

ETHz AFEA CHEAT SHEET

FEM & ELEMENT

$$Ku = f$$

K : stiffness matrix

u : displacement

f : external force

- Stiffness Matrix

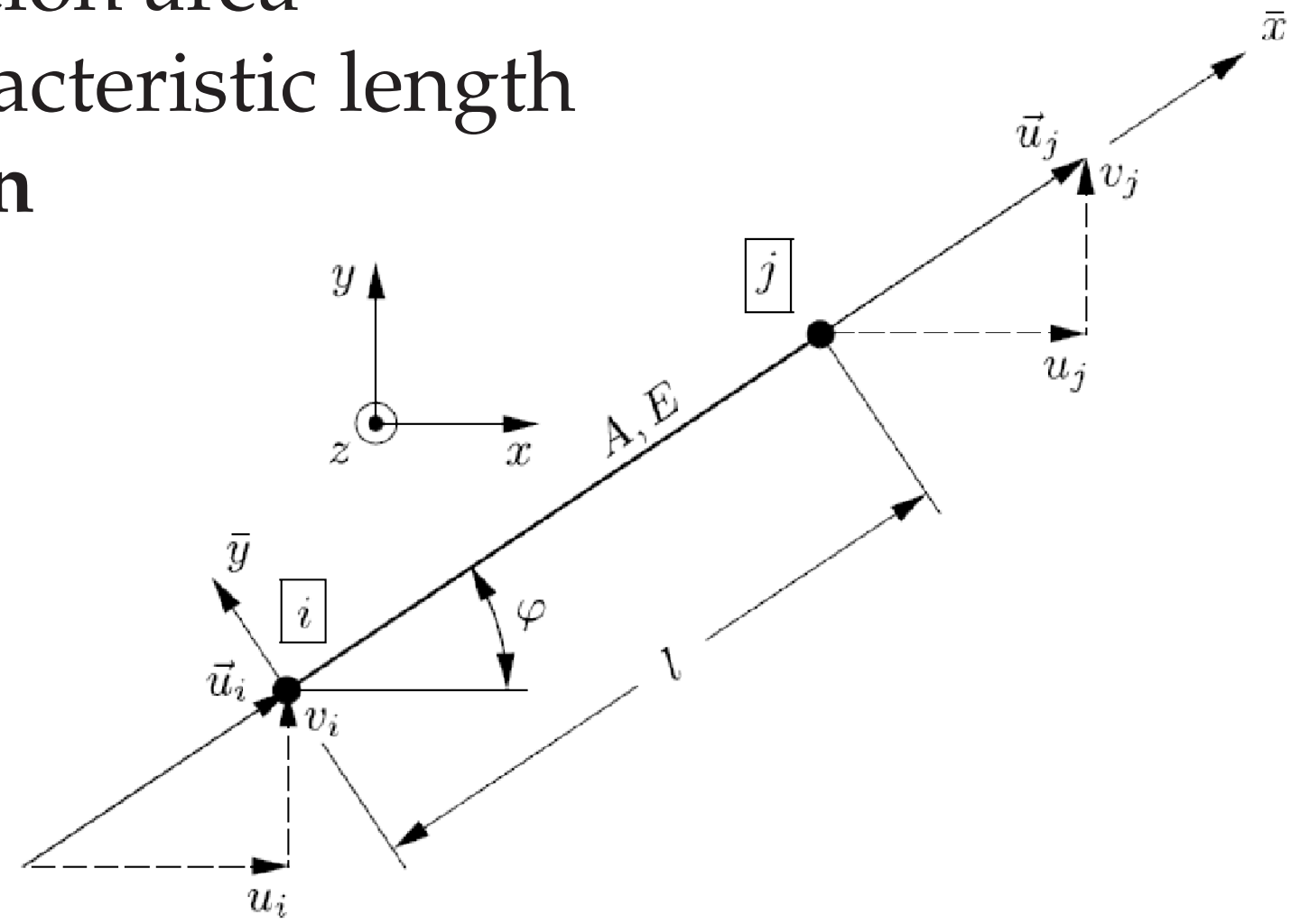
$$K = \frac{EA}{l} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

E : Young's Module

A : section area

l : characteristic length

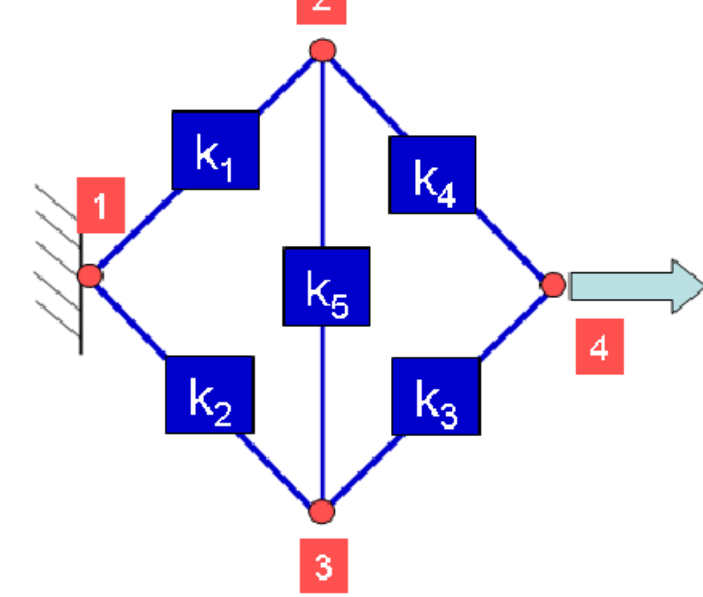
rotation



$$\hat{K} = \begin{bmatrix} \cos\theta & \sin\theta & 0 & 0 \\ 0 & 0 & \cos\theta & \sin\theta \end{bmatrix}^T K \begin{bmatrix} \cos\theta & \sin\theta & 0 & 0 \\ 0 & 0 & \cos\theta & \sin\theta \end{bmatrix}$$

