

Date
16/10/2017

Lecture 21

-1-

$f(x)$ periodic
 $\checkmark F.S$

Fourier Integral

Fourier series — periodic

non-periodic f^n

Fourier series \nearrow extend -

Q) What can be done to extend the method of Fourier series to non-periodic f^n .
We begin with a special f^n .

e.g.; $f_L(x)$ of period $2L$

& see what happens
to $f_L(x)$ as $L \rightarrow \infty$.

$f(x)$)

Square wave

Let us consider the periodic square wave $f_L(x)$ of period

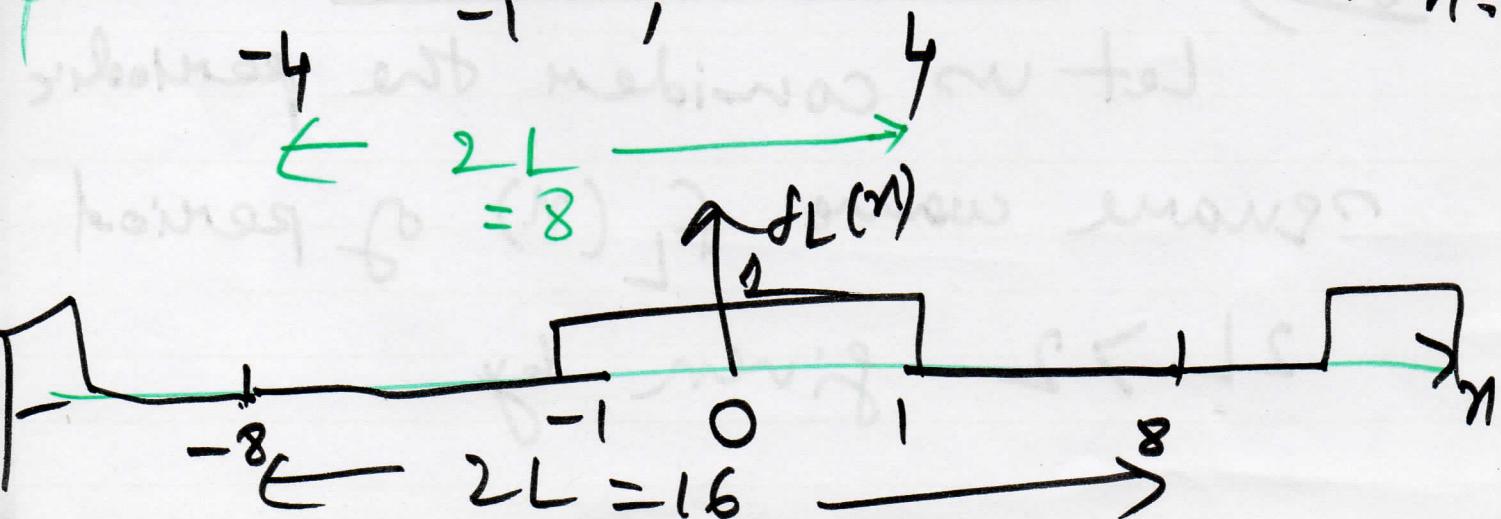
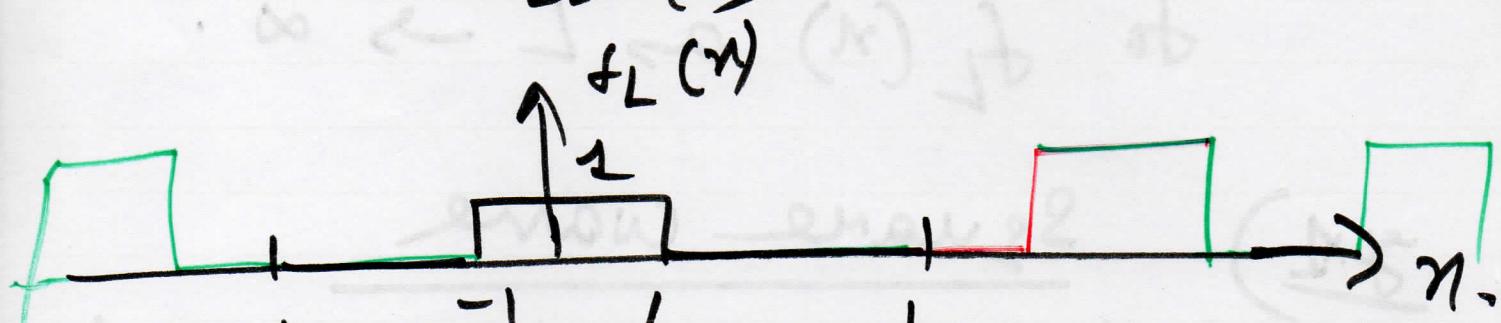
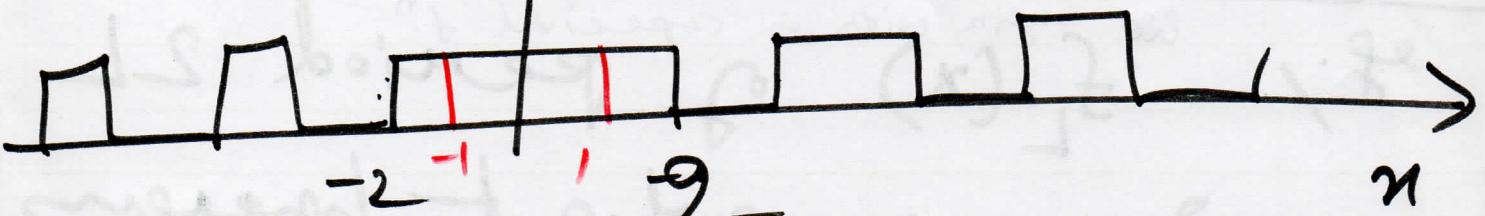
$2L > 2$ given by

