

Physically-based Simulation in Computer Graphics

Summary of slides

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FINITE ELEMENTS

1.1 Poisson's equation

$$(1) \quad -\Delta u(x, y) = f(x, y)$$

1.2 Finite Difference

Example 1.1 (1D problem).

$$-u''(x) = f(x), \quad x \in \Omega = (0, 1) \quad (2)$$

$$u(0) = u(1) = 0 \quad (3)$$

$$u[i] = u(i \cdot h), \quad i \in (0, \dots, n) \quad (4)$$

$$\begin{aligned} u''(x) &= \frac{u'[i] - u'[i-1]}{h} \\ &= \frac{u[i+1] - 2u[i] + u[i-1]}{h^2} \end{aligned} \quad (5)$$

$$h^2 f[i] = u[i-1] + 2u[i] + u[i+1] \quad (6)$$

1.3 Finite Elements

Example 1.2 (1D problem).

$$-u''(x) = f(x), \quad x \in \Omega = (0, 1) \quad (7a)$$

$$u(0) = u(1) = 0 \quad (7b)$$

Assume PDF is satisfied pointwise, then

$$-\int_{\Omega} dx u''(x) = \int_{\Omega} dx f(x) \quad (7c)$$

and also

$$-\int_{\Omega} dx u''(x) \cdot v(x) = \int_{\Omega} dx f(x) \cdot v(x) \quad (7d)$$

for arbitrary functions $v : \Omega \rightarrow \mathbb{R}$.

Example 1.3. • Assume v sufficiently smooth and

$$v(0) = v(1) = 0. \quad (8)$$

• Integration by parts

$$\int_a^b dx f'(x)g(x) = [f(x) \cdot g(x)]_a^b - \int_a^b dx f(x)g'(x). \quad (9)$$

$$-\int_{\Omega} dx u''(x) \cdot v(x) = \int_{\Omega} dx f(x) \cdot v(x) \quad (10a)$$

$$\int_{\Omega} dx u'(x) \cdot v'(x) = \int_{\Omega} dx f(x) \cdot v(x) + [u'(x) \cdot v(x)]_0^1 \quad (10b)$$

with the imposed boundary conditions of $v(x)$ this leads to

$$\int_{\Omega} dx u'(x) \cdot v'(x) = \int_{\Omega} dx f(x) \cdot v(x) \quad (10c)$$

Weak form because of the weaker continuity requirements.

1.3.1 Ritz-Galerkin Approach

So far, $u(x)$ and $f(x)$ continuous function, but now choose finite-dimensional subspace for $u(x)$ and $f(x)$. Solve problem in *weak* form in subspace (projection). \rightarrow Discretize u

- Sample with nodal positions x_i
- Nodal coefficients u_i
- Basis functions $N_i(x)$

$$\leadsto u(x) = \sum_i u_i N_i(x) \quad (11)$$

Discretize both u and v on n -dimensional space

$$\begin{aligned} u(x) &= \sum_{i=1}^N u_i N_i(x) \\ \leadsto \frac{\partial}{\partial x} u(x) &= \sum_{i=1}^n u_i \frac{\partial}{\partial x} N_i(x) \end{aligned}$$

insert into *weak* formulation

$$\begin{aligned} \int_{\Omega} dx u'(x) v'(x) &= \int_{\Omega} dx f(x) v(x) \\ \leadsto \int_{\Omega} dx \sum_{i=1}^n u_i \frac{\partial N_i}{\partial x} \cdot \sum_{j=1}^n v_j \frac{\partial N_j}{\partial x} &= \int_{\Omega} dx f \cdot \sum_{j=1}^n v_j N_j(x) \\ \leadsto \sum_{i,j=1}^n u_i v_j \int_{\Omega} dx \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} &= \sum_{j=1}^n v_j \int_{\Omega} dx f(x) N_j(x) \\ \leadsto \sum_{j=1}^n v_j \left[\sum_{i=1}^n u_i \int_{\Omega} dx \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} - \int_{\Omega} dx f(x) N_j(x) \right] &= 0, \end{aligned}$$

with an arbitrary v_j (assumed)

$$\sum_{i=1}^n u_i \int_{\Omega} dx \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} - \int_{\Omega} dx f(x) N_j(x) = 0, \quad \forall i, \dots, n.$$

Now we have n linear equations for n unknowns, which can be written as

$$K \mathbf{u} = \mathbf{f} \quad (12a)$$

$$\underbrace{\begin{pmatrix} K_{11} & \cdots & K_{1n} \\ \vdots & \ddots & \vdots \\ K_{n1} & \cdots & K_{nn} \end{pmatrix}}_K \underbrace{\begin{pmatrix} u_1 \\ \vdots \\ u_n \end{pmatrix}}_{\mathbf{u}} = \underbrace{\begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix}}_{\mathbf{f}} \quad (12b)$$

$$K_{ij} = \int_{\Omega} dx \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} \quad (13a)$$

$$f_i = \int_{\Omega} dx f(x) N_i(x) \quad (13b)$$

This matrix is

- symmetric (definition of K_{ij})
- positive-definite (elliptic PDE)
- sparse (N_i have compact support)

\leadsto use sparse solver, e.g., conjugate gradients

Choice of function space (types of Finite Elements)

$$u(x) = \sum_i u_i N_i(x) \quad (14)$$

- smooth enough \leadsto once differentiable
- simple \leadsto polynomial functions
- interpolation $\leadsto N_i(\mathbf{x}_j) = \delta_{ij}$
- compact support \leadsto defined piecewise on simple geometry

use piecewise linear basis functions \leadsto piece wise linear approximation $u(x)$.

