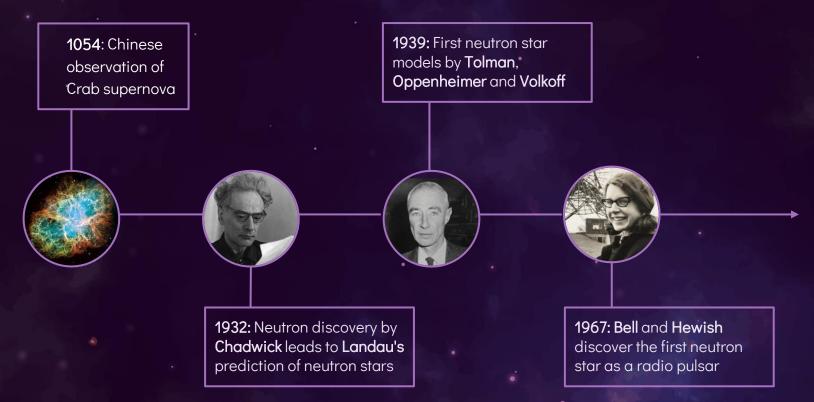


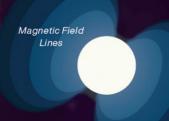
The discovery



Neutron star types

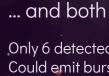
Magnetar

Extremely high magnetic fields: $10^{13} - 10^{15}$ Gauss Long rotation periods: 5 - 12 seconds



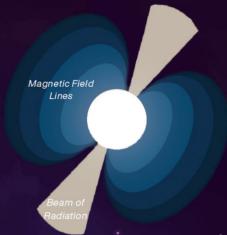
Pulsar

Emit twin beams of radiation from magnetic poles.
Short rotation periods:
down to milliseconds



Only 6 detected so far. Could emit bursts at irregular intervals





15 – 30 km

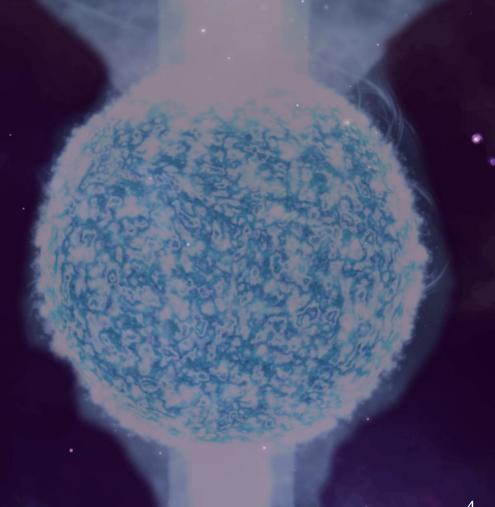
Diameter

 $\geq 10 M_{\odot}$

Original star mass

1-2 M_O

Total mass



Non-relativistic equilibrium equations

Gravitational force for an element of volume:

$$F_G = -\frac{Gm(r)}{r^2} \, \rho(r)$$

- $G = 6.67 \times 10^{-11} \text{m}^3 / \text{kgs}^2$ gravitational constant
- ρ = mass density
- $m(r) = 4\pi \int_0^r \rho(r') r'^2 dr'$

Non-relativistic equilibrium equations

The structure equations are

$$\frac{dm}{dr} = 4\pi r^2 \rho(r) = \frac{4\pi r^2 \epsilon(r)}{c^2}$$

$$\frac{dP}{dr} = -\frac{Gm(r)}{r^2}\rho(r) = -\frac{Gm(r)}{r^2}\frac{\epsilon(r)}{c^2}$$

- $\epsilon = \rho c^2$ energy density
- Valid when $\frac{V(r)}{mc^2} = \frac{Gm}{rc^2} \ll 1$



Tolman – Oppenheimer – Volkoff equations

$$\frac{dm}{dr} = 4\pi r^2 \rho(r)$$

$$\frac{dP}{dr} = -\frac{Gm(r)}{r^2} \rho(r) \frac{\left[1 + \frac{P(r)}{\rho(r)c^2}\right] \left[1 + \frac{4\pi r^3 P(r)}{m(r)c^2}\right]}{\left[1 - \frac{2Gm(r)}{rc^2}\right]}$$

- Account for special and general relativity corrections
- Relativity <u>enhances</u> gravity!

Solving the equations



$$r=0 \to m(0)=0$$

Final conditions

$$r = R \to P(R) = 0$$
$$\to M = m(R)$$

...What about ρ ?

