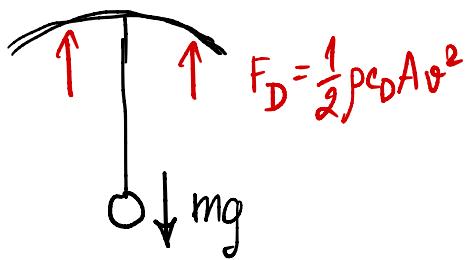


ผลลัพธ์ของความเร็ว $v(t)$ ซึ่งมีแนวโน้มลดลง

ตาม Newton's :

$$\vec{SF} = m \frac{d\vec{v}}{dt}$$



$$mg - \frac{1}{2} \rho C_D A v^2 = m \frac{dv}{dt}$$

$$\text{ถ้า } \alpha = \frac{1}{2} \rho C_D A :$$

$$g - \frac{\alpha}{m} v^2 = \frac{dv}{dt}$$

$$\text{ดูอย่าง } \frac{m}{\alpha} \int_{v(0)}^{v(t)} \frac{1}{mg - v^2} dv$$

$$\int_{t=0}^t 1 dt = \Delta t = \int_{v(0)}^{v(t)} \frac{1}{g - \frac{\alpha}{m} v^2} dv$$

$$\text{จากผลลัพธ์ } \int \frac{1}{a^2 - u^2} du = \frac{1}{2a} \ln \left| \frac{u+a}{u-a} \right| + C = \frac{1}{a} \operatorname{arctanh} \left(\frac{u}{a} \right) + C = \frac{1}{a} \operatorname{arccoth} \left(\frac{u}{a} \right) + C$$

$$\text{ดังนั้น } \frac{u}{a} \in [-1, 1] \rightarrow \frac{u^2}{a^2} \in [0, 1]$$

$$\text{เมื่อ } u = v \text{ และ } a = \sqrt{\frac{mg}{\alpha}}$$

$$\text{หมายเหตุ } 0 \leq \frac{\alpha v^2}{mg} \leq 1 \quad \begin{array}{l} \text{กรณีเร่งต้านทาน} \\ \text{กรณีห้าม} \end{array}$$

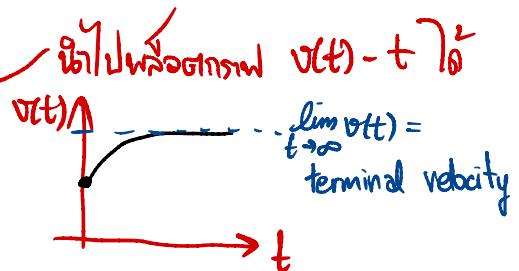
ดังนั้น $v(t)$ ต้องเป็นฟังก์ชันที่ $v(t) \rightarrow 0$ เมื่อ $t \rightarrow \infty$ / เนื่องจาก $v(t) \rightarrow 0$

