MA2011-W2 Dynamical Systems

clear all, disp(date)

13-Mar-2023

Zero-Order Systems

Here, by zero-order systems, we mean systems for which the output y(t) at any given time t only depends on the <u>current value</u> of the input x(t) but <u>not on its history</u>. These systems can be symply modeled by a, possibly nonlinear, *algebraic equation*, e.g.

$$y(t) = f(x(t), t)$$

The function $f(\cdot,t)$ might or might not depend explicitly on time t, in which case we denote the systems as *time-invariant*.

$$y(t) = f(x(t))$$

The simplest case is that of linear systems, i.e. those which can be simply described by a linear equation

$$y(t) = a(t) \cdot x(t) + b(t)$$

The easiest case is when the system is *Linear and Time-Invariant (LTI)*:

$$y(t) = a \cdot x(t) + b$$

i.e. when a and b are constant.

First-Order Systems (LTI) systems

Very often, the output of a system will depend also on the history of the input. Many systems of interest can be described via Ordinary Differential Equations (ODE). Besides its current input u(t), the state of the system x(t) will also depend on its time derivatives $\frac{d}{dt}x(t), \frac{d^2}{dt^2}x(t), \cdots, \frac{d^n}{dt^n}x(t)$, up to the n-th order.

<u>First-oder systems only involve derivatives up the first order,</u> i.e. it will only contain $\frac{d}{dt}x(t)$ (denoted as $\dot{x}(t)$, for short).

A linear, time-invariant first-order system can in general be described as

$$\dot{x}(t) + \frac{x(t)}{\tau} = u(t)$$

$$x(0) = x_0$$

where

- τ is the time constant;
- *u*(*t*) is the forcing input;

• x_0 is the initial condition.

In MATLAB, this can be easily written as

```
syms x(t) u(t) syms tau x0 real assume (tau > 0)  %% NOTE this assumption myEQ = diff(x) == -x/tau + u  %% fist order ODE myEQ(t) = \frac{\partial}{\partial t} x(t) = u(t) - \frac{x(t)}{\tau} myIC = x(0) == x(0)
```

dsolve(myEQ, myIC)

ans = $e^{-\frac{t}{\tau}} \left(x_0 + \int_0^t e^{y/\tau} u(y) dy \right)$

A few words on Dynamical Analogies - the heart of mechatronics

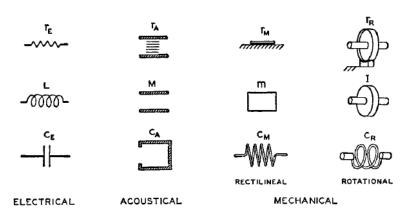


Fig. 2.1. Graphical representation of the three basic elements in electrical, mechanical rectilineal, mechanical rotational and acoustical systems.

$$r_E = {
m electrical \ re-sistance}$$
 $r_A = {
m acoustical \ re-sistance}$ $r_A = {
m acoustical \ re-sistance}$ $r_B = {
m mechanical \ re-tilineal}$ $r_B = {
m mechanical \ re-tili$

figure from [Olson 1958]

• Olson, H. F. (1958). Dynamical analogies (Vol. 2729). Princeton: Van Nostrand.

	electrical	mechanical (lin.)	mechanical (rot.)
flow	current [A]: i	lin. vel. [m/s]: \dot{x}	ang. vel [rad/s]: $\dot{ heta}$
effort	voltage [V]: Δv	force [N]: f	torque [Nm]: $ au$
power [W]	$\Delta v \cdot i$	$f\cdot \dot{x}$	$ au\cdot\dot{ heta}$
energy [J]	$\int \Delta v \cdot i dt$	$\int f \cdot \dot{x} dt$	$\int au \cdot \dot{ heta} dt$
Elements			
resistance	$\Delta v = R i$	$f=b\dot{x}$	$ au=eta\dot{ heta}$
inertia	$\Delta v = L rac{di}{dt}$	$f=m\ddot{x}$	$ au = I\ddot{ heta}$
compliance	$\Delta v = rac{1}{C}\int idt \equiv rac{q}{C}$	$f=\int k\dot{x}dt\equiv \ kx$	$ au = \int \kappa \dot{ heta} dt \equiv \ \kappa heta$
En. storage			
kinetic energy	$E = \frac{1}{2}Li^2$	$E=rac{1}{2}m\dot{x}^2$	$E=rac{1}{2}I\dot{ heta}^2$
potential energy	$E=rac{1}{2}C\Delta v^2$	$E=rac{1}{2}k\Delta x^2$	$E=rac{1}{2}\kappa\Delta heta^2$

