

# 1 April 6th, 2020

## 1.1 Sturm-Liouville

**Definition 1.1.** An ODE of the form:

$$\frac{d}{dx} (s(x)\phi'(x)) + q(x)\phi(x) + \lambda w(x)\phi(x) = 0, \alpha < x < \beta.$$

along with the boundary conditions:

$$c_1\phi(\alpha) + c_2\phi'(\alpha) = 0.$$

$$d_1\phi(\beta) + d_2\phi'(\beta) = 0.$$

is called a **regular Sturm Liouville Problem** (denoted RSLP) if the following conditions hold.

1.  $s(x), s'(x), q(x), w(x)$  are all continuous functions in the open interval  $\alpha < x < \beta$
2.  $s(x) > 0$  and  $w(x) > 0$  for all  $\alpha < x < \beta$
3.  $c_1^2 + c_2^2 > 0$  and  $d_1^2 + d_2^2 > 0$ , i.e. can't have both  $c_1$  and  $c_2$  equal zero, same for  $d_1, d_2$
4. The  $\lambda$  occurs only in the ODE as indicated by  $\lambda w(x)\phi(x)$ .

Recall that any linear 2nd order homogeneous ODE of the form:

$$a_2(x)\phi''(x) + a_1(x)\phi'(x) + a_0(x)\phi(x) + \lambda b(x)\phi(x) = 0.$$

can be placed in the form:

$$\frac{d}{dx} (s(x)\phi'(x)) + q(x)\phi(x) + \lambda w(x)\phi(x) = 0, \alpha < x < \beta.$$

by setting

$$s(x) = e^{\int \frac{a_1(x)}{a_2(x)} dx}.$$

$$q(x) = \frac{a_0(x)s(x)}{a_2(x)}.$$

$$w(x) = \frac{b(x)s(x)}{a_2(x)}.$$

**Example 1.2**

Consider:

$$x\phi''(x) + 2x\phi(x) + \phi(x) + \lambda x^2\phi(x) = 0, \quad 0 < x < 1.$$

we have:

$$s(x) = e^{\int \frac{2x}{x} dx} = e^{2x}.$$

$$q(x) = \frac{1}{x}e^{2x}.$$

$$w(x) = \frac{x^2}{x}e^{2x} = xe^{2x}.$$

Meaning that the equation can be written in the form of:

$$\frac{d}{dx} (e^{2x}\phi'(x)) + \frac{1}{x}e^{2x}\phi(x) + \lambda xe^{2x}\phi(x) = 0, \quad 0 < x < 1.$$

**Remark 1.3** — The form above is called the **self-adjoint form**.

## 1.2 Properties of Regular Sturm-Liouville Problems

- There exist an infinite number of  $\lambda$ 's that lead to non-zero solutions  $\phi(x)$  to the ODE and boundary conditions. These  $\lambda$ 's which can be ordered from smallest to largest are called the **eigenvalues** of the RSLP. Moreover

$$\lim_{n \rightarrow \infty} \lambda_n = +\infty.$$

This means if you solve a RLSP and found that  $\lambda_n = \frac{n}{n+1}$  then there is a problem, since  $\lim_{n \rightarrow \infty} \frac{n}{n+1} = 1 \neq +\infty$ .

- If these  $\lambda$ 's are ordered and if the non-zero functions  $\phi_n(x)$  satisfy:

$$\frac{d}{dx} (s(x)\phi'_n(x)) + q(x)\phi_n(x) + \lambda_n\omega(x)\phi_n(x) = 0.$$

and it satisfies the boundary conditions, then  $\phi_n(x)$  is called the eigenfunction associated with the eigenvalue  $\lambda_n$  and  $\phi_n(x)$  goes through zero exactly  $n - 1$  times in the open interval  $\alpha < x < \beta$ . As such  $\phi_1(x)$  does not go through zero,  $\phi_2(x)$  goes through zero once, and so on.

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$$\phi_m \cdot \phi_n = \begin{cases} 0, & n \neq m \\ > 0, & n = m \end{cases}.$$

where

$$f \cdot g = \int_{\alpha}^{\beta} f(x)g(x)w(x) dx.$$

In other words there is a dot product that can be defined on the functions  $\phi_n$ .

- The set of eigenfunctions

$$\{\phi_1(x), \phi_2(x), \dots\}.$$

is called a complete set of basic functions so that if  $f(x)$  is any piecewise continuous function in the interval  $\alpha < x < \beta$ , then we may expand  $f(x)$  as:

$$f(x) = \sum_{n=1}^{\infty} a_n \phi_n(x).$$

with

$$a_n = \frac{\phi_n \cdot f}{\phi_n \cdot \phi_n}.$$

moreover, the sum converges to:

$$\frac{f(x^+) + f(x^-)}{2}.$$

Where:

$$f(a^+) = \lim_{x \rightarrow a^+} f(x) \quad f(a^-) = \lim_{x \rightarrow a^-} f(x).$$

Note that piecewise continuous means that there can only be a finite number of hole or jump discontinuities but not essential discontinuities.

**Remark 1.4** — This means that the regular Sturm Liouville series will fill in all of the hole discontinuities and where there is a jump discontinuity, it will converge to the midpoint of the jump.

**Definition 1.5.** A function  $f(x)$  is called piecewise continuous in a finite interval  $\alpha < x < \beta$  if it has at most a finite number of hole or jumps in  $\alpha < x < \beta$ .

**Example 1.6**

Consider:

$$\phi''(x) + \lambda\phi(x) = 0, 0 < x < 1.$$

$$\phi(0) = 1 \quad \phi(1) = 0.$$

Thus we have:

$$s(x) = e^{\int \frac{0}{1} dx} = 1.$$

$$q(x) = \frac{(0)(1)}{1} = 0.$$

$$w(x) = \frac{(1)(1)}{1} = 1.$$

Note that this is a RSLP. To solve this problem, notice that this is a problem with constant coefficients, giving us:

$$\phi(x) = \begin{cases} A \cosh(x\sqrt{-\lambda}) + B \sinh(x\sqrt{-\lambda}), & \lambda < 0 \\ A + Bx, & \lambda = 0 \\ A \cos(x\sqrt{\lambda}) + B \sin(x\sqrt{\lambda}), & \lambda > 0 \end{cases}.$$

If we consider  $\lambda < 0$ , because of the boundary conditions, we get  $A = 0$ , and  $B = 0$ , meaning that  $\lambda < 0$  gives us  $\phi(x) = 0$ , meaning that there are no negative eigenvalues.

For  $\lambda = 0$ , we have:

$$\phi(x) = A = 0 \implies \phi(x) = Bx.$$

$$\phi(1) = B = 0 \implies \phi(x) = 0.$$

meaning that zero is not an eigenvalue.

For  $\lambda > 0$ , we have:

$$\phi(0) = A = 0 \implies \phi(x) = B \sin(x\sqrt{\lambda}).$$

$$\phi(1) = B \sin(x\sqrt{\lambda}) \implies \sqrt{\lambda} = n\pi \implies \lambda = (n\pi)^2 = \lambda_n.$$

Note that we will always get a multiplicative constant when we try to calculate  $\phi_n$ , thus we can set  $B_n = 1$  giving us:

$$\phi_n(x) = \sin(n\pi x).$$

**Remark 1.7** — Note that for the above example,  $\phi_1(x) = \sin(\pi x)$  does not go through zero on the open interval between 0 and 1. Similarly  $\phi_2(x) = \sin(2\pi x)$  goes through zero once at  $x = \frac{1}{2}$ , etc.