Before that, white dwarfs

- Low- or medium-mass star at the end of its lifetime, $R \approx 10^4 \text{km}$
- Not enough thermal pressure → electrons pushed closer together
- Electrons in the lowest energy levels
- Pauli exclusion principle → star is stabilised against gravity

Fermi gas model for electrons

Number of states dn available at momentum k per unit volume

$$dn = \frac{d^3k}{(2\pi\hbar)^3} = \frac{4\pi k^2 dk}{(2\pi\hbar)^3}$$

Integrating

$$n = \frac{8\pi}{(2\pi\hbar)^3} \int_0^{k_F} k^2 dk = \frac{k_F^3}{3\pi^2\hbar^3}$$

Where $k_F c$ is the Fermi energy level, i.e., maximum electron energy

Fermi gas model for electrons

The electron energy density ϵ is found similarly

$$\epsilon_{el}(k_F) = \frac{8\pi}{(2\pi\hbar)^3} \int_0^{k_F} \sqrt{k^2 c^2 + m_e^2 c^4} \, k^2 dk$$

The total energy density includes the rest mass density (primarily given by nucleons)

$$\epsilon = nm_N \frac{A}{Z} + \epsilon_{el}(k_F)$$

Fermi gas model for electrons

From thermodynamics

$$P = -\left[\frac{\partial U}{\partial V}\right]_{T=0} = n^2 \frac{d\left(\frac{\epsilon}{n}\right)}{dn} = n \frac{d\epsilon}{dn} - \epsilon$$

From the calculations we get:

Relativistic case $(k_F \gg m_e)$:

$$P(k_F) \approx K_{rel} \epsilon^{\frac{4}{3}}$$

Non-relativistic case $(k_F \ll m_e)$:

$$P(k_F) \approx K_{nonrel} \epsilon^{\frac{5}{3}}$$

Pure neutron star, Fermi gas EoS

Same as for white dwarfs, but with m_n instead of m_e

Relativistic case $(k_F \gg m_n)$:

$$P(k_F) \approx K_{rel}\epsilon$$

Non-relativistic case $(k_F \ll m_n)$:

$$P(k_F) \approx K_{nonrel} \epsilon^{\frac{5}{3}}$$

For arbitrary relativity, we can fit the energy density with

$$\epsilon(P) = A_{NR}P^{\frac{3}{5}} + A_RP$$

Numerical solutions

Variables rescaling

Introducing:

- $M_s = 1.98892 \times 10^{30} \text{ kg solar mass}$
- $R_S = \frac{2GM_S}{c^2}$ sun Schwarzschild radius
- $\bullet \quad \gamma = \frac{M_S c^2}{R_S^3}$

The quantities present in the equations can be rescaled as:

$$M = \overline{M}M_{S}, \qquad P = \overline{P}\gamma, \qquad r = \overline{r}R_{S}, \qquad \rho = \frac{\overline{\rho}\gamma}{c^{2}}$$

where \overline{M} , \overline{P} , \overline{r} and $\overline{\rho}$ are dimensionless variables

Numerical solutions

Equations to be solved

Non relativistic equations

$$\frac{d\overline{m}}{d\overline{r}} = 4\pi \overline{r}^2 \overline{\rho}(\overline{r})$$

$$\frac{d\overline{P}}{d\overline{r}} = -\frac{\overline{m}(\overline{r})}{2\overline{r}^2}\overline{\rho}(\overline{r})$$

Relativistic equations

$$\frac{d\overline{m}}{d\overline{r}} = 4\pi \overline{r}^2 \overline{\rho}(\overline{r})$$

$$\frac{\mathrm{d}\overline{P}}{\mathrm{d}\overline{r}} = -\frac{\overline{m}(\overline{r})}{2}\overline{\rho}(\overline{r})\left[1 + \frac{\overline{P}(\overline{r})}{\overline{\rho}(\overline{r})}\right]\left[1 + \frac{4\pi\overline{r}^{3}\overline{P}(\overline{r})}{\overline{m}(\overline{r})}\right]\left[\frac{1}{\overline{r}^{2} - \overline{m}(\overline{r})\overline{r}}\right]$$

Pure neutron star, Fermi gas EoS

Using the phenomenological EoS

$$\bar{\rho}(\bar{P}) = C_{NR}\bar{P}^{\frac{3}{5}} + C_R\bar{P}$$

The values of C_{NR} and C_R which have been used are the ones suggested in [1]:

$$\bar{\rho}(\bar{P}) = 0.871\bar{P}^{\frac{3}{5}} + 2.867\bar{P}$$

Fourth order Runge-Kutta method

• Given an ordinary differential equation with initial conditions

$$\frac{dy(t)}{dt} = f(t, y), \qquad y(t_0) = y_0$$

- General predictor-corrector method:
 - Compute slope at $t_i \rightarrow f(t_i, y_i)$
 - Prediction $y_{i+1} \approx y(t_i) + hf(t_i, y_i)$
 - Compute slope at $t_{i+1} \rightarrow f(t_{i+1}, y_{i+1})$
 - Corrected prediction $y_{i+1} \approx y(t_i) + h/2(f(t_i, y_i) + f(t_{i+1}, y_{i+1}))$
 - **...**