$$= \frac{EA}{l} \begin{bmatrix} \cos^2 & \sin\cos & -\cos^2 & -\sin\cos \\ \sin\cos & \sin^2 & -\sin\cos & -\sin^2 \\ -\cos^2 & -\sin\cos & \cos^2 & \sin\cos \\ -\sin\cos & \sin^2 & \sin\cos & \sin^2 \end{bmatrix}$$

k_{11}	k_{12}	k_{13}	k_{14}	k_{15}	k_{16}	0	0	$\hat{U}_x^{(1)}$	$\hat{U}_y^{(1)}$
k_{21}	k_{22}	k_{23}	k_{24}	k_{25}	k_{26}	0	0	$\hat{U}_x^{(2)}$	$\hat{U}_y^{(2)}$
k_{31}	k_{32}	k_{33}	k_{34}	k_{35}	k_{36}	k_{37}	k_{38}	$\hat{U}_x^{(3)}$	$\hat{U}_y^{(3)}$
k_{41}	k_{42}	k_{43}	k_{44}	k_{45}	k_{46}	k_{47}	k_{48}	$\hat{U}_x^{(4)}$	$\hat{U}_y^{(4)}$
k_{51}	k_{52}	k_{53}	k_{54}	k_{55}	k_{56}	k_{57}	k_{58}	$\hat{U}_x^{(5)}$	$\hat{U}_y^{(5)}$
k_{61}	k_{62}	k_{63}	k_{64}	k_{65}	k_{66}	k_{67}	k_{68}	$\hat{U}_x^{(6)}$	$\hat{U}_y^{(6)}$
0	0	k_{73}	k_{73}	k_{73}	k_{73}	k_{73}	k_{73}	$\hat{U}_x^{(7)}$	$\hat{U}_y^{(7)}$
0	0	k_{83}	k_{84}	k_{85}	k_{86}	k_{87}	k_{88}	$\hat{U}_x^{(8)}$	$\hat{U}_y^{(8)}$



- Continuum Elements

Purpose: Given the position of a point r_j in local coordinate, compute the strain tensor $\varepsilon_{ii}, \varepsilon_{i_1 i_2}$ in global coordinate.

$$J_{ji} = \frac{\partial x_i(r_j)}{\partial r_j} = \frac{\partial \sum_k h^k(r_j) \hat{x}_i^k}{\partial r_j}$$

Interpolate(B matrix)

$$\varepsilon_{ii} = \sum_k \frac{\partial h^k(r_j)}{\partial r_j} J_{ji}^{-1} \hat{u}_i^k$$

$$2\varepsilon_{i_1 i_2} = \sum_k \frac{\partial h^k(r_j)}{\partial r_j} J_{ji_1}^{-1} \hat{u}_{i_1}^k + \frac{\partial h^k(r_j)}{\partial r_j} J_{ji_2}^{-1} \hat{u}_{i_2}^k$$

\hat{u}_i^k : nodal k displacement in global coordinate

\hat{x}_i^k : nodal k position in global coordinate

x_i : random point position in global coordinate

r_j : random point position in local coordinate

$h^k(r_j)$: nodal k interpolate weight $\sum_k h^k(r_j) = 1$

ε_{ii} : diagonal of strain tensor ε

$\varepsilon_{i_1 i_2}$: triangle part of strain tensor ε

Element	Triangular	Tetrahedral	Hexahedral (Brick)
Dimension	2D	3D	3D
Node	3	4	6

ELEMENT

- Continuum Elements

Continuity

- C-0: element boundaries displacements
- C-1: Continuous 1st derivatives

Failed

If the element is distorted or folded there will be no 1-to-1 relationship between natural and global coordinates and the Jacobian will be singular

Quadrilateral

$$\hat{u} = [\hat{u}_x^1 \hat{u}_y^1 \hat{u}_x^2 \hat{u}_y^2 \hat{u}_x^3 \hat{u}_y^3 \hat{u}_x^4 \hat{u}_y^4]^T$$

$$h^1 = 1/4(1 - r_1)(1 - r_2) \quad h^2 = 1/4(1 + r_1)(1 - r_2)$$

$$h^3 = 1/4(1 + r_1)(1 + r_2) \quad h^4 = 1/4(1 - r_1)(1 + r_2)$$

$$\frac{\partial h^k}{\partial x_i} = \frac{\partial h^k}{\partial r_j} J_{ji}^{-1}$$

$$B = \begin{bmatrix} \frac{\partial h^1}{\partial x_1} & 0 & \frac{\partial h^2}{\partial x_1} & 0 & \frac{\partial h^3}{\partial x_1} & 0 & \frac{\partial h^4}{\partial x_1} & 0 \\ 0 & \frac{\partial h^1}{\partial x_2} & 0 & \frac{\partial h^2}{\partial x_2} & 0 & \frac{\partial h^3}{\partial x_2} & 0 & \frac{\partial h^4}{\partial x_2} \\ \frac{\partial h^1}{\partial x_1} & \frac{\partial h^1}{\partial x_2} & \frac{\partial h^2}{\partial x_1} & \frac{\partial h^2}{\partial x_2} & \frac{\partial h^3}{\partial x_1} & \frac{\partial h^3}{\partial x_2} & \frac{\partial h^4}{\partial x_1} & \frac{\partial h^4}{\partial x_2} \end{bmatrix}$$

$$\begin{bmatrix} \varepsilon_{11} & \varepsilon_{22} & \varepsilon_{12} + \varepsilon_{21} \end{bmatrix}^T = B \hat{u}$$

