

Date
17/11/2017

Lecture 28

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a) Solve $\frac{\partial^2 U}{\partial t} = \frac{\partial^2 U}{\partial x^2}$, $n > 0$, $t > 0$

subject to the cond'n's

$$U(0, t) = 0, U(n, 0) = \begin{cases} 1, & 0 < n < 1 \\ 0, & n \geq 1. \end{cases}$$

$U(n, t)$ is bounded.

To find $U(n, t) = ?$

Sol) i) - Taking the Fourier sine transform of both sides of the given p.d.e w.r.t n , we get

$$\mathcal{F}_S \left[\frac{\partial^2 U}{\partial t} \right] = \mathcal{F}_S \left[\frac{\partial^2 U}{\partial x^2} \right]$$

Recall

$$\hat{f}_s(\omega) = \sqrt{\frac{2}{\pi}} \int_0^\infty f(n) \sin \omega n \, dn$$

$$f(n) = \sqrt{\frac{2}{\pi}} \int_0^\infty \hat{f}_s(\omega) \sin \omega n \, d\omega.$$

$$\sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial V}{\partial t} \sin(2\pi n) \, dx$$

$$= \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 V}{\partial n^2} \sin(2\pi n) \, dx$$

Then if

$$u = u(\lambda, t) = \hat{f}_s \{ V(x, t) \}$$

$$= \sqrt{\frac{2}{\pi}} \int_0^\infty V(x, t) \sin \lambda x \, dx.$$

This becomes $\rightarrow (*)$

$$\frac{\partial^2}{\partial t^2} \left(\sqrt{\frac{2}{\pi}} \int_0^\infty V(x, t) \sin \lambda x \, dx \right) = \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 V}{\partial n^2} \sin \lambda n \, dn$$

$$\Rightarrow \frac{du}{dt} = \sqrt{\frac{2}{\pi}} \left[\frac{\partial V}{\partial x} \sin \lambda n \right]_0^\infty - \lambda \int_0^\infty \frac{\partial^2 V}{\partial n^2} C(\lambda n) \, dn$$

F^0 (why!)

$$\Rightarrow \frac{du}{dt} = \sqrt{\frac{2}{\pi}} \left[-\lambda \left[U_{\infty} \sin(\lambda x) \right] \Big|_{x=0} + \lambda^2 \int_0^\infty U \sin(\lambda x) dx \right]$$

$$\Rightarrow \frac{du}{dt} = \sqrt{\frac{2}{\pi}} \rightarrow \frac{\underline{U(0,t)}}{(=0)} - \lambda^2 u$$

$$\Rightarrow \frac{du}{dt} + \lambda^2 u = 0 \rightarrow (2)$$

assuming $\rightarrow U \propto \frac{\partial U}{\partial x} \rightarrow 0$
 $\text{as } x \rightarrow \infty$

Now,

$$\mathcal{F}_S[U(x,0)] = u(\lambda, 0)$$

$$= \sqrt{\frac{2}{\pi}} \int_0^\infty \underline{U(x,0) \sin(\lambda x)} dx$$

$$= \sqrt{\frac{2}{\pi}} \int_0^\infty 1 \cdot \sin(\lambda x) dx$$

$$\begin{aligned} \frac{dy}{dt} &= -\lambda^2 u \\ \Rightarrow \frac{dy}{u} &= -\lambda^2 dt \\ (\text{on } 1) \quad & \\ \Rightarrow \ln u &= -\lambda^2 t + \text{Int C} \\ \Rightarrow u(\lambda, t) &= C e^{-\lambda^2 t} \end{aligned}$$

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$$\therefore u(\lambda, 0) = \sqrt{\frac{2}{\pi}} \int_0^1 \sin(\lambda n) dn$$

$$= \sqrt{\frac{2}{\pi}} \left[-\frac{\cos(\lambda n)}{\lambda} \right]_0^1$$

$$= \sqrt{\frac{2}{\pi}} \frac{(1 - \cos \lambda)}{\lambda} \rightarrow (3)$$

We have

$$u(\lambda, t) = C e^{-\lambda^2 t}$$

$$\therefore u(\lambda, 0) = C \cdot 1 = C$$

$$\Rightarrow C = \sqrt{\frac{2}{\pi}} \frac{(1 - \cos \lambda)}{\lambda} \quad (\text{from (3)})$$

$$\therefore u(x, t) = C \cdot e^{-\lambda^2 t}$$
$$= \boxed{\sqrt{\frac{2}{\pi}} \frac{(1 - \cos \lambda)}{\lambda} e^{-\lambda^2 t}}$$

Then taking the inverse Fourier sine transform,
we find the reqd. soln

$$\begin{aligned} \mathcal{F}(u, t) &= \sqrt{\frac{2}{\pi}} \int_0^\infty u(\lambda, t) \sin(\lambda x) dx \\ &= \sqrt{\frac{2}{\pi}} \int_0^\infty \left(\frac{1 - e^{-\lambda^2 t}}{\lambda} \right) e^{-\lambda^2 t} \sin(\lambda x) dx. \end{aligned}$$

Physically, this can (solve?) represent the temperature in a solid.



