SCALAR PROXY

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1. Introduction

We will work most with the scalar proxy given by the lagrangian,

(1.1)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi - \frac{1}{2M^{2}}\Box\phi\Box\phi - \frac{\kappa}{2}\Box\phi\phi^{2}$$

The idea here is reintegrate the higher derivative term, in order to obtain a lower derivative term, but in terms of additional fields. This is easily done by,

1

(1.2)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi + \Box\phi\eta + \frac{M^{2}}{2}\eta^{2} - \frac{\kappa}{2}\Box\phi\phi^{2}$$

The new lagrangian has mixed propagator terms, to diagonalize it is also easy, we just open in terms of $\phi = h - \eta$,

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \partial_{\mu}h\partial^{\mu}\eta - \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2} + \eta\Box(h-\eta) + \frac{M^{2}}{2}\eta^{2}$$

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \partial_{\mu}h\partial^{\mu}\eta - \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2} - \partial_{\mu}\eta\partial^{\mu}(h-\eta) + \frac{M^{2}}{2}\eta^{2}$$

(1.5)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta + \frac{M^{2}}{2}\eta^{2} - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2}$$

The Feynman rules are easily red as,

•
$$h - - - - - h = \frac{1}{i} \frac{1}{p^2}$$

$$\bullet \qquad \eta = -\frac{1}{\mathrm{i}} \frac{1}{p^2 + M^2}$$

•
$$h_1 - \cdots = i\kappa (p_1^2 + p_2^2 + p_3^2)$$

•
$$h_1 - - - \prec (p_1^2 + p_2^2 + p_3^2)$$

•
$$h_1 = i\kappa (p_1^2 + p_2^2 + p_3^2)$$

•
$$\eta_1 = -i\kappa (p_1^2 + p_2^2 + p_3^2)$$

Which can also be seen directly from the Feynman rules of the ϕ field,

$$\bullet \qquad \phi = \frac{1}{i} \frac{1}{p^2 + \frac{p^4}{M^2}}$$

So that the four point amplitude can be computed by,

$$(1.6) \qquad \begin{array}{c} \phi_{2} \\ P \\ \phi_{3} \\ \phi_{4} \end{array} = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] + \langle 2P \rangle [2P] + \langle P1 \rangle [P1]) (\langle 34 \rangle [34] - \langle 4P \rangle [4P] - \langle P3 \rangle [P3]) \\ (1.7) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - \langle P2 \rangle [2P] - \langle P1 \rangle [1P]) (\langle 34 \rangle [34] + \langle P4 \rangle [4P] + \langle P3 \rangle [3P]) \\ (1.8) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] + \langle P|1 + 2|P]) (\langle 34 \rangle [34] - \langle P|3 + 4|P]) \\ (1.9) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - \langle P|P|P]) (\langle 34 \rangle [34] - \langle P|P|P]) \\ (1.10) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - 2P^{2}) (\langle 34 \rangle [34] - 2P^{2}) \\ (1.11) \\ = -\mathrm{i} \frac{(\kappa M)^{2}}{\mathrm{s}(M^{2} - \mathrm{s})} (\langle 12 \rangle [12] + 2\mathrm{s}) (\langle 34 \rangle [34] + 2\mathrm{s}) \end{array}$$

It's trivial to read the t and u channels from this expression,

(1.12)
$$\phi_{2} \qquad \phi_{3} = -i\frac{(\kappa M)^{2}}{t(M^{2} - t)} (\langle 23 \rangle [23] + 2t) (\langle 41 \rangle [41] + 2t)$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4} \qquad (\langle 24 \rangle [24] + 2u) (\langle 31 \rangle [31] + 2u)$$

$$\phi_{1} \qquad \phi_{2} \qquad \phi_{3} \qquad \phi_{4} \qquad \phi_{5} \qquad \phi_{6} \qquad \phi_{7} \qquad \phi_{8} \qquad$$