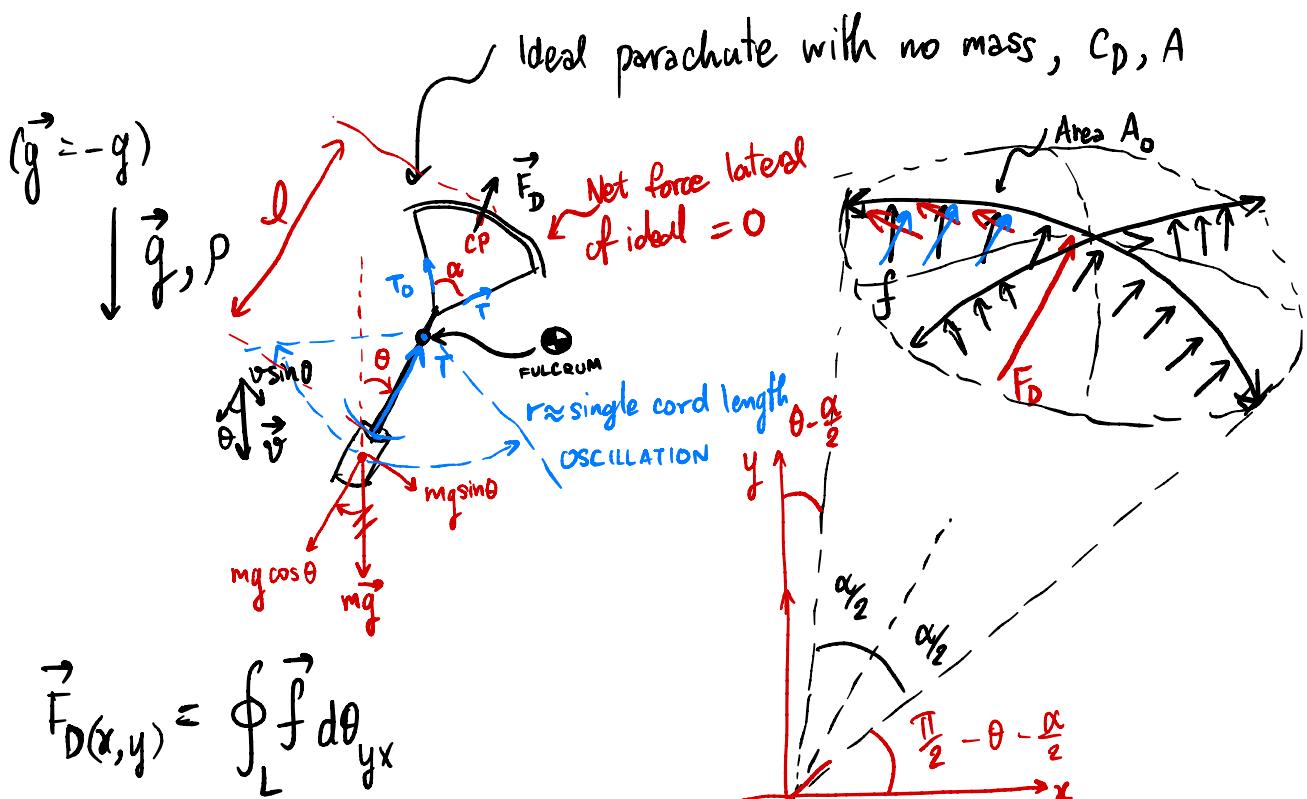
$$\text{จึง } t = \Delta t = \sqrt{\frac{m}{\alpha g}} \operatorname{arctanh} \left(\sqrt{\frac{\alpha}{mg}} v \right) \Big|_{v=v(0)}^{v(t)}$$

$$\sqrt{\frac{dg}{m}} \Delta t = \operatorname{arctanh} \left(\sqrt{\frac{\alpha}{mg}} v(t) \right) - \operatorname{arctanh} \left(\sqrt{\frac{\alpha}{mg}} v(0) \right)$$

$$\text{ดังนั้น } v(t) = \sqrt{\frac{mg}{\alpha}} \tanh \left(\operatorname{arctanh} \left(\sqrt{\frac{\alpha}{mg}} v(0) \right) + \sqrt{\frac{dg}{m}} t \right) ; \alpha = \frac{1}{2} \rho C_D A$$

$$\therefore v(t) = v_0 \pm \sqrt{\frac{mg}{\alpha}} \tanh \left(\sqrt{\frac{dg}{m}} t \right)$$





$$\vec{F}_D(x, y, z) = \iint_A \vec{f} d\theta_{yx} d\theta_{yz}$$

with integration boundary of $\left[\theta - \frac{\alpha}{2}, \frac{\pi}{2} - \theta - \frac{\alpha}{2}\right]$

while f is composite force of aerodynamic drag $F.D. \equiv \frac{1}{2} \rho c_D A v^2$

$$\therefore f(v) = \frac{1}{2} \rho c_D A_{xz} v^2 = \alpha_{xz} v^2.$$

Newton's 2nd:

$$\sum \vec{F} = m \frac{d\vec{v}}{dt}$$

$$\vec{F}_D + \vec{mg} = m \frac{d\vec{v}}{dt}$$

$$\vec{F}_D(x, y, z)(\vec{v}) + \vec{mg} = m \frac{d\vec{v}}{dt}$$

$$\int_0^t \frac{1}{m} dt = \int_{v_0}^v \frac{1}{F_D(x, y, z)(\vec{v}) + \vec{mg}} d\vec{v}$$

