SCALAR PROXY

VICENTE V. FIGUEIRA

1. Introduction

We will work most with the scalar proxy given by the lagrangian,

(1.1)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi - \frac{1}{2M^{2}}\Box\phi\Box\phi - \frac{\kappa}{2}\Box\phi\phi^{2}$$

The idea here is reintegrate the higher derivative term, in order to obtain a lower derivative term, but in terms of additional fields. This is easily done by,

1

(1.2)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi + \Box\phi\eta + \frac{M^{2}}{2}\eta^{2} - \frac{\kappa}{2}\Box\phi\phi^{2}$$

The new lagrangian has mixed propagator terms, to diagonalize it is also easy, we just open in terms of $\phi = h - \eta$,

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \partial_{\mu}h\partial^{\mu}\eta - \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2} + \eta\Box(h-\eta) + \frac{M^{2}}{2}\eta^{2}$$

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \partial_{\mu}h\partial^{\mu}\eta - \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2} - \partial_{\mu}\eta\partial^{\mu}(h-\eta) + \frac{M^{2}}{2}\eta^{2}$$

(1.5)
$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \frac{1}{2}\partial_{\mu}\eta\partial^{\mu}\eta + \frac{M^{2}}{2}\eta^{2} - \frac{\kappa}{2}\Box(h-\eta)(h-\eta)^{2}$$

The Feynman rules are easily red as,

•
$$h - - - - - h = \frac{1}{i} \frac{1}{p^2}$$

$$\bullet \qquad \eta = -\frac{1}{\mathrm{i}} \frac{1}{p^2 + M^2}$$

•
$$h_1 - \cdots = i\kappa (p_1^2 + p_2^2 + p_3^2)$$

•
$$h_1 - - - \prec (p_1^2 + p_2^2 + p_3^2)$$

•
$$h_1 = i\kappa (p_1^2 + p_2^2 + p_3^2)$$

•
$$\eta_1 = -i\kappa (p_1^2 + p_2^2 + p_3^2)$$

Which can also be seen directly from the Feynman rules of the ϕ field,

$$\bullet \qquad \phi = \frac{1}{i} \frac{1}{p^2 + \frac{p^4}{M^2}}$$

So that the four point amplitude can be computed by,

$$(1.6) \qquad \begin{array}{c} \phi_{2} \\ P \\ \phi_{3} \\ \phi_{4} \end{array} = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] + \langle 2P \rangle [2P] + \langle P1 \rangle [P1]) (\langle 34 \rangle [34] - \langle 4P \rangle [4P] - \langle P3 \rangle [P3]) \\ (1.7) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - \langle P2 \rangle [2P] - \langle P1 \rangle [1P]) (\langle 34 \rangle [34] + \langle P4 \rangle [4P] + \langle P3 \rangle [3P]) \\ (1.8) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] + \langle P|1 + 2|P]) (\langle 34 \rangle [34] - \langle P|3 + 4|P]) \\ (1.9) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - \langle P|P|P]) (\langle 34 \rangle [34] - \langle P|P|P]) \\ (1.10) \\ = \frac{1}{\mathrm{i}} (-\mathrm{i}\kappa)^{2} \frac{1}{P^{2} + \frac{P^{4}}{M^{2}}} (\langle 12 \rangle [12] - 2P^{2}) (\langle 34 \rangle [34] - 2P^{2}) \\ (1.11) \\ = -\mathrm{i} \frac{(\kappa M)^{2}}{\mathrm{s}(M^{2} - \mathrm{s})} (\langle 12 \rangle [12] + 2\mathrm{s}) (\langle 34 \rangle [34] + 2\mathrm{s}) \end{array}$$

It's trivial to read the t and u channels from this expression,

(1.12)
$$\phi_{2} \qquad \phi_{3} = -i\frac{(\kappa M)^{2}}{t(M^{2} - t)} (\langle 23 \rangle [23] + 2t) (\langle 41 \rangle [41] + 2t)$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4} \qquad (\langle 24 \rangle [24] + 2u) (\langle 31 \rangle [31] + 2u)$$

$$\phi_{1} \qquad \phi_{2} \qquad \phi_{3} \qquad \phi_{4} \qquad \phi_{5} \qquad \phi_{6} \qquad \phi_{7} \qquad \phi_{8} \qquad$$

So that the full 4-point amplitude is,

$$\begin{array}{ll}
\phi_{2} & \phi_{3} \\
\phi_{2} & \phi_{3}
\end{array}$$

$$= -i \frac{(\kappa M)^{2}}{stu(M^{2} - s)(M^{2} - t)(M^{2} - u)} \left[(\langle 12 \rangle [12] + 2s)(\langle 34 \rangle [34] + 2s)tu(M^{2} - t)(M^{2} - u) + (\langle 23 \rangle [23] + 2t)(\langle 41 \rangle [41] + 2t) \right]$$

$$\phi_{1} & \phi_{4}$$

Let us specialize when 1, 2 are massless and 3, 4 are massive, then,

2. Conformal Toy Model

Consider the following Lagrangian,

$$\mathcal{L} = -\frac{1}{2}\Box\phi\Box\phi - \frac{g}{2}\phi^2\Box\phi - \frac{g^2}{8}\phi^4 + m^2\left(-\frac{1}{2}\partial_\mu\phi\partial^\mu\phi + \frac{g}{3!}\phi^3\right)$$

Notice the form of the Lagrangian,

$$\mathcal{L} = -\frac{1}{2} \left(\Box \phi + \frac{g}{2} \phi^2 \right)^2 + m^2 \left(-\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{g}{3!} \phi^3 \right)$$

It possesses the Feynman rules,

Let's compute the self energy,

(2.1)
$$i\Pi(p^2) = \dots + \dots + \dots + \dots$$

(2.2)
$$i\Pi^{(1)} = -\frac{3}{2}ig^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{1}{i} \frac{1}{\ell^2} \frac{1}{\ell^2 + m^2}$$

(2.3)
$$i\Pi^{(1)} = -\frac{3}{2}g^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{1}{\ell^2} \frac{1}{\ell^2 + m^2}$$

(2.4)
$$i\Pi^{(1)} = -\frac{3}{2}g^2 \frac{i}{(4\pi)^{\frac{D}{2}}\Gamma(\frac{D}{2})} (m^2)^{\frac{D}{2}-2} \frac{\Gamma(2-\frac{D}{2})\Gamma(\frac{D}{2}-1)}{\Gamma(1)}$$

(2.5)
$$i\Pi^{(1)} = -\frac{3}{2}ig^2 \frac{(m^2)^{-\epsilon}\Gamma(\epsilon)\Gamma(1-\epsilon)}{(4\pi)^{2-\epsilon}\Gamma(2-\epsilon)}$$

(2.6)
$$i\Pi^{(2)} = \frac{1}{2} (ig)^2 \int \frac{d^D \ell}{(2\pi)^D} \frac{1}{i^2} \frac{1}{\ell^2 (\ell+p)^2} \frac{\left(m^2 + \ell^2 + p^2 + (\ell+p)^2\right)^2}{\ell^2 + m^2} \frac{1}{(\ell+p)^2 + m^2}$$

For the mass renormalization we can take p = 0,

(2.7)
$$i\Pi^{(2)} = \frac{1}{2}g^2 \int \frac{\mathrm{d}^D \ell}{(2\pi)^D} \frac{\left(m^2 + 2\ell^2\right)^2}{\ell^4 (\ell^2 + m^2)^2}$$

Let's compute the four point amplitude for this theory,

(2.8)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{\left(p_{1}^{2} + p_{2}^{2} + (p_{1} + p_{2})^{2} + m^{2}\right) \left(p_{3}^{2} + p_{4}^{2} + (p_{3} + p_{4})^{2} + m^{2}\right)}{i(p_{1} + p_{2})^{2} \left((p_{1} + p_{2})^{2} + m^{2}\right)}$$