Linear simplicial elements Basis functions are uniquely defined through

- Geometry \mathbf{x}_j and
- interpolations requirement $N_i(\mathbf{x}_j) = \delta_{ij}$

$$\begin{aligned} \mathbf{x}_j &= x_j, \quad \text{in 1D,} \\ &= (x_j, y_j)^t, \quad \text{in 2D,} \\ &= (x_j, y_j, z_j)^t, \quad \text{in 3D.} \end{aligned} \quad (15a)$$

- Basis functions are linear

$$N_i(x, y) = a_i x + b_i y + c \quad (16)$$

- Due to $N_i(\mathbf{x}_j) = \delta_{ij}$ we have

$$\begin{pmatrix} x_1 & y_1 & 1 \\ x_2 & y_2 & 1 \\ x_3 & y_3 & 1 \end{pmatrix} \begin{pmatrix} a_i \\ b_i \\ c_i \end{pmatrix} = \begin{pmatrix} \delta_{1i} \\ \delta_{2i} \\ \delta_{3i} \end{pmatrix} \quad (17a)$$

$$\begin{pmatrix} a_i \\ b_i \\ c_i \end{pmatrix} = \begin{pmatrix} x_1 & y_1 & 1 \\ x_2 & y_2 & 1 \\ x_3 & y_3 & 1 \end{pmatrix}^{-1} \begin{pmatrix} \delta_{1i} \\ \delta_{2i} \\ \delta_{3i} \end{pmatrix} \quad (17b)$$

What is a finite element A finite element is a triple consisting of

- a closed subset $\Omega_e \subset \mathbb{R}^d$
- a set of n basis function $N_i : \Omega_E \rightarrow \mathbb{R}$
- a set of n associated nodal variables \mathbf{x}_i

Summary of problem

- 1D Poisson problem given as

$$u''(x) = f(x) \quad (18)$$

- Weak form + galerkin approach gives linear system

$$Ku = f \quad (19a)$$

with

$$K_{ij} = \int_{\Omega} dx \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x}, \quad (19b)$$

and

$$\mathbf{f}_i = \int_{\Omega} dx f(x) N_i(x). \quad (19c)$$

- Use simple elements with linear basis functions

What is left is

- tessate domain into elements
- evaluate integrals (basis functions & derivatives)
- assemble the system matrix and right hand side

1.4 Exercise 2

Task: use linear triangle elements

- evaluate integrals
- assemble the global system matrix and right hand side

Goal: evaluate

$$\mathbf{f}_i = \int_{\Omega} dx dy f(x) N_i(x) \quad (20)$$

- N_i extends over all $n_{e,i}$ elements incident to \mathbf{x}_i
- N_i is zero outside $\Omega_i \subset \Omega$
- Consider restrictions N_i^e onto elements

It holds

$$\int_{\Omega} dx f(x) N_i(x) = \sum_{\ell=1}^{n_{\ell,i}} \int_{\Omega_{\ell}} dx f(x) N_i^{\ell}(x), \quad \text{with } \Omega_i = \bigcup_{\ell} \Omega_{\ell} \quad (21a)$$

evaluate integrals over Ω_{ℓ} with *quadratic rule*

$$\int_{\Omega_{\ell}} dx dy f(x) N_i^e(x) \approx f(x_q, y_q) \cdot N_i^{\ell}(x_q, y_q) \cdot A_{\ell}, \quad (21b)$$

where A_{ℓ} area of element ℓ and (x_{ℓ}, y_{ℓ}) the single quadrature point at bary-center

Goal: evaluate

$$\begin{aligned} \mathbf{K}_{ij} &= \int_{\Omega} dx dy \int_{\Omega_{\ell}} dx dy N_i^{\ell} N_j^{\ell} \\ &> 0 \end{aligned} \quad (22a)$$

i.e., \exists an element containing vertices \mathbf{x}_i and \mathbf{x}_j . Let S_{ij} denote a set of all these elements, then

$$K_{ij} = \sum_{\ell \in S_{ij}} \int_{\Omega_{\ell}} dx dy \frac{\partial N_i^{\ell}}{\partial x} \frac{\partial N_j^{\ell}}{\partial x} + \frac{\partial N_i^{\ell}}{\partial y} \frac{\partial N_j^{\ell}}{\partial y}. \quad (22b)$$

This is an element-centered implementation. For each element ℓ , compute basis function derivatives

$$\frac{\partial N_k^\ell}{\partial \mathbf{x}}, \quad k = 1, \dots, 3 \quad (23a)$$

form products and integrate

$$\frac{\partial N_k^\ell}{\partial \mathbf{x}} \frac{\partial N_m^\ell}{\partial \mathbf{x}}, \quad k = 1, \dots, 3; m = k, \dots, 3, \quad (23b)$$

which is a constant on element and add to global matrix (local vs. global numbering)

$$K_{ij} += A_\ell \left(\frac{\partial N_k^\ell}{\partial \mathbf{x}} \frac{\partial N_m^\ell}{\partial \mathbf{x}} + \frac{\partial N_k^\ell}{\partial \mathbf{y}} \frac{\partial N_m^\ell}{\partial \mathbf{y}} \right) \quad (23c)$$

Advantages of FEM (over Finite Differences)

- can easily deal with comple geometries
- solution in nodes is naturally interpolated inside elements using basis functions
- weaker smoothness requirements on solution
- weak form is natural for many applications