+y}, \quad a=1, b=1.$

$$F(a, b) = 0$$

$$u_p = \frac{e^{x+y}}{(D^2 - D')} = \frac{e^{x+y}}{0} ?$$

Theorem: If  $f(x,y) = e^{ax+by} \phi(x,y)$  then  $\frac{1}{F(D, D')} (e^{ax+by} \phi(x,y))$

$$= e^{ax+by} \cdot \frac{1}{F(D+a, D'+b)} \phi(x,y)$$

$$\rightarrow u_p = \frac{1}{(D^2 - D')} e^{x+y} \cdot \frac{1}{((D+1)^2 - (D'+1))} \phi(x,y)$$

$$= e^{x+y} \cdot \frac{1}{(D^2 + 2D - D')}$$

$$= e^{x+y} \cdot \frac{1}{(-D') \left( 1 - \frac{2D + D^2}{D'} \right)}$$

$$= e^{x+y} \left( -\frac{1}{D'} \right) \left( 1 + \frac{2D + D^2}{D'} + \dots \right)$$

$$= e^{x+y} \left( -\frac{1}{D'} \right) (1)$$

$$= -y e^{x+y} \leftarrow$$

$$(D^2 + DD' + D + D' + 1) u = e^{-2x} (x^2 + y^2)$$

$$u_p = e^{-2x} \cdot \frac{1}{(D-2)^2 + (D-2)D' + D-2 + D' + 1} (x^2 + y^2),$$

$$= e^{-2x} \cdot \frac{1}{D^2 + DD' - 3D - D' + 3} (x^2 + y^2)$$

$$= e^{-2x} \cdot \frac{1}{3} \left[ 1 + \frac{(D^2 + DD' - 3D - D')}{3} + \frac{(D^2 + DD' - 3D - D')^2}{9} + \dots \right] (x^2 + y^2).$$

$$= \frac{e^{-2x}}{3} \left[ x^2 + y^2 - \frac{1}{3} (2 - 6x - 2y) + \frac{1}{9} (18 + 2) \right].$$

$$= \frac{e^{-2x}}{27} (9x^2 + 9y^2 + 18x + 6y + 14).$$

# periodic functions-

$$(D^2 + 2DD' + D'^2) u = \cos(x+2y).$$

$$\text{Let } u_p = c_1 \cos(x+2y) + c_2 \sin(x+2y).$$

$$D^2 u = -c_1 \cos(x+2y) - c_2 \sin(x+2y)$$

$$D'^2 u = -4c_1 \cos(x+2y) - 4c_2 \sin(x+2y)$$

$$DD' u = -2c_1 \cos(x+2y) - 2c_2 \sin(x+2y).$$

Substitute in the eq<sup>n</sup> and equate coeff. of independent func<sup>n</sup>.

$$-9c_1 \cos(x+2y) - 9c_2 \sin(x+2y) = \cos(x+2y)$$

$$\therefore c_2 = 0, c_1 = -1/9.$$

$$\therefore u_p = -\frac{1}{9} \cos(x+2y)$$

$$\textcircled{1} \quad (D^2 - D') u = x \sin y.$$

$$\textcircled{2} \quad (D^2 - D') u = e^y \cos s(2x + 3y)$$

D+ Q G 6 1000

$$q_j^2 - b_j = 0 \quad b_j > q_j^2$$

$$u = \sum_{j=1}^{\infty} a_j e^{q_j x + q_j^2 y}$$

$$u_p = \frac{1}{(D^2 - D')} x \sin y = \frac{1}{(1 + (D^2 - D' - i))} x \sin y = \frac{1}{(1 - (D' + i - D^2))} x \sin y$$

$$= \left[ 1 + (D' + i - D^2) + \frac{(D' + i)^2}{(D^2 - D')} + \dots \right] x \sin y$$

$$= x \cos y + x \sin y +$$

$$= \operatorname{Im} \left[ \frac{1}{(D^2 - D')} x e^{iy} \right]$$

$$= \operatorname{Im} \left[ e^{iy} \cdot \frac{x}{(D^2 - D' - i)} \right]$$

$$= \operatorname{Im} \left[ e^{iy} \cdot \frac{x}{(1 - (D' + i - D^2))} \right]$$

$$= \operatorname{Im} \left[ e^{iy} \cdot \frac{x}{(1 + (D' + i - D^2))} \right]$$

$$= \operatorname{Im} \left[ e^{iy} \cdot \frac{x}{i(D^2 - D' - 1)} \right]$$

$$= \operatorname{Im} \left[ \frac{e^{iy}}{i} \frac{x}{(1 - (\frac{D^2 - D'}{i})')} \right]$$

$$= \operatorname{Im} \left[ (-1) \frac{e^{iy}}{i} (x) \right]$$

$$u_p = x \cos y$$

$$u_e + u_p$$

## Classification of 2nd order PDE and Integration

①

→ Some motivation

i) Conservation of vector field

$$\nabla \cdot \vec{V} = 0 \quad \text{--- ①}$$

$$\text{if rotational } \vec{\nabla} \times \vec{V} = 0, \quad \vec{V} = \nabla \phi. \quad \text{--- ②}$$

$$\text{① or ②} \rightarrow \nabla^2 \phi = 0.$$

$$2D \quad \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0.$$

when water drop falls in steady state water

$$\text{flux } \vec{j} = -D \nabla c$$

$$\text{total flux } \vec{\nabla} \cdot \vec{j} = \nabla (-D \nabla c).$$

$$\frac{\partial c}{\partial A} + D \nabla^2 c = 0 \quad \text{diffusion / unsteady condition}$$

Consider a general linear 2nd order PDE with variable coefficients.

$$a(x,y) u_{xx} + b(x,y) u_{xy} + c(x,y) u_{yy} + d(x,y) u_x + e(x,y) u_y + f(x,y) u(x,y) = g(x,y) \quad \text{--- ③}$$

$$\text{Ex: 1) } x^2 u_{xx} + 2xy u_{xy} + y^2 u_{yy} + x u_x - y u_y + u = \ln x.$$

$$2) \quad x u_{xx} - u_{xy} + y^2 u_{yy} = 0$$

aim: to integrate ③

$$u_{xx} + u = 0.$$

If ③ can be to one of the following forms.

$$u_{xy} = F(x, y, u, u_x, u_y)$$

$$u_{xy} =$$

$$u_{yy} =$$

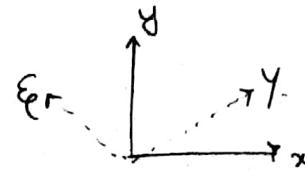
→ i.e. we need to transform  $(x, y)$  coordinate to  $(\xi, \eta)$  coordinates; we reduce the 3 2nd order derivatives terms in ③ to above form (having one 2nd order

### Motivation

$$au_x + bu_y = 0, \quad ; \quad a, b \in \mathbb{R}, \quad ; \quad a^2 + b^2 \neq 0.$$

$$(a, b) \cdot (u_x, u_y) = 0.$$

$$\nabla \cdot \nabla u = 0$$



We need to transform to another  $(E_r, Y)$

such that only  $u_{E_r}$  or  $u_Y$  remains in the PDE.

$$(E_r, Y) = (ax + by, bx - ay)$$

because

$$\frac{dx}{a} = \frac{dy}{b} = \frac{du}{0} \quad \Rightarrow \quad u = c_1, \quad bx - ay = c_2. \\ \therefore u = f(bx - ay)$$

$$\therefore E_r = ax + by \quad Y = bx - ay$$

$$u = f(E_r, Y) \quad (\text{my assumption})$$

$$u_x = u_{E_r} E_{rx} + u_Y Y_x = u_{E_r} \cdot a + u_Y \cdot b$$

$$u_y = u_{E_r} E_{ry} + u_Y Y_y = u_{E_r} b - u_Y \cdot a$$

$$\therefore a u_x + b u_y = 0 \quad \Rightarrow \quad a(a u_{E_r} + b u_Y) + b(b u_{E_r} - a u_Y) = 0.$$

$$\therefore (a^2 + b^2) u_{E_r} = 0 \quad \Rightarrow \quad u_{E_r} = 0 \quad \therefore a^2 + b^2 \neq 0.$$

$$\therefore u = f(Y)$$

$$u = f(bx - ay).$$

We use this motivation for eq" having multiples 2nd order terms to having only one by choosing appropriate  $(E_r, z)$  coordinate axis.

$$(x, y) \rightarrow (E_r, \eta) \quad J = \begin{vmatrix} E_{rx} & E_{ry} \\ \eta_x & \eta_y \end{vmatrix} \neq \{0, \infty\}$$

So after transforming, we get a PDE of below form

$$u_{E_r \eta} = \dots$$

$$u_{E_r E_r} = \dots$$

$$u_{\eta \eta} = \dots$$

consider  $\epsilon_p = \epsilon_p(x, y)$ ,  $n = n(x, y)$

$$u_x = u_{cp} \epsilon_{px} + u_n n_x, \quad u_y = u_{cp} \epsilon_{py} + u_n n_y$$

$$u_{xx} = u_{cp} (\epsilon_{nn}) + \epsilon_x (u_{cp} \epsilon_{px} + u_n n_x) + u_n n_{xx} + n_x (u_n \epsilon_{px} + u_{nn} n_x).$$

$$u_{yy} = "y$$

$$u_{yy} = "y.$$

eqn ④

$$a u_{xx} + b u_{xy} + c u_{yy} + \dots = g(x, y).$$

reduces to

$$(a \epsilon_x^2 + b \epsilon_x \epsilon_y + c \epsilon_y^2) u_{cp} + (2a(\epsilon_x n_n) + b(\epsilon_x n_y + \epsilon_y n_x) + 2c(\epsilon_y n_y)) u_n + (a n_x^2 + b n_x n_y + c n_y^2) u_{nn} + F(u, \epsilon_p, n, u_{cp}, u_n) = 0. \quad \text{--- (5)}$$

⑤ can be written as

$$\tilde{A} u_{cp} + \tilde{B} u_n + \tilde{C} u_{nn} + F = 0$$

$$A = a \epsilon_x^2 + b \epsilon_x \epsilon_y + c \epsilon_y^2$$

$$B =$$

$$B =$$

$$\text{note that } (\tilde{B}^2 - 4\tilde{A}\tilde{C}) = (b^2 - 4ac)(\epsilon_x n_y - \epsilon_y n_x)^2 \\ = (b^2 - 4ac)J^2$$

If  $J \neq 0$ , discriminant is invariant.