DC Analysis (response to constant inputs)

Very often, a system is subjected to a constant input

$$\dot{x}(t) + \frac{x(t)}{\tau} = u_0$$
$$x(0) = x_0$$

$$\begin{aligned} & \text{myEQO(t)} &= \\ & \frac{\partial}{\partial t} \; x(t) = u_0 - \frac{x(t)}{\tau} \end{aligned}$$

DC_Sol =

$$\tau u_0 + e^{-\frac{t}{\tau}} (x_0 - \tau u_0)$$

Let's see what happens when $t \to \infty$, let's define

$$x_{\infty} := \lim_{t \to \infty} x(t)$$

```
limit(DC_Sol, t, inf)

ans = \tau u_0
```

From this, we shall rewrite the first order equation as

$$\dot{x}(t) = \frac{x_{\infty} - x(t)}{\tau}$$
$$x(0) = x_0$$

```
\begin{aligned} &\text{syms x\_inf real} \\ &\text{myEQ0} = &\text{subs(myEQ, u(t), x\_inf/tau)} \\ &\text{myEQ0(t)} = \\ &\frac{\partial}{\partial t} \; x(t) = \frac{x_{\text{inf}}}{\tau} - \frac{x(t)}{\tau} \\ &\text{dsolve(myEQ0, myIC)} \\ &\text{ans} = \\ &x_{\text{inf}} + \mathrm{e}^{-\frac{t}{\tau}} \left( x_0 - x_{\text{inf}} \right) \end{aligned}
```

and we wish to know how it will respond. In this case, rather than a generic forcing input, we should consider a constant input u(t) and a system initially in its zero state, i.e. x(0) = 0

```
mylegend = {};
for tau_list = [0.1 0.2 1 2 5 10]
    SOL = dsolve (subs(myEQ0, {x_inf, tau}, {5, tau_list}), 'x(0)==0')
    fplot(SOL, [0 10]); grid on; hold on
    mylegend= {mylegend{:}, ['\tau = ' num2str(tau_list) ' sec']};
    xlabel('time [sec]'); ylabel('x(t) [V]')
    title('x_\infty = 5V')
end
```

```
SOL = 5 - 5e^{-10t}

SOL = 5 - 5e^{-5t}

SOL = 5 - 5e^{-t}

SOL = 5 - 5e^{-t}

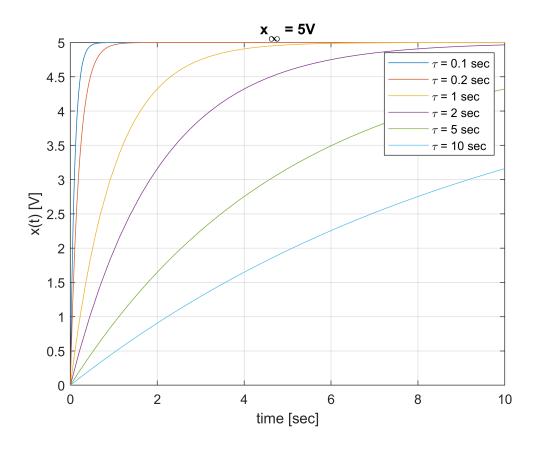
SOL = 5 - 5e^{-t}

SOL = 5 - 5e^{-t}
```

$$5 - 5e^{-\frac{t}{5}}$$

SOL = $5 - 5e^{-\frac{t}{10}}$

legend(mylegend);