- Structural Elements

Truss: Stresses only in the axes direction No bending resistance

Beam: With bending resistance

Plate & Shell

- thickness is 0
- perpendicular to the mid-surface

Kirchhof Shell

- section perpendicular to the mid-surface
- infinite shear-stiffness: for thin shell it's allowed, for thick shell, deviations to the real behavior occur

- C1 continuity required
- Displacements and rotations are coupled

Mindlin Shell

- section may rotate, but still straight
- could be used for thick structures
- strong bending: Compression-strains go to infinity(aspect-ratio matters)
- C0 continuity required
- Displacements and rotations are independent of each other

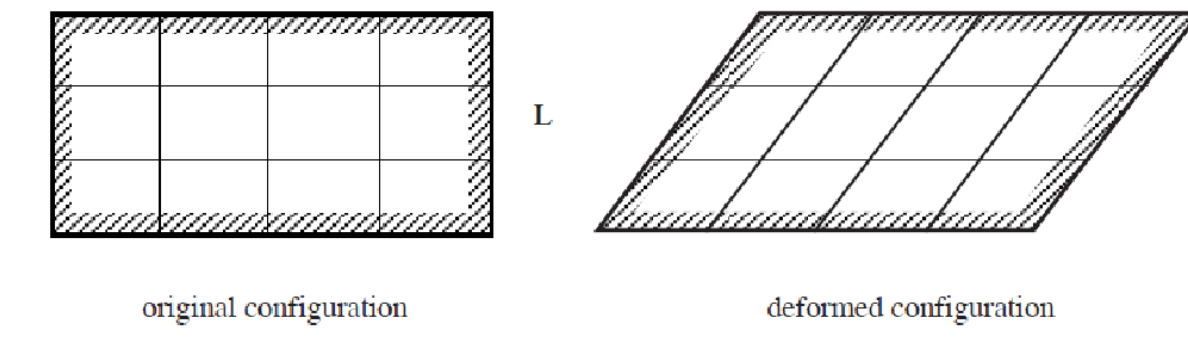
- Numerical Issues

Volumetric Locking: high order elements + reduced integration(reduce computation)

Shear Locking: Reduced integration avoid shear-locking effects, may cause Hourglassing, controlled by adding artificial shear stiffness

Hourglass Control: avoid unphysical deformation of the structure (avoiding zero energy modes)

CONTINUUM MECHANICS



lagrangian(T) and eulerian(B) configuration

$$x = X + u \quad dx = F dX$$

X : is the origin position

x : transform position

u : displacement

F : deformation gradient (not symmetric), $F = RU = VR$

Right Cauchy Green Tensor: $C = F^T F = U^T U$

Left Cauchy Green Tensor: $B = F F^T = V V^T$

Engineering stress: $\sigma = \frac{F_t}{A_0}$

Cauchy stress: $\sigma = \frac{F_t}{A_t}$

Green Langrange Strain Tensor

$$E = \frac{1}{2}(F^T F - I)$$

$$E_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i} + u_{k,i} u_{k,j})$$

first Piola-Kirchhoff stress

$$P = J \sigma F^{-T} \quad J = \det(F)$$

- not symmetric
- seldom used in constitutive models
- axial stress same as Engineering stress

second Piola-Kirchhoff stress

$$S = J F^{-1} \sigma F^{-T}$$

- symmetric
- frequently used in constitutive models
- not a clear physical meaning
- Corotational stress(objective)
- energy conjugate to Green Lagrangian strain tensor(E)

Spin tensor

L : velocity gradient, $L = \frac{dv}{dx} = \dot{F} F^{-1}$

W : asymmetric part of L , $W = \frac{1}{2}(L - L^T)$, $W' = Q W Q^T + \dot{Q} Q^{-1}$

objective: $\sigma' = Q \sigma Q^T$, Q is orthogonal

Jaumann stress rate

$$\dot{\sigma} = \dot{\sigma} + \sigma W - W \sigma$$

take $\dot{\sigma}'$ and use $\dot{Q} = W' Q - Q W$ and $\dot{Q}^T = -Q^T W' + Q W$

Truesdell stress rate

$$\dot{\sigma} = \dot{\sigma} - L \sigma - \sigma L^T + \text{tr}(L) \sigma$$

ELASTO-PLASTIC

$$d\varepsilon = d\varepsilon^e + d\varepsilon^p \quad d\sigma = \mathbb{C}^e d\varepsilon^e$$