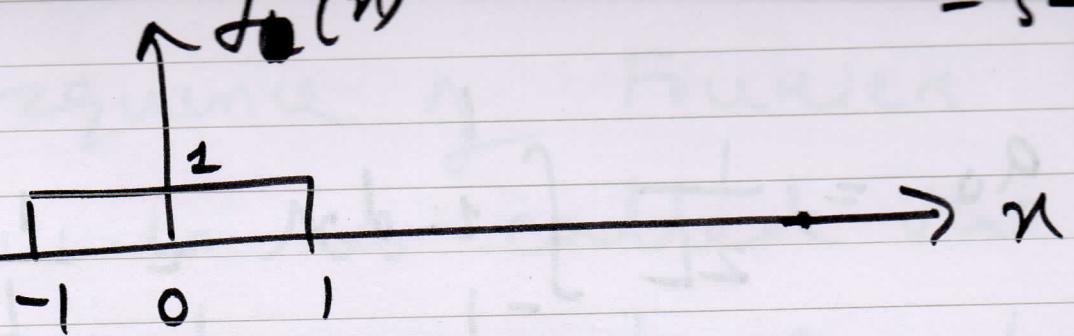
$$f_L(n) = \begin{cases} 0, & \text{if } -L < n < -1 \\ 1, & \text{if } -1 < n < 1 \\ 0, & \text{if } 1 < n < L \end{cases}$$

Σ Lt $f_L(n) = \begin{cases} 1, & \text{if } -1 < n < 1 \\ 0, & \text{otherwise} \end{cases}$

$f_L(n)$

Waveform of $f_L(n)$





~~Fig 1 :-~~ shows that f^n for $2L = 4, 8, 16$ as well as the non-periodic f^n , $f(x)$, which we obtain from $f_L(x)$ if we let $L \rightarrow \infty$.

We now explore what happens to the Fourier co-efficients of f_L as L increases.

Since f_L is even (why?)

$$\boxed{b_n = 0} \text{ for all } n.$$

For a_n , the Euler formulas give

$$a_0 = \frac{1}{2L} \int_{-1}^1 1 \cdot dn = \frac{1}{L}$$

$$a_n = \frac{1}{L} \int_{-1}^1 1 \cdot \cos \frac{n\pi x}{L} dn$$

$$= \frac{2}{L} \int_0^1 \cos \left(\frac{n\pi x}{L} \right) dn \quad \begin{cases} f_L(n) = 1 \\ -1 < n < 1. \end{cases}$$

$$= \frac{2}{L} \frac{L}{n\pi} \cdot \left[\sin \frac{n\pi x}{L} \right]_0^1$$

$$= \frac{2}{L} \cdot \frac{\sin \frac{n\pi}{L}}{(n\pi)}$$

\therefore Let $a_n = 0$ (how?)