First let's consider all legs massless,

(2.9)
$$\phi_{2} \xrightarrow{\phi_{3}} = ig^{2} \frac{(-s+m^{2})(-s+m^{2})}{(-s)(-s+m^{2})} = -ig^{2} \frac{(-s+m^{2})}{s}$$

So,

(2.10)
$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{\left(-s+m^{2}\right)}{s} - ig^{2} \frac{\left(-t+m^{2}\right)}{t} - ig^{2} \frac{\left(-u+m^{2}\right)}{u} - 3ig^{2}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{\left(-s+m^{2}\right)}{s} - ig^{2} \frac{\left(-t+m^{2}\right)}{t} - ig^{2} \frac{\left(-u+m^{2}\right)}{u} - ig^{2} \frac{s}{s} - ig^{2} \frac{t}{t} - ig^{2} \frac{u}{u}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2}m^{2} \left(\frac{1}{s} + \frac{1}{t} + \frac{1}{u}\right) = ig^{2}m^{2} \left(\frac{1}{\langle 12\rangle[12]} + \frac{1}{\langle 14\rangle[14]} + \frac{1}{\langle 13\rangle[13]} \right)$$

$$\phi_{1} \qquad \phi_{3}$$

$$(2.12)$$

Para uma perna massiva, ϕ_4 ,

(2.13)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{(-s+m^{2})(-s)}{i(-s)(-s+m^{2})} = ig^{2}$$

$$\phi_{4} \qquad \phi_{4}$$

So,

(2.14)
$$\phi_{2} \qquad \phi_{3} = ig^{2} + ig^{2} + ig^{2} - 3ig^{2} = 0$$

$$\phi_{1} \qquad \phi_{4}$$

Para duas pernas massivas, $\phi_{3.4}$,

(2.15)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{\left(-s + m^{2}\right)\left(-s - m^{2}\right)}{i(-s)\left(-s + m^{2}\right)} = ig^{2} \frac{s + m^{2}}{s}$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{3} \qquad \phi_{3} = (ig)^{2} \frac{\left(-t\right)\left(-t\right)}{i(-t)\left(-t + m^{2}\right)} = -ig^{2} \frac{t}{-t + m^{2}}$$

(2.17)
$$\phi_{2} \qquad \phi_{4} = -ig^{2} \frac{u}{-u + m^{2}}$$

$$\phi_{1} \qquad \phi_{3} \qquad \phi_{3}$$

So,

$$(2.18) \qquad \phi_{2} \qquad \phi_{3} \qquad (2.18) \qquad = ig^{2} \frac{s+m^{2}}{s} - ig^{2} \frac{t}{-t+m^{2}} - ig^{2} \frac{u}{-u+m^{2}} - 3ig^{2}$$

$$(2.19) \qquad \phi_{3} \qquad = ig^{2} \frac{s+m^{2}}{s} - ig^{2} \frac{t}{-t+m^{2}} - ig^{2} \frac{u}{-u+m^{2}} - ig^{2} \frac{s}{s} - ig^{2} \frac{-t+m^{2}}{-t+m^{2}} - ig^{2} \frac{-u+m^{2}}{-u+m^{2}}$$

$$(2.19) \qquad \phi_{4} \qquad \phi_{4} \qquad \phi_{5} \qquad (2.20) \qquad = -ig^{2}m^{2} \left(-\frac{1}{s} + \frac{1}{-t+m^{2}} + \frac{1}{-u+m^{2}} \right) = -ig^{2}m^{2} \left(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 14 \rangle [14]} + \frac{1}{\langle 13 \rangle [13]} \right)$$

Para uma perna sem massa ϕ_1 ,

(2.21)
$$\phi_{2} \qquad \phi_{3} = (ig)^{2} \frac{(-s)(-s-m^{2})}{i(-s)(-s+m^{2})} = -ig^{2} \frac{s+m^{2}}{-s+m^{2}}$$

$$\phi_{2} \qquad \phi_{3} \qquad \phi_{3} = (ig)^{2} \frac{(-t)(-t-m^{2})}{i(-t)(-t+m^{2})} = -ig^{2} \frac{t+m^{2}}{-t+m^{2}}$$

$$\phi_{1} \qquad \phi_{4} \qquad \phi_{4} \qquad \phi_{2} \qquad \phi_{4}$$

(2.23)
$$\phi_{2} \qquad \phi_{4} = (ig)^{2} \frac{(-u)(-u-m^{2})}{i(-u)(-u+m^{2})} = -ig^{2} \frac{u+m^{2}}{-u+m^{2}}$$

$$\phi_{1} \qquad \phi_{3}$$

(2.24)

So,

$$(2.25) \qquad \phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{s+m^{2}}{-s+m^{2}} - ig^{2} \frac{t+m^{2}}{-t+m^{2}} - ig^{2} \frac{u+m^{2}}{-u+m^{2}} - 3ig^{2}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2} \frac{s+m^{2}}{-s+m^{2}} - ig^{2} \frac{t+m^{2}}{-t+m^{2}} - ig^{2} \frac{u+m^{2}}{-u+m^{2}} - ig^{2} \frac{-s+m^{2}}{-s+m^{2}} - ig^{2} \frac{-t+m^{2}}{-t+m^{2}} - ig^{2} \frac{-u+m^{2}}{-u+m^{2}}$$

$$\phi_{1} \qquad \phi_{4}$$

$$\phi_{2} \qquad \phi_{3}$$

$$= -ig^{2}m^{2} \left(\frac{1}{-s+m^{2}} + \frac{1}{-t+m^{2}} + \frac{1}{-u+m^{2}} \right) = -ig^{2}m^{2} \left(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 14 \rangle [14]} + \frac{1}{\langle 13 \rangle [13]} \right)$$

Cut comparison, only massless legs

$$(2.28) \qquad l+3+4 \qquad (2.28)$$

$$\lim_{l \to \infty} \frac{1}{l} = \lim_{l \to \infty} \frac{ig^2 m^2 (igm^2)^2}{(im^2)^3} \left(\frac{1}{\langle 12 \rangle [12]} - \frac{1}{\langle 1l \rangle [1l]} - \frac{1}{\langle 2l \rangle [2l]} \right)$$

to solve for the cuts, $l^2 = (l+3+4)^2 = (l+4)^2 = 0$,

$$(2.29) l^2 = 0 \Rightarrow l = -|l\rangle[l]$$

(2.30)
$$0 = (l+4)^2 = \langle l4 \rangle [l4] = 0 \Rightarrow |l| = |4|$$

$$(2.31) 0 = (l+3+4)^2 = \langle lP_{34}\rangle[lP_{34}] + (3+4)^2 = \langle l|3+4|l] + \langle 34\rangle[34] = \langle l|3+4|4] + \langle 34\rangle[34]$$

$$(2.32) \qquad \langle 43 \rangle [34] = -\langle l3 \rangle [34] \Rightarrow |l\rangle = -|4\rangle + z|3\rangle$$

$$(2.33) l = -(-|4\rangle + z|3\rangle)[4]$$

The cuts are solved by this. Hence,

$$\begin{split} &=g^4\bigg(\frac{1}{\langle 12\rangle[12]}-\frac{1}{\langle 1l\rangle[1l]}-\frac{1}{\langle 2l\rangle[2l]}\bigg)\\ &=g^4\bigg(\frac{1}{\langle 12\rangle[12]}-\frac{1}{(-\langle 14\rangle+z\langle 13\rangle)[14]}-\frac{1}{(-\langle 24\rangle+z\langle 23\rangle)[24]}\bigg) \end{split}$$

Now for internal massive lines,

$$(2.34) \qquad l+3+4 \qquad \qquad l+3+4 \qquad \qquad = \frac{-\mathrm{i}g^2m^2\left(-\mathrm{i}gm^2\right)^2}{\left(-\mathrm{i}m^2\right)^3} \left(\frac{1}{\langle 12\rangle[12]} - \frac{1}{\langle 1l\rangle[1l]} - \frac{1}{\langle 2l\rangle[2l]}\right)$$