Let  $\chi = \epsilon_p \text{ or } n$ .

$$\text{consider } a \chi_x^2 + b \chi_x \chi_y + c \chi_y^2 = 0$$

$$\text{i.e. } a \left( \frac{\chi_x}{\chi_y} \right)^2 + b \left( \frac{\chi_x}{\chi_y} \right) + c = 0. \quad \text{--- (6)}$$

or  $m \Rightarrow \chi = \text{const}$

$$d\chi = \chi_x dx + \chi_y dy = 0 \quad \text{--- (7)}$$

$$\frac{dy}{dx} = -\frac{x_x}{x_y} \quad \text{--- (2)}$$

$\textcircled{1}$  and  $\textcircled{2}$   $a \left( \frac{dy}{dx} \right)^2 - b \frac{dy}{dx} + c = 0.$

$$\frac{dy}{dx} = \frac{b \pm \sqrt{b^2 - 4ac}}{2a}$$

Soln if the above PDE is curve

$$\xi = \text{const.}, \eta = \text{const.}$$

We are killing  $\tilde{A}$  and  $\tilde{C}$  provided  $\tilde{B}$  is non zero.

$\xi = \text{const.}$  &  $\eta = \text{const.}$  are called "characteristic curves"

$$\xi = \xi(x, y) = \phi(x, y)$$

$$\eta = \eta(x, y) = \psi(x, y) \quad \text{reduce given PDE to}$$

$$u_{\xi\eta} = F \quad \text{provided } \tilde{B} \neq 0.$$

$\Rightarrow$  PDE of form  $\boxed{u_{\xi\eta} = F}$  are called canonical forms.

Q1  $u_{xx} + 2u_{xy} + u_{yy} + u_x = 0.$

$$a=1 \quad b=2 \quad c=1$$

Q2  $y^2 u_{xx} - x^2 u_{yy} = 0$

$$a=y^2 \quad b=0 \quad c=-x^2$$

$$\frac{dy}{dx} = \pm \frac{\sqrt{4x^2y^2}}{2y^2} = \pm \frac{2xy}{2y^2} = \pm \frac{x}{y}.$$

$$\therefore \xi = \xi(x, y) = y^2 - x^2 = C_1$$

$$\eta = \eta(x, y) = y^2 + x^2 = C_2.$$

$$u_x = u_{\xi\eta}(-2x) + u_\eta(2x) = -2xu_{\xi\eta} + 2xu_\eta.$$

$$u_{xx} = -2(u_{\xi\eta} + x(u_{\xi\xi\eta}(-2x) + u_{\xi\eta\eta}(2x))) + 2(u_\eta + x(u_{\eta\xi\eta}(-2x) + u_{\eta\eta\eta}(2x)))$$

$$u_y = u_{\xi\eta}(2y) + u_\eta(2y).$$

Characteristic curve: A smooth curve  $\vec{x}(k) \in \mathbb{R}^2$  is called a characteristic curve for the 2nd order PDE. (★) if its tangent vector

$\vec{x} = (x, y)^T$  satisfy the quadratic  $ay^2 - by + cy = 0$

$$a u_{xx} + b u_{xy} + c u_{yy} + d u_x + e u_y + f u = g.$$

$$\Delta = b^2 - 4ac.$$

$$\text{Q } (\xi, \eta) \Rightarrow a\xi^2 + b\xi\eta + c\eta^2 + \dots + fu = 0.$$

$$a\left(\frac{dy}{dx}\right)^2 - b\left(\frac{dy}{dx}\right) + c = 0$$

- if  $b^2 - 4ac > 0$  + 2 real solns

$$\xi = \xi(x, y), \eta = \eta(x, y)$$

PDE reduces to  $u_{\xi\eta} = F$

characteristic curve.

and PDE of hyperbolic type.

- if  $b^2 - 4ac = 0$  then one real root parabolic type.

- if  $b^2 - 4ac < 0$ , 2 imaginary roots, ~~integrating~~ elliptic type.

(I). Choice 1) if we choose ( $b^2 - 4ac > 0$ )

$$\xi_x = + \frac{b + \sqrt{b^2 - 4ac}}{2a}, \xi_y = 0.$$

then  $\tilde{A} = 0$

$$\eta_x = + \frac{b - \sqrt{b^2 - 4ac}}{2a}, \eta_y = 0.$$

then  $\tilde{C} = 0$ .

PDE is hyperbolic type.

$$\tilde{B} = - \frac{\Delta \xi_x \eta_y}{a} \neq 0.$$

(★) reduces to

$$\boxed{u_{\xi\eta} = F}$$

Canonical form.

$$4u_{xx} - 2u_{xy} - 3u_{yy} = 0$$

$$a=1, b=-2, c=-3$$

$$b^2-4ac = 4+12 = 16 > 0$$

∴ the given PDE is hyperbolic.

∴ we expect two real characteristics

∴ consider

$$\frac{dy}{dx} = \frac{b \pm \sqrt{b^2 - 4ac}}{2a} = \frac{-2 \pm 4}{2} = 1, -3$$

$$x - y = c_1$$

$$3x + y = c_2$$

$$\therefore \xi_p = x - y, \eta = 3x + y$$

$$u_x = u_{\xi_p} \cdot 1 + u_{\eta} \cdot 3$$

$$u_y = u_{\xi_p}(-1) + u_{\eta}(1) = -u_{\xi_p} + u_{\eta}$$

$$u_{xx} = u_{\xi_p \xi_p} + u_{\xi_p \eta}(3) + u_{\eta \eta}(3)$$

$$u_{yy} =$$

$$u_{xy} =$$

$$16u_{\xi_p \eta} = 0 \Rightarrow u_{\xi_p \eta} = 0$$

$$u_{\xi_p} = f(\xi_p)$$

$$u = g(\xi_p) + h(\eta)$$

$$u = g(x-y) + h(3x+y)$$

II) Choice 2:  $\Delta = b^2 - 4ac = 0$  (parabolic)

$$\frac{dy}{dx} = \frac{b \pm \sqrt{b^2 - 4ac}}{2a} = \frac{b}{2a}$$

$$N = a\xi_p x^2 + b\xi_p x \xi_y + c\xi_y^2$$

$$= a \left( \xi_p x + \frac{b}{2a} \xi_y \right)^2 \quad \frac{\xi_p}{\xi_y} = -\frac{b}{2a}$$

$$N = 0$$

$$u_y = u_{\xi_p}(2y) + u_{\eta}(2y)$$

choose  $\eta$  such that  $\frac{h_n}{h_y} \neq -\frac{b}{2a} = \underline{\underline{c}} \neq 0$

(ii)  $u_{xx} + 2u_{xy} + u_{yy} + u_x = 0$

$$a=1, b=2, c=1$$

$$b^2 - 4ac = 0, \quad \frac{dy}{dx} = \frac{b}{2a} = 1$$

$$x-y = c_1 \Rightarrow \xi_p(x, y) = x-y$$

along which  $\tilde{A}=0$ .

choose  $\eta$  such that  $(\xi_p, \eta)$  are independent

let  $\eta = y$  (we can take  $x+y$  also, take  $y$  to make calculation easy)

$$u_x = u_{\xi_p} + u_n(0) = u_{\xi_p}, \quad u_{xx} = u_{\xi_p \xi_p}$$

$$u_y = -u_{\xi_p} + u_n, \quad u_{xy} = u_{\xi_p \xi_p}(-1) + u_{\xi_p n}(1)$$

$$u_{yy} = u_{\xi_p \xi_p} - 2u_{\xi_p n} + u_{nn}$$

$$\therefore u_{\xi_p \xi_p} - 2u_{\xi_p n} + 2u_{nn} + u_{\xi_p \xi_p} - 2u_{\xi_p n} + u_{nn} + u_{\xi_p} = 0$$

$$\boxed{u_{nn} + u_{\xi_p} = 0}$$

Canonical form.

$$u_n = - \int u_{\xi_p} d\eta + g(\xi_p)$$

$$u = - \iint u_{\xi_p} d\eta d\eta + ng(\xi_p) + h(\xi_p)$$

choice 3)  $\Delta = b^2 - 4ac < 0$  (elliptic)

$$\frac{dy}{dx} = \frac{b \pm \sqrt{b^2 - 4ac}}{2a} \quad \text{gives two linearly independent soln.}$$

$\xi_p, \bar{n}$  but are imaginary

$$\xi_p = \bar{\xi}_p + \bar{n}, \quad \eta = (\bar{\xi}_p - \bar{n})i \quad \text{are real}$$

(contd)

$$a) \quad u_{xx} + 2u_{xy} + u_{yy} = 0$$

$$a=1, \quad b=2, \quad c=2$$

$$b^2 - 4ac = -4$$

$$\frac{dy}{dx} = \frac{1+i}{2-i} = 1+i$$

$$\bar{c}_p = y - (1+i)x \quad \bar{n} = y - (1-i)x$$

$$c_p = 2(y-x), \quad n = 2x$$

$$c_p = 2y - 2x, \quad n = 2x.$$

$$u_x = u_{cp}(-2) + u_n(2).$$

$$u_{xx} =$$

$$u_y = u_{cp}(2) + u_n(0)$$

$$u_{xy} = -2$$

$$\underline{u_{cp} + u_{nn} = 0}$$

Parabolic-

diffusion process

$$c = c(\bar{x}, A)$$

chemical / nutrient conduction process.