$L \rightarrow \infty$

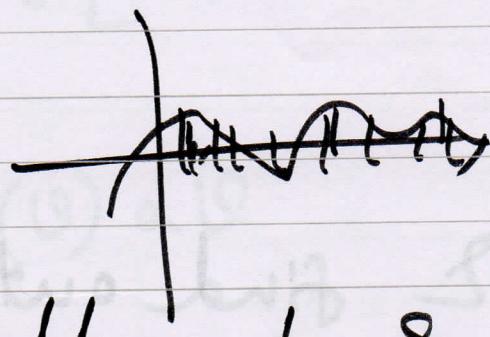
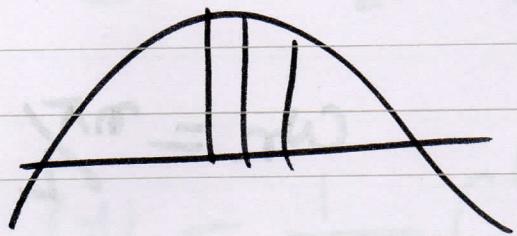
Ex

i.e., $|a_n| \rightarrow 0 \Rightarrow L \rightarrow \infty$

A sequence of Fourier coefficients is called the amplitude spectrum of f

because $|a_n|$ is the max^m amplitude of the wave
 $a_n \cos(n\pi x/L)$

i.e., these amplitudes a_n



will eventually be everywhere dense on the axis & will decrease to zero.

From Fourier series to the Fourier Integral

We now consider any periodic function $f_L(x)$ of period $2L$ that can be represented by a Fourier series.

$$f_L(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos \omega_n x + b_n \sin \omega_n x)$$
$$\omega_n = \frac{n\pi}{L}.$$

2 find out what happens if we let $L \rightarrow \infty$.

Here, $\cos \omega_n x \approx \sin \omega_n x$ with ω_n is no longer restricted to integer multiples.

$$\omega = \omega_n = n\pi/L \text{ or } \pi/L$$

for all values of n .

We want to see, what form such an integral might have.
we now insert the values
of a_0, a_n & b_n & denote
the variable of integration
by ϑ , the Fourier
series of $f_L(\vartheta)$ becomes

$$f_L(\vartheta) = \frac{1}{2L} \int_{-L}^L f_L(u) du + \frac{1}{L} \sum_{n=1}^{\infty} \left[c_n \cos \omega_n \vartheta \int_{-L}^L f_L(u) \cos(\omega_n u) du + s_n \sin \omega_n \vartheta \int_{-L}^L f_L(u) \sin(\omega_n u) du \right]$$

We now let

$$\begin{aligned}\Delta \omega &= \omega_{n+1} - \omega_n \\ &= \frac{(n+1)\pi}{L} - \frac{n\pi}{L} \\ &= \pi/L\end{aligned}$$

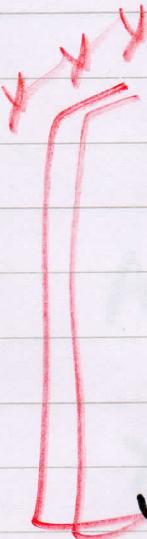
Then $\frac{1}{L} = \frac{\Delta \omega}{\pi}$ & we may
write the Fourier series
in the form

$$\begin{aligned}f_L(x) &= \frac{1}{2L} \int_{-L}^L f_L(u) du \\ &\quad + \frac{1}{\pi} \sum_{n=1}^{\infty} \left[(\cos \omega_n x) \Delta \omega \int_{-L}^L f_L(u) \cos(\omega_n u) du \right. \\ &\quad \left. + (\sin \omega_n x) \Delta \omega \int_{-L}^L f_L(u) \sin(\omega_n u) du \right] \rightarrow (1)\end{aligned}$$

This representation is valid for any fixed L arbitrarily large, but finite.

we now let $L \rightarrow \infty$

& assume that the resulting $n m$ -periodic

 f^n $f(n) = \lim_{L \rightarrow \infty} f_L(n)$

is absolutely integrable

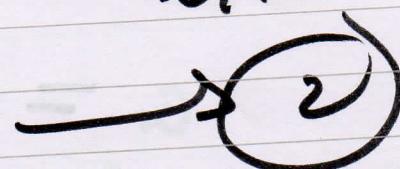
on the x -axis

i.e., the following (finite!) limits exist:

$$\lim_{a \rightarrow -\infty} \int_a^0 |f(n)| dn + \lim_{b \rightarrow \infty} \int_0^b |f(n)| dn$$

(written as $\int_{-\infty}^{\infty} |f(n)| dn$)

exist.