ε : strain tensor, composed of elastic strain ε^e and plastic strain ε^p

\mathbb{C}^e : 4-th order elasticity tensor

in uniaxial case, $\bar{\sigma}$ is equivalent stress

$\phi = \bar{\sigma}(\sigma) - \sigma_y < 0$ elastic

$\phi = \bar{\sigma}(\sigma) - \sigma_y = 0$ plastic

Deviatoric Stress: $s = \sigma - \frac{1}{3} \text{tr}(\sigma) I$

$$J_1 = \text{tr}(s) = 0 \quad J_2 = \frac{1}{2} s : s \quad J_3 = \det(s)$$

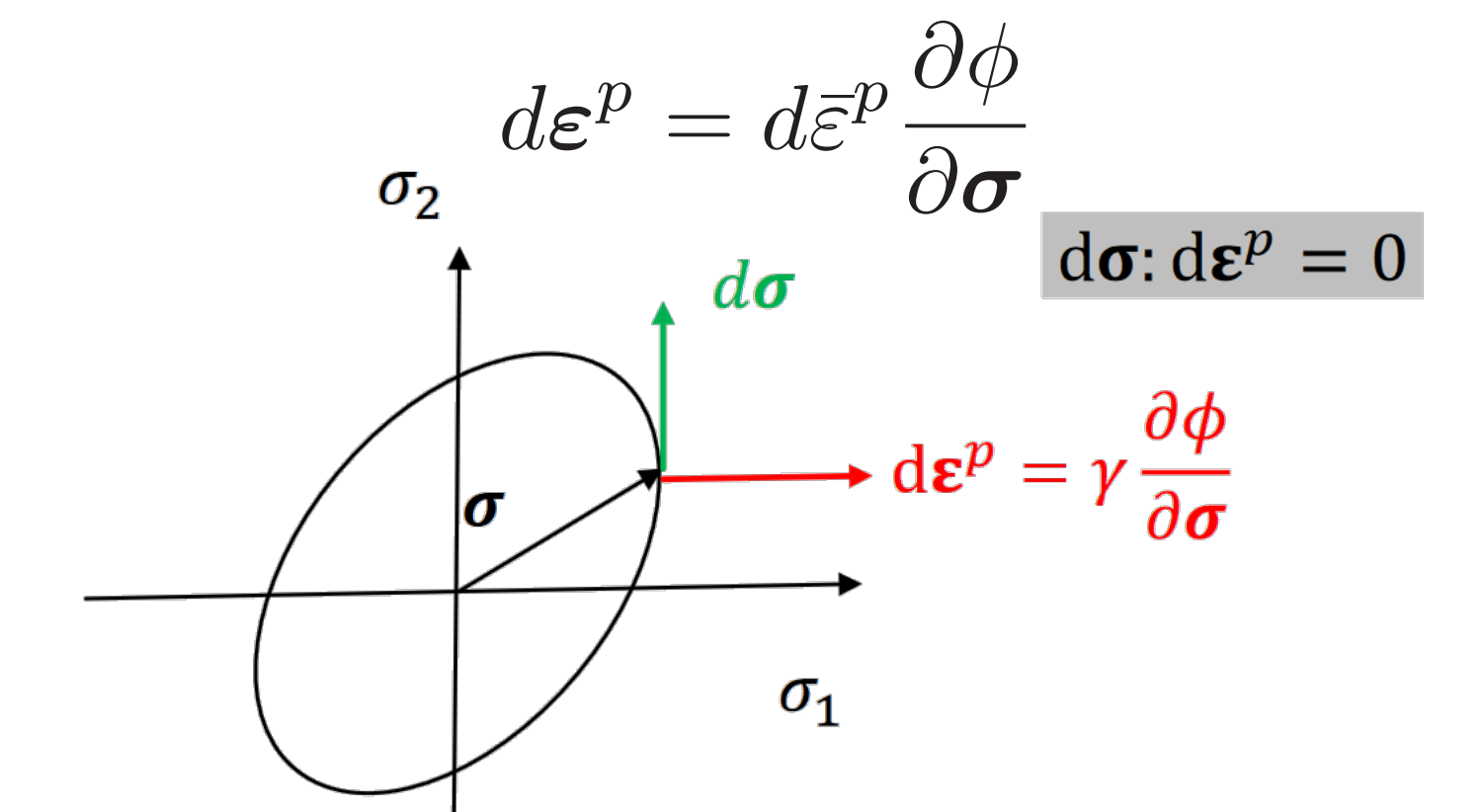
Von Mises stress

$$\bar{\sigma} = \sqrt{3J_2} = \sqrt{\frac{1}{2} \sum_{i \neq j} [(\sigma_{ii} - \sigma_{jj})^2 + 6\sigma_{ij}^2]}$$

stress tensor only depend on the elastic strain
Equivalent Plastic Strain

$$d\bar{\varepsilon}^p = \sqrt{\frac{2}{3}} d\varepsilon^p : d\varepsilon^p$$

Normality Rule



plastic strain tensor is proportional to the normal of the yield surface

Hardening

$$\phi = \bar{\sigma}(\sigma) - \sigma_y(\bar{\varepsilon}^p) = 0$$

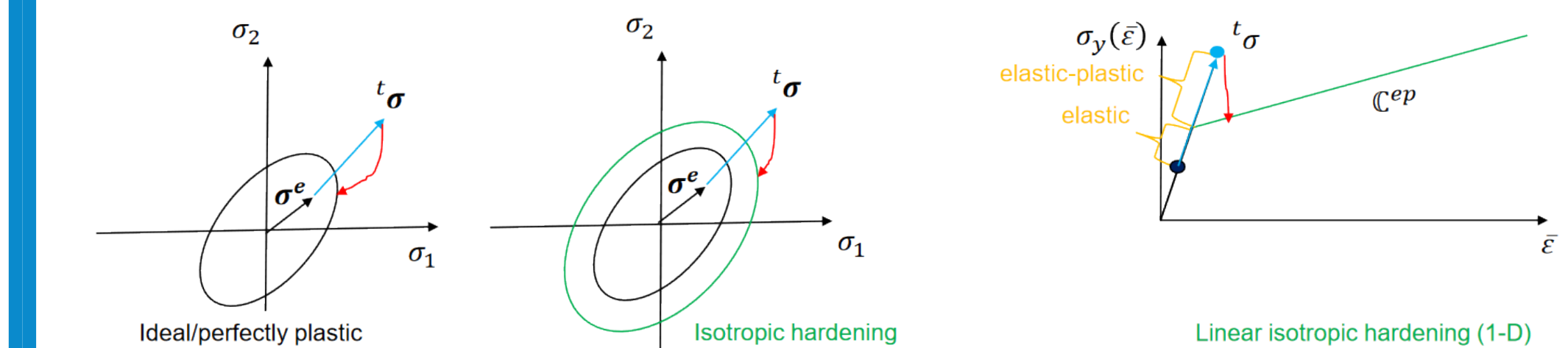
Consistency Condition

$$\frac{\partial \bar{\sigma}}{\partial \sigma} d\sigma - \frac{\partial \sigma_y}{\partial \bar{\varepsilon}^p} d\bar{\varepsilon}^p = 0$$