$$T = T(\bar{x}, A)$$

temperature.



$$\vec{j} = -k \nabla T$$

$$\vec{j} = -\nabla(k \nabla T) = -k \nabla^2 T \quad (k \text{ const})$$

Rate of change of temp. = total flux accumulated

$$\frac{\partial T}{\partial A} = k \nabla^2 T$$

$$m \cdot 1D : \frac{\partial T}{\partial A} = k \frac{\partial^2 T}{\partial x^2} \quad (\text{gives parabola})$$

Such that  $J \neq 0$

$$(x, y)$$

so

$$dxdy = J dx dy$$

Steady conduction

$$\nabla^2 T = 0$$

(elliptic)

$$\nabla \cdot \vec{V} = 0 \quad \vec{V} = \nabla \phi$$

$$\nabla \cdot \vec{V} = 0$$

$$\nabla^2 \phi = 0$$

wave eq<sup>n</sup>

$$\frac{\partial^2 u}{\partial A^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

# Sol<sup>n</sup> of 2<sup>nd</sup> order PDE

Motivation to separation of variables.

$$2u_{xx} - u_{yy} = 0 \quad (\text{linear, homogeneous})$$

$$u = u(x, y) \quad \text{let } u \text{ be } X(x)Y(y)$$

$$u_x = x' Y \quad ; \quad u_y = x Y'$$

$$\therefore u_{xx} = x'' Y \quad ; \quad u_{yy} = x Y''$$

$$u_x = \frac{\partial u}{\partial x}$$

$$u_y = \frac{\partial u}{\partial y}$$

$$2x'' Y - x Y' = 0. \quad \textcircled{2}$$

looking for a non trivial sol<sup>n</sup>  $\neq xY \neq 0$

hence divide by  $xY$ .

$$1 - \frac{2x''}{x} - \frac{Y'}{Y} = 0. \quad \rightarrow \frac{2x''}{x} = \frac{Y'}{Y}.$$

$\uparrow$   $\uparrow$   
f(x) only f(y) only.

$$\therefore \frac{2x''}{x} = \frac{Y'}{Y} = d.$$

$$\rightarrow x'' - \frac{d}{2} x = 0, \quad Y'' - dY = 0 \quad (\text{eigen value problems})$$

why eigen value  
in general

$$A x = d x.$$
  
 $\uparrow$   
transform.

here,

$$L x = d x$$
  
 $\uparrow \quad \uparrow$   
operator eigen value

$$u_{xx} + 2u_{xy} + u_{yy} = 0$$

$$\underline{d=0 \Rightarrow x''=0} \quad \underline{x(x) = c_1 x + c_2}$$

$$\underline{d>0 \Rightarrow x'' - \frac{d}{2}x = 0} \quad \underline{x(x) = c_1 e^{\sqrt{d/2}x} + c_2 e^{-\sqrt{d/2}x}}$$

$$\underline{d<0 \Rightarrow d=-m^2}$$

$$x'' + \frac{m^2}{2}x = 0$$

$$x(x) = c_1 \cos\left(\frac{m}{\sqrt{2}}x\right) + c_2 \sin\left(\frac{m}{\sqrt{2}}x\right)$$

$$\underline{d=0 \Rightarrow y'=0} \quad \underline{y(y)=d_1}$$

$$\underline{d>0 \Rightarrow y' = dy}$$

for  $d=0$

$$u(x,y) = (c_1 x + c_2)d_1 = \underline{a_1 x + a_2}$$

$$\underline{d>0 \Rightarrow y' = dy}$$

$$y = d_1 e^{dy}$$

$$\underline{d \leq 0} \quad y = d_2 e^{-m^2 y}$$

$$u(x,y) = (d_1 e^{dy}) (c_1 e^{\sqrt{d/2}x} + c_2 e^{-\sqrt{d/2}x})$$

on eigenvalues

$$u_{xx} - u_y = 0 \quad , \quad 0 < x < L.$$

$$u(x,0) = f(x) \quad ; \quad u(0,y) = 0 \quad ; \quad u(L,y) = 0.$$

$$u = xy \quad \frac{x''}{x} = \frac{y'}{y} = d.$$

$$d=0 \quad x(x) = c_1 x + c_2$$

$$c_2 = 0 \quad \therefore u(0,y) = 0$$

$$x(0) \cdot y(y) = 0$$

$$x(0) = 0 \quad \therefore y(y) \neq 0.$$

$$u(L,y) = 0 \cdot x(L) \cdot y(y).$$

$$= c_L = 0$$

$$\underline{c_L = 0}$$

$\therefore$  trivial soln for  $d=0$ .

$\lambda > 0$

$$x(x) = C_1 e^{\sqrt{\lambda}x} + C_2 e^{-\sqrt{\lambda}x}$$

$$x(0) = 0 \quad \therefore C_1 + C_2 = 0$$

$$x(L) = 0 \Rightarrow C_1 e^{\sqrt{\lambda}L} + C_2 e^{-\sqrt{\lambda}L} = 0$$

$$C_1 = 0, C_2 = 0$$

$\lambda < 0$

$\lambda = -m^2$

$$x(x) = C_1 \cos mx + C_2 \sin mx$$

$$x(0) = 0 \Rightarrow C_1 = 0$$

$$x(L) = C_2 \sin mL$$

$$x(L) = 0 \quad \therefore C_2 \sin mL = 0$$

We are looking for non trivial sol<sup>n</sup>

$$\therefore C_2 \neq 0, \quad mL = n\pi$$

$$m = \frac{n\pi}{L}, \quad n = 1, 2, \dots$$

$$\lambda = -m^2 = -\frac{n^2\pi^2}{L^2}$$

( $\lambda = 0$  is excluded so  $n$  starts from 1, 2, ...)

$$\therefore \lambda_n = -\left(\frac{n\pi}{L}\right)^2$$

( $\because$  there are multiple eigen values).

eigen functions corresponding to  $\lambda_n$   
 $\sin\left(\frac{n\pi}{L}x\right)$ .

there's no sense to compute  $y(y)$  for  $\lambda > 0$

( $\because x(x)$  is zero &  $u = xy$ )

$$\lambda < 0, \quad d = -m^2$$

$$y' - dy = 0 \quad \therefore y = C_1 \cos mx + C_2 \sin mx$$

$$\therefore y = d_1 e^{-mx}$$

Note if  $u(0,y) = 1$

$$\Rightarrow x(0) \cdot y(y) = 1$$

doesn't mean  $x(0) = 1$ .

instead we have infinite sol<sup>n</sup>.

$$u = xy.$$

$$= C_2 \sin\left(\frac{n\pi x}{L}\right) d_1 e^{-\left(\frac{n\pi}{L}\right)^2 y}.$$

$$u_n(x,y) = A_n \sin\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 y}.$$

because each eigen value yields a sol<sup>n</sup> u.

∴ the general sol<sup>n</sup> is

$$u(x,y) = \sum_{n=1}^{\infty} u_n(x,y)$$

$$u(x,y) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 y}$$

$$u(x,0) = f(x)$$

$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) \quad (\text{fourier series})$$

Recall  $f(x) = \sum_{n=0}^{\infty} a_n (x-a)^n$ , then

$$a_n = \frac{f^{(n)}(a)}{n!} \quad \text{Taylor's series}$$

along those lines, we wish to expand a function  $f(x)$  as  $[-\pi, \pi]$  which is periodic (i.e.  $f(x+\pi) = f(x)$ ) as  $f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$

I need to determine  $a_0, a_n, b_n$  ( $n \geq 1$ ) assume that  $\star$  can be integrated term by term.

Simply integrate  $\star$  from  $-\pi$  to  $\pi$

$$\int_{-\pi}^{\pi} f(x) dx = \int_{-\pi}^{\pi} a_0 dx + \sum_{n=1}^{\infty} a_n \int_{-\pi}^{\pi} \cos nx dx + b_n \int_{-\pi}^{\pi} \sin nx dx$$

$$\int_{-\pi}^{\pi} f(x) dx = a_0 \int_{-\pi}^{\pi} dx = a_0 \cdot 2\pi.$$

$$\underline{a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx.}$$

$$f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

Multiply  $\star$  with  $\cos mx$  and  $\int_{-\pi}^{\pi}$

$$\int_{-\pi}^{\pi} f(x) \cos mx dx = a_0 \int_{-\pi}^{\pi} \cos mx dx + \sum_{n=1}^{\infty} a_n \int_{-\pi}^{\pi} \cos nx \cos mx dx + \sum_{n=1}^{\infty} b_n \int_{-\pi}^{\pi} \sin nx \cos mx dx$$

$$\int_{-\pi}^{\pi} \sin nx \cos mx dx = 0. \quad \text{if } m, n.$$

$$\int_{-\pi}^{\pi} f(x) \cos nx dx = a_n \pi.$$

$$+ a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx dx.$$

$$b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin nx dx.$$

$f$  piecewise continuous on  $(-\pi, \pi)$  and periodic

$$\text{Eg. } f(x) = \begin{cases} 0 & -\pi \leq x < 0 \\ 1 & 0 \leq x < \pi \end{cases}$$

$$f(x+2\pi) = f(x).$$

$$a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx = \frac{1}{2\pi} \left[ \int_{-\pi}^0 0 dx + \int_0^{\pi} 1 dx \right] = \frac{\pi}{2\pi} = \underline{\underline{\frac{1}{2}}}.$$

$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos nx dx$$

$$= \frac{1}{\pi} \left[ \int_0^{\pi} 1 \cdot \cos nx dx \right]$$

$$= \frac{1}{\pi} \left[ \frac{\sin nx}{n} \right]_0^{\pi} = 0$$

$$b_n = \frac{1}{\pi} \int_0^{\pi} \sin nx dx$$

$$= -\frac{1}{\pi} \left[ \frac{\cos nx}{n} \right]_0^{\pi} = -\frac{1}{n\pi} (\cos n\pi - 1)$$

$$= -\frac{1}{n\pi} (\cos n\pi - 1) = 0, n \text{ is even}$$

$$\frac{2}{n\pi}, n \text{ is odd}$$

$$\therefore f(x) = \frac{1}{2} + \frac{2}{\pi} \sin x + \frac{2}{3\pi} \sin 3x + \frac{2}{5\pi} \sin 5x \dots$$

heat conduction in a rod (wire) of finite length  $0 \leq x \leq L$ .

assumptions: homogenous material thin so that heat distribution is uniform

$$x=0 \quad x=L$$

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < L$$

$$u = u(x, t)$$

defined as an Initial Boundary Value Problem.

init. condition--  $u(x, 0) = f(x)$ .

fixed ends (fixed temperature)

homogeneous  
b.c.

$$u(0, t) = 0 \quad \forall t > 0.$$

insulated ends

$$\frac{\partial u}{\partial x}(0, t) = 0 \quad \forall t > 0.$$

$$\frac{\partial u}{\partial x}(L, t) = 0$$

I.B.V.P.