$$\frac{t}{m} = \int_{v_0}^v \left\{ mg - \alpha v^2 \right\}^{-1} d\vec{v}$$

$$\frac{t}{m} = \frac{1}{\alpha mg} \operatorname{arctanh} \left(\sqrt{\frac{\alpha}{mg}} v \right) \Big|_{v_0}^v$$

$$2t \sqrt{\frac{\alpha}{m}} = \ln \left[\frac{1 + \sqrt{\frac{\alpha}{mg}} v}{1 - \sqrt{\frac{\alpha}{mg}} v} \right] - \ln \left[\frac{1 + \sqrt{\frac{\alpha}{mg}} v_0}{1 - \sqrt{\frac{\alpha}{mg}} v_0} \right]$$

$$2t \sqrt{\frac{\alpha}{m}} = \ln \left[\frac{(1+k\alpha)(A)}{(1-k\alpha)(B)} \right]$$

$$(1+k\alpha)A = Be^{2t \sqrt{\frac{\alpha}{m}}} (1-k\alpha)$$

$$A + Ak\alpha = Be^{kt} - Be^{kt}k\alpha$$

$$v(Ak + Be^{kt}k) = Be^{kt} - A$$

$$v = \frac{Be^{kt} - A}{k(Be^{kt} + A)}$$

DESCENT : $\therefore v(t) = \sqrt{\frac{mg}{\alpha}} \frac{(1 + \sqrt{\frac{\alpha}{mg}} v_0) e^{2 \sqrt{\frac{\alpha}{m}} t} - (1 - \sqrt{\frac{\alpha}{mg}} v_0)}{(1 + \sqrt{\frac{\alpha}{mg}} v_0) e^{2 \sqrt{\frac{\alpha}{m}} t} + (1 - \sqrt{\frac{\alpha}{mg}} v_0)}$

$v \rightarrow v \cos \theta$

$v \rightarrow v \sin \theta$

$\alpha \rightarrow \alpha_{xz}$

Oscillation

$$\sum F = m \frac{d^2x}{dt^2} ; \theta = \frac{x}{l}$$

$$mg \sin \theta = - mr \frac{d^2\theta}{dt^2}$$

$$\frac{d^2\theta}{dt^2} + \frac{g}{r} \sin \theta = 0$$

* Given that descending mass with parachute has approximately inertial frame of reference as the velocity is slowly increasing.

In oscillation, realistic parachute calc. method is required.

The F_D in lateral component must be included:

$$mg \sin \theta + F_D + mr \frac{d^2\theta}{dt^2} = 0 ; F_D = \underbrace{\frac{1}{2} \rho C_D A_{xy} v^2}_{\alpha_{xy}} \sin^2 \theta$$

$$mr \frac{d^2\theta}{dt^2} + mg \sin \theta + \alpha_{xy} v^2 \sin^2 \theta = 0$$

$$\frac{d^2\theta}{dt^2} + \frac{g}{r} \sin \theta + \frac{\alpha v^2}{mr} \sin^2 \theta = 0$$

Thus, the Period of oscillation will follow small-angle approximation

if $\theta \ll 1$ and $\sin \theta > \sin^2 \theta$. The $\frac{\alpha v^2}{mr} \sin^2 \theta$ term will be negligible:

$$T = 2\pi \sqrt{\frac{r}{g}} \sum_{n=0}^{\infty} \left(\frac{(2n)!}{(2^n n!)^2} \right)^2 \sin^2 \theta$$

$$T \approx 2\pi \sqrt{\frac{r}{g}} \quad (\text{small-angle approximation})$$

But the term v^2 can be formulated as $\left(\frac{dx}{dt} \right)^2 = \left(r \frac{d\theta}{dt} \right)^2$

$$mr \frac{d^2\theta}{dt^2} + mg \sin \theta + \alpha \cdot \left((l-r) \frac{d\theta}{dt} \right)^2 \sin^2 \theta = 0$$