So that the full 4-point amplitude is,

$$\begin{array}{ll}
\phi_{2} & \phi_{3} \\
\phi_{2} & \phi_{3}
\end{array}$$

$$= -i \frac{(\kappa M)^{2}}{stu(M^{2} - s)(M^{2} - t)(M^{2} - u)} \left[(\langle 12 \rangle [12] + 2s)(\langle 34 \rangle [34] + 2s)tu(M^{2} - t)(M^{2} - u) + (\langle 23 \rangle [23] + 2t)(\langle 41 \rangle [41] + 2t) \right]$$

$$\phi_{1} & \phi_{4}$$

Let us specialize when 1, 2 are massless and 3, 4 are massive, then,

2. Conformal Toy Model

Consider the following Lagrangian,

$$\mathcal{L} = -\frac{1}{2}\Box\phi\Box\phi - \frac{g}{2}\phi^2\Box\phi - \frac{g^2}{8}\phi^4 + m^2\left(-\frac{1}{2}\partial_\mu\phi\partial^\mu\phi + \frac{g}{3!}\phi^3\right)$$

Notice the form of the Lagrangian,

$$\mathcal{L} = -\frac{1}{2} \left(\Box \phi + \frac{g}{2} \phi^2 \right)^2 + m^2 \left(-\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{g}{3!} \phi^3 \right)$$

It possesses the Feynman rules,

Let's compute the self energy,

(2.1)
$$i\Pi(p^2) = \dots + \dots + \dots + \dots$$

(2.2)
$$i\Pi^{(1)} = -\frac{3}{2}ig^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{1}{i} \frac{1}{\ell^2} \frac{1}{\ell^2 + m^2}$$

(2.3)
$$i\Pi^{(1)} = -\frac{3}{2}g^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{1}{\ell^2} \frac{1}{\ell^2 + m^2}$$

(2.4)
$$i\Pi^{(1)} = -\frac{3}{2}g^2 \frac{i}{(4\pi)^{\frac{D}{2}}\Gamma(\frac{D}{2})} (m^2)^{\frac{D}{2}-2} \frac{\Gamma(2-\frac{D}{2})\Gamma(\frac{D}{2}-1)}{\Gamma(1)}$$

(2.5)
$$i\Pi^{(1)} = -\frac{3}{2}ig^2 \frac{(m^2)^{-\epsilon}\Gamma(\epsilon)\Gamma(1-\epsilon)}{(4\pi)^{2-\epsilon}\Gamma(2-\epsilon)}$$

(2.6)
$$i\Pi^{(2)} = \frac{1}{2} (ig)^2 \int \frac{d^D \ell}{(2\pi)^D} \frac{1}{i^2} \frac{1}{\ell^2 (\ell+p)^2} \frac{\left(m^2 + \ell^2 + p^2 + (\ell+p)^2\right)^2}{\ell^2 + m^2} \frac{1}{(\ell+p)^2 + m^2}$$

For the mass renormalization we can take p = 0,

(2.7)
$$i\Pi^{(2)} = \frac{1}{2}g^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{\left(m^2 + 2\ell^2\right)^2}{\ell^4 (\ell^2 + m^2)^2}$$

Let's compute the four point amplitude for this theory,

(2.8)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{\left(p_{1}^{2} + p_{2}^{2} + (p_{1} + p_{2})^{2} + m^{2}\right) \left(p_{3}^{2} + p_{4}^{2} + (p_{3} + p_{4})^{2} + m^{2}\right)}{i(p_{1} + p_{2})^{2} \left((p_{1} + p_{2})^{2} + m^{2}\right)}$$

First let's consider all legs massless,

(2.9)
$$\phi_{2} \xrightarrow{\phi_{3}} = ig^{2} \frac{(-s+m^{2})(-s+m^{2})}{(-s)(-s+m^{2})} = -ig^{2} \frac{(-s+m^{2})}{s}$$

So,

(2.10)
$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{\left(-s+m^{2}\right)}{s} - ig^{2} \frac{\left(-t+m^{2}\right)}{t} - ig^{2} \frac{\left(-u+m^{2}\right)}{u} - 3ig^{2}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{\left(-s+m^{2}\right)}{s} - ig^{2} \frac{\left(-t+m^{2}\right)}{t} - ig^{2} \frac{\left(-u+m^{2}\right)}{u} - ig^{2} \frac{s}{s} - ig^{2} \frac{t}{t} - ig^{2} \frac{u}{u}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2}m^{2} \left(\frac{1}{s} + \frac{1}{t} + \frac{1}{u}\right) = ig^{2}m^{2} \left(\frac{1}{\langle 12\rangle[12]} + \frac{1}{\langle 14\rangle[14]} + \frac{1}{\langle 13\rangle[13]} \right)$$