With the cuts being, $l^2 = (l+3+4)^2 = (l+4)^2 = -m^2$,

$$(2.35) 0 = (l+4)^2 - l^2 = 2l \cdot p_4$$

$$(2.36) 0 = (l+4+3)^2 - l^2 = 2l \cdot (4+3) + (4+3)^2 = 2l \cdot p_3 + (4+3)^2$$

As ansatz, $l = |4\rangle[4| + \alpha|4\rangle[3| + \beta|3\rangle[4|$ satisfy both conditions above. The remaining condition is,

$$(2.37) l^2 = -m^2$$

$$(2.38) -\alpha\beta[43]\langle 43\rangle = -m^2 \Rightarrow \alpha = \frac{m^2}{\beta\langle 34\rangle[34]}$$

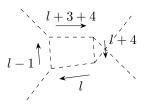
Setting now $-\beta = z$,

$$(2.39) l = |4\rangle[4| - \frac{m^2}{z\langle 34\rangle[34]}|4\rangle[3| - z|3\rangle[4|$$

The value of the diagram is,

$$\begin{split} &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{\langle 1l \rangle [l1]} + \frac{1}{\langle 2l \rangle [l2]} \bigg) \\ &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} + \frac{1}{-\langle 1|l|1]} + \frac{1}{-\langle 2|l|2]} \bigg) \\ &=g^4 \bigg(\frac{1}{\langle 12 \rangle [12]} - \frac{1}{\langle 14 \rangle [41] - \frac{m^2}{z\langle 34 \rangle [34]} \langle 14 \rangle [31] - z\langle 13 \rangle [41]} - \frac{1}{\langle 2|l|2]} \bigg) \end{split}$$

The explicit cut loop amplitude is,



Triple cut has no improvement, what about a double cut,

$$(2.40) \qquad p_{1} \qquad p_{3} \qquad p_{3} \qquad = \frac{\left(ig^{2}m^{2}\right)^{2}}{\left(im^{2}\right)^{2}} \left(\frac{1}{\langle 12\rangle[12]} - \frac{1}{\langle 1\ell\rangle[1\ell]} - \frac{1}{\langle 2\ell\rangle[2\ell]}\right) \left(\frac{1}{\langle 34\rangle[34]} + \frac{1}{\langle 3\ell\rangle[3\ell]} + \frac{1}{\langle 4\ell\rangle[4\ell]}\right) \qquad p_{4} \qquad p_{4} \qquad p_{5} \qquad p_{4} \qquad p_{5} \qquad p_{6} \qquad p_{6} \qquad p_{7} \qquad p_{7} \qquad p_{8} \qquad p_{7} \qquad p_{8} \qquad p_{8} \qquad p_{9} \qquad$$

Five point amplitude,

Let's consider the special case of all massless,

Combining this graph with,

$$3 - \frac{4}{1} = \frac{-i3g^{2}ig(m^{2} + (p_{1} + p_{2})^{2})}{i(p_{1} + p_{2})^{2}((p_{1} + p_{2})^{2} + m^{2})} = \frac{-3ig^{3}}{(p_{1} + p_{2})^{2}}$$

We get,

Now we have to sum the contributions of 1 being in the middle,

Which will be,

$$\frac{\mathrm{i}g^3}{\left(p_2+p_3\right)^2}\frac{m^2+\left(p_2+p_3\right)^2+\left(p_4+p_5\right)^2}{\left(p_4+p_5\right)^2}-\frac{3\mathrm{i}g^3}{\left(p_2+p_3\right)^2}-\frac{3\mathrm{i}g^3}{\left(p_4+p_5\right)^2}+\left(3\leftrightarrow 4,5\right)$$

Summing all the contributions we have,

$$\begin{split} &=\frac{\mathrm{i}g^{3}m^{2}}{(p_{1}+p_{2})^{2}}\left[\frac{1}{(p_{3}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}\right]+(2\leftrightarrow3,4,5)\\ &+\mathrm{i}g^{3}\left[\frac{1}{(p_{3}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{2}+p_{4})^{2}}+\frac{1}{(p_{2}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{3}+p_{5})^{2}}+\frac{1}{(p_{5}+p_{2})^{2}}+\frac{1}{(p_{5}+p_{2})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{4})^{2}}+\frac{1}{(p_{5}+p_{5})^{2}}+\frac{1}{(p_$$

By residue, any amplitude with just one massive external on-shell leg is zero. For two massive external on-shell legs, let's take as massive 1, 2,

Combining this graph with,

$$3 - \frac{4}{5}$$

$$2 - 3ig^{3} \frac{(p_{1} + p_{2})^{2} - m^{2}}{(p_{1} + p_{2})^{2} \left(m^{2} + (p_{1} + p_{2})^{2}\right)}$$

so,

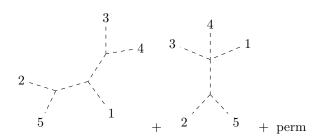
$$\begin{array}{ll}
4 \\
5 \\
2 \\
1 \\
 &= ig^3 \frac{\left((p_1 + p_2)^2 - m^2\right)}{\left(p_1 + p_2\right)^2 \left(m^2 + (p_1 + p_2)^2\right)} \left[\frac{\left(m^2 + (p_1 + p_2)^2 + (p_3 + p_4)^2\right)}{\left(p_3 + p_4\right)^2} - \frac{\left(p_3 + p_4\right)^2}{\left(p_3 + p_4\right)^2} + (5 \leftrightarrow 3, 4) \right] \\
 &= ig^3 \frac{\left((p_1 + p_2)^2 - m^2\right)}{\left(p_1 + p_2\right)^2} \left[\frac{1}{\left(p_3 + p_4\right)^2} + \frac{1}{\left(p_5 + p_4\right)^2} + \frac{1}{\left(p_3 + p_5\right)^2} \right]
\end{array}$$

The other contributions are,

$$2 \xrightarrow{\frac{4}{5}} 5$$

$$3 \xrightarrow{1} = \frac{ig^3(p_1 + p_3)^2}{\left(m^2 + (p_1 + p_3)^2\right)} \left[\frac{1}{m^2 + (p_2 + p_4)^2} + \frac{1}{m^2 + (p_2 + p_5)^2} + \frac{1}{(p_4 + p_5)^2} \right]$$
In the contributions of 1 being in the middle,

Now we have to sum the contributions of 1 being in the middle



which are,

$$=ig^{3}\frac{\left(p_{2}+p_{3}\right)^{2}+\left(p_{4}+p_{5}\right)^{2}}{\left(m^{2}+\left(p_{2}+p_{3}\right)^{2}\right)\left(p_{4}+p_{5}\right)^{2}}-\frac{3ig^{3}}{m^{2}+\left(p_{2}+p_{3}\right)^{2}}-\frac{3ig^{3}}{\left(p_{4}+p_{5}\right)^{2}}+\left(3\leftrightarrow4,5\right)$$