To manually sketch exponential responses given x_0, x_∞, τ it is useful to evaluate the time derivative of the explonential at time t = 0, i.e.

$$\dot{x}(0) = \frac{x_{\infty} - x_0}{\tau}$$

$$\left(\left(\frac{\partial}{\partial t} x(t) \right) \Big|_{t=0} \right) = \frac{x_{\text{inf}}}{\tau} - \frac{x(0)}{\tau}$$

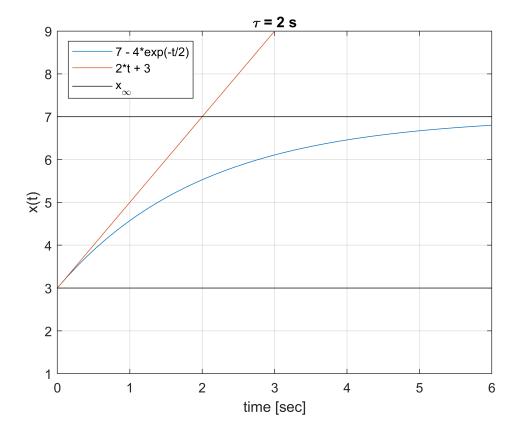
as this provides the slope of the tangent line

```
MyExp = subs( dsolve(myEQ0, 'x(0)==x0'), {x0, x_inf, tau}, {3, 7, tau_val})

MyExp = 
7 - 4e^{-\frac{t}{2}}
MyTaylor = taylor (MyExp, t, 'Order', 2)
```

```
MyTaylor = 2t + 3
```

```
figure
fplot(MyExp, range); grid on, hold on
fplot(MyTaylor, range)
fplot(x_inf_val, 'k', range)
fplot(x_0_val, 'k', range)
ylim([x_0_val-2 x_inf_val+2])
title(['\tau = ' num2str(tau_val) ' s'])
legend(char(MyExp), char(MyTaylor), 'x_\infty', 'Location', 'NorthWest')
xlabel('time [sec]'); ylabel('x(t)')
```



AC Steady-State analysis (steady-state response to sinusoidal inputs)

AC-analysis is based on the fact that given a <u>purely sinusoidal input</u>, e.g. $\cos(\omega t)$, after a long time, we expect to see a purely sinusoidal output, e.g. $A_{\omega}\cos(\omega t + \phi_{\omega})$

 ullet at the same frequency ω

• but with likely different (<u>frequency-dependent</u>) Amplitude A_{ω} and phase ϕ_{ω}

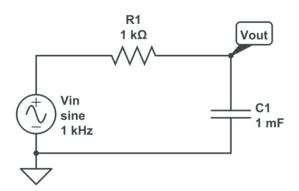
$$\cos(\omega t) \Rightarrow \|system\| \Rightarrow A_{\omega}\cos(\omega t + \phi_{\omega})$$

The **objective of AC analysis** is to determine Amplitude A_{ω} and phase ϕ_{ω} for every input frequency ω

This can be accomplished in basically two ways

- via trigonometric manipulation (hard way)
- via complex exponentials (easy way)

Let's look at a simple case, a 1st order RC circuit



The hard way - via trigonometry

 $V_{in} = \cos(\omega t)$

 $V_{-}out = A\cos(\phi + \omega t)$

Define device equations

$$i = -\frac{A \omega \sin(\phi + \omega t)}{1000}$$

$$V_R = -A \omega \sin(\phi + \omega t)$$

Solve equations $\forall t$, in fact, only for a couple of values, then we can verify that solutions are valid at all times.