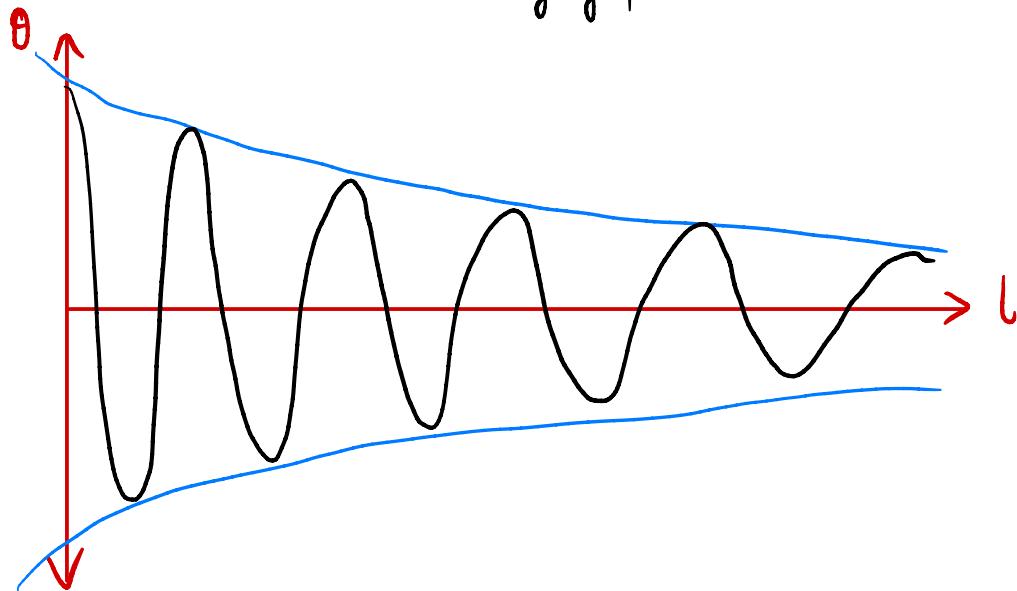
$$\therefore \frac{d^2\theta}{dt^2} + \frac{g}{r} \sin \theta + \alpha(l-r)^2 \left(\frac{d\theta}{dt} \right)^2 \sin^2 \theta = 0$$

$\alpha \rightarrow \alpha_{xy}$
OSCILLATION

This equation describes the motion of oscillation of a descending mass object (payload) attached to parachute.

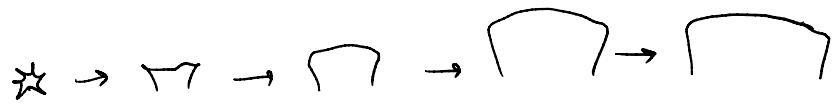
The first and second term represents oscillation with full amplitude $\theta_0 / r(1-\cos\theta_0)$, $\omega^2 = \frac{g}{r}$, and period $2\pi\sqrt{\frac{r}{g}}$.

The third term describes decaying pattern of the oscillation.

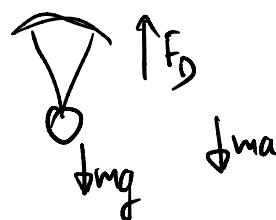
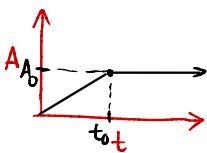


Which is similar to damped oscillation with solution of $\theta(t) = A e^{-\gamma t} \cos(\omega t)$.

PARACHUTE OPENING : Linear Approximation



$$① A = \begin{cases} A_0 \left(\frac{t}{t_0} \right) & 0 \leq t \leq t_0 \\ A_0 & t > t_0 \end{cases}$$



$$② A = A_0 \left(\frac{1 - e^{-kt}}{1 + e^{-kt}} \right)$$

$$③ A = A_0 \tanh(kt)$$

$$④ A = A_0 (1 - e^{-kt})$$

$$mg - \frac{1}{2} \rho c_D A(t) v^2 = m \frac{dv}{dt}$$

$$\text{From } A(t) = A_0 t$$

$$\rightarrow mg - kt v^2 = m \frac{dv}{dt}$$

$$g - \frac{k}{m} \cdot t \cdot v^2 = \frac{dv}{dt}$$

$$\frac{g}{t} - \frac{k}{m} v^2 = \frac{dv}{dt}$$

$$u - v = \frac{dv}{du}$$

Double Pendulum with 2 Parachute

Consider a 3D system :

* m_1 and m_2 is approx. point mass

* using coordinates $\begin{matrix} z \\ y \\ x \end{matrix}$
and ZYX vectors
(Tait-Bryan) $V(\Psi, \theta, \phi)$