Syllabus for the end - Sem

Marks: 50 , Time: 3 hours

- 1) Laplace Transform &
its properties, application
(^{especially} F. B. E.)
to boundary value problems,
(solving p.d.e's).
D.O.D. eigh \Rightarrow system of
O.D.E's
- 2) Fourier series, Dirichlet's
criterion,
half-range, even-odd,
Parseval's theorem,
complex form
for Fourier series
Fourier series \Rightarrow Fourier F.T.
Fourier Integral representation,
convolution.
Fourier sine or Cosine Integral
representation, Parseval's R.

Fourier Transform, Fourier
Sine or Cosine Transform,

Various properties, — convolution,
Dirac delta (etc.)

1st shifting, 2nd shifting

Solving Boundary value Problem
(i.e., solving P.D.E.)

(Q1) The steady state temperature
in a semi-infinite plate
is determined from

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, \quad 0 < x < \pi, \quad y > 0$$

solve for
 $u(x, y)$:-

$$u(0, y) = 0, \quad u(\pi, y) = e^{-y}, \quad y > 0$$

$$\boxed{\left. \frac{\partial u}{\partial y} \right|_{y=0} = 0, \quad 0 < x < \pi}$$

Solⁿ.- The domain of the variable y & the prescribed condⁿ at $y=0$ indicate that the Fourier cosine transform

is suitable for the problem.

We define

$$\left[\mathcal{F}_c \{u(x,y)\} = \sqrt{\frac{2}{\pi}} \int_0^{\infty} u(x,y) \cos(\alpha y) dy \right. \\ \left. = V(x,\alpha) \right]$$

$$\therefore \mathcal{F}_c \left\{ \frac{\partial^2 u}{\partial x^2} \right\} + \mathcal{F}_c \left\{ \frac{\partial^2 u}{\partial y^2} \right\} = \tilde{f}(0)$$

$$\Rightarrow \frac{d^2}{dx^2} V - \alpha^2 V(x,\alpha) - \sqrt{\frac{2}{\pi}} \underset{(=0)}{u_y(x,0)} = 0$$

$$\boxed{\frac{u_y}{y} = 2u}$$

We know,

$$\mathcal{F}_c \{ f''(x) \} = -\omega^2 \mathcal{F}_c \{ f(x) \}$$
$$-\sqrt{\frac{2}{\pi}} f'(0)$$

$$\Rightarrow \frac{d^2 V}{dx^2} - \omega^2 V = 0$$

Since, the domain
of x is a finite
interval, we

choose to write
the soln of the
o.d.e. as

$$\left[\begin{array}{l} A \cdot e^{ix} \\ m^2 - \omega^2 = 0 \\ \Rightarrow m = \pm \omega \end{array} \right]$$

$$V(x, \omega) = C_1 \cosh(\omega x) + C_2 \sinh(\omega x) \quad \rightarrow (1)$$

$$\text{Now, } \mathcal{F}_c \{ u(0, y) \} = F_c(0)$$
$$\Rightarrow V(0, \omega) = 0$$

we know,

$$\mathcal{F}_c \{ f''(x) \} = -\omega^2 \mathcal{F}_c \{ f(x) \}$$
$$-\sqrt{\frac{2}{\pi}} f'(0)$$

$$\Rightarrow \frac{d^2 V}{dx^2} - \omega^2 V = 0$$

Since, the domain
 $[0, l]$ is a finite
interval, we

choose to write
the solⁿ of the
o.d.e.

$$\begin{aligned} A \cdot E \\ m^2 - \omega^2 = 0 \\ \Rightarrow m = \pm \omega \end{aligned}$$

$$V(x, \omega) = C_1 \cosh(\omega x) + C_2 \sinh(\omega x)$$

$$\text{Now, } \mathcal{F}_c \{ V(0, \omega) \} = F_c(0) \rightarrow \text{S(1)}$$
$$\Rightarrow V(0, \omega) = 0$$

$$2 \mathcal{F}_c \{u(\pi, y)\} = \mathcal{F}_c \{e^{-y}\}$$

$$\Rightarrow U(\pi, \alpha) = \sqrt{\frac{2}{\pi}} \left(\frac{1}{1+\alpha^2} \right)$$

$$\Rightarrow 0 = c_1 \cdot 1 + c_2 \cdot 0$$

$\Rightarrow c_1 = 0$

we know,

$$\mathcal{F}_c \{e^{-bx}\}$$

$$= \sqrt{\frac{2}{\pi}} \int_0^\infty e^{-bn} \cos(kn) dn$$

$$\therefore U(n, \alpha)$$

$$= c_2 \sinh(bn)$$

$$\therefore \underbrace{U(\pi, \alpha)}_{n=\pi} = c_2 \sinh(b\pi)$$

$$\Rightarrow c_2 = \frac{U(\pi, \alpha)}{\sinh(b\pi)} = \sqrt{\frac{2}{\pi}} \frac{1}{(1+\alpha^2) \sinh(b\pi)}$$

$$\therefore U(n, \alpha) = \sqrt{\frac{2}{\pi}} \cdot \frac{\sinh(bn)}{(1+\alpha^2) \sinh(b\pi)}$$

$$\therefore u(x, y) = \mathcal{F}_C^{-1}[f(x, \zeta)]$$

$$= \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\sinh(kx)}{(1+\zeta^2) \sinh(\zeta x)} \cos(\zeta y) d\zeta.$$

(Solve it if possible)

Note :- Had $u(x, 0)$ been given in the above example rather than $u_y(x, 0)$, then the sine transform could have been appropriate.