Fourth order Runge-Kutta method

Given an ordinary differential equation with initial conditions

$$\frac{dy(t)}{dt} = f(t, y), \qquad y(t_0) = y_0$$

•
$$y_{i+1} = y_i + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$
, where

$$k_1 = f(y_i, t_i)$$

•
$$k_1 = f(y_i, t_i)$$

• $k_2 = f(y_i + \frac{k_1}{2}, t_i + \frac{h}{2})$

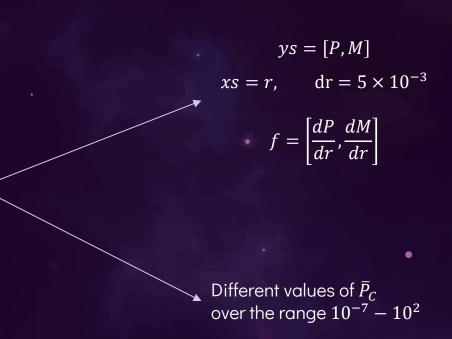
•
$$k_3 = f\left(y_i + \frac{k_2}{2}, t_i + \frac{h}{2}\right)$$

$$k_4 = f(y_i + k_3, t_i + h)$$

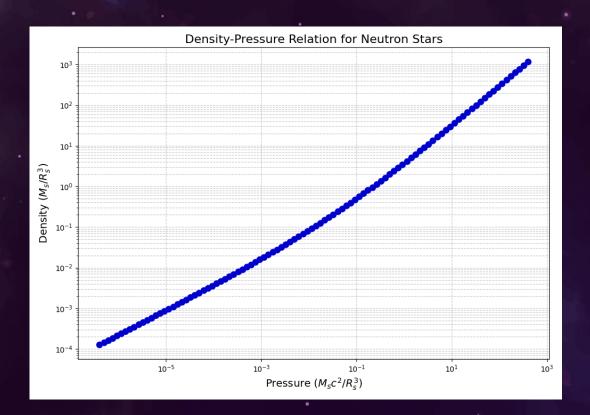
Error of the order $O(h^4)$

RK4 in TOV equations

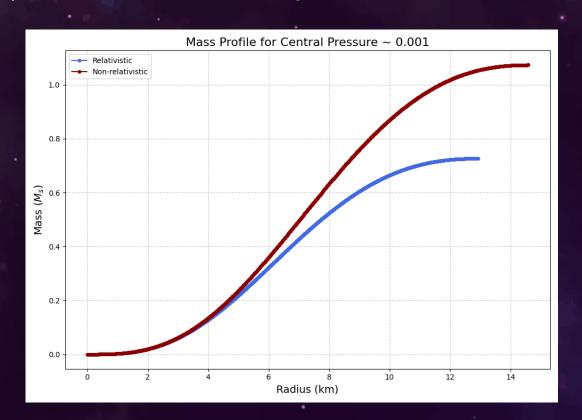
```
def rk4_gen(f, r0, y0, step, max_step):
  xs = r0 + np.arange(max_step) * step
  ys = []
  yvals = y0.copy()
  for x in xs:
    if yvals[0] < 0: # Stop when pressure crosses 0
    break
    ys.append(yvals.copy())
    k0 = step * f(x, yvals)
    k1 = step * f(x + step / 2, yvals + k0 / 2)
    k2 = step * f(x + step / 2, yvals + k1 / 2)
    k3 = step * f(x + step, yvals + k2)
    yvals += (k0 + 2 * k1 + 2 * k2 + k3) / 6
  return np.array(xs[:len(ys)]), np.array(ys)
```