 $KVL = \cos(\omega t) = A\cos(\phi + \omega t) - A\omega\sin(\phi + \omega t)$

expand (KVL)

ans = $\cos(\omega t) = A\cos(\omega t)\cos(\phi) - A\sin(\omega t)\sin(\phi) - A\omega\cos(\omega t)\sin(\phi) - A\omega\sin(\omega t)\cos(\phi)$

myEquations = [subs(KVL, t, 0); expand(subs(KVL, t, 1/omega*pi/2))]

myEquations =

$$\begin{pmatrix} 1 = A\cos(\phi) - A\omega\sin(\phi) \\ 0 = -A\sin(\phi) - A\omega\cos(\phi) \end{pmatrix}$$

myVariables = [A; phi]

myVariables =

 $\begin{pmatrix} A \\ \phi \end{pmatrix}$

SOL = solve(myEquations, myVariables)

SOL = struct with fields:
 A: [1×1 sym]
 phi: [1×1 sym]

SOL.A, SOL.phi

ans =

$$\frac{1}{\sqrt{\omega^2+1}}$$

ans =

$$-2 \operatorname{atan} \left(\frac{\sqrt{\omega^2 + 1} - 1}{\omega} \right)$$

NOTE: that $-2\tan^{-1}\left(\frac{\sqrt{\omega^2+1}-1}{\omega}\right) \equiv -\tan^{-1}\omega$

However, this is a very cumbersome (too much trigonometry) way to proceed!!!

The easy way via the 'complex trick'

again, start from the time-domain equations, for simplicity let's set $\tau = 1$

myEQ(t) =

$$\frac{\partial}{\partial t} x(t) = u(t) - x(t)$$

Now, let's consider complex exponentials as inputs

$$\left. e^{j\omega t} \Rightarrow \left| \stackrel{LTI}{SYS} \right| \Rightarrow A_{\omega} e^{j(\omega t + \phi_{\omega})} = \underbrace{A_{\omega} e^{j\phi}}_{H(j\omega)} e^{j\omega t}$$

```
syms H_0 complex

syms omega real

j=sqrt(-1);

myEQ\_AC = subs(myEQ, \{u(t), x(t)\}, \{exp(j*omega *t), H_0 * exp(j*omega *t)\})

myEQ\_AC(t) = H_0 \omega e^{\omega t i} i = e^{\omega t i} - H_0 e^{\omega t i}

eqn\_cmplx = simplify(myEQ\_AC)

eqn\_cmplx(t) = H_0 (1 + \omega i) = 1

% simplify (real(AC_Sol_complex) - AC_Sol)

Sol_myEQ_AC = solve(eqn_cmplx, H_0)

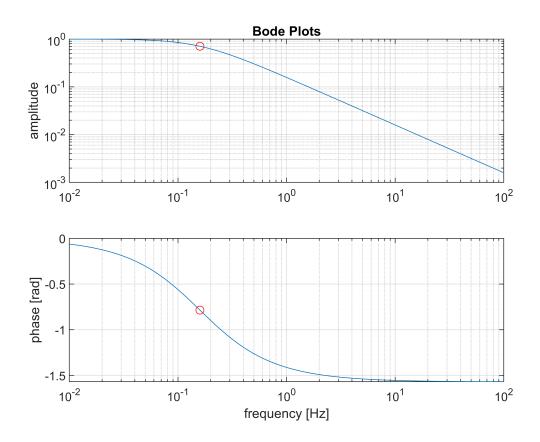
Sol_myEQ_AC = \frac{1}{1 + \omega i}
```

Bode Plots

```
freq = logspace(-2, 2, 100);
H = subs(Sol_myEQ_AC,{ tau, omega}, {1, 2*pi*freq} );
figure
subplot(2,1,1)
loglog(freq, abs(H)); grid on; ylabel('amplitude'); title ('Bode Plots')
hold on
subplot(2,1,2)
semilogx(freq, angle(H)); grid on; xlabel('frequency [Hz]'); ylabel('phase [rad]')
hold on
```

Cut-off frequency: the frequency at which the output is reduced by -3dB (or $1/\sqrt{2}$), clearly for a first order system