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < L.$$

i.e.  $u(x, 0) = f(x)$ .

b.c.  $u(0, t) = 0 \quad \forall t > 0.$

$u(L, t) = 0$

let  $u(x, t) = X(x) T(t)$

$$\underline{X T' = k T X''}$$

$$\frac{X'}{X''} = \frac{1}{k} \frac{T'}{T} = \text{const.} = d.$$

$$\underline{k = 1}$$

$$X'' - \lambda X = 0$$

$$u(0,t) = X(0)T(t) = 0$$

$$\underline{X(0) = 0}$$

$$u(L,t) = X(L)T(t) = 0$$

$$\underline{\underline{X(L) = 0}}$$

$\lambda = 0$

$$X(x) = ax + b$$

$$X(0) = 0, \quad X(L) = 0 \quad \rightarrow a = 0 \\ b = 0.$$

trivial solution.

$$\lambda > 0 \quad X(x) = c_1 e^{\sqrt{\lambda}x} + c_2 e^{-\sqrt{\lambda}x}$$

$$X(0) = 0 = c_1 + c_2$$

$$X(L) = 0 = c_1 e^{\sqrt{\lambda}L} + c_2 e^{-\sqrt{\lambda}L}$$

$$\underline{c_1 = 0, c_2 = 0}$$

trivial soln.

$$\lambda < 0, \quad \underline{\lambda = -m^2} \quad (\text{say})$$

$$X(x) = c_1 \cos mx + c_2 \sin mx.$$

$$X(0) = 0 \rightarrow c_1 = 0$$

$$X(L) = 0 \rightarrow c_2 \sin mL \quad X(L) = 0 = c_2 \sin mL.$$

non-trivial soln

$$c_2 \neq 0 \quad \sin mL = 0$$

$$\underline{m = \frac{n\pi}{L}} \quad n = 1, 2, \dots$$

$$\text{for } \lambda < 0, \quad = -m^2,$$

$$\frac{T'}{T} = d = -m^2$$

$$\rightarrow T = d_1 e^{-m^2 t} = d_1 e^{-(n\pi/L)^2 t}$$

$$u_n(x,t) = X_n(t) T_n(x)$$

$$= C_n \sin\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 t} \quad \text{Cn is constant}$$

$$\therefore u(x,t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{L}\right) e^{-\left(\frac{n\pi}{L}\right)^2 t}$$

need to determine  $C_n$  from i.e.  $u(x,0) = f(x)$ .

$$f(x) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{L}\right).$$

Now

$$\int_0^L f(x) \sin \frac{m\pi x}{L} dx = \sum_{n=1}^{\infty} C_n \int_0^L \sin \frac{m\pi x}{L} \sin \frac{n\pi x}{L} dx$$

$$+ C_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx.$$

$$\text{Suppose } f(x) = (1-x)$$

$$C_n = \frac{2}{L} \int_0^L (1-x) \sin\left(\frac{n\pi x}{L}\right) dx.$$

$$= \frac{2}{L} \int_0^L \sin\left(\frac{n\pi x}{L}\right) dx - \underbrace{\frac{2}{L} \int_0^L x \sin\left(\frac{n\pi x}{L}\right) dx}.$$

$$\int_0^L x \sin\left(\frac{n\pi x}{L}\right) dx = \left[ -\frac{Lx \cos\left(\frac{n\pi x}{L}\right)}{n\pi} \right]_0^L + \frac{L}{n\pi} \int_0^L \cos\left(\frac{n\pi x}{L}\right) dx$$

$$\begin{aligned} \int_0^L \sin\left(\frac{n\pi x}{L}\right) dx &= \frac{L}{n\pi} \left[ -\cos\left(\frac{n\pi x}{L}\right) \right]_0^L = -\frac{L^2}{n\pi} \cos(n\pi) + \frac{L}{n\pi} \left[ \frac{L}{n\pi} \sin(n\pi) \right] \\ &\quad + \frac{L}{n\pi} \left( -\frac{\cos n\pi}{n\pi} + 1 \right) = -\frac{L^2}{n\pi} \cos(n\pi) \end{aligned}$$

$$\Rightarrow \frac{2}{L} \left[ -\frac{L^2}{n\pi} \cos(n\pi) + \frac{L}{n\pi} \left( -\frac{\cos n\pi}{n\pi} + 1 \right) \right]$$

solve

## # Heat conduction equation, - insulated ends.

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < L$$

$$u(x, 0) = f(x)$$

$$\frac{\partial u}{\partial x}(0, t) = 0 \quad ; \quad \frac{\partial u}{\partial x}(L, t) = 0 \quad \forall t > 0$$

$$\text{look for } u(x, t) = X(x) T(t)$$

$$X T' = X'' T \quad \left. \begin{array}{l} X'' - \lambda X = 0 \\ T' - k \lambda T = 0 \end{array} \right\}$$

$$\underline{b-c} \quad \frac{\partial u}{\partial x} = \frac{\partial}{\partial x}(X T) = X' T$$

$$\frac{\partial u}{\partial x}(0, t) = X'(0) T(t) = 0$$

$$\frac{\partial u}{\partial x}(L, t) = 0 \quad \left. \begin{array}{l} X'(0) = 0 \\ X'(L) = 0 \end{array} \right\}$$

$$\underline{n=0} \quad X'' - \lambda X = 0 \quad \& \quad X(x) = ax + b$$

$$\begin{aligned} X'(0) &= 0 \Rightarrow a \\ X'(L) &= 0 \quad (\text{satisfied}) \end{aligned} \quad \left. \begin{array}{l} \\ \end{array} \right\}$$

$\therefore b$  is arbitrary.

$$\therefore X(x) = b. \quad \text{non-trivial soln.}$$

$$\underline{d>0} \quad X(x) = C_1 e^{\sqrt{\lambda} x} + C_2 e^{-\sqrt{\lambda} x}$$

$$X'(x) = \sqrt{\lambda} (C_1 e^{\sqrt{\lambda} x} + C_2 e^{-\sqrt{\lambda} x})$$

$$X'(0) = 0 \Rightarrow \sqrt{\lambda} (C_1 - C_2) = 0$$

$$x'(L)=0 = \sqrt{d} (c_1 e^{\sqrt{d}L} - c_2 e^{-\sqrt{d}L}).$$

$$\therefore c_1=0, c_2=0.$$

∴ trivial soln

$$\underline{d \neq 0} \quad x(x,t), d = -m^2$$

$$x(x) = c_1 \cos mx + c_2 \sin mx.$$

$$x' = m(-c_1 \sin mx + c_2 \cos mx)$$

$$x' = -mc_1 \sin mx + mc_2 \cos mx.$$

$$x'(0) = 0 = c_2$$

$$x'(L) = 0 \Rightarrow mc_1 \sin mL$$

$$+ m = \frac{n\pi}{L}, n=1, 2, \dots$$

$$\therefore d_n = -\left(\frac{n\pi}{L}\right)^2, n=1, 2, \dots$$

$$x_n = \cos \frac{n\pi}{L} x$$

$$\therefore u_n(x, t) = x_n T_n$$

$$\underline{d=0} \quad T' - dT = 0 \quad \therefore T = d,$$

$$\underline{d \neq 0} \quad T' + m^2 T = 0 \quad \therefore T = d_2 e^{-m^2 t}$$

$$u_n(x, t) = x_n T_n + x_n T_n$$

$$(d=0) \quad (d \neq 0)$$

$$\therefore u_n(x, t) = b d_1 + c_n \cos \left( \frac{n\pi x}{L} \right) e^{-\left(\frac{n\pi}{L}\right)^2 t}$$

$$u(x, t) = A_0 + \sum_{n=1}^{\infty} c_n \cos \left( \frac{n\pi x}{L} \right) e^{-\left(\frac{n\pi}{L}\right)^2 t}$$

$$= \sum_{n=0}^{\infty} A_n \cos \left( \frac{n\pi x}{L} \right) e^{-\left(\frac{n\pi}{L}\right)^2 t}$$

$$\underline{x} \quad u(x, 0) = f(x)$$

$$u(x) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right).$$

$$\therefore A_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$

Let  $f(x) = x$

$$= \frac{2}{L} \int_0^L x \cos\left(\frac{n\pi x}{L}\right) dx$$

1)  $\underline{u(0,t) = 0}$

$$\frac{\partial u}{\partial x}(L, t) = 0$$

$$u(x, 0) = \begin{cases} 1, & 0 < x < L/2 \\ 2, & L/2 < x < L \end{cases}$$

2).  $\frac{\partial u}{\partial x}(0, t) = 0$

$$\frac{\partial u}{\partial x}(L, t) = 0$$

$$u(x, 0) = \begin{cases} x, & 0 < x < L/2 \\ 1-x, & L/2 < x < L \end{cases}$$

2)  $x'' - \lambda x = 0$   
 $T'' - \lambda T = 0$

~~for~~  $x'(0, t) = 0$  } init.  
 $x'(L, t) = 0$ .