Then $\frac{1}{L} \rightarrow 0$ & the value

of the first term in the

R.H.S of eqn (1) approaches

zero. (i.e., $\int_{-\infty}^{\infty} |f_L(u)| du$ ^{converges} exists.)

Also, $\delta\omega = \pi/L \rightarrow 0$ as $L \rightarrow \infty$

∴ it seems plausible

(how?)

that the infinite series

in (1) becomes an integral

from 0 to ∞ which represents
 $f(x)$ namely.

$$\text{i.e., } \lim_{\Delta\omega \rightarrow 0} \left\{ \sum_{n=1}^{\infty} F(\omega_n) \Delta\omega \right\}$$

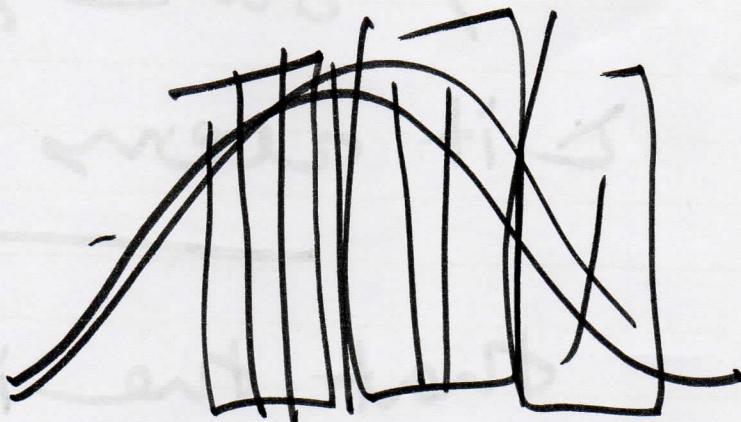
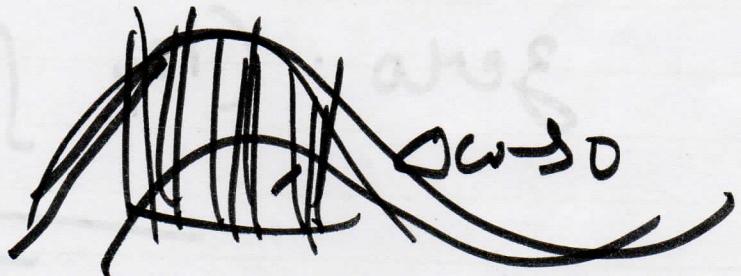
$$= \int_0^\infty F(\omega) d\omega$$

which resembles

a Riemann sum of a definite integral.

(Ex)
how?

Note :- Periodic f^n 's
are represented
by a Fourier series
& non-periodic f^n 's
are represented
in the form of an
integral i.e; by
Fourier integral.



together we formed (1) in
charge of which is a non
-periodic (2)

$$f(n) = \text{dt} f_L(n)$$

L → N

$$= \frac{1}{\pi} \int_0^{\omega} \left[\cos(\omega n) \int_{-\omega}^{\omega} f(v) \cos(\omega v) dv \right.$$

$$\left. + \sin(\omega n) \int_{-\omega}^{\omega} f(v) \sin(\omega v) dv \right] dw$$

→ (3)

If we introduce the notation

$$A(\omega) = \frac{1}{\pi} \int_{-\omega}^{\omega} f(v) \cos(\omega v) dv$$

$$B(\omega) = \frac{1}{\pi} \int_{-\omega}^{\omega} f(v) \sin(\omega v) dv$$

→ (4)

We can write this in the form

$$f(x) = \int_0^{\infty} [A(\omega) \cos(\omega x) + B(\omega) \sin(\omega x)] d\omega \rightarrow (5)$$

This is called a representation of $f(x)$

by a Fourier integral.

~~n-1~~ / (Sufficient Condition)
(Fourier integral)

If $f(x)$ is piecewise continuous in every finite interval Σ has a

R.H. Derivative \neq L.H.

Derivative at every point

\Leftrightarrow if the integral (2)

exists i.e., $\int_{-\infty}^{\infty} |f(x)| dx$ exists

then $f(x)$ can be represented by a Fourier Integral (3) with A & B given by (4).

At a point where $f(x)$ is discontinuous, the value of the Fourier Integral equals the average of the left-hand & right-hand limits of $f(x)$ at that point.