elastic-plastic masterial modulus

$$\mathbb{C}^{ep} = \mathbb{C}^e \left(1 - \frac{\frac{\partial \bar{\sigma}}{\partial \sigma} : \mathbb{C}^e : \frac{\partial \bar{\sigma}}{\partial \sigma}}{\frac{\partial \sigma_y}{\partial \bar{\varepsilon}^p} + \frac{\partial \bar{\sigma}}{\partial \sigma} : \mathbb{C}^e : \frac{\partial \bar{\sigma}}{\partial \sigma}} \right)$$

return mapping algorithm



ETHZ AFEA CHEAT SHEET

HYPER-ELASTIC

incompressible material

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}' + \frac{1}{3}\boldsymbol{\varepsilon}_V \mathbf{I}$$

$$\boldsymbol{\sigma} = \kappa \boldsymbol{\varepsilon}_V \mathbf{I} + 2G\boldsymbol{\varepsilon}' = -p\mathbf{I} + 2G\boldsymbol{\varepsilon}'$$

$\boldsymbol{\varepsilon}'$: deviatoric strain tensor

$\boldsymbol{\varepsilon}_V$: volumetric strain

κ : Bulk modulus, $\kappa = \frac{E}{3(1-2\nu)}$

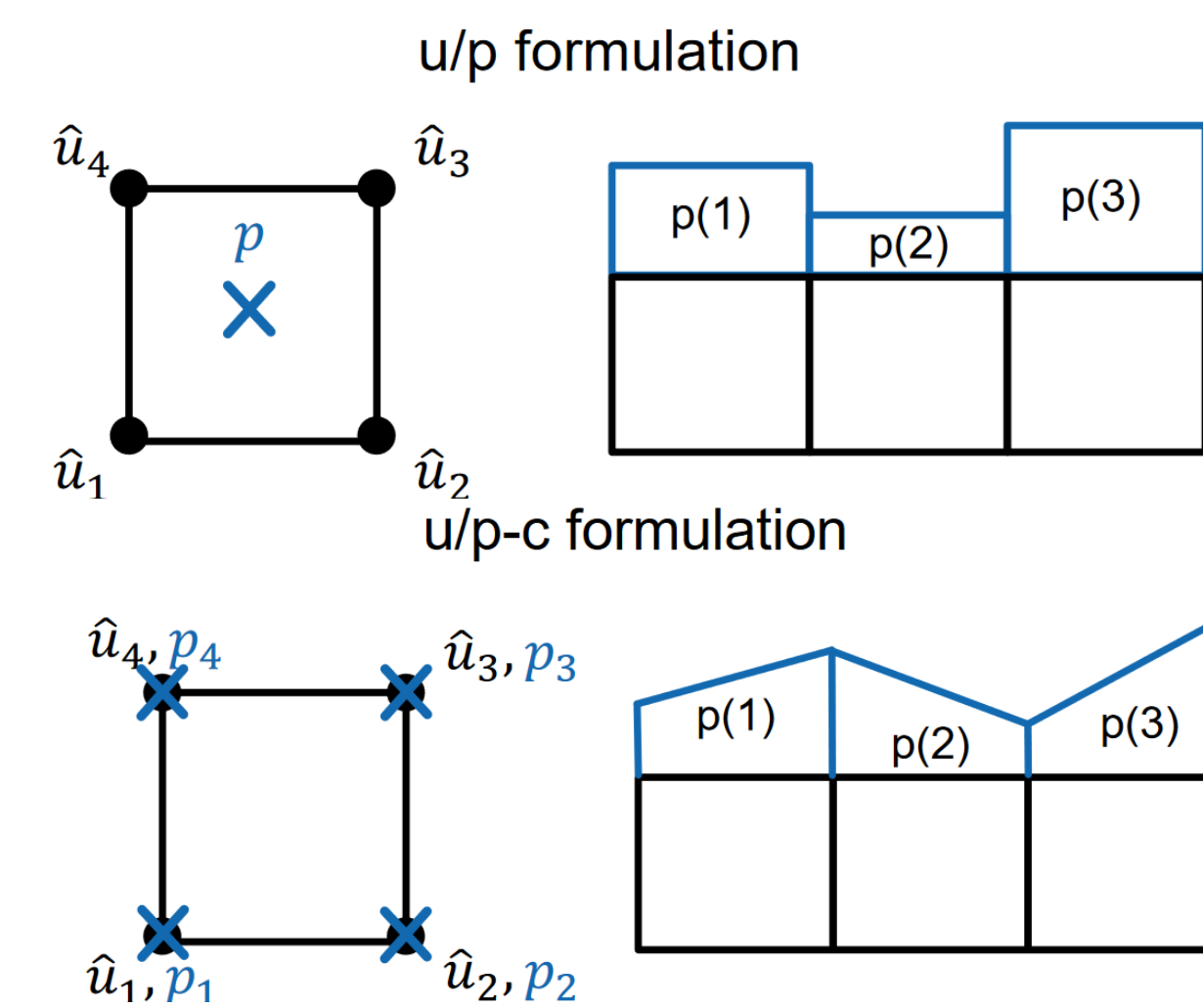
G : Shear modulus, $G = \frac{E}{1+\nu}$

p : pressure, $p = -\frac{1}{3}\text{Tr}(\boldsymbol{\sigma}) = -\kappa \boldsymbol{\varepsilon}_V$

$$\begin{bmatrix} K_{uu} & K_{up} \\ K_{pu} & K_{pp} \end{bmatrix} \begin{bmatrix} \hat{u} \\ \hat{p} \end{bmatrix} = \begin{bmatrix} f^{ext} \\ 0 \end{bmatrix}$$

if material total incompressible $\kappa \rightarrow \infty$, $K_{pp} \rightarrow 0$