$$\phi_{1} \qquad \phi_{3}$$

$$(2.12)$$

Para uma perna massiva, ϕ_4 ,

(2.13)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{(-s+m^{2})(-s)}{i(-s)(-s+m^{2})} = ig^{2}$$

$$\phi_{4} \qquad \phi_{4}$$

So,

(2.14)
$$\phi_{2} \qquad \phi_{3} = ig^{2} + ig^{2} + ig^{2} - 3ig^{2} = 0$$

$$\phi_{1} \qquad \phi_{4}$$

Para duas pernas massivas, $\phi_{3.4}$,

(2.15)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{\left(-s + m^{2}\right)\left(-s - m^{2}\right)}{i(-s)\left(-s + m^{2}\right)} = ig^{2} \frac{s + m^{2}}{s}$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{3} \qquad \phi_{3} = (ig)^{2} \frac{\left(-t\right)\left(-t\right)}{i(-t)\left(-t + m^{2}\right)} = -ig^{2} \frac{t}{-t + m^{2}}$$

(2.17)
$$\phi_{2} \qquad \phi_{4} = -ig^{2} \frac{u}{-u + m^{2}}$$

$$\phi_{1} \qquad \phi_{3} \qquad \phi_{3}$$

So,

$$(2.18) \qquad \phi_{2} \qquad \phi_{3} \qquad (2.18) \qquad = ig^{2} \frac{s+m^{2}}{s} - ig^{2} \frac{t}{-t+m^{2}} - ig^{2} \frac{u}{-u+m^{2}} - 3ig^{2}$$

$$(2.19) \qquad \phi_{3} \qquad = ig^{2} \frac{s+m^{2}}{s} - ig^{2} \frac{t}{-t+m^{2}} - ig^{2} \frac{u}{-u+m^{2}} - ig^{2} \frac{s}{s} - ig^{2} \frac{-t+m^{2}}{-t+m^{2}} - ig^{2} \frac{-u+m^{2}}{-u+m^{2}}$$

$$(2.19) \qquad \phi_{4} \qquad \phi_{4} \qquad \phi_{5} \qquad (2.20) \qquad = -ig^{2}m^{2} \left(-\frac{1}{s} + \frac{1}{-t+m^{2}} + \frac{1}{-u+m^{2}} \right) = -ig^{2}m^{2} \left(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 14 \rangle [14]} + \frac{1}{\langle 13 \rangle [13]} \right)$$

Para uma perna sem massa ϕ_1 ,

(2.21)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{(-s)(-s-m^{2})}{i(-s)(-s+m^{2})} = -ig^{2} \frac{s+m^{2}}{-s+m^{2}}$$

$$\phi_{2} \qquad \phi_{3} \qquad \phi_{3} = (ig)^{2} \frac{(-t)(-t-m^{2})}{i(-t)(-t+m^{2})} = -ig^{2} \frac{t+m^{2}}{-t+m^{2}}$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4}$$

(2.23)
$$\phi_{2} \qquad \phi_{4} = (ig)^{2} \frac{(-u)(-u-m^{2})}{i(-u)(-u+m^{2})} = -ig^{2} \frac{u+m^{2}}{-u+m^{2}}$$

$$\phi_{1} \qquad \phi_{3}$$

(2.24)

So,

$$(2.25) \qquad \phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{s+m^{2}}{-s+m^{2}} - ig^{2} \frac{t+m^{2}}{-t+m^{2}} - ig^{2} \frac{u+m^{2}}{-u+m^{2}} - 3ig^{2}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{s+m^{2}}{-s+m^{2}} - ig^{2} \frac{t+m^{2}}{-t+m^{2}} - ig^{2} \frac{u+m^{2}}{-u+m^{2}} - ig^{2} \frac{-s+m^{2}}{-s+m^{2}} - ig^{2} \frac{-t+m^{2}}{-t+m^{2}} - ig^{2} \frac{-u+m^{2}}{-u+m^{2}}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2}m^{2} \left(\frac{1}{-s+m^{2}} + \frac{1}{-t+m^{2}} + \frac{1}{-u+m^{2}} \right) = -ig^{2}m^{2} \left(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 14 \rangle [14]} + \frac{1}{\langle 13 \rangle [13]} \right)$$