So, summing all the contributions

$$\begin{split} &= \mathrm{i} g^3 \frac{\left((p_1 + p_2)^2 - m^2\right)}{(p_1 + p_2)^2} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \frac{\mathrm{i} g^3 (p_1 + p_3)^2}{\left(m^2 + (p_1 + p_3)^2\right)} \left[\frac{1}{m^2 + (p_2 + p_4)^2} + \frac{1}{m^2 + (p_2 + p_5)^2} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5) \\ &\quad + \mathrm{i} g^3 \frac{(p_2 + p_3)^2 + (p_4 + p_5)^2}{\left(m^2 + (p_2 + p_3)^2\right) (p_4 + p_5)^2} - \frac{3\mathrm{i} g^3}{m^2 + (p_2 + p_3)^2} - \frac{3\mathrm{i} g^3}{(p_4 + p_5)^2} + (3 \leftrightarrow 4, 5) \\ &= -\mathrm{i} g^3 \frac{m^2}{(p_1 + p_2)^2} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \mathrm{i} g^3 \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \mathrm{i} g^3 \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5) \\ &\quad - \mathrm{i} g^3 \frac{m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{2p_2 \cdot p_4} + \frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5) \\ &\quad + \mathrm{i} g^3 \frac{-m^2 + 2p_2 \cdot p_3 + (p_4 + p_5)^2}{(2p_2 \cdot p_3)(p_4 + p_5)^2} - \frac{3\mathrm{i} g^3}{2p_2 \cdot p_3} - \frac{3\mathrm{i} g^3}{(p_4 + p_5)^2} + (3 \leftrightarrow 4, 5) \\ &\quad = -\frac{\mathrm{i} g^3 m^2}{(p_1 + p_2)^2} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \mathrm{i} g^3 \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \mathrm{i} g^3 \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \mathrm{i} g^3 \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad - \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad - \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{2p_2 \cdot p_4} + \frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5) \\ &\quad - \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad - \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{2p_2 \cdot p_4} + \frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5) \\ &\quad - \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} + \frac{1}{(p_3 + p_5)^2} \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} + \frac{1}{(p_3 + p_5)^2} \right] \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}$$

$$-\frac{\mathrm{i}g^3m^2}{(2p_2 \cdot p_3)(p_4 + p_5)^2} + (3 \leftrightarrow 4, 5)$$

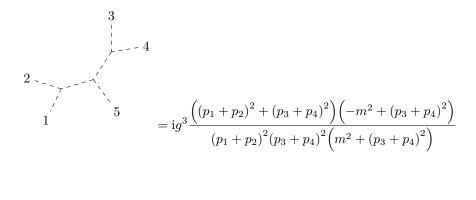
$$-\frac{2\mathrm{i}g^3}{2p_2 \cdot p_3} - \frac{2\mathrm{i}g^3}{(p_4 + p_5)^2} - \frac{2\mathrm{i}g^3}{2p_2 \cdot p_4} - \frac{2\mathrm{i}g^3}{(p_3 + p_5)^2} - \frac{2\mathrm{i}g^3}{2p_2 \cdot p_5} - \frac{2\mathrm{i}g^3}{(p_4 + p_3)^2}$$

$$= -\frac{\mathrm{i}g^3m^2}{(p_1 + p_2)^2} \left[\frac{1}{(p_3 + p_4)^2} + \frac{1}{(p_5 + p_4)^2} + \frac{1}{(p_3 + p_5)^2} \right]$$

$$-\frac{\mathrm{i}g^3m^2}{(2p_1 \cdot p_3)} \left[\frac{1}{2p_2 \cdot p_4} + \frac{1}{2p_2 \cdot p_5} + \frac{1}{(p_4 + p_5)^2} \right] + (3 \leftrightarrow 4, 5)$$

$$-\frac{\mathrm{i}g^3m^2}{(2p_2 \cdot p_3)(p_4 + p_5)^2} + (3 \leftrightarrow 4, 5)$$

Is almost the same of the all massless, but with a different denominator in the 1, 2 channel. Now with three massive legs, being 3, 4, 5,



With,

$$3 - \frac{4}{1} = \frac{-i3g^{2}ig(m^{2} + (p_{1} + p_{2})^{2})}{i(p_{1} + p_{2})^{2}((p_{1} + p_{2})^{2} + m^{2})} = \frac{-3ig^{3}}{(p_{1} + p_{2})^{2}}$$

So,

$$\begin{array}{ll}
3 & \stackrel{4}{\longrightarrow} 5 \\
2 & 1 & = ig^{3} \frac{\left((p_{1} + p_{2})^{2} + (p_{3} + p_{4})^{2}\right)\left(-m^{2} + (p_{3} + p_{4})^{2}\right)}{\left(p_{1} + p_{2}\right)^{2}\left(p_{3} + p_{4}\right)^{2}\left(m^{2} + (p_{3} + p_{4})^{2}\right)} - \frac{3ig^{3}}{\left(p_{1} + p_{2}\right)^{2}} \\
& = \frac{ig^{3}}{\left(p_{1} + p_{2}\right)^{2}} \left[\frac{\left((p_{1} + p_{2})^{2} + (p_{3} + p_{4})^{2}\right)\left(-m^{2} + (p_{3} + p_{4})^{2}\right)}{\left(p_{3} + p_{4}\right)^{2}\left(m^{2} + (p_{3} + p_{4})^{2}\right)} - \frac{\left(p_{3} + p_{4}\right)^{2}\left(m^{2} + (p_{3} + p_{4})^{2}\right)}{\left(p_{3} + p_{4}\right)^{2}\left(m^{2} + (p_{3} + p_{4})^{2}\right)} \right] + (5 \leftrightarrow 3, 4) \\
& = \frac{ig^{3}}{\left(p_{1} + p_{2}\right)^{2}} \left[\frac{\left(p_{1} + p_{2}\right)^{2}\left(-m^{2} + (p_{3} + p_{4})^{2}\right) - 2m^{2}(p_{3} + p_{4})^{2}}{\left(p_{3} + p_{4}\right)^{2}\left(m^{2} + (p_{3} + p_{4})^{2}\right)} \right] + (5 \leftrightarrow 3, 4)
\end{array}$$

The other topology is,

$$\begin{array}{ll}
3 \\
& = ig^{3} \frac{\left(m^{2} + p_{3}^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + p_{5}^{2} + (p_{1} + p_{3})^{2} + (p_{2} + p_{4})^{2}\right)\left(m^{2} + p_{4}^{2} + (p_{2} + p_{4})^{2}\right)}{\left(p_{1} + p_{3}\right)^{2}\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(p_{2} + p_{4}\right)^{2}\left(m^{2} + (p_{2} + p_{4})^{2}\right)} \\
& = ig^{3} \frac{\left((p_{1} + p_{3})^{2} + (p_{2} + p_{4})^{2}\right)}{\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{2} + p_{4})^{2}\right)} \\
& = ig^{3} \frac{\left(m^{2} + p_{3}^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{1} + p_{3})^{2} + (p_{5} + p_{4})^{2}\right)}{\left(p_{1} + p_{3}\right)^{2}\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(p_{5} + p_{4}\right)^{2}\left(m^{2} + (p_{5} + p_{4})^{2}\right)} \\
& = ig^{3} \frac{\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{1} + p_{3})^{2} + (p_{5} + p_{4})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)}{\left(p_{1} + p_{3}\right)^{2}\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(p_{5} + p_{4}\right)^{2}\left(m^{2} + (p_{5} + p_{4})^{2}\right)} \\
& = ig^{3} \frac{\left(m^{2} + (p_{1} + p_{3})^{2} + (p_{5} + p_{4})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)}{\left(m^{2} + (p_{1} + p_{3})^{2} + (p_{5} + p_{4})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)} \\
& = ig^{3} \frac{\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)\left(m^{2} + (p_{5} + p_{4})^{2}\right)}{\left(p_{1} + (p_{1} + p_{3})^{2}\right)\left(m^{2} + (p_$$