Mixed Formulations



Right Cauchy stress tensor: $\mathbf{C} = \mathbf{F}^\top \mathbf{F} = \text{diag}(\lambda_i^2)$

$$I_1 = \text{Tr}(\mathbf{C}) = \sum_i \lambda_i^2$$

$$I_2 = \frac{1}{2}(\text{Tr}(\mathbf{C})^2 - \text{Tr}(\mathbf{C}^2))$$

$$I_3 = \det(\mathbf{C})$$

Neo-Hook Model

$$W = \frac{1}{2}G(I_1 - 3)$$

- tensile test same as first Piola-Kirchhoff stress
- accurate only for small strain
- only one parameter G

Mooney Model

$$W = C_1(I_1 - 3) + C_2 \left(\frac{1}{\lambda_1^2} + \frac{1}{\lambda_2^2} + \frac{1}{\lambda_3^2} - 3 \right)$$

assumption

- incompressible and isotropic
- obeys Hooke's law under simple shear

Other Model

Mooney Rivlin Model: $W = C_1(I_1 - 3) + C_2(I_2 - 3)$

Ogden Model: $W = \sum_{k=1}^N \frac{\mu_k}{\alpha_k} (\lambda_1^{\alpha_k} + \lambda_2^{\alpha_k} + \lambda_3^{\alpha_k} - 3)$

IMPLICIT QUASI-STATIC FEM

$$K(u)\Delta u = \Delta F$$

Full Newton-Raphson

$$K^i \Delta u^{i+1} = \Delta f^i$$

quadratic convergence rate.

Modified Newton-Raphson

$$K^0 \Delta u^{i+1} = \Delta f^i$$

Quasi Newton-Methods

secant approximation

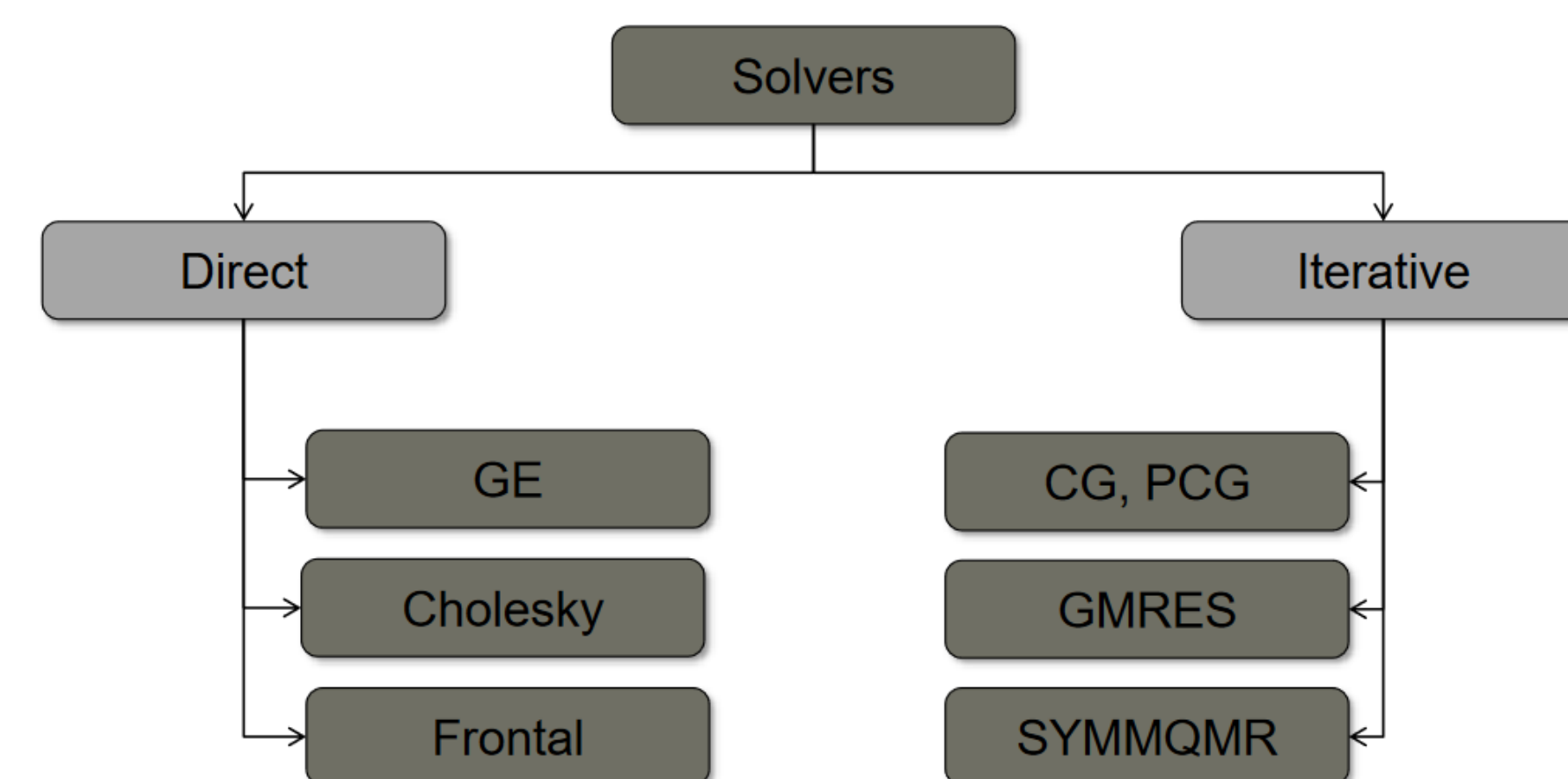
BFGS method

Convergence Criteria

convergence check for each iteration

- Displacement: $\frac{\|\Delta u^{i+1}\|_2}{\|u^i + \Delta u^i\|_2} \leq \text{TOL}_{dsp}$
- Residual: $\frac{\|u^{i+\Delta t} f_{ext} - u^i f_{int}\|_2}{\|u^{i+\Delta t} f_{ext} - u^i f_{int}\|_2} \leq \text{TOL}_{res}$