Solving for x

$\lambda = 0$   $x = \underline{b} ax + b$

b is arbitrary

$$x(x) = \underline{b}$$

$\lambda > 0$  no sol.

$\lambda < 0$   $x_n = \cos\left(\frac{n\pi x}{L}\right)$

1)  $\lambda = 0$   $T = d_1$

$\lambda < 0$   $T = d_2 e^{-m^2 t}$

$$u_n(x, t) = \underline{A_n} \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) e^{-\frac{(n\pi)^2 t}{L}}$$

$$u(x,t) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) = \begin{cases} x & , 0 < x < L/2 \\ 1-x & , L/2 < x < L \end{cases}$$

$$A_n = \frac{2}{L} \int_0^{L/2} x \cos\left(\frac{n\pi x}{L}\right) dx + \frac{2}{L} \int_{L/2}^L (1-x) \cos\left(\frac{n\pi x}{L}\right) dx.$$

$$= \frac{2}{L} \left[ \int_0^{L/2} x \cos\left(\frac{n\pi x}{L}\right) dx + \int_{L/2}^L (1-x) \cos\left(\frac{n\pi x}{L}\right) dx \right].$$

$$\int_0^{L/2} x \cos\left(\frac{n\pi x}{L}\right) dx = \left[ \frac{L}{n\pi} \cdot x \sin\left(\frac{n\pi x}{L}\right) \right]_0^{L/2} - \int_0^{L/2} \frac{L}{n\pi} \sin\left(\frac{n\pi x}{L}\right) dx$$

$$= \frac{L^2}{2n\pi} \sin\left(\frac{n\pi}{2}\right) + \frac{L^2}{(n\pi)^2} \left[ \cos\left(\frac{n\pi x}{L}\right) \right]_0^{L/2}$$

1). b.c.  $x(0)=0$   $\underline{x'(L)}=0$

$$x'' - dx = 0$$

$$\underline{d=0} \quad x = ax + b \quad \underline{d \neq 0} \quad \underline{d \neq 0} \quad c_1 + c_2 = 0$$

$$x'(L) = 0 \quad \rightarrow a = 0 \quad c_1 e^{J\lambda L} - c_2 e^{-J\lambda L} = 0$$

$$x = b$$

$$\underline{c_1 = c_2 = 0}$$

$$\underline{d \neq 0} \quad x(x) = c_1 \cos mx + c_2 \sin mx.$$

$$\underline{c_1 = 0} ; m c_2 \sin mL = 0$$

$\rightarrow m$  arbitrary.

$$x(x) = \underline{\underline{c_2 \sin mx}}.$$

(I). Determine the region in which the given eq<sup>n</sup> is hyperbolic, parabolic, elliptic. Identify the transformation in respective regions that transforms the given eq<sup>n</sup> to canonical form. Reduce it to corresponding canonical form.

(a)  $\alpha u_{xx} + u_{yy} = x^2$

(b)  $u_{xx} + y^2 u_{yy} = y$

(c)  $x^2 u_{xx} - 2xy u_{xy} + y^2 u_{yy} = e^x$

(d)  $u_{xx} - \sqrt{y} u_{xy} + x u_{yy} = \cos(x^2 - 2y), \quad y > 0$

## # Exercises on classification -

The type of PDE is a local property.

An equation can change its behaviour in different regions of a plane. Can change its behaviour depending on a parameter.

Eg 1)

Tricomi's eq<sup>n</sup>

$$y \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

$$\begin{array}{l} a=y \\ b=0 \\ c=1 \end{array}$$

$$b^2 - 4ac = -4y$$

$y=0$  parabolic

$y < 0$  hyperbolic

$y > 0$  elliptic

Eg 2)

$$\frac{1}{1-m^2} \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = 0$$

$$b^2 - 4ac = -\frac{1}{1-m^2}$$

$m^2 > 1$  hyperbolic

$m^2 < 1$  elliptic

Eg 3)

Reduce to canonical form  $y^2 u_{xx} - 2xy u_{xy} + x^2 u_{yy} = \frac{1}{xy} (y^3 u_x + x^2 u_y)$

$$a=y^2 \quad b=-2xy \quad c=x^2$$

$$b^2 - 4ac = 4x^2 y^2 - 4x^2 y^2 = 0. \text{ parabolic.}$$

$$\frac{dy}{dx} = \frac{b \pm \sqrt{b^2 - 4ac}}{2a} = \frac{-2xy}{2y^2} = -\frac{x}{y}.$$

$\rightarrow \xi = x^2 + y^2$  only one characteristic (real).

Let  $\eta = y$

$$u_x = u_q \cdot 2x + u_{\eta} \cdot 0 = 2x \cdot u_q.$$

$$u_y = 2(u_q + x \cdot 2x u_{qq})$$

$$u_{xy} = 2x(u_{qq} \cdot 2y + u_{q\eta} \cdot 1).$$

$$u_y = u_q \cdot 2y + u_n \cdot 1$$

$$u_{yy} = u_{qq} \cdot 2y + \cancel{u_q \cdot 2} + \cancel{u_{q\eta} \cdot 2y}.$$

$$2(u_q + y u_{qq} \cdot 2y + y u_{q\eta} + u_{qp} \cdot 2y + u_{nn})$$

$$u_{nn} - \frac{1}{\eta} u_n = 0 \quad \text{Canonical / standard form.}$$

$$\eta u_{nn} - u_n = 0.$$

$$\frac{\partial}{\partial \eta} (\eta u_n - 2u) = 0.$$

$$\eta u_n - 2u = f(\frac{e_1}{y})$$

~~rather do integration by parts.~~

$$\cancel{u_y} \quad \frac{1}{\eta} u_{nn} - \frac{1}{\eta} u_n = 0.$$

$$\frac{\partial}{\partial \eta} \left( \frac{1}{\eta} u_n \right) = 0.$$

$$u_n = \eta f(\frac{e_1}{y})$$

$$u = \eta^2 f(\xi) + g(\eta)$$

~~Hyperbolic~~

$$= y^2 f(x^2 + y^2) + g(x^2 + y^2).$$

Eq 4)  $u_{xx} + \alpha^2 u_{yy} = 0.$

$$a=1 \quad b=0 \quad c=n^2$$

$$b^2 - 4ac = -4\alpha^2$$

$\xrightarrow{n=0}$  parabolic  
 $\xrightarrow{\alpha \neq 0}$  elliptic

$$\frac{dy}{dx} = \pm \frac{\sqrt{-4\alpha^2}}{2} = \pm i\alpha.$$

$$iy = \pm \frac{i\alpha^2}{2} \Rightarrow c_{1,2}.$$

8  $\xi_p(x,y) = y + \frac{i\alpha^2}{2}$        $\eta(x,y) = y - \frac{i\alpha^2}{2}$       imaginary

Need to switch to a plane where the given PDE admits

(this is a characteristic of elliptical form)

a canonical form

$$\alpha = \frac{\xi_p + \eta}{2}, \quad \beta = \frac{\xi_p - \eta}{2i}$$

$$= y = \frac{\alpha^2}{2}.$$

$$\text{so } \underline{\alpha = y, \beta = \frac{\alpha^2}{2}}$$

remove 2 as it can be absorbed in the family

no point working with the factor.

$$u_x = u_\alpha \cdot 0 + u_\beta \cdot 2\alpha = 2u_\beta \alpha$$

$$u_{xx} = 2u_{\beta\beta} + (2\alpha)^2 u_{\beta\beta}$$

$$u_{xy} = 2u_{\beta\alpha} \cdot 2\alpha$$

$$u_{yy} = u_{\alpha\alpha}$$

$$u_{xx} + u^2 u_{yy} = 0.$$

$$2u_\beta + u x^2 u_{\beta\beta} + u^2 u_{\alpha\alpha} = 0.$$

$$+ u_{\alpha\alpha} + 4u u_{\beta\beta} = \frac{-1}{2x^2} u_\beta = \frac{-1}{2\beta} u_\beta.$$

$$\boxed{+ u_{\alpha\alpha} + u_{\beta\beta} = \frac{-1}{2\beta} u_\beta.}$$

Canonical

(5) Reduce  
 $u_{xx} + u_{yy} - 2u_{yy} + 1 = 0.$

in  $0 \leq x \leq 1, y \geq 0.$

to canonical form and solve with  $u = \frac{dy}{dx} = n$  on  $y=0.$

$$a=1, b=1, c=-2$$

$$b^2 - 4ac = 9 > 0. \quad \underline{\text{Hyperbolic}}$$

$$\therefore \frac{dy}{dx} = \frac{1 \pm \sqrt{9}}{2} = \frac{1 \pm 3}{2} = \underline{2, -1}.$$

$$+ \underline{\underline{\eta}} = y - 2x \quad \underline{\underline{\zeta}} = y + x$$

$$+ n =$$

$$y =$$

$$u_x = u_{\zeta\eta}(-2) + u_n(1) = -2u_{\zeta\eta} + u_n.$$

$$u_y = u_{\zeta\eta} + u_n.$$

$$u_{yy} = -2(u_{\zeta\zeta} + u_{\eta\eta}) + u_{\zeta\zeta} + u_{\eta\eta}.$$

$$u_{yy} = u_{\epsilon\epsilon\eta\eta} + u_{\eta\eta\eta\eta} + 2u_{\epsilon\eta\eta\eta}$$

$$u_{xx} = -2u_{\epsilon\epsilon\epsilon\eta} \cdot (-2)$$

$$u_{xx} + u_{yy} - 2u_{yxy} + 1 = 0$$

$$-2u_{\epsilon\eta\eta\eta} + u_{\eta\eta\eta\eta} + u_{\epsilon\epsilon\epsilon\eta} \cdot (-2)$$

