

# Supplementary Material for: “Realization of metallic state in 1T-TaS<sub>2</sub> with persisting long-range coherence of charge density wave”

Xin-Yang Zhu,<sup>1</sup> Shi Wang,<sup>1</sup> Zhen-Yu Jia,<sup>1</sup> Li Zhu,<sup>1</sup> Qi-Yuan Li,<sup>1</sup> Wei-Min Zhao,<sup>1</sup> Cheng-Long Xue,<sup>1</sup>  
Yong-Jie Xu,<sup>1</sup> Zhen Ma,<sup>1</sup> Jinsheng Wen,<sup>1,2</sup> Shun-Li Yu,<sup>1,2</sup> Jian-Xin Li,<sup>1,2</sup> and Shao-Chun Li<sup>1,2</sup>

<sup>1</sup>*National Laboratory of Solid State Microstructures and Department of Physics, Nanjing University, Nanjing 210093, China*

<sup>2</sup>*Collaborative Innovation Center of Advanced Microstructures, Nanjing University, Nanjing 210093, China*

(Dated: October 29, 2019)

Here, we show the details of the cluster perturbation theory for the theoretical calculations of the single-particle spectral weight in the main text.

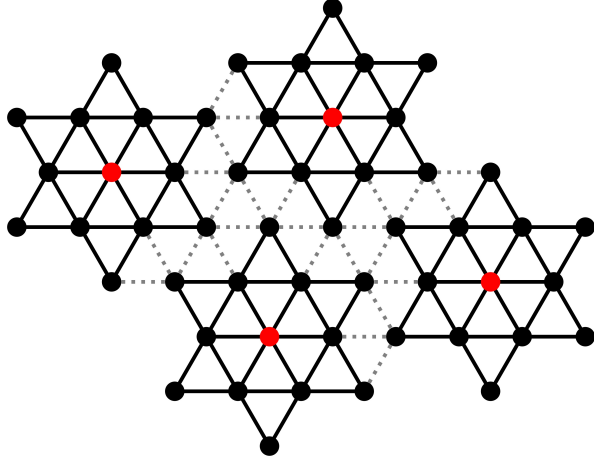


FIG. 1. Tiling of the triangular lattice with 13-site clusters. The nearest-neighbor hoppings within a cluster are represented by solid lines, and intercluster hoppings by dashed lines. Here, each site represents a David star in 1T-TaS<sub>2</sub>, and the red sites are the David stars where the K<sup>+</sup> cations are located in the theoretical calculations.

In the main text, the computational technique for our calculation of the single-particle spectral weight  $A_i(\omega)$  on site  $i$  is based on the cluster perturbation theory (CPT), which has been successfully applied to many strongly correlated systems [1–7]. As shown in Fig. 1, in the CPT calculations, we divide the original lattice into identical 13-site clusters which constitute a superlattice. Then, the lattice Hamiltonian is written as

$$H = H' + V, \quad (1)$$

where  $H'$  is the cluster Hamiltonian, obtained by severing the hopping terms between different clusters, and  $V$  contains the intercluster hoppings. The CPT method is based on a strong-coupling perturbation expansion of the one-body operators  $V$  linking the individual clusters[1]. At the lowest order of this expansion, the Green's function  $\mathbf{G}$  of the original lattice can be expressed (in matrix

form) as

$$\mathbf{G}(\tilde{\mathbf{k}}, \omega) = \mathbf{G}'(\omega)[1 - \mathbf{V}(\tilde{\mathbf{k}})\mathbf{G}'(\omega)]^{-1}, \quad (2)$$

where  $\mathbf{G}'$  is the cluster Green's function, and  $\tilde{\mathbf{k}}$  is the wavevector in the Brillouin zone (BZ) of the superlattice.  $\mathbf{G}'$  is independent of  $\tilde{\mathbf{k}}$ , while  $\mathbf{V}$  is frequency independent. The cluster Green's function  $\mathbf{G}'$  is calculated by the exact diagonalization method with the Lanczos algorithm, and it has the following expression,

$$G'_{\mu\nu}(\omega) = \langle 0 | c_\mu \frac{1}{\omega - H + E_0 + i\eta} c_\nu^\dagger | 0 \rangle + \langle 0 | c_\nu^\dagger \frac{1}{\omega + H - E_0 + i\eta} c_\mu | 0 \rangle, \quad (3)$$

where  $E_0$  is the energy of the ground state  $|0\rangle$ , and  $\eta$  is a small real number introduced in the calculation to shift the poles of the Green's function away from the real axis. Here,  $\mu$  and  $\nu$  denote both the site and spin degrees of freedom in a cluster. The spectral function of the local single-particle excitation is given by

$$A_{\mu\mu}(\omega) = -\frac{1}{\pi} \sum_{\tilde{\mathbf{k}}} G_{\mu\mu}(\tilde{\mathbf{k}}, \omega). \quad (4)$$

- 
- [1] D. S  n  chal, D. Perez, and M. Pioro-Ladri  re, *Phys. Rev. Lett.* **84**, 522 (2000).
  - [2] M. G. Zacher, R. Eder, E. Arrigoni, and W. Hanke, *Phys. Rev. Lett.* **85**, 2585 (2000).
  - [3] D. S  n  chal and A.-M. S. Tremblay, *Phys. Rev. Lett.* **92**, 126401 (2004).
  - [4] S.-L. Yu, X. C. Xie, and J.-X. Li, *Phys. Rev. Lett.* **107**, 010401 (2011).
  - [5] J. Kang, S.-L. Yu, T. Xiang, and J.-X. Li, *Phys. Rev. B* **84**, 064520 (2011).
  - [6] W.-H. Leong, S.-L. Yu, T. Xiang, and J.-X. Li, *Phys. Rev. B* **90**, 245102 (2014).
  - [7] S.-L. Yu, W. Wang, Z.-Y. Dong, Z.-J. Yao, and J.-X. Li, *Phys. Rev. B* **98**, 134410 (2018).