Linear System Solver



CG solver should be definite

Direct Solver

based on Gauss – Jordan elimination algorithm

Advantage

- predicable number of operations and acc
- fast for small to midsize systems

Disadvantage

- large storage requirements
- slow performance for large system
- hard to parallel

Iterative Solver

Advantage

- smaller storage
- few operations
- found up to desired accuracy
- easy to parallel

Disadvantage

- Convergence not guaranteed
- badly conditioned system slow convergence

Sparse Storage

- Banded Storage: along diagonal
- Skyline Storage: along column near diagonal
- CSR and CSC (no 0 value)

DYNAMIC FEM

time dependent problem (time integration)

- implicit dynamic FEM
- explicit dynamic FEM

$$M\ddot{u} + C\dot{u} + Ku = F$$

Explicit Dynamic FEM

$$u^{t+\Delta t} = \hat{M}^{-1} \hat{f}$$

$$\ddot{u} = \frac{u^{t+\Delta t} - 2u^t + u^{t-\Delta t}}{\Delta t^2}$$

$$\dot{u} = \frac{u^{t+\Delta t} - u^{t-\Delta t}}{2\Delta t}$$

\hat{M} : effective mass(diagonal)

\hat{f} : effective load

$$\Delta t_{crit} = \frac{L_{char}}{c}$$

$$T_{tot}^{CPU} \propto \frac{t_{process}}{\Delta t_{critic}}$$

init: $u^{-\Delta t} = u^0 - \dot{u}^0 \Delta t + \ddot{u}^0 \frac{\Delta t^2}{2}$

accelerate computation

- time scaling**: increase velocity
- mass scaling**: increase density

Bar	$c = \sqrt{\frac{E}{\rho}}$	$l_{char} = L$
Shell	$c = \sqrt{\frac{E}{(1-\nu^2)\rho}}$	quadraliteral $l_{char} = \frac{A}{\min(l_{max}, d_{max})}$ trianguera $l_{char} = \frac{A}{h_{min}}$
3D Solid	$c = \sqrt{\frac{E(1-\nu)}{(1+\nu)(1-2\nu)\rho}}$	Hexadral $l_{char} = \frac{V}{A_{max}}$ Tetrahedral $l_{char} = h_{min}$

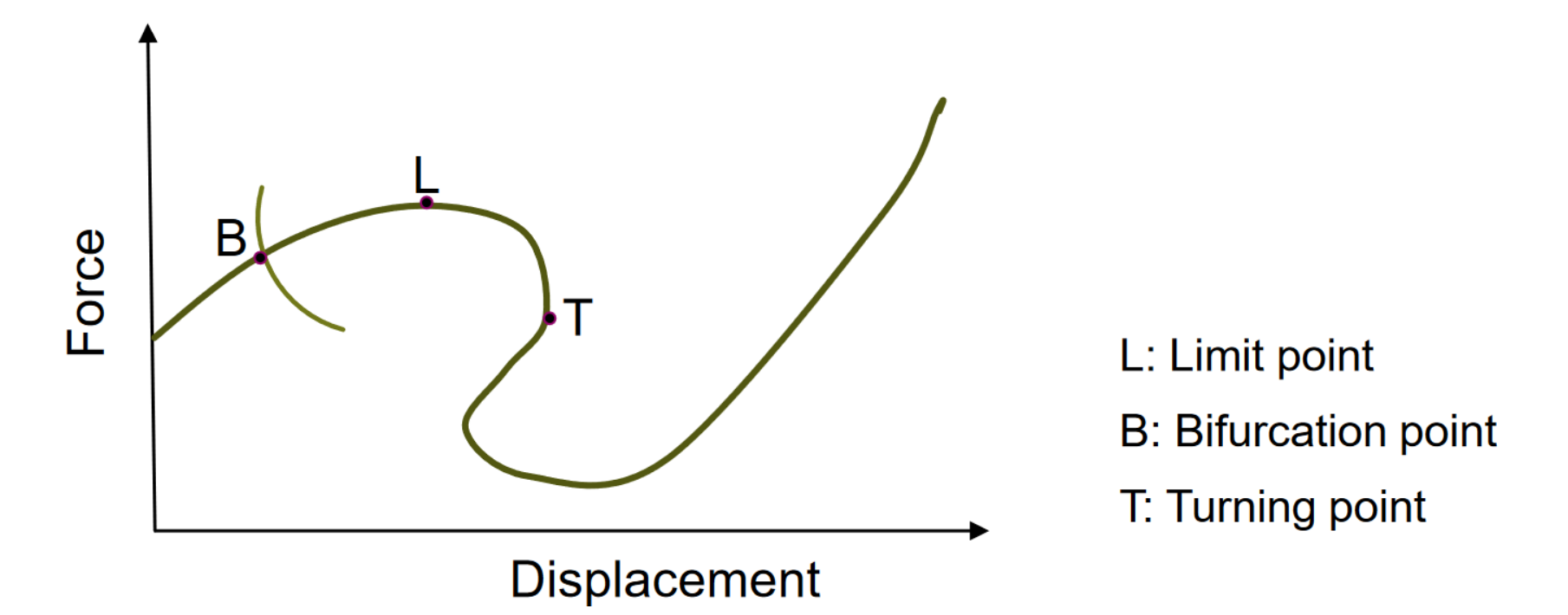
explicit method not suitable for

- slow processed
- spring back analysis
- incompressible material
- small plastic zone but large deformation