$$= \underline{u_{\epsilon\epsilon\epsilon\eta}} - \underline{2u_{\epsilon\eta\eta\eta}} + \underline{u_{\eta\eta\eta\eta}} - \underline{2u_{\epsilon\epsilon\eta\eta}}$$

$$\therefore \underline{u_{\epsilon\eta\eta\eta}} = \frac{1}{g} . \quad \underline{\text{Canonical form}}$$

$$u_{\epsilon\eta\eta\eta} = \frac{1}{g} \eta + f(\epsilon) + f(\eta).$$

$$u = \underline{\frac{1}{g} \eta \epsilon + g(\epsilon) + h(\eta)}.$$

$$u = \frac{1}{g} (y-2x)(y+x) + g(y-2x) + h(y+x).$$

$$x = \underline{-\frac{2}{g} x^2 + g(-2x) + h(x)}$$

① Solve

$$u_t = g u_{xx}, \quad 0 < x < 2$$

$t > 0$

$$u(0, t) = 0$$

$$\frac{\partial u}{\partial x}(2, t) = 0$$

$$u(x, 0) = x$$

~~boundary conditions~~

$$u(0, t) = 0$$

$$u(x, 0) = 0 \quad x$$

$$x'' - \lambda x = 0$$

for  $\lambda = 0$

$$x = ax + b$$

$$x'(0) = 0 = a$$

$$x'(2) = 0 = a$$

$$\Rightarrow a = 0$$

$$b = \text{arbitrary } 0 \quad \text{as } x(0) = 0.$$

for  $\lambda > 0$

$$x(x) = C_1 e^{\sqrt{\lambda} x} + C_2 e^{-\sqrt{\lambda} x}$$

$$x'(x) = C_1 \sqrt{\lambda} e^{\sqrt{\lambda} x} + -C_2 \sqrt{\lambda} e^{-\sqrt{\lambda} x}$$

$$x'(0) = C_1 \sqrt{\lambda} e^{2\sqrt{\lambda}} - C_2 \sqrt{\lambda} e^{-2\sqrt{\lambda}}$$

$$\cancel{x'(0)} = C_1 - C_2 = 0 \quad (x(0)) = 0.$$

$$+ \quad \underline{C_1 = C_2 = 0}$$

for  $\lambda < 0$

$$\lambda = -m^2$$

$$x(x) = C_1 \cos mx + C_2 \sin mx.$$

$$x'(x) = -C_1 m \sin mx + C_2 m \cos mx$$

$$x'(0) = C_2 m \cancel{\cos 0}$$

$$\cancel{+ C_2 = 0}$$

$$x'(2) = 0$$

$$- \cancel{C_1 \sin(2m)} \quad C_2 m \cos(2m) = 0.$$

$$2m = (2n+1)\frac{\pi}{2}$$

$$m = \frac{(2n+1)\pi}{4} ; n = 0, 1, 2, \dots$$

$$T' - \lambda T = 0$$

for  $\lambda > 0$

$$Y(x) = 0$$

so no T needed

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$$T' + m^2 T = 0$$

$$T = \underline{\underline{d_2 e^{-m^2 t}}}$$

sint explanation

let  $u(x,t) = \underbrace{x(x)}_{\text{func of } x} \cdot \underbrace{T(t)}_{\text{func of } t}$

$$\text{PDE} = \frac{\partial u}{\partial t} = g \frac{\partial^2 u}{\partial x^2}$$

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial t} (x(x) T(t)) = x(x) \frac{\partial T}{\partial t} = x(x) T'(t)$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial^2}{\partial x^2} (x(x) T(t)) = \frac{\partial^2 x}{\partial x^2} (T(t))$$

$$= x''(x) T(t)$$

$$+ x(x) T'(t) = g x''(x) T(t)$$

$$\rightarrow f_x \cdot f_x = f_x \cdot f_t$$

$$\rightarrow \frac{T'}{g T} = \frac{x''}{x} = \lambda \quad \text{const.}$$

$$\boxed{x'' - \lambda x = 0 \quad ; \quad T' - g \lambda T = 0}$$

$$\underline{b.c.} \quad u(0,t) = 0 \quad \rightarrow x(0) T(t) = 0.$$

for non trivial soln

$$T(t) \neq 0 \quad \rightarrow \boxed{x(0) = 0}$$

$$\frac{\partial u}{\partial x}(2,t) = 0 \rightarrow x'(2) T(t) = 0$$

$$\boxed{x'(2) = 0}$$

to solve.

$$x'' - dx = 0, \quad x(0) = 0, \quad x'(2) = 0.$$

$d=0$   $\rightarrow x'' = 0 \rightarrow x(x) = ax + b$

$$x(0) = 0 \rightarrow b = 0$$

$$x'(2) = 0 \rightarrow \underline{a = 0}.$$

$\rightarrow$  trivial solution.

$d > 0$   $\rightarrow x(x) = c_1 e^{\sqrt{d}x} + c_2 e^{-\sqrt{d}x}$

$$x(0) = 0 \rightarrow c_1 + c_2 = 0$$

$$x'(2) = 0 \rightarrow d(c_1 e^{2\sqrt{d}} - c_2 e^{-2\sqrt{d}})$$

$$\rightarrow \underline{c_1 = 0, c_2 = 0} \rightarrow$$

trivial solution.

$d < 0$  let  $\underline{d = -m^2}$

$$x(x) = c_1 \cos mx + c_2 \sin mx.$$

$$x(0) = 0 \rightarrow \underline{c_1 = 0}$$

$$x'(x) = c_2 \sin mx$$

$$x'(2) = m c_2 \cos 2m$$

$$x'(2) = 0 = m c_2 \cos 2m.$$

non trivial solns  $\rightarrow c_2 \neq 0$

$$2m = (2n+1)\pi/2$$

$$m = (2n+1)\pi/4$$

$$\therefore \text{Eigenvalues : } \underline{\lambda_n = \left[ (2n+1)\pi/4 \right]^2}$$

eigen functions :  $x_n = \frac{\cos((2n+1)\pi x)}{2}$

$$\sin \left( \frac{(2n+1)\pi}{4} x \right)$$

b)  $T' = g \Delta T = 0$

$$T = B e^{g \Delta T} = B e^{-g \left[ \frac{(2n+1)\pi}{4} \right]^2 t}$$

$$u(x,t) = \sum u_n \\ = \sum_{n=0}^{\infty} A_n \sin \left( \frac{(2n+1)\pi}{4} x \right) e^{-g \left[ \frac{(2n+1)\pi}{4} \right]^2 t}$$

$$u(x,0) = x$$

$$x = \sum_{n=0}^{\infty} A_n \sin \left( \frac{(2n+1)\pi}{4} x \right).$$

$$A_n = \frac{2}{2} \int_0^2 x \sin \left( \frac{(2n+1)\pi}{4} x \right) dx.$$

This was one sense end insulated and other end free inflow case

This can be extended to more dimensions

$$\frac{\partial u}{\partial t} = \nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}.$$

### Wave Equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}, \quad x \in \mathbb{R}$$

$-\infty < x < \infty$

$t > 0$

$$u = u(x, t)$$

↳ displacement

initial data

initial displacement  $u(x, 0) = f(x)$

initial velocity  $\frac{du}{dt}(x, 0) = g(x)$

(infinite string)  $x \in \mathbb{R}$ .

↓  
this is one of the ~~the~~ 3 kinds of possible cases

Note:

infinite string  $\rightarrow$  vanishing disturbances at far field at  $x \rightarrow \infty$ .

hyperbolic  $\frac{c^2 \partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial t^2} = 0$

$A = c^2 \quad B = 0 \quad C = -1$

$b^2 - 4AC = 4 \underline{c^2} > 0$

$$\frac{dt}{dx} = \frac{B \pm \sqrt{B^2 - 4AC}}{2A} = \frac{2c}{2c^2} = \pm \frac{1}{c}$$

$dt = \pm \frac{dx}{c}$

$x - ct = c_1 = \underline{\xi_1(x, t)}$

$x + ct = c_2 = \underline{\eta_1(x, t)}$

two real characteristics.

$u_x = u_{\xi_1} + u_{\eta_1}$

$u_t = -cu_{\xi_1} + cu_{\eta_1}$

$\underline{u_{xx} = u_{\xi_1\xi_1} + u_{\eta_1\eta_1} + 2u_{\xi_1\eta_1}}$

$\underline{u_{tt} = c^2 u_{\xi_1\xi_1} - c^2 u_{\xi_1\eta_1} + c^2 u_{\eta_1\eta_1} - c^2 u_{\xi_1\eta_1}}$

$\underline{= c^2 u_{\xi_1\xi_1} + c^2 u_{\eta_1\eta_1} - 2c^2 u_{\xi_1\eta_1}}$

$c^2 u_{nn} - u_{tt} = 0$

$c^2 u_{\xi_1\xi_1} + c^2 u_{\eta_1\eta_1} + 2c^2 u_{\xi_1\eta_1} - c^2 u_{\xi_1\xi_1} - c^2 u_{\eta_1\eta_1} + 2c^2 u_{\xi_1\eta_1} = 0$

$$\underline{4\mu c^2 u_{\eta\eta} = 0.}$$

(Canonical form)

$$\star \boxed{u_{\eta\eta} = 0}$$

$$\star u_{\eta\eta} = f(\eta)$$

$$\star u = \underline{g(\eta) + h(\eta)}$$

$$\therefore u(x,t) = \underline{f(x-ct) + g(x+ct)}$$

$$u(x,t) = f(x+ct) + g(x-ct)$$

$$u(x,t) = \underline{\phi(x+ct) + \psi(x-ct)}$$

Cond's: 1).  $u(x,0) = f(x) \rightarrow \phi(x) + \psi(x)$  (A)  $f(x)$  = initial displacement