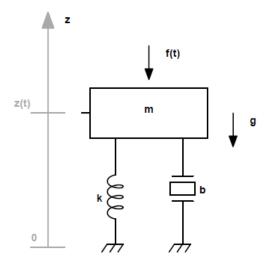
```
subplot(2,1,1)
loglog(omega_0/2/pi, abs(H_0), 'ro')
subplot(2,1,2)
semilogx(omega_0/2/pi, angle(H_0), 'ro')
```



2nd order dynamics

Newton's law

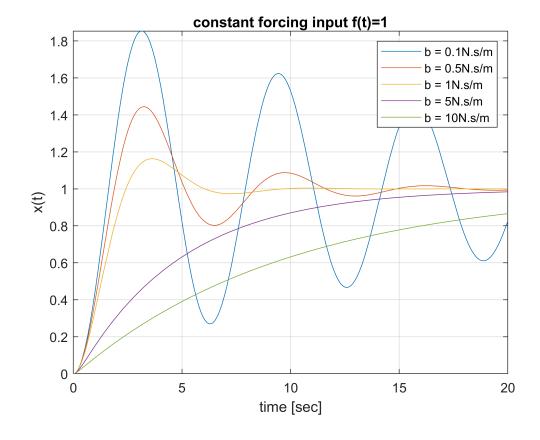
Consider a mass-springer-damper case:



```
syms m b k real syms f(t) x(t) eqn_Newton = f(t) == m*diff(x(t), 2)+b*diff(x(t)) + k*x(t) eqn_Newton = f(t) = m \frac{\partial^2}{\partial t^2} x(t) + b \frac{\partial}{\partial t} x(t) + k x(t)
```

Time Response to constant forcing Input

```
figure;
mylegend = {};
for b_var = [0.1, 0.5, 1, 5, 10]
    mylegend = {mylegend{:}, ['b = ' num2str(b_var) 'N.s/m']};
    eqn_Newton_DC = subs(eqn_Newton, {f(t), m, b, k}, {1, 1, b_var, 1});
    Dx = diff(x);
    cond = [x(0)==0, Dx(0) == 0];
    Sol_Newton_time = dsolve(eqn_Newton_DC, cond);
    fplot(Sol_Newton_time, [0 20]); grid on; hold on
end
xlabel('time [sec]'); ylabel('x(t)'); title ('constant forcing input f(t)=1')
legend(mylegend)
```



AC Analysis

```
syms omega real syms F_0 X_0 complex
```

```
AC_Exp_Newton = subs(eqn_Newton, {f(t), x(t)}, {F_0*exp(j*omega*t), X_0*exp(j*omega*t)})  \text{AC}_{\text{Exp}}_{\text{Newton}} = F_0 \, \mathrm{e}^{\omega \, t \, \mathrm{i}} = X_0 \, k \, \mathrm{e}^{\omega \, t \, \mathrm{i}} + X_0 \, b \, \omega \, \mathrm{e}^{\omega \, t \, \mathrm{i}} \, \mathrm{i} - X_0 \, m \, \omega^2 \, \mathrm{e}^{\omega \, t \, \mathrm{i}}
```

```
ans = X_0 m \omega^2 + F_0 = X_0 (k + b \omega i)
```

simplify(AC_Exp_Newton)

```
Sol_ODE2= solve(simplify (AC_Exp_Newton) , X_0)
```

```
Sol_ODE2 = \frac{F_0}{-m \omega^2 + i b \omega + k}
```

```
freq = logspace(-2, 1, 100);
H = subs(Sol_ODE2,{F_0, m, b, k, omega}, {1, 1, .1, 1, 2*pi*freq} );
figure
subplot(2,1,1)
loglog(freq, abs(H)); grid on; ylabel('amplitude'); title ('Bode Plots')
subplot(2,1,2)
semilogx(freq, angle(H)); grid on; xlabel('frequency [Hz]'); ylabel('phase [rad]')
```