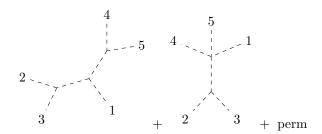
Also,

So,

$$\begin{array}{ll}
3 & 1 & = \frac{ig^3}{m^2 + (p_1 + p_3)^2} \left[\frac{\left((p_1 + p_3)^2 + (p_2 + p_4)^2\right) - m^2 - (p_2 + p_4)^2}{\left(m^2 + (p_2 + p_4)^2\right)} + \frac{\left((p_1 + p_3)^2 + (p_2 + p_5)^2\right) - m^2 - (p_2 + p_5)^2}{\left(m^2 + (p_2 + p_5)^2\right)} \right] \\
& + \frac{ig^3}{m^2 + (p_1 + p_3)^2} \left[\frac{\left(m^2 + (p_1 + p_3)^2 + (p_4 + p_5)^2\right) \left(-m^2 + (p_4 + p_5)^2\right) - (p_4 + p_5)^2 \left(m^2 + (p_4 + p_5)^2\right)}{\left(p_4 + p_5\right)^2 \left(m^2 + (p_4 + p_5)^2\right)} \right] \\
& = \frac{ig^3 \left((p_1 + p_3)^2 - m^2\right)}{m^2 + (p_1 + p_3)^2} \left[\frac{1}{m^2 + (p_2 + p_4)^2} + \frac{1}{m^2 + (p_2 + p_5)^2} \right]
\end{array}$$

$$\begin{split} &+\frac{\mathrm{i}g^3}{m^2+(p_1+p_3)^2}\left[\frac{-m^4-m^2(p_1+p_3)^2+(p_1+p_3)^2(p_4+p_5)^2-m^2(p_4+p_5)^2}{(p_4+p_5)^2\Big(m^2+(p_4+p_5)^2\Big)}\right]\\ &=\frac{\mathrm{i}g^3\Big(\big(p_1+p_3\big)^2-m^2\Big)}{m^2+\big(p_1+p_3\big)^2}\left[\frac{1}{m^2+(p_2+p_4)^2}+\frac{1}{m^2+(p_2+p_5)^2}+\frac{1}{m^2+(p_4+p_5)^2}\right]+\big(3\leftrightarrow4,5\big)\\ &-\frac{\mathrm{i}g^3m^2}{\big(p_4+p_5\big)^2\Big(m^2+(p_4+p_5)^2\Big)}+\big(3\leftrightarrow4,5\big) \end{split}$$

Now, with 1 in the middle,



Which is,

$$=\frac{\mathrm{i}g^{3}\Big(m^{2}+p_{3}^{2}+(p_{2}+p_{3})^{2}\Big)\Big(m^{2}+p_{1}^{2}+(p_{2}+p_{3})^{2}+(p_{5}+p_{4})^{2}\Big)\Big(m^{2}+p_{5}^{2}+p_{4}^{2}+(p_{5}+p_{4})^{2}\Big)}{(p_{2}+p_{3})^{2}\Big(m^{2}+(p_{2}+p_{3})^{2}\Big)(p_{4}+p_{5})^{2}\Big(m^{2}+(p_{4}+p_{5})^{2}\Big)}-\frac{3\mathrm{i}g^{3}}{m^{2}+(p_{2}+p_{3})^{2}}-\frac{3\mathrm{i}g^{3}\Big(-m^{2}+(p_{4}+p_{5})^{2}\Big)}{(p_{4}+p_{5})^{2}\Big(m^{2}+(p_{5}+p_{4})^{2}\Big)}-\frac{3\mathrm{i}g^{3}\Big(-m^{2}+(p_{4}+p_{5})^{2}\Big)}{m^{2}+(p_{2}+p_{3})^{2}}-\frac{3\mathrm{i}g^{3}\Big(-m^{2}+(p_{4}+p_{5})^{2}\Big)}{(p_{4}+p_{5})^{2}\Big(m^{2}+(p_{4}+p_{5})^{2}\Big)}+(3\leftrightarrow4,5)$$

Summing all the contributions we have,

$$\begin{split} & ig^3 \frac{\left(-m^2 + (p_5 + p_4)^2\right)}{\left(p_5 + p_4\right)^2\left(m^2 + (p_5 + p_4)^2\right)} + (3 \leftrightarrow 4, 5) \\ & - \frac{2ig^3m^2}{\left(p_1 + p_2\right)^2} \frac{1}{\left(m^2 + (p_5 + p_4)^2\right)} + (3 \leftrightarrow 4, 5) \\ & + \frac{ig^3\left(\left(p_1 + p_3\right)^2 - m^2\right)}{m^2 + \left(p_1 + p_3\right)^2} \left[\frac{1}{m^2 + \left(p_2 + p_4\right)^2} + \frac{1}{m^2 + \left(p_2 + p_5\right)^2} + \frac{1}{m^2 + \left(p_4 + p_5\right)^2}\right] + (3 \leftrightarrow 4, 5) \\ & - \frac{ig^3m^2}{\left(p_4 + p_5\right)^2\left(m^2 + \left(p_4 + p_5\right)^2\right)} + (3 \leftrightarrow 4, 5) \\ & + \frac{ig^3\left(m^2 + \left(p_2 + p_3\right)^2 + \left(p_5 + p_4\right)^2\right)\left(-m^2 + \left(p_5 + p_4\right)^2\right)}{\left(m^2 + \left(p_2 + p_3\right)^2\right)\left(p_4 + p_5\right)^2\left(m^2 + \left(p_4 + p_5\right)^2\right)} + (3 \leftrightarrow 4, 5) \\ & - \frac{3ig^3}{m^2 + \left(p_2 + p_3\right)^2} - \frac{3ig^3\left(-m^2 + \left(p_4 + p_5\right)^2\right)}{\left(p_4 + p_5\right)^2\left(m^2 + \left(p_4 + p_5\right)^2\right)} + (3 \leftrightarrow 4, 5) \\ & - \frac{2ig^3m^2}{\left(p_1 + p_2\right)^2} \frac{1}{\left(m^2 + \left(p_5 + p_4\right)^2\right)} + (3 \leftrightarrow 4, 5) \\ & + \frac{ig^3\left(\left(p_1 + p_3\right)^2 - m^2\right)}{m^2 + \left(p_1 + p_3\right)^2} \left[\frac{1}{m^2 + \left(p_2 + p_4\right)^2} + \frac{1}{m^2 + \left(p_2 + p_5\right)^2} + \frac{1}{m^2 + \left(p_4 + p_5\right)^2}\right] + (3 \leftrightarrow 4, 5) \\ & - \frac{ig^3m^2}{\left(p_4 + p_5\right)^2\left(m^2 + \left(p_4 + p_5\right)^2\right)} + (3 \leftrightarrow 4, 5) \end{split}$$