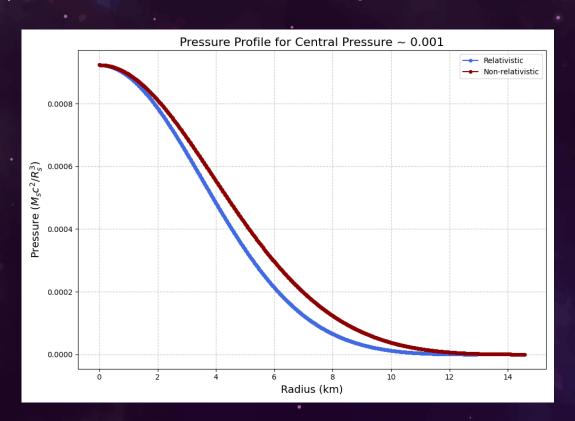
Results: EOS



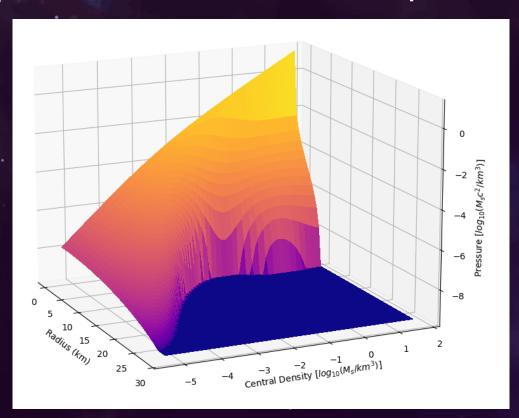
Results: Mass profile



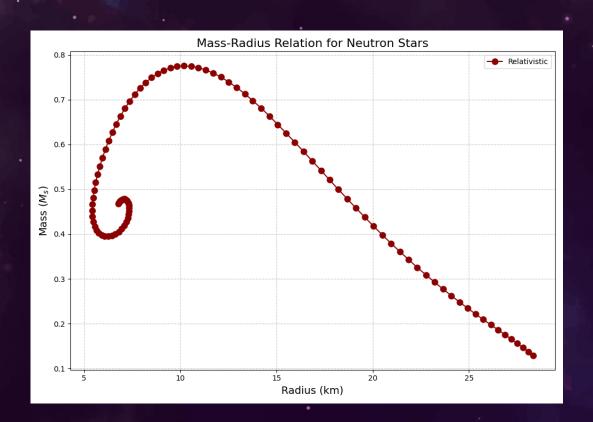
Results: Pressure profile



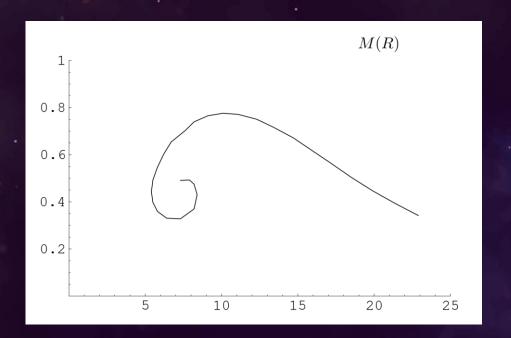
Results: Pressure 3D profile



Results: Mass-Radius relation

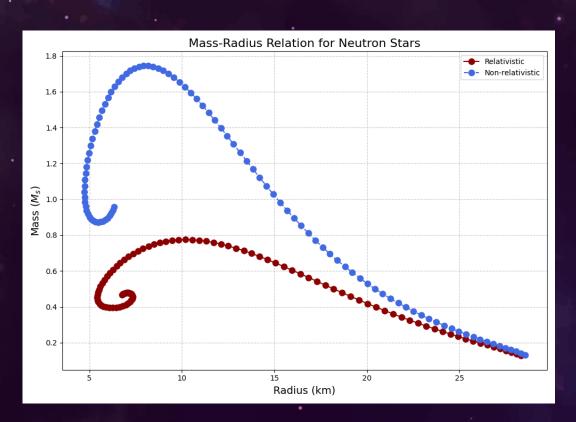


Results: Mass-Radius relation



Richard R. Silbar and Sanjay Reddy. Neutron stars for undergraduates. American Journal of Physics, 72(7):892-905, 2004. doi: 10.1119/1.1703544. URL https://arxiv.org/abs/nucl-th/0309041

Results: Mass-Radius relation, comparison



Execution time test

Time perfromances evaluated for both RK4 method and solve_ivp Python function

For $dr = 5 \times 10^{-3}$ average time:

- Custom RK4: 170 ms
- solve_ivp: 138 ms, but worse results

For $dr = 1 \times 10^{-2}$ average time:

- Custom RK4: 138 ms
- solve_ivp: 150 ms

Code structure



- <u>constants.py</u>: Contains all the physical constants and all the parameters
 of the simulation.
- <u>eos.py</u>: Contains the equation of state employed in the project
- <u>tov_solver.py</u>: Contains the equations to be solved and the fourth-order Runge-Kutta method.
- tov_calculations.py: Contains the function that manages the computation of the solutions.
- <u>utils.py</u>: Contains utility functions for operations across the project.
- <u>time_performance.py</u>: Contains a function to estimate execution time of the RK4 method
- <u>outputs/</u>: Directory where result plots are saved.

Thank you!

Nicholas Pieretti

Theoretical and Numerical aspects of Nuclear Physics