2).  $u_t(x,0) = c\phi'(x) - c\psi'(x) = g(x)$  (B)  $g(x)$  = initial velocity

integrate (B) from  $x_0$  to  $x$ .

$$\int_{x_0}^x g(J) dJ = c \int_{x_0}^x \phi'(J) dJ - c \int_{x_0}^x \psi'(J) dJ$$

$$= c(\phi(x) - \psi(x)) - c(\phi(x_0) - \psi(x_0))$$

~~$$-c\phi(x) + c\phi(x_0)$$~~

$$= 2c\phi'(x) - c(f(x)) - c(\phi(x_0) - \psi(x_0))$$

$$= 2c\phi'(x) - c f(x) - \tilde{k}$$

$$\phi(x) + \frac{1}{2} f(x) + \frac{1}{2c} \int_{x_0}^x g(J) dJ + \frac{\tilde{k}}{2c}$$

$$\tilde{k} = \underline{c(\phi(x_0) - \psi(x_0))}$$

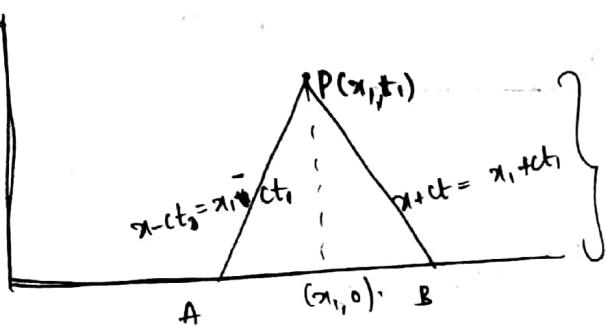
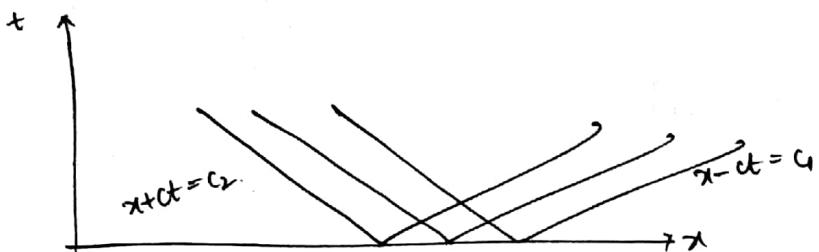
$$\Psi(x) = \frac{1}{2} f(x) - \frac{1}{2c} \int_{x_0}^x g(\xi) d\xi - \frac{k}{2c}$$

$$u(x,t) = \frac{1}{2} f(x+ct) + \frac{1}{2c} \int_{x_0}^{x+ct} g(\xi) d\xi + \frac{k}{2c} + \frac{1}{2} f(x-ct) - \frac{1}{2} \int_{x_0}^{x-ct} g(\xi) d\xi$$

$$u(x,t) = \frac{1}{2} [f(x+ct) + f(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} g(\xi) d\xi$$

(in exams, derive canonical form  
and solve)

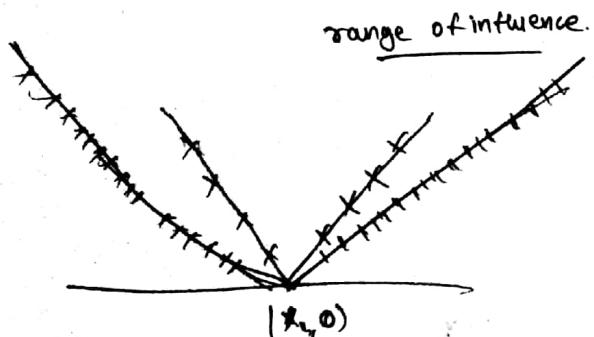
d' Alembert's solution.



this is the complete scope for getting to point P.

$$u(x_1, t_1) = u_p = \frac{1}{2} [f(A) + g(B)] + \frac{1}{2c} \int_A^B g(\xi) d\xi$$

AB: domain of dependence.



d : Atember's Soln

$$u(x,t) = \frac{1}{2} [f(x+ct) + f(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} g(\tau) d\tau.$$

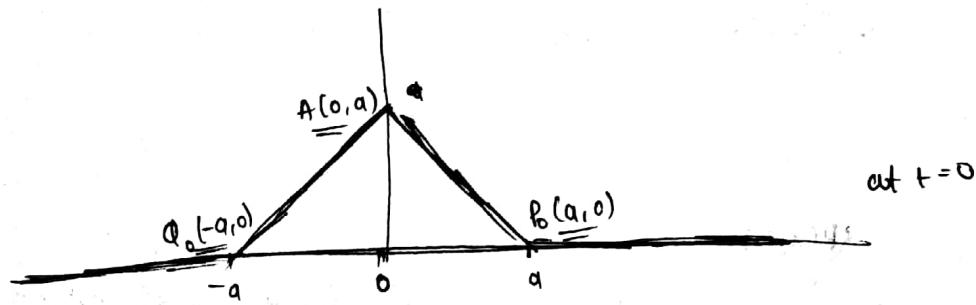
case-1)  $g(x)=0$ . initial velocity = 0.

Let

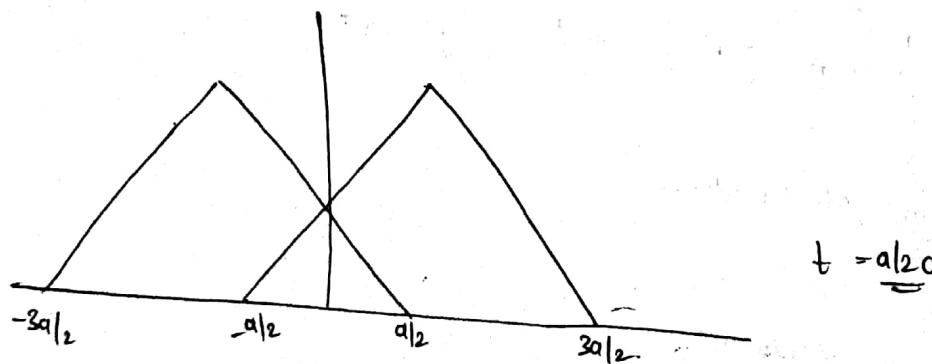
$$f(x) =$$

$$\begin{cases} 0 & , -\alpha \leq x \leq 0 \\ x+a & , -a \leq x \leq 0 \\ a-x & , 0 \leq x \leq a \\ 0 & , a \leq x < \infty. \end{cases}$$

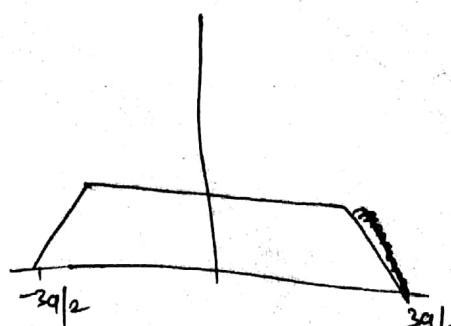
$$\therefore u(x,t) = \frac{1}{2} [f(x+ct) + f(x-ct)].$$



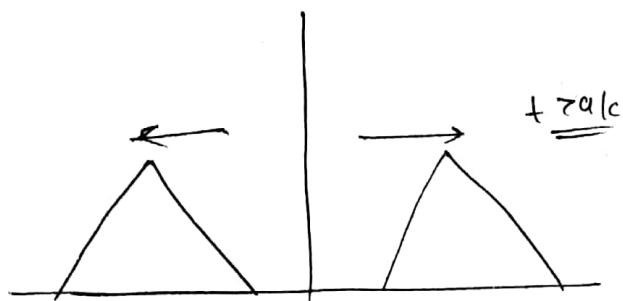
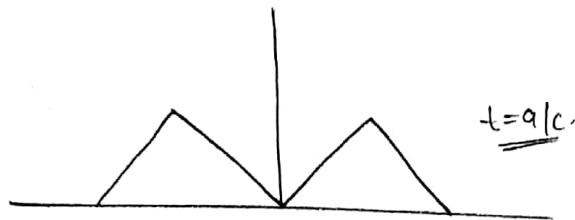
$$u(x, t=a/2c) = \frac{1}{2} [f(x+a/2) + f(x-a/2)] + \frac{1}{2}$$



=



$$u(x,t=a/c) = \frac{1}{2} [f(x+a/c) + f(x-a/c)]$$



Case-2

$$f(x) = 0$$

$$g(x) \neq 0$$

$$u(x,t) = \frac{1}{2c} \int_{x-ct}^{x+ct} g(J) dJ$$

$$\text{let } g(x) = x \quad \text{then } u(x,t) = xt.$$

problem:

$$u_{tt} = g u_{xx}, \quad x \in \mathbb{R}, t > 0.$$

$$u(x,0) = \sin x = f(x).$$

$$c = 3$$

$$u_t(x,0) = \omega \sin x = g(x).$$

$$u(x,t) = \frac{1}{2} [\sin(3t+x) + \sin(x-3t)] + \frac{1}{6} \int_{x-3t}^{x+3t} \cos J dJ.$$

$$= \frac{1}{2} [2 \sin(\cancel{x}) - 2 \sin(\cancel{x}) \cos(3t)]$$

$$= \underline{\sin(x) \cos(3t)}$$

$$[\sin J]_{x-3t}^{x+3t}$$

$$+ \frac{1}{6} [\sin(x+3t) - \sin(x-3t)]$$

$$\frac{1}{6} \left\{ \sin x \cos 3t + \cos x \sin 3t - \sin x \cos 3t + \cos x \sin 3t \right\}$$

$$= \frac{1}{3} (\cos x \sin 3t)$$

$$= \sin x \cos 3t + \frac{1}{3} \cos x \sin 3t$$

# Wave eq^n finite string

$$u_{tt} = c^2 u_{xx} \quad \text{--- (1)} \quad x \in [0, L].$$

$$u(0, t) = 0$$

fixed ends.

$$\text{BC: } u(L, t) = 0$$

$$u(x, 0) = f(x)$$

IC:

$$u_t(x, 0) = g(x).$$

Look for separable soln  $u = X(x) T(t)$ .