Equivalent forms of Fourier

Integral Theorem.

Fourier's integral theorem can also be written in the forms

$$\textcircled{1} \quad f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(u) e^{ix(u-k)} du dk$$

$\checkmark x \checkmark u \checkmark k$
 $\checkmark E^x$

$$\textcircled{2} \quad (\text{Complex form of Fourier integral})$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(u) e^{ix(u-k)} du dk$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} C(k) e^{-ikx} dk, \quad \text{where } C(k) = \int_{-\infty}^{\infty} f(u) e^{iku} du$$

$$f(n) = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{inx} dx \times \int_{-\pi}^{\pi} f(u) e^{-iuy} du$$

where it is understood
that if $f(n)$ is not
continuous at n , the
L.H.S must be replaced
by $\frac{f(n+0) + f(n-0)}{2}$.

These results can be
simplified somewhat if
 $f(n)$ is either odd or even.

$$(I) f(n) = \frac{2}{\pi} \int_0^\infty \sin(nx) dx \times \int_0^\infty f(u) \sin(uy) du$$

$f(n)$ is odd

$$f(n) = \frac{2}{\pi} \int_0^\infty \cos(kn) dx \int_0^\infty f(u) \cos(ku) du$$



if $f(n)$ is even.

Note :- Fourier integrals are used mainly in solving O.D.E's & P.D.E's. However, they also help in evaluating integrals as well.

§ Fourier Cosine & Sine Integrals

For an even or odd f^n ,
the Fourier integral
becomes simpler, just as
in the case of Fourier
series.

Indeed, if $f(n)$ is an
even f^n , then

$$\text{II } B(\omega) = 0 \quad \underline{\underline{\text{in } e_j^n(u)}}$$

$$A(\omega) = \frac{2}{\pi} \int_0^\infty f(v) \cos(\omega v) dv$$

The Fourier integral (5) then reduces to the Fourier cosine integral

$$f(n) = \int_0^\infty A(\omega) \cos(\omega n) d\omega$$

(if even).

slly, if $f(n)$ is odd, then

in (4), we have

$$A(\omega) = 0$$

$$B(\omega) = \frac{2}{\pi} \int_0^\infty f(v) \sin(\omega v) dv$$

The Fourier integral (5) then reduces to the Fourier integral

$$f(n) = \int_0^\omega B(\omega) \sin(\omega n) d\omega$$

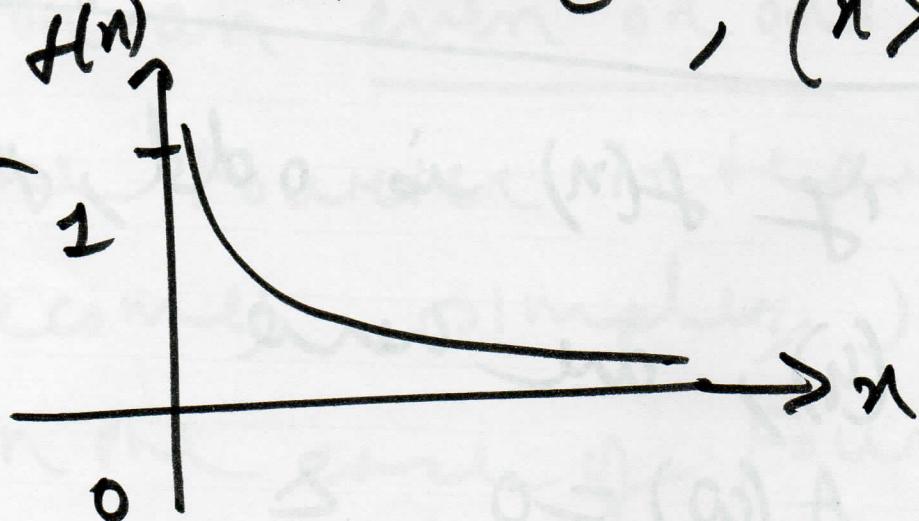
(f odd)