Cut comparison, only massless legs

$$(2.28) \qquad l+3+4 \qquad (2.28)$$

$$\lim_{l \to \infty} \frac{1}{l} = \lim_{l \to \infty} \frac{ig^2 m^2 (igm^2)^2}{(im^2)^3} \left(\frac{1}{\langle 12 \rangle [12]} - \frac{1}{\langle 1l \rangle [1l]} - \frac{1}{\langle 2l \rangle [2l]} \right)$$

to solve for the cuts, $l^2 = (l+3+4)^2 = (l+4)^2 = 0$,

$$(2.29) l^2 = 0 \Rightarrow l = -|l\rangle[l]$$

(2.30)
$$0 = (l+4)^2 = \langle l4 \rangle [l4] = 0 \Rightarrow |l| = |4|$$

$$(2.31) 0 = (l+3+4)^2 = \langle lP_{34}\rangle[lP_{34}] + (3+4)^2 = \langle l|3+4|l] + \langle 34\rangle[34] = \langle l|3+4|4] + \langle 34\rangle[34]$$

$$(2.32) \qquad \langle 43 \rangle [34] = -\langle l3 \rangle [34] \Rightarrow |l\rangle = -|4\rangle + z|3\rangle$$

$$(2.33) l = -(-|4\rangle + z|3\rangle)[4]$$

The cuts are solved by this. Hence,

$$\begin{split} &=g^4\bigg(\frac{1}{\langle 12\rangle[12]}-\frac{1}{\langle 1l\rangle[1l]}-\frac{1}{\langle 2l\rangle[2l]}\bigg)\\ &=g^4\bigg(\frac{1}{\langle 12\rangle[12]}-\frac{1}{(-\langle 14\rangle+z\langle 13\rangle)[14]}-\frac{1}{(-\langle 24\rangle+z\langle 23\rangle)[24]}\bigg) \end{split}$$

Now for internal massive lines,

$$(2.34) \qquad l+3+4 \qquad \qquad l+3+4 \qquad \qquad = \frac{-\mathrm{i}g^2m^2\left(-\mathrm{i}gm^2\right)^2}{\left(-\mathrm{i}m^2\right)^3} \left(\frac{1}{\langle 12\rangle[12]} - \frac{1}{\langle 1l\rangle[1l]} - \frac{1}{\langle 2l\rangle[2l]}\right)$$

With the cuts being, $l^2 = (l+3+4)^2 = (l+4)^2 = -m^2$,

$$(2.35) 0 = (l+4)^2 - l^2 = 2l \cdot p_4$$

$$(2.36) 0 = (l+4+3)^2 - l^2 = 2l \cdot (4+3) + (4+3)^2 = 2l \cdot p_3 + (4+3)^2$$

As ansatz, $l = |4\rangle[4| + \alpha|4\rangle[3| + \beta|3\rangle[4|$ satisfy both conditions above. The remaining condition is,

$$(2.37) l^2 = -m^2$$

$$(2.38) -\alpha\beta[43]\langle 43\rangle = -m^2 \Rightarrow \alpha = \frac{m^2}{\beta\langle 34\rangle[34]}$$

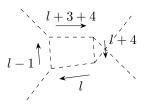
Setting now $-\beta = z$,

$$(2.39) l = |4\rangle[4| - \frac{m^2}{z\langle 34\rangle[34]}|4\rangle[3| - z|3\rangle[4|$$

The value of the diagram is,

$$\begin{split} &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 1l \rangle [l1]} + \frac{1}{\langle 2l \rangle [l2]} \bigg) \\ &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{-\langle 1|l|1]} + \frac{1}{-\langle 2|l|2]} \bigg) \\ &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} - \frac{1}{\langle 14 \rangle [41] - \frac{m^2}{z\langle 34 \rangle [34]} \langle 14 \rangle [31] - z\langle 13 \rangle [41]} - \frac{1}{\langle 2|l|2]} \bigg) \end{split}$$

The explicit cut loop amplitude is,



Triple cut has no improvement, what about a double cut,

$$(2.40) \qquad p_{1} \qquad p_{3} \qquad p_{3} \qquad = \frac{\left(ig^{2}m^{2}\right)^{2}}{\left(im^{2}\right)^{2}} \left(\frac{1}{\langle 12\rangle[12]} - \frac{1}{\langle 1\ell\rangle[1\ell]} - \frac{1}{\langle 2\ell\rangle[2\ell]}\right) \left(\frac{1}{\langle 34\rangle[34]} + \frac{1}{\langle 3\ell\rangle[3\ell]} + \frac{1}{\langle 4\ell\rangle[4\ell]}\right) \qquad p_{4} \qquad p_{4} \qquad p_{5} \qquad p_{6} \qquad p_{6} \qquad p_{7} \qquad p_{7} \qquad p_{8} \qquad p_{8} \qquad p_{8} \qquad p_{9} \qquad$$