--- (2)

(2) in (1).

$$X T'' = c^2 X'' T$$

$$\Rightarrow \frac{X'}{X} = \frac{T''}{c^2 T} = \lambda \quad (\text{say}).$$

pair of 2nd order ODEs:

$$X'' - \lambda X = 0 \quad ; \quad T'' - c^2 \lambda T = 0$$

$$\text{BC: } \underline{X(0)=0} \quad \underline{X(L)=0}$$

$\lambda = 0, d > 0$  + trivial soln.

$$\lambda < 0: \quad X(x) = C_1 \cos mx + C_2 \sin mx.$$

$$X(0) = \underline{\cos 0} \quad C_1 = 0$$

$$X(L) = C_2 \sin(mL)$$

$$mL = n\pi \quad n=1, 2, \dots$$

$$m = \frac{n\pi}{L}$$

$$\therefore \text{eigen values. } d_n = -n^2 = -\left(\frac{n\pi}{L}\right)^2$$

$$\text{eigen functions } x_n: \sin\left(\frac{n\pi x}{L}\right)$$

correspondingly, (2<0)

$$T_{n0} = A T^{2n} + C^2 \left(\frac{n\pi}{L}\right)^2 T = 0.$$

$$\therefore T_n(0,t) = d_n \cos\left(\frac{n\pi c}{L} t\right) + e_n \sin\left(\frac{n\pi c}{L} t\right).$$

$$u(x,t) = \sum_{n=1}^{\infty} \left( A_n \cos \frac{n\pi c}{L} t + B_n \sin \frac{n\pi c}{L} t \right) \sin \frac{n\pi x}{L}.$$

$$\text{IL} \quad u(x,0) = f(x) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right).$$

$$u_t(x,0) = g(x) = \sum_{n=1}^{\infty} B_n \left(\frac{n\pi c}{L}\right) \sin\left(\frac{n\pi x}{L}\right).$$

$$\therefore \int_0^L f(x) \sin \frac{n\pi x}{L} dx = \sum_{n=1}^{\infty} A_n \int_0^L \sin \frac{n\pi x}{L} \sin \frac{n\pi x}{L} dx.$$

L/2 (by orthogonality property)

$$\therefore A_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

$$\text{By } B_n = \frac{2}{L} \cdot \frac{L}{n\pi c} \int_0^L g(x) \sin \frac{n\pi x}{L} dx.$$

$$\text{Let } f(x) = 1 \quad g(x) = x$$

$$B_n = \frac{2}{n\pi c} \int_0^L x \sin \frac{n\pi x}{L} dx.$$

$$A_n = \frac{2}{L} \left[ \frac{x \cos(n\pi x)}{n\pi} \right]_0^L$$

$$\therefore \underline{\underline{A_n}} = \frac{2}{n\pi} \left[ -\underline{\underline{\cos(n\pi x)}} + 1 \right] = \frac{2}{n\pi} \underline{\underline{1 - (-1)^n}}.$$

get the answer

Remark:

Suppose the initial data is

$$u(x,0) = \frac{1}{3} \sin 5\pi x$$

$$u_t(x,0) = 6 \sin 10\pi x - 100 \sin 3\pi x$$

in case of  $\phi, f, g$ ,  
algebraic terms only  
proceed through the  
previous process as only  
algebraic terms can be  
expressed as Fourier series.

$$u(x,t) = \sum_{n=1}^{\infty} \left( A_n \cos \frac{n\pi c t}{L} + B_n \sin \frac{n\pi c t}{L} \right) \sin \frac{n\pi x}{L}$$

$$u(x,0) = \frac{1}{3} \sin 5\pi x = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{L}$$

general solution

$$\Rightarrow A_5 = \frac{1}{3}; A_n = 0 \text{ if } n \neq 5.$$

$$u_t(x,0) = 6 \sin 10\pi x - 100 \sin 3\pi x$$

$$= \sum_{n=1}^{\infty} \left( B_n \cdot \frac{n\pi c}{L} \right) \sin \frac{n\pi x}{L}$$

for  $n=3$

$$-100 = B_3 \cdot \frac{3\pi c}{L}$$

$$B_3 = -\frac{L \cdot 100}{3\pi c}$$

for  $n=10$

$$B_{10} \cdot \frac{10\pi c}{L} = 6$$

$$B_{10} = \frac{6L}{10\pi c}$$

and all others zero.

$$u(x,t) = \sum_{n=1}^{\infty} K_n \cos \frac{n\pi c}{L} (t + \delta_n) \sin \frac{n\pi x}{L}$$

harmonic motion

with amplitude  $K_n \sin \frac{n\pi x}{L}$

each motion is a standing wave with nodes at  $\sin \frac{n\pi x}{L} = 0$ .

max. amplitude at  $\sin \frac{n\pi x}{L} = \pm 1$ .

for a fixed t,

$$u(x,t) = \sum_{n=1}^{\infty} C_n(t) \sin \frac{n\pi x}{L}$$

$$\underline{C_n(t) = \alpha_n \cos \omega_n (t + \delta_n)}$$

$$\underline{\omega_n = \frac{n\pi c}{L}}$$

free ends:

$$\frac{\partial u}{\partial x}(0,t) = 0 \quad \frac{\partial u}{\partial x}(L,t) = 0$$

one end fixed other free

$$u(0,t) = 0 \quad \frac{\partial u}{\partial x}(L,t) = 0$$

$$u(x,0) = x$$

$$u_t(x,0) = x^2$$

$$\cancel{x'' - dx = 0} \quad T'' - c^2 d T = 0$$

↓

$$\cancel{d=0}$$

$$x = ax + b$$

$$\underline{d \geq 0}$$

$$x(x) = C_1 e^{\sqrt{d}x} + C_2 e^{-\sqrt{d}x}$$

$$x(0) = 0 \Rightarrow b = 0$$

$$x(0) = C_1 + C_2 = 0$$

$$x'(L) = a = 0$$

$$x'(L) = C_1 \sqrt{d} e^{\sqrt{d}L} - C_2 \sqrt{d} e^{-\sqrt{d}L}$$

\* trivial

$$C_1 e^{\sqrt{d}L}$$

$$\underline{C_1 = C_2 = 0}$$

$$\cancel{d < 0}$$

$$x(x) = C_1 \cos mx + C_2 \sin(mx)$$

$$x(0) = C_1 = 0$$

$$x'(0) = -m C_2 \sin(0) + m C_2 \cos(0) = m C_2$$

$$\cos mL = 0 \\ mL = (2n+1)\frac{\pi}{2} \quad m = \frac{(2n+1)\pi}{L}$$

$$u(x,t) = \sum_{n=1}^{\infty} \left( A_n \cos\left(\frac{(2n+1)\pi}{2L} ct\right) + B_n \sin\left(\frac{(2n+1)\pi}{2L} ct\right) \right) \sin\left(\frac{(2n+1)\pi}{2L} x\right).$$

$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{(2n+1)\pi}{2L} x\right) = x.$$

$$A_n = \frac{2}{L} \int_0^L x \sin\left(\frac{(2n+1)\pi}{2L} x\right) dx.$$

$$= \frac{2}{L} \left\{ 2L \left[ \frac{x \cos\left(\frac{(2n+1)\pi}{2L} x\right)}{\frac{(2n+1)\pi}{2L}} \right] \Big|_0^L - \int_0^L 2L \cdot \frac{\cos\left(\frac{(2n+1)\pi}{2L} x\right)}{\frac{(2n+1)\pi}{2L}} dx \right\}.$$

$$= \frac{2}{\cancel{x}} \left[ \frac{2L}{(2n+1)\pi} \cos\left(\frac{(2n+1)\pi}{2L} \cdot 0\right) + \frac{2L}{(2n+1)\pi} \cdot \frac{2L}{(2n+1)\pi} \sin\left(\frac{(2n+1)\pi}{2L} \cdot L\right) \right]$$

$$A_n = \frac{8L}{((2n+1)\pi)^2} (-1)^n$$

$$u_t(x,0) = \sum B_n \cdot \frac{2L}{(2n+1)\pi} \sin\left(\frac{(2n+1)\pi}{2L} x\right) = x^2.$$

$$B_n = \frac{2}{L} \int_0^L \left( \frac{(2n+1)\pi}{2L} x^2 \right) \sin\left(\frac{(2n+1)\pi}{2L} x\right) dx$$

$$B_n = \frac{(2n+1)\pi}{L^2} \int_0^L x^2 \sin\left(\frac{(2n+1)\pi}{2L}x\right) dx.$$

$$\int_0^L x^2 \sin\left(\frac{(2n+1)\pi}{2L}x\right) dx = \left[ -2L \frac{x^2 \cos\left(\frac{(2n+1)\pi}{2L}x\right)}{(2n+1)\pi} \right]_0^L + 2 \int_0^L x \frac{2L}{(2n+1)\pi} \cos\left(\frac{(2n+1)\pi}{2L}x\right) dx \\ = \frac{4L}{(2n+1)\pi} \int_0^L x \cos\left(\frac{(2n+1)\pi}{2L}x\right) dx$$

$$\int_0^L x \cos\left(\frac{(2n+1)\pi}{2L}x\right) dx = \left[ 2L \frac{x \sin\left(\frac{(2n+1)\pi}{2L}x\right)}{(2n+1)\pi} \right]_0^L - \int_0^L \frac{2L}{(2n+1)\pi} \sin\left(\frac{(2n+1)\pi}{2L}x\right) dx \\ = \frac{2L \cdot L \cdot \sin\left(\frac{(2n+1)\pi}{2}\right)}{(2n+1)\pi} - \frac{2L}{(2n+1)\pi} \cdot \frac{2L}{(2n+1)\pi}$$