~~Ex 1~~ / Find the Fourier cosine

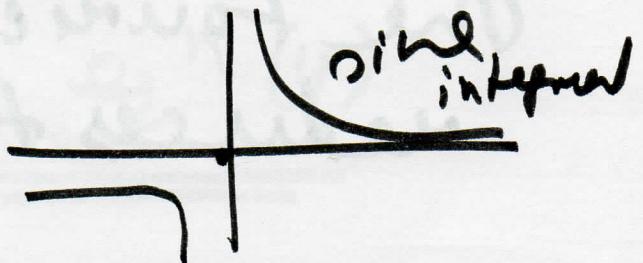
of sine integrals

$$f(n) = e^{-kn}, (n > 0, k > 0)$$

Sol:



cosine integral



sine integral

∴ - we have

$$A(\omega) = \frac{2}{\pi} \int_0^\infty f(v) c_\omega(\omega v) dv$$

$$= \frac{2}{\pi} \int_0^\infty e^{-kv} \underbrace{c_\omega(\omega v)}_{\text{Now, integration by parts}} dv$$

Now, integration by parts,

$$\int e^{-kv} c_\omega(\omega v) dv$$

$$= \frac{-k}{k^2 + \omega^2} e^{-kv} \left(\frac{-\omega}{k} \sin \omega v + c_\omega \omega v \right)$$

$$= \left[\frac{-k}{k^2 + \omega^2} e^{-kv} \left(\frac{\omega}{k} \sin(\omega v) + c_\omega(\omega v) \right) \right]_0^{\infty}$$

(Ex)

$$= -\frac{k}{k^2 + \omega^2} [0 - 1] = \left(\frac{k}{k^2 + \omega^2} \right)$$

Thus

$$A(\omega) = \frac{2k\pi}{k^2 + \omega^2}$$

∴ The Fourier cosine Integral representation is

$$f(n) = \int_0^\infty A(\omega) \cos(\omega n) d\omega$$

$$\Rightarrow e^{-kn} = \frac{\int_0^\infty \cos \omega n}{k^2 + \omega^2} d\omega$$

$$\Rightarrow \left[\int_0^\infty \frac{\cos \omega n}{k^2 + \omega^2} d\omega = \frac{\pi}{2k} e^{-kn} \right]$$

\int_1^∞ $e^{-kn} \cdot (n > 0, k > 0)$
Cosine integral.

b) silly, are here

$$B(\omega) = \frac{2}{\pi} \int_0^\infty f(v) \sin(\omega v) dv$$

$$= \frac{2}{\pi} \int_0^\infty e^{-kv} \sin(\omega v) dv$$

~~✓ Sec'y.~~
~~rough work~~

Integrating by parts.

$$\int e^{-kv} \sin(\omega v) dv = -\frac{\omega}{k^2 + \omega^2} e^{-kv} \left(\frac{k}{\omega} \sin(\omega v) + C \cos(\omega v) \right)$$

$$= \left[-\frac{\omega}{(k^2 + \omega^2)} \cdot e^{-kv} \left(\frac{k}{\omega} \sin(\omega v) + C \cos(\omega v) \right) \right]_0^\infty$$

$$= -\frac{\omega}{k^2 + \omega^2} \left[0 - 1 \right] = \frac{\omega}{k^2 + \omega^2}$$

- 2

Thus,
$$B(\omega) = \frac{2\omega/k}{k^2 + \omega^2}$$

\therefore the Fourier sine integral representation is

$$f(n) = \int_0^\infty B(\omega) \sin(\omega n) d\omega$$

$$\Rightarrow c_n^{-k} = \frac{2}{\pi} \int_0^\infty \frac{\omega \sin(\omega n)}{k^2 + \omega^2} d\omega$$

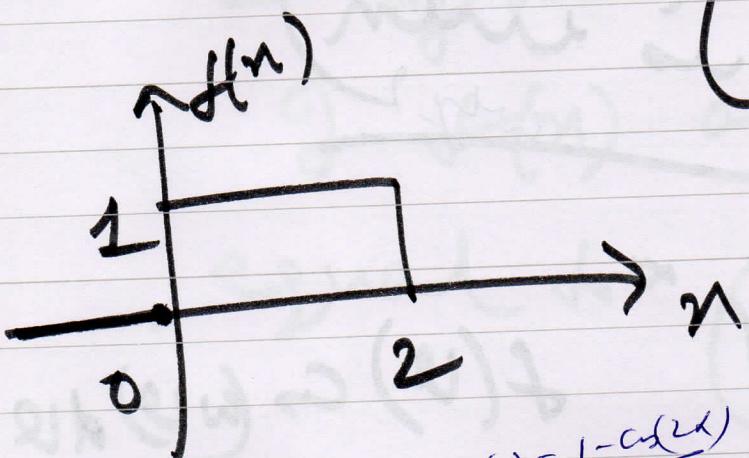
$$\Rightarrow \boxed{\int_0^\infty \frac{\omega \sin(\omega n)}{k^2 + \omega^2} d\omega \stackrel{dw = \frac{\pi}{2} e^{-kn}}{=} \frac{\pi}{2} e^{-kn}}$$

$n > 0, k > 0$

The above integrals are called the ⁽¹⁾ Laplace integrals.