	QS Implicit	QS Explicit
Equilibrium	✓	×
Time step	arbitrary	$\Delta t < \Delta t_{char}$
Convergence	stable	conditionally stable
Boundary	RB suppressed	free
Contact	iterated	not iterated
Rezoning	$t_{fine} = t_{course}$	$\Delta t_{critic}^{new} = \frac{1}{2} \Delta t_{critic}^{old}$
Stiffness matrix	evaluated	not computed
Parallelization	complex	easy

INSTABILITY & BOUNDARY & COU

Instability



Arc-Length method

$$\mathbf{r}(\mathbf{u}, \lambda) := \hat{f}^{int}(\mathbf{u}(t)) - \lambda(t) \mathbf{f}_0^{ext} = 0$$

$$g(\mathbf{u}(t), \lambda(t)) := \sqrt{\|\dot{\mathbf{u}}\|^2 + |\dot{\lambda}|^2} - \dot{s} = 0$$

\dot{s} is the chosen value

$n + 1$ unknowns

Contact

Lagrange Multiplier Method

$$\begin{bmatrix} \mathbf{K} & \mathbf{e}_i \\ \mathbf{e}_i^\top & 0 \end{bmatrix} \begin{bmatrix} \mathbf{U} \\ \lambda \end{bmatrix} = \begin{bmatrix} \mathbf{R} \\ \mathbf{U}_i^* \end{bmatrix}$$

loss band structure of K

λ is additional variable

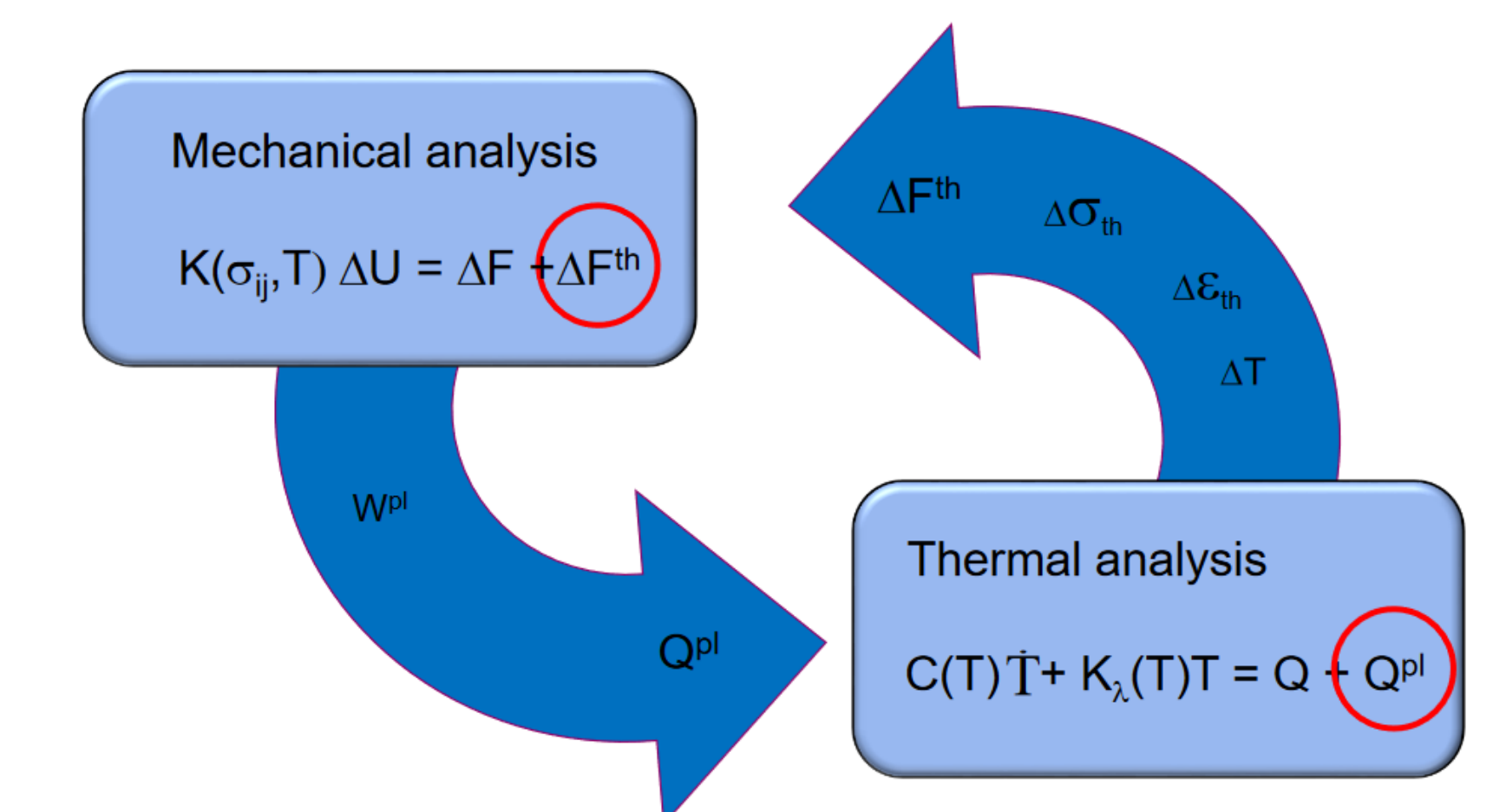
Penalty method

$$(\mathbf{K} + \alpha \mathbf{Z}^\top \mathbf{Z}) \mathbf{U} = \mathbf{R} + \alpha \mathbf{Z}^\top \mathbf{V}$$

boundary condition: $\mathbf{Z} \mathbf{U} = \mathbf{V}$

\mathbf{U} is nodal displacement

Thermo coupled



Heating source

- plastic work
- friction
- heat transfer