$$\begin{split} &+\frac{ig^3(m^2+(p_2+p_3)^2+(p_5+p_4)^2)(m^2+(p_4+p_5)^2)}{(m^2+(p_2+p_3)^2)(p_4+p_5)^2(m^2+(p_4+p_5)^2)} + (3 \leftrightarrow 4, 5)}{(m^2+(p_2+p_3)^2)(p_4+p_5)^2(m^2+(p_4+p_5)^2)} - \frac{2ig^3m}{(m^2+(p_2+p_5)^2)} + (3 \leftrightarrow 4, 5)} \\ &-\frac{2ig^3m^2}{(p_1+p_2)^2} \frac{1}{(m^2+(p_2+p_4)^2)} + (3 \leftrightarrow 4, 5)}{(m^2+(p_2+p_3)^2)(m^2+(p_2+p_5)^2)} + \frac{1}{m^2+(p_4+p_5)^2}} \\ &-\frac{2ig^3m^2}{m^2+(p_1+p_2)^2} \frac{1}{(m^2+(p_2+p_4)^2)} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_4+p_5)^2} \\ &+ ig^2 \left[\frac{1}{m^2+(p_2+p_4)^2} + \frac{1}{m^2+(p_2+p_5)^2} + \frac{1}{m^2+(p_4+p_5)^2} \right] \\ &+ ig^3 \left[\frac{1}{m^2+(p_2+p_4)^2} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_4+p_5)^2} \right] \\ &+ ig^3 \left[\frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_4+p_5)^2} \right] \\ &+ ig^3 \left[\frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_4+p_5)^2} \right] \\ &+ ig^3 \left[\frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_4+p_5)^2} \right] \\ &+ \frac{ig^3(m^2+(p_2+p_3)^2)(p_4+p_3)^2(m^2+(p_4+p_5)^2)}{(m^2+(p_2+p_3)^2)(p_4+p_3)^2(m^2+(p_4+p_5)^2)} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{m^2+(p_2+p_3)^2} + \frac{1}{(p_4+p_3)^2(m^2+(p_4+p_5)^2)} - \frac{2ig^3}{m^2+(p_2+p_3)^2} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{(m^2+(p_2+p_3)^2)(p_4+p_3)^2(m^2+(p_4+p_3)^2)} + \frac{1}{m^2+(p_2+p_3)^2} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{m^2+(p_2+p_3)^2} + \frac{ig^3m^2}{(p_4+p_2)^2(m^2+(p_4+p_3)^2)} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{m^2+(p_2+p_3)^2(p_4+p_3)^2(m^2+(p_4+p_3)^2)} - \frac{ig^3}{(m^2+(p_2+p_3)^2)} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{m^2+(p_2+p_3)^2(p_4+p_3)^2(m^2+(p_4+p_3)^2)} + \frac{1}{m^2+(p_2+p_3)^2} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{m^2+(p_2+p_3)^2(p_2+p_3)^2(p_2+p_3)^2(m^2+(p_2+p_3)^2)} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{(p_2+p_3)^2(m^2+(p_2+p_3)^2)} + \frac{1}{m^2+(p_2+p_3)^2} + \frac{1}{m^2+(p_2+p_3)^2} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{(p_2+p_3)^2(m^2+(p_2+p_3)^2)(p_4+p_3)^2(m^2+(p_2+p_3)^2)} + (3 \leftrightarrow 4, 5) \\ &-\frac{2ig^3m^2}{(p_2+p_3)^2(m^2+(p_2+p_3)^2)(p_2+p_3)^2(m^2+(p_2+p_3)^2)} + \frac{ig^3}{m^2+(p_2+p_3)^2(m^2+(p_2+p_3)^2)} + (3 \leftrightarrow 4, 5) \\ &+\frac{2ig^3m^2}{(m^2+(p_2+p_3)^2)(p_2+p_3)^2(m^2+(p_2+p_3)^2)} + \frac{ig^3}{m^2+(p_2+p_3)^2} + \frac{ig^3}{m^2+(p_2+p_3)^2} + \frac{ig^3}{m^2+(p_2+p_3)^$$

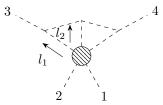
$$+ ig^{3} \frac{-(p_{4} + p_{5})^{2}m^{2} - (p_{4} + p_{5})^{4} + m^{4} + m^{2}(p_{2} + p_{3})^{2} - (p_{4} + p_{5})^{2}m^{2} - (p_{2} + p_{3})^{2}(p_{4} + p_{5})^{2}}{\left(m^{2} + (p_{2} + p_{3})^{2}\right)\left(p_{4} + p_{5}\right)^{2}\left(m^{2} + (p_{4} + p_{5})^{2}\right)} + (3 \leftrightarrow 4, 5)$$

$$- \frac{2ig^{3}m^{2}}{(p_{1} + p_{2})^{2}} \frac{1}{\left(m^{2} + (p_{5} + p_{4})^{2}\right)} + (3 \leftrightarrow 4, 5)$$

$$- \frac{2ig^{3}m^{2}}{m^{2} + (p_{1} + p_{3})^{2}} \left[\frac{1}{m^{2} + (p_{2} + p_{4})^{2}} + \frac{1}{m^{2} + (p_{2} + p_{5})^{2}} + \frac{1}{m^{2} + (p_{4} + p_{5})^{2}}\right] + (3 \leftrightarrow 4, 5)$$

$$- 2ig^{3}m^{2} \frac{1}{\left(m^{2} + (p_{2} + p_{3})^{2}\right)\left(m^{2} + (p_{4} + p_{5})^{2}\right)} + (3 \leftrightarrow 4, 5)$$

Now, let's do the cuts, consider a two loop four point amplitude with five cuts,



$$= -\frac{g^3}{m^4} [\mathcal{A}_5(1, 2, l_{1h}, l_{2h}, l_{3h}) - \mathcal{A}_5(1, 2, l_{1h}, l_{2\eta}, l_{3\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2h}, l_{3\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{3h}) + 2\mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{3\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}, l_{2\eta}) - \mathcal{A}_5(1, 2, l_{1\eta}, l_{2\eta}, l_{2\eta},$$

3. Cut solutions

Of course in each amplitude we have different cut solutions. Now let us solve them,

3.1. all massless. The cut condition is,

$$k_1^2 = k_2^2 = (3 - k_1)^2 = (3 - k_1 - k_2)^2 = (3 + 4 - k_1 - k_2)^2 = 0$$

The first and third condition enforces $k_1 = -|k_1|\langle 3|$. But the fourth and fifth conditions enforces $3 - k_1 - k_2 = n$, with $n \cdot 4 = 0$ & $n^2 = 0$. Lastly, the second condition imposes $(3 - k_1 - n)^2 = -23 \cdot n + 2k_1 \cdot n = 0$, that is,

$$[3n]\langle n3\rangle = [k_1n]\langle n3\rangle$$

which has two solutions, $|n| = |k_1| - |3| \& |n\rangle = z|4\rangle$ or $|n\rangle = |3\rangle \& |n| = z|4|$. When working with scalar particles it's better to choose the first solution, as this avoids singularities in denominators such as $(k_1 \cdot k_2)^{-1}$. Hence, the solution we're going to choose is,

$$\begin{cases} k_1 &= -|k_1|\langle 3| \\ k_2 &= -|3|\langle 3| + |k_1|\langle 3| + z(|k_1| - |3|)\langle 4| \end{cases}$$

3.2. massive legs first topology. Our approach to massive legs is to shift the solution with massless, in order to obtain a well behaved solution in the $m^2 \to 0$ limit. For this topology the cut constrains are,

$$l_1^2 = l_2^2 = (3 - l_1)^2 = -m^2 \& (3 - l_1 - l_2)^2 = (3 + 4 - l_1 - l_2)^2 = 0$$

The ideia here is to define, $l_i = k_i + \alpha_i q_i$ (no sum), with $q_i^2 = 0$ and $\alpha_i = -m^2 (2k_i \cdot q_i)^{-1}$, then, q_i, k_i are not allowed to have any dependence on m^2 . The first and second constrains are already satisfied. The third one gives,

$$-23 \cdot l_1 = 0 \rightarrow 3 \cdot (k_1 + \alpha_1 q_1) = 0 \rightarrow 3 \cdot q_1 = 0$$

As $|q_1\rangle = |3\rangle$ is forbidden, $|q_1| = |3|$. The fourth and fifth constrains imposes,

$$\begin{cases} -n \cdot (\alpha_1 q_1 + \alpha_2 q_2) + \alpha_1 \alpha_2 q_1 \cdot q_2 &= 0\\ 4 \cdot (\alpha_1 q_1 + \alpha_2 q_2) &= 0 \end{cases}$$

This imposes actually $q_1 \cdot q_2 = 0$, for this to be true we have to options, either $|q_2| = |3|$, or $|q_2\rangle = |q_1\rangle$. If we choose the first, we can shift k_1 by 3 such to make $|q_1\rangle = |4\rangle$, this imposes further $|q_2\rangle = |4\rangle$. Hence, a possible solution is,

$$q_1 = q_2 = -|3|\langle 4|$$

3.3. massive legs second topology. The constrains now are slightly different,

$$l_1^2 = (3 - l_1)^2 = 0 \& l_2^2 = (3 - l_1 - l_2)^2 = (3 + 4 - l_1 - l_2)^2 = -m^2$$

Which has as solution $q_2 = -|4|\langle 3|$

3.4. massive legs third topology. Now the constrain is difficult to solve,

$$l_2^2 = 0 \& l_1^2 = (3 - l_1)^2 = (3 - l_1 - l_2)^2 = (3 + 4 - l_1 - l_2)^2 = -m^2$$