Q) Find the Fourier integral representation of the piece-wise continuous f^n

$$f(n) = \begin{cases} 0, & n < 0 \\ 1, & 0 < n < 2 \\ 0, & n > 2 \end{cases}$$



Hence evaluate

$$\text{Hint: } A(\alpha) = \frac{\sin 2\alpha}{\alpha}, \quad B(k) = \frac{1 - \cos(2k)}{k}$$

$$f(n) = \frac{2}{\pi} \int_0^\infty \frac{\sin \alpha}{\alpha} \cos \left(\alpha(n-1) \right) d\alpha$$

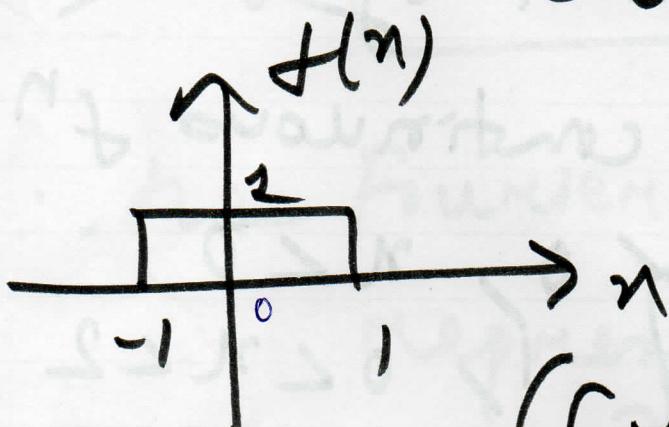
$$\int_0^\infty \frac{\sin x}{x} dx \quad (= \pi/2)$$

~~Explain~~ single pulse, sine integral

Q) Find the Fourier integral representation of the function f^r

$$f(x) = \begin{cases} 1, & \text{if } |x| < 1 \\ 0, & \text{if } |x| > 1 \end{cases}$$

ie, $-1 < x < 1$



Ex/ Find Fourier integral representation
for
 $f(x) = \begin{cases} 1, & |x| < a \\ 0, & |x| > a \end{cases}$

(find even
why?)

We have

$$A(\omega) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \cos(\omega v) dv$$

$$= \frac{1}{\pi} \int_{-1}^{1} 1 \cdot \cos(\omega v) dv$$

$$= \left[\frac{\sin(\omega v)}{\pi v} \right]_{-1}^1 = \boxed{\frac{2 \sin \omega}{\pi \omega}}$$

$$B(\omega) = 0$$

$$\therefore f(x) = \frac{2}{\pi} \int_0^\infty \frac{\sin \omega x \sin \omega}{\omega} d\omega$$

$$\Rightarrow \int_0^\infty \frac{\omega \cos \omega x \sin \omega}{\omega} d\omega = \frac{\pi}{2} f(n) \rightarrow x)$$

The average of the left & right hand limits of $f(n)$ at $n=1$ is

equal to $\left(\frac{1+0}{2} = \frac{\pi}{2} \right)$

~~xxx~~ ∵ from (1) & $\lim_{n \rightarrow 1} ?$
we get

$$\boxed{\int_0^\infty \frac{\omega \cos \omega x \sin \omega}{\omega} d\omega} = \begin{cases} \frac{\pi}{2} \cdot 1 = \frac{\pi}{2}, & 0 \leq x < 1 \\ \frac{\pi}{2} \cdot \frac{\pi}{2} = \frac{\pi^2}{4} \text{ if } x = 1. \\ \frac{\pi}{2} \cdot 0 = 0 \text{ if } x > 1 \end{cases}$$

This integral is called
Dirichlet's discontinuous
function.

Let $n = 0$,

then $\int_0^\infty \frac{\sin w}{w} dw = \frac{\pi}{2}$

This integral is the limit of the so-called sine integral.

sine integral

$$\text{Si}(u) = \int_0^u \frac{\sin w}{w} dw$$

as $u \rightarrow \infty$.

X