Five point amplitude,

Let's consider the special case of all massless,

Combining this graph with,

$$3 - \frac{4}{1} = \frac{-i3g^{2}ig(m^{2} + (p_{1} + p_{2})^{2})}{i(p_{1} + p_{2})^{2}((p_{1} + p_{2})^{2} + m^{2})} = \frac{-3ig^{3}}{(p_{1} + p_{2})^{2}}$$

We get,

Now we have to sum the contributions of 1 being in the middle,

Which will be,

$$\frac{\mathrm{i}g^3}{\left(p_2+p_3\right)^2}\frac{m^2+\left(p_2+p_3\right)^2+\left(p_4+p_5\right)^2}{\left(p_4+p_5\right)^2}-\frac{3\mathrm{i}g^3}{\left(p_2+p_3\right)^2}-\frac{3\mathrm{i}g^3}{\left(p_4+p_5\right)^2}+\left(3\leftrightarrow 4,5\right)$$

Summing all the contributions we have,

$$\begin{split} &=\frac{\mathrm{i}g^{3}m^{2}}{(p_{1}+p_{2})^{2}}\left[\frac{1}{(p_{3}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}\right]+(2\leftrightarrow3,4,5)\\ &+\mathrm{i}g^{3}\left[\frac{1}{(p_{3}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{2}+p_{4})^{2}}+\frac{1}{(p_{2}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{2})^{2}}+\frac{1}{(p_{5}+p_{2})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{5})^{2}}+\frac{1}{(p_$$

By residue, any amplitude with just one massive external on-shell leg is zero. For two massive external on-shell legs, let's take as massive 1, 2,

Combining this graph with,

$$3 - \frac{4}{5}$$

$$2 - 3ig^{3} \frac{(p_{1} + p_{2})^{2} - m^{2}}{(p_{1} + p_{2})^{2} \left(m^{2} + (p_{1} + p_{2})^{2}\right)}$$

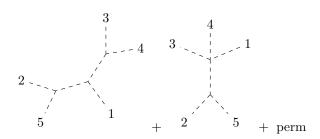
so,

$$\begin{array}{ll}
4 \\
5 \\
2 \\
1 \\
 &= ig^3 \frac{\left((p_1 + p_2)^2 - m^2\right)}{\left(p_1 + p_2\right)^2 \left(m^2 + (p_1 + p_2)^2\right)} \left[\frac{\left(m^2 + (p_1 + p_2)^2 + (p_3 + p_4)^2\right)}{\left(p_3 + p_4\right)^2} - \frac{\left(p_3 + p_4\right)^2}{\left(p_3 + p_4\right)^2} + (5 \leftrightarrow 3, 4) \right] \\
 &= ig^3 \frac{\left((p_1 + p_2)^2 - m^2\right)}{\left(p_1 + p_2\right)^2} \left[\frac{1}{\left(p_3 + p_4\right)^2} + \frac{1}{\left(p_5 + p_4\right)^2} + \frac{1}{\left(p_3 + p_5\right)^2} \right]
\end{array}$$

The other contributions are,

$$3 \qquad 1 = \frac{\mathrm{i} g^3 (p_1 + p_3)^2}{\left(m^2 + (p_1 + p_3)^2\right)} \left[\frac{1}{m^2 + (p_2 + p_4)^2} + \frac{1}{m^2 + (p_2 + p_5)^2} + \frac{1}{(p_4 + p_5)^2} \right]$$
 um the contributions of 1 being in the middle,

Now we have to sum the contributions of 1 being in the middle,



which are,

$$=ig^{3}\frac{\left(p_{2}+p_{3}\right)^{2}+\left(p_{4}+p_{5}\right)^{2}}{\left(m^{2}+\left(p_{2}+p_{3}\right)^{2}\right)\left(p_{4}+p_{5}\right)^{2}}-\frac{3ig^{3}}{m^{2}+\left(p_{2}+p_{3}\right)^{2}}-\frac{3ig^{3}}{\left(p_{4}+p_{5}\right)^{2}}+\left(3\leftrightarrow4,5\right)$$