The second and third constrains give, $l_1 = -|k_1|\langle 3| - \alpha|3|\langle l_1|$. Now, the fourth and fifth constrains gives,

$$3 - l_1 - l_2 = -z(|k_1| - |3|)\langle 4| + \beta |4|\langle n|$$

With of course $\beta = -\frac{m^2}{z\langle 4n\rangle[4|(|k_1|-|3])}$. At last the second constrain gives,

$$\begin{split} l_2 &= -|3] \langle 3| + |k_1] \langle 3| + \alpha |3] \langle l_1| + z(|k_1] - |3] \rangle \langle 4| - \beta |4] \langle n| \\ l_2^2 &= 0 = -z \langle 34 \rangle [k_1 3] + \beta \langle 3n \rangle [43] + \alpha \langle 3l_1 \rangle [3k_1] - z \langle 34 \rangle [3k_1] - \beta \langle 3n \rangle [4k_1] + \alpha z \langle l_1 4 \rangle [k_1 3] - \alpha \beta \langle l_1 n \rangle [43] - \beta z \langle 4n \rangle [4|(|k_1] - |3]) \\ 0 &= \beta \langle 3n \rangle [43] + m^2 - \beta \langle 3n \rangle [4k_1] + \alpha z \langle l_1 4 \rangle [k_1 3] - \alpha \beta \langle l_1 n \rangle [43] + m^2 \\ -2m^2 &= \beta \langle 3n \rangle [4|(|3] - |k_1]) + \alpha z \langle l_1 4 \rangle [k_1 3] - \alpha \beta \langle l_1 n \rangle [43] \\ -2m^2 &= -\frac{m^2}{z \langle 4n \rangle [4|(|k_1] - |3])} \langle 3n \rangle [4|(|3] - |k_1]) + \frac{m^2}{\langle 3l_1 \rangle [3k_1]} z \langle l_1 4 \rangle [k_1 3] - \alpha \beta \langle l_1 n \rangle [43] \end{split}$$

For the quadratic term in m^2 to vanish is necessary $\langle l_1 n \rangle = 0 \rightarrow |l_1\rangle \propto |n\rangle$, thus,

$$-2m^{2} = \frac{m^{2}}{z\langle 4n\rangle}\langle 3n\rangle + \frac{m^{2}}{\langle 3l_{1}\rangle}z\langle 4l_{1}\rangle$$
$$-2 = \frac{1}{z\langle 4n\rangle}\langle 3n\rangle + \frac{1}{\langle 3n\rangle}z\langle 4n\rangle \to \langle 3n\rangle = -z\langle 4n\rangle$$

The best parametrization is $|n\rangle = |l_1\rangle = |4\rangle - \frac{1}{z}|3\rangle$.

3.5. massive legs fourth topology. Now the constrain is the hardest to solve,

$$l_1^2 = l_2^2 = (3 - l_1)^2 = (3 - l_1 - l_2)^2 = (3 + 4 - l_1 - l_2)^2 = -m^2$$

Happily, most of the work was done already in the last solution, we just have to change:

$$l_{2} = -|3]\langle 3| + |k_{1}]\langle 3| + \alpha|3]\langle l_{1}| + z(|k_{1}] - |3]\rangle\langle 4| - \beta|4]\langle n|$$

$$l_{2}^{2} = -m^{2} \rightarrow m^{2} = -z\langle 34\rangle[k_{1}3] + \beta\langle 3n\rangle[43] + \alpha\langle 3l_{1}\rangle[3k_{1}] - z\langle 34\rangle[3k_{1}] - \beta\langle 3n\rangle[4k_{1}] + \alpha z\langle l_{1}4\rangle[k_{1}3] - \alpha\beta\langle l_{1}n\rangle[43] - \beta z\langle 4n\rangle[4|(|k_{1}] - |3]\rangle$$

$$m^{2} = \beta\langle 3n\rangle[43] + m^{2} - \beta\langle 3n\rangle[4k_{1}] + \alpha z\langle l_{1}4\rangle[k_{1}3] - \alpha\beta\langle l_{1}n\rangle[43] + m^{2}$$

$$-m^{2} = \beta\langle 3n\rangle[4|(|3] - |k_{1}]) + \alpha z\langle l_{1}4\rangle[k_{1}3] - \alpha\beta\langle l_{1}n\rangle[43]$$

$$-m^{2} = -\frac{m^{2}}{z\langle 4n\rangle[4|(|k_{1}] - |3]\rangle}\langle 3n\rangle[4|(|3] - |k_{1}]) + \frac{m^{2}}{\langle 3l_{1}\rangle[3k_{1}]}z\langle l_{1}4\rangle[k_{1}3] - \alpha\beta\langle l_{1}n\rangle[43]$$

Again we fix $|l_1\rangle \propto |n\rangle$. The solution then is given by,

$$-m^{2} = \frac{m^{2}}{z\langle 4n\rangle}\langle 3n\rangle + \frac{m^{2}}{\langle 3l_{1}\rangle}z\langle 4l_{1}\rangle$$
$$-1 = \frac{1}{z\langle 4n\rangle}\langle 3n\rangle + \frac{1}{\langle 3n\rangle}z\langle 4n\rangle \rightarrow \langle 3n\rangle = -z\left(\frac{1}{2} + \frac{\mathrm{i}}{2}\sqrt{3}\right)\langle 4n\rangle = -z\mathrm{e}^{\frac{1}{3}\pi\mathrm{i}}\langle 4n\rangle$$

4. Amplitudes evaluated in the cuts

4.1. all massless. The expression for the amplitude is,

$$\begin{split} &= \frac{\mathrm{i} g^3 m^2}{2 p_1 \cdot p_2} \left[\frac{1}{2 k_1 \cdot k_2} + \frac{1}{2 k_1 \cdot k_3} + \frac{1}{2 k_2 \cdot k_3} \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}{2 p_1 \cdot k_1} \left[\frac{1}{2 p_2 \cdot k_2} + \frac{1}{2 p_2 \cdot k_3} + \frac{1}{2 k_3 \cdot k_2} \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}{2 p_1 \cdot k_2} \left[\frac{1}{2 k_1 \cdot p_2} + \frac{1}{2 k_1 \cdot k_3} + \frac{1}{2 k_3 \cdot p_2} \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}{2 p_1 \cdot k_3} \left[\frac{1}{2 k_1 \cdot k_2} + \frac{1}{2 k_1 \cdot p_2} + \frac{1}{2 p_2 \cdot k_2} \right] \\ &\quad + \frac{\mathrm{i} g^3 m^2}{2 p_2 \cdot k_1 2 k_2 \cdot k_3} + \frac{\mathrm{i} g^3 m^2}{2 p_2 \cdot k_2 2 k_1 \cdot k_3} + \frac{\mathrm{i} g^3 m^2}{2 p_2 \cdot k_3 2 k_2 \cdot k_1} \end{split}$$

 $4.2.1.\ l_1\ and\ l_2\ massive.$

$$\begin{split} &= -\frac{\mathrm{i} g^3 m^2}{-2 m^2 + 2 l_1 \cdot l_2} \left[\frac{1}{2 k_3 \cdot p_1} + \frac{1}{2 p_2 \cdot p_1} + \frac{1}{2 k_3 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 l_1 \cdot p_1)} \left[\frac{1}{2 l_2 \cdot k_3} + \frac{1}{2 l_2 \cdot p_2} + \frac{1}{2 k_3 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 l_1 \cdot p_2)} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 l_2 \cdot k_3} + \frac{1}{2 p_1 \cdot k_3} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 l_1 \cdot k_3)} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 l_2 \cdot p_2} + \frac{1}{2 p_1 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{2 l_2 \cdot k_3 2 p_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_2 \cdot p_1 2 k_3 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_2 \cdot p_2 2 p_1 \cdot k_3} \\ &= - \frac{\mathrm{i} g^3 m^2}{2 k_1 \cdot k_2} \left[\frac{1}{2 k_3 \cdot p_1} + \frac{1}{2 p_2 \cdot p_1} + \frac{1}{2 k_3 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 l_1 \cdot p_1)} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 l_2 \cdot p_2} + \frac{1}{2 l_3 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 l_1 \cdot p_2)} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 k_2 \cdot k_3} + \frac{1}{2 p_1 \cdot k_3} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{(2 k_1 \cdot k_3)} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 l_2 \cdot p_2} + \frac{1}{2 p_1 \cdot p_2} \right] \\ &- \frac{\mathrm{i} g^3 m^2}{2 k_2 \cdot k_3 2 p_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_2 \cdot p_1 2 k_3 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_2 \cdot p_2 2 p_1 \cdot k_3} \end{split}$$

We used that $l_3 = k_3, l_2 \cdot k_3 = k_2 \cdot k_3, l_1 \cdot k_3 = k_1 \cdot k_3, 2l_1 \cdot l_2 = 2k_1 \cdot k_2 + 2m^2$.

4.2.2. l_2 and l_3 massive.

$$= -\frac{\mathrm{i} g^3 m^2}{-2m^2 + 2l_3 \cdot l_2} \left[\frac{1}{2k_1 \cdot p_1} + \frac{1}{2p_2 \cdot p_1} + \frac{1}{2k_1 \cdot p_2} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2l_3 \cdot p_1)} \left[\frac{1}{2l_2 \cdot k_1} + \frac{1}{2l_2 \cdot p_2} + \frac{1}{2k_1 \cdot p_2} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2l_3 \cdot p_2)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2l_2 \cdot k_1} + \frac{1}{2p_1 \cdot k_1} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2l_3 \cdot k_1)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2l_2 \cdot p_2} + \frac{1}{2p_1 \cdot p_2} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{2l_2 \cdot k_1 2p_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2l_2 \cdot p_1 2k_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2l_2 \cdot p_2 2p_1 \cdot k_1}$$

$$= -\frac{\mathrm{i} g^3 m^2}{2k_3 \cdot k_2} \left[\frac{1}{2k_1 \cdot p_1} + \frac{1}{2p_2 \cdot p_1} + \frac{1}{2k_1 \cdot p_2} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2k_3 \cdot p_1)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2l_2 \cdot p_2} + \frac{1}{2k_1 \cdot p_2} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2k_3 \cdot p_2)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2k_2 \cdot k_1} + \frac{1}{2p_1 \cdot k_1} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2k_3 \cdot k_1)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2l_2 \cdot p_2} + \frac{1}{2p_1 \cdot k_1} \right]$$

$$-\frac{\mathrm{i} g^3 m^2}{(2k_3 \cdot k_1)} \left[\frac{1}{2l_2 \cdot p_1} + \frac{1}{2l_2 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2l_2 \cdot p_1 2k_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2l_2 \cdot p_2 2p_1 \cdot k_1} \right]$$

We used that $l_1 = k_1, l_2 \cdot k_1 = k_2 \cdot k_1, l_3 \cdot k_1 = k_3 \cdot k_1, 2l_2 \cdot l_3 = 2k_2 \cdot k_3 + 2m^2$.

4.2.3. l_1 and l_3 massive.

$$=-\frac{\mathrm{i} g^3 m^2}{-2 m^2+2 l_3 \cdot l_1} \left[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 p_2 \cdot p_1} + \frac{1}{2 l_2 \cdot p_2} \right] \\ -\frac{\mathrm{i} g^3 m^2}{(2 l_3 \cdot p_1)} \left[\frac{1}{2 l_2 \cdot l_1} + \frac{1}{2 l_1 \cdot p_2} + \frac{1}{2 l_1 \cdot p_2} \right]$$

$$\begin{split} &-\frac{\mathrm{i} g^3 m^2}{(2 l_3 \cdot p_2)} \bigg[\frac{1}{2 l_2 \cdot p_1} + \frac{1}{2 l_2 \cdot l_1} + \frac{1}{2 p_1 \cdot l_1} \bigg] \\ &-\frac{\mathrm{i} g^3 m^2}{(2 l_3 \cdot l_2)} \bigg[\frac{1}{2 l_1 \cdot p_1} + \frac{1}{2 l_1 \cdot p_2} + \frac{1}{2 p_1 \cdot p_2} \bigg] \\ &-\frac{\mathrm{i} g^3 m^2}{2 l_1 \cdot l_2 2 p_1 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_1 \cdot p_1 2 l_2 \cdot p_2} - \frac{\mathrm{i} g^3 m^2}{2 l_1 \cdot p_2 2 p_1 \cdot l_1} \end{split}$$

4.2.4. l_1, l_2 and l_3 massive.

5.
$$DF^2$$
 Theory

The $(DF)^2$ + YM theory is given by the following Lagrangian,

$$\mathcal{L} = -\frac{1}{2} D_{\mu} F^{a\mu\nu} D_{\alpha} F_{a\ \nu}^{\ \alpha} + \frac{1}{3} f_{abc} F^{a\ \mu}_{\ \mu}{}^{\nu} F^{b\ \alpha}_{\ \nu} F^{c\ \mu}_{\ \alpha} - \frac{1}{2} D_{\mu} \phi^{I} D^{\mu} \phi_{I} + \frac{g}{2} C^{Iab} \phi_{I} F_{a\mu\nu} F_{b}^{\ \mu\nu} + \frac{g}{6} d^{IJK} \phi_{I} \phi_{J} \phi_{K} - \frac{m^{2}}{2} \phi_{I} \phi^{I} - \frac{m^{2}}{4} F_{a\mu\nu} F^{a\mu\nu}$$
 Where of course,

$$F^{a}_{\mu\nu} = \partial_{\mu}A^{a}_{\nu} - \partial_{\nu}A^{a}_{\mu} + gf^{a}_{bc}A^{b}_{\mu}A^{c}_{\nu}$$
$$(D_{\alpha}F^{\mu\nu})^{a} = \partial_{\alpha}F^{a}_{\mu\nu} + gf^{a}_{bc}A^{b}_{\alpha}F^{c}_{\mu\nu}$$
$$(D_{\alpha}\phi)^{I} = \partial_{\alpha}\phi^{I} - igT^{a}_{R}{}^{I}{}_{I}A_{a\alpha}\phi^{J}$$

We also have to incorporate the gauge fixing part,

$$\mathcal{L}_{GF} = -\frac{1}{2}\partial_{\mu}A^{a\mu} \left(-\Box + m^{2}\right)\partial_{\nu}A_{a}^{\nu}$$

Let us collect all the quadratic pieces,

$$\begin{split} -\frac{1}{2}D_{\mu}F^{a\mu\nu}D_{\alpha}F_{a\nu}^{\alpha} &= -\frac{1}{2}\big(\partial_{\mu}F^{a\mu\nu} + gf_{bc}^{}A^{}_{\mu}F^{c\mu\nu}\big)\big(\partial_{\alpha}F_{a\nu}^{\alpha} + gf_{ade}A^{}_{\alpha}F^{e\alpha}_{\nu}\big) \\ &= -\frac{1}{2}(\partial_{\mu}F^{a\mu\nu})(\partial_{\alpha}F_{a\nu}^{\alpha}) \\ &= -\frac{1}{2}\big(\partial_{\mu}\big(\partial^{\mu}A^{a\nu} - \partial^{\nu}A^{a\mu} + gf_{bc}^{}A^{b\mu}A^{c\nu}\big)\big)\big(\partial_{\alpha}\big(\partial^{\alpha}A_{a\nu} - \partial_{\nu}A_{a}^{\alpha} + gf_{ade}A^{}_{\alpha}A^{e}_{\nu}\big)\big